# SU(2) symmetry in underdoped cuprates





X. Montiel (IPhT,CEA-Saclay)



Th. Kloss (IPhT, CEA-Saclay) K.B. Efetov (Bochum)

H. Meier (Yale)

M. Einenkel (Bochum)

C. Pépin (IPhT, CEA-Saclay)

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V. Silva de Carvalho (Goiania, Brazil)

#### <sup>89</sup>Y NMR Evidence for a Fermi-Liquid Behavior in YBa<sub>2</sub>Cu<sub>3</sub>O<sub>6+x</sub>

H. Alloul, T. Ohno,<sup>(a)</sup> and P. Mendels Physique des Solides, Université de Paris-Sud, 91405 Orsay, France (Received 15 May 1989)

We report NMR shift  $\Delta K$  and  $T_1$  data of <sup>19</sup>Y taken from 77 to 300 K in YBa<sub>2</sub>Cu<sub>3</sub>O<sub>h+z</sub> for 0.35 < x < 1, from the insulating to the metallic state. A Korringa law and therefore a Fermi-liquid picture is found to apply for the spin part K, of  $\Delta K$ . The spin contribution  $\chi_c(x, T)$  to  $\chi_m$  is singled out, as the T variation of  $\Delta K$  scales linearly with the macroscopic susceptibility  $\chi_m$ . This implies that Cu(3d) and O(2p) holes do not have independent degrees of freedom. Their hybridization, which has a  $\sigma$  character, hardly varies with doping. These results put severe constraints on theoretical models of high- $T_c$  cuprates.

PACS numbers: 74.70.Vy, 75.20.En, 76.60.Cq, 76.60.Es



FIG. 1. The shift  $\Delta K$  of the <sup>89</sup>Y line, referenced to YCl<sub>3</sub> plotted vs T, from 77 to 300 K. The lines are guides to the eye.

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Cyr-Choignière, preprint 2015





#### BSCCO (opt. doped)

1.0



0.15

Cyr-Choignière, preprint 2015





#### **Emergent symmetries for the pseudo-gap**



Sachdev et al (2013) Efetov, Meier, CP (2013)



SU(2) symmetry related to the SU(2) symmetry of the superexchange hamiltonian and gauge SU(2) symmetry

$$U_{ij} = \begin{pmatrix} -\chi_{ij}^* & \Delta_{ij} \\ \Delta_{ij}^* & \chi_{ij} \end{pmatrix}$$

$$\chi_{ij}\delta_{\alpha\beta} = 2\langle f_{i\alpha}^{\dagger}f_{j\beta}\rangle, \quad \chi_{ij} = \chi_{ji}^{*},$$

$$\Delta_{ij}\epsilon_{\alpha\beta} = 2\langle f_{i\alpha}f_{j\beta}\rangle, \quad \Delta_{ij} = \Delta_{ji}.$$



Sachdev et al (2013) Kotliar and Liu (1988) Lee, Wen, Nagaosa, RMP (2006)



















## **Exact realization : neglecting Coulomb interactions**



$$H = \bigvee \left[ -\sum_{\langle ij \rangle,\sigma} t_{ij} c_{i\sigma}^{\dagger} c_{i\sigma} + J \sum_{\langle ij \rangle} \left( \mathbf{S}_{i} \cdot \mathbf{S}_{j} - \frac{1}{4} n_{i} n_{j} \right) \right] \bigvee$$

P: projection on no double occupancy

## **Exact realization of SU(2) symmetry** AFM QCP in d=2

K.B.Efetov, H.Meier, C.P. Nat. Phys. 9, (2013)

$$\mathcal{L} = \chi^{\dagger} \left( \partial_{\tau} + \varepsilon (-\mathrm{i}\hbar\nabla) + \lambda \vec{\phi}\vec{\sigma} \right) \chi$$

$$\langle \phi^i_{\omega,\mathbf{k}} \phi^j_{-\omega,-\mathbf{k}} \rangle \propto \frac{\delta_{ij}}{(\omega/v_s)^2 + (\mathbf{k} - \mathbf{Q})^2 + a}$$

A Abanov, A. Chubukov, Schmalian RMP 2003 Belitz, Kirkpatrick, Vojta, RMP 2005 J Rech, CP, A Chubukov, PRB 2006

Eliashberg theory : neglect vertices

Chubukov, Morr (2003 ...)

$$i\Sigma(\omega_n) = \frac{1}{k \cdot \omega_n}$$

$$\chi_0^{-1} \Pi(q, \Omega_m) = \frac{q \cdot \nu_m}{2} \sqrt{1}$$

M. Metlitsky and S. Sachdev (2010)





 $\delta \ll 1$ 

#### Composite order parameter

$$c_{\mathbf{p}}^{\mathrm{pp}}\left\langle \left(i\sigma_{2}\right)_{\alpha\beta}\psi_{\alpha,\mathbf{p}}\psi_{\beta,-\mathbf{p}}\right\rangle + c_{\mathbf{p}}^{\mathrm{ph}}\left\langle \delta_{\alpha\beta}\psi_{\alpha,\mathbf{p}}\psi_{\beta,-\mathbf{p}}^{*}\right\rangle,$$

SU(2) symmetry and fluctuations

$$u = \begin{pmatrix} \Delta_{-} & \Delta_{+} \\ -\Delta_{+}^{*} & \Delta_{-}^{*} \end{pmatrix} \quad \text{with} \quad |\Delta_{+}|^{2} + |\Delta_{-}|^{2} = 1$$



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## Gap equations around the QCP are universal

$$b(\varepsilon) = \frac{3\lambda^2}{4N\nu}T\sum_{\varepsilon'}\frac{\bar{D}(\varepsilon-\varepsilon')b(\varepsilon')}{\sqrt{f^2(\varepsilon')+b^2(\varepsilon')}},$$
$$\bar{D}(\omega) = \frac{1}{\sqrt{\gamma\Omega(\omega)+a}}$$

linear dispersion at the hot-spots





d-wave symmetry





Curvature breaks the SU(2) symmetry :  ${}^{49}$ C dome  ${}^{80}$ 

How much is it bro



120

λ

160



#### **Confirmation by other techniques ?**



2 loops RG

#### In search for the collective mode...



X. Montiel, T. Kloss, Y. Gallais, A. Sacuto, CP, preprint



#### **Controversy with the CDW wave vector**



Y. Wang, A. Chubukov (2014)



Pre-emptive order breaking TR



S. Sachdev, D. Chowdhury (2014)



Atkinson, A. Kampf (2013)



Secondary order at the tip of the arcs

Co-existence : PG + Qx, Qy CDW



## Conclusions

• Charge orders are a key players in cuprate physics: natural competitor of superconductivity.

• SU(2) symmetry present in the under-doped region of the phase diagram

• Pseudo-gap with SU(2) symmetry and charge orders are precursors of the AFM order

• One can stabilize axial-CDW in co-existence with composite Peierls-SC phase

• SU(2) rotation of axial CDW=PWD

