Méthodes qualitatives pour les problèmes de diffraction inverses en électromagnétisme Main topic: Transmission Eigenfrequencies for Dielectrics and Their Use in The Identification Problem

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Joint work with

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Radar, Sonar, Medical Imaging, Non destructive testing, ···



Inverse problem: Determine the geometry (imaging) and some physical properties (identification) of inclusions from the knowledge of diffracted fields (associated with several incident waves).

- nonlinear problem
- unstable with respect to measurement error (ill-posed problem)

• uniqueness is not guaranteed with only several measurements?

None of the existing numerical methods can efficiently treat the problem in its general setting.

Radar, Sonar, Medical Imaging, Non destructive testing, ···



Inverse problem: Determine the geometry (imaging) and some physical properties (identification) of inclusions from the knowledge of diffracted fields (associated with several incident waves).

Commonly used approaches: (linearization of the inverse problem)

- The Kirchhoff approximation (high frequencies)
- The Born approximation (weak scatterer, low frequencies, small inclusions)

 \Rightarrow Provide a satisfying solution for a large variety of practical problems

Radar, Sonar, Medical Imaging, Non destructive testing, ···



In our case: We consider problems for which the linearization is not possible (strongly non linear problems)

- inclusions with high contrast, frequency in the resonant regime
- important multiple scattering effects (closely spaced objects)
- \Rightarrow We use multistatic data at fixed or multiple frequencies.

Goal: Get reliable qualitative information with few a priori information.

Radar, Sonar, Medical Imaging, Non destructive testing, ···



Typical targeted applications:

- Imaging of urban infrastructures, landmines, ... (GPR)
- Microwave imaging of malignancies or blood content tissues
- Nondestructive testing of complex media

Physical (nonhomogeneous, anisotropic) properties are not known (a priori).

• Part I. Beginning of the story ...

Sampling Methods for Imaging problems

- Imaging of anisotropic inclusions using electromagnetic waves
- Interior transmission problem
- Part II. The story continues...

Qualitative identification procedures

- Transmission eigenvalues (or eigenfrequencies)!
- Application to the identification problem and to nondestructive testings

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Sampling Methods for Imaging problems

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Mathematical Model: Forward problem

- Constant permeability $\mu = \mu_0$.
- ϵ and σ can be matrices. $\epsilon = \epsilon_0$ and $\sigma = 0$ in $\mathbb{R}^3 \setminus D$.

• Inclusion D is characterized by the index N

$$N(x)=\frac{\epsilon(x)}{\epsilon_0}+i\frac{\sigma(x)}{\epsilon_0\omega},\ N(x)=I \ \mbox{in}\ \ \mathbb{R}^3\setminus D$$

Wave number: $k^2 = \epsilon_0 \, \mu_0 \, \omega^2$

$$\begin{cases} (i) \quad \operatorname{curl}\operatorname{curl} E - k^2 N E = p\delta(x - x_0) \quad \text{in } \mathbb{R}^3, \\ (ii) \quad E^s = E - E^i \in H_{loc}(\operatorname{curl}, \mathbb{R}^3) \\ (iii) \quad \lim_{r \to \infty} (\operatorname{curl} E^s \times x - ik \, r \, E^s) = 0 \end{cases}$$

Incident electromagnetic field E^i :

$$\operatorname{curl}\operatorname{curl} E^i - k^2 E^i = p\delta(x - x_0) \text{ in } \mathbb{R}^3$$

Mathematical Model: Inverse problem



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 $\mathbb{G}(x; x_0)$: the Green tensor of the background medium.

Mathematical Model: Inverse problem



 $\mathbb{G}(x; x_0)$: the Green tensor of the background medium.

Inverse problem: Determine D and N from the knowledge of $E^{s}(x; x_{0}, p) \times \nu$, for all $x_{0} \in \Sigma$, $p \in T_{x_{0}}$, and all $x \in \Gamma$.

Theorem (*Cakoni-Colton*, 2003): Assume that the complement of D is connected, and (N - I) is positive definite on D, then D is uniquely determined.

However, N cannot be uniquely determined, even if measurements are available for a range of frequencies $k \in [k_1, k_2]$.

Retrieve the geometry: Sampling Methods

Linear Sampling Method, Factorization, Probe Method, MUSIC, Reciprocity Gap

A quick development in recent years after the introduction of the Linear Sampling Method by *D. Colton* and *A. Kirsch* in 1996...

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- Algorithms are not based on a linearization assumption
- Do not require a forward solver and a priori knowledge on the physical parameters
- Quick methods (trivially parallel)!

Limitations:

- Use multi-static data (at a fixed frequency)
- Precision strongly dependent on the number of measurements

• Only the geometrical information is provided

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Tutorial Monographs:

- Cakoni-Colton ('06) LSM for the scalar problem
- Grinberg-Kirsch ('08) The Factorization method

Associate with a sampling point z of the probed domain a criterion $\mathcal{G}(z)$ that indicates whether z is in the interior or the exterior of the scatterer.

Inversion scheme: consider a collection of points $(z_i)_{i=1,\dots,N}$ meshing the probed region and evaluate $\mathcal{G}(z_i)$ for each z_i .



The computation of \mathcal{G} needs to be quick!

Proposed methods differ according to the choice and the way of computing of \mathcal{G} .

The Linear Sampling Method (LSM)



 $\mathbb{G}(x; x_0)$: the Green tensor of the background medium.

Near field equation of LSM: seek $g_z \in \tilde{H}_{curl}^{-\frac{1}{2}}(\Sigma)$ such that for all $x \in \Gamma$,

$$(F g_{\boldsymbol{z}})(x) := \int_{\Sigma} \nu(x) \times \underline{E}^{\boldsymbol{s}}(x; x_0, g_{\boldsymbol{z}}(x_0)) \, ds(x_0) \simeq \nu(x) \times \mathbb{G}(x; \boldsymbol{z}) \, q$$

Fg scattered wave associated with: $(\mathcal{H}g)(x) := \int_{\Sigma} \mathbb{G}(x,y)g(y) \, ds(y)$

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Criterion: $\mathcal{G}(z) = ||g_z||$ has large values if $z \notin D$ as compared with $z \in D$.

$$(F g_z)(x) \simeq \nu(x) \times \mathbb{G}(x; z) q \quad x \in \Gamma,$$

for $z \in D$, is equivalent to finding E and E_0 solutions to

$$\begin{cases} \operatorname{curl}\operatorname{curl} E - k^2 N E = 0 & \operatorname{in} D, \\ \operatorname{curl}\operatorname{curl} E_0 - k^2 E_0 = 0 & \operatorname{in} D. \end{cases}$$
(1)

$$\begin{cases} (\underline{E} - \underline{E}_0) \times \nu = (\mathbb{G}(\cdot; z) q) \times \nu & \text{on } \partial D, \\ (\operatorname{curl} \underline{E} - \operatorname{curl} \underline{E}_0) \times \nu = \operatorname{curl}(\mathbb{G}(\cdot; z) q) \times \nu & \text{on } \partial D. \end{cases}$$
(2)

Problem (2)-(1) is called the Interior Transmission Problem.

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Space of solutions: $(E, E_0) \in L^2(D)^3$ satisfy (1) in the distributional sense; $E - E_0 \in H(\operatorname{curl}, D)$; $\operatorname{curl}(E - E_0) \in H(\operatorname{curl}, D)$ and satisfy (2)

Interior Transmission Problem

Theorem (H. 04): Assume that $(N-1)^{-1}$ is bounded and that $\mathcal{R}e(N-1)^{-1}$ is a positive definite matrix field on D then ITP is of Fredholm type.

Definition: Transmission eigenvalues (eigenfrequencies) are the values of k for which uniqueness of ITP solutions fails.

Observations:

- Uniqueness for ITP \longrightarrow Injectivity of $F \iff 1$ is not an eigenvalue of the scattering matrix)
- The LSM fails for k being a transmission eigenvalue.

Theorem (continued):

- If $\mathcal{I}m(N) > 0$ in D then the solution is unique.
- If not, then the set of k for which uniqueness fails is at most discrete.



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Aperture (sources-measurements)

Reconstruction

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Aperture (sources-measurements) Reconstruction

Numerical examples for dielectrics with experimental data from the Fresnel institute : Corocco et al '09

A qualitative approach for the identification problem

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The identification problem: an original approach Cakoni-Colton-H. (2008-2009)

Example of geometry reconstruction using LSM (M'B Fares)





Goal: Obtain information on N, assuming that D (or an approximation of it) is known.

The idea is to use the transmission eigenvalues...

Interior Transmission Problem: equivalent formulation

Assume that $(N-1)^{-1}$ is bounded a.e. in D.

Define: $\mathbf{u} = \mathbf{E} - \mathbf{E}_0$.

 $\mathbf{u} \in \mathcal{U}(D) := \{\mathbf{u} \in H(\operatorname{curl}, D); \ \operatorname{curl} \mathbf{u} \in H(\operatorname{curl}, D)\},\$

satisfies

$$\begin{aligned} \operatorname{curl} \operatorname{curl} (N-1)^{-1} \operatorname{curl} \operatorname{curl} \mathbf{u} &- k^2 \operatorname{curl} \operatorname{curl} (N-1)^{-1} \mathbf{u} \\ &- k^2 N (N-1)^{-1} \operatorname{curl} \operatorname{curl} \mathbf{u} + k^4 N (N-1)^{-1} \mathbf{u} &= 0 \text{ in } D, \end{aligned}$$

and

$$\left\{ \begin{array}{ll} \mathbf{u} \times \nu = (\mathbb{G}(\cdot; \mathbf{z}) \, q) \times \nu, & \text{on } \partial D, \\ (\operatorname{curl} \mathbf{u}) \times \nu = \operatorname{curl}(\mathbb{G}(\cdot; \mathbf{z}) \, q) \times \nu & \text{on } \partial D. \end{array} \right.$$

Reciprocally:

$$E = (N-1)^{-1} (\operatorname{curl} \operatorname{curl} \mathbf{u} - k^2 \mathbf{u})/k^2$$
$$E_0 = (N-1)^{-1} (\operatorname{curl} \operatorname{curl} \mathbf{u} - k^2 N \mathbf{u})/k^2$$

Solvability of the Interior Transmission Problem

The key point is to observe that the variational formulation can be written in the form

$$\mathcal{A}_k(\mathbf{u},\mathbf{u}') - k^2 \mathcal{B}(\mathbf{u},\mathbf{u}') = 0 \quad \text{for all } \mathbf{u}' \in \mathcal{U}_0(D), \tag{3}$$

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with $\mathcal{U}_0(D) := \{ \mathbf{u} \in H_0(\operatorname{curl}, D); \operatorname{curl} \mathbf{u} \in H_0(\operatorname{curl}, D) \}$ and

$$\begin{aligned} \mathcal{A}_k(\mathbf{u},\mathbf{u}') &= \left((N-1)^{-1} (\operatorname{curlcurl} \mathbf{u} - k^2 \mathbf{u}), \left(\operatorname{curlcurl} \mathbf{u}' - k^2 \mathbf{u}' \right) \right)_D \\ &+ k^4 \left(\mathbf{u}, \mathbf{u}' \right)_D \\ \mathcal{B}(\mathbf{u},\mathbf{u}') &= \left(\operatorname{curl} \mathbf{u}, \operatorname{curl} \mathbf{u}' \right)_D \end{aligned}$$

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Definition: Transmission eigenvalues (eigenfrequencies) are the values of k for which uniqueness of ITP solutions fails; i.e. $\exists \mathbf{u} \in \mathcal{U}_0(D)$ non trivial solution of (3)

 \longrightarrow Quadratic eigenvalue problem.

It can be shown that it is also a linear **non self-adjoint** eigenvalue problem.

Remark: In the case of perfect conductors these special frequencies correspond to Maxwell's eigenfrequencies of the cavity D, i.e the k > 0 such that:

$$\begin{cases} \exists \mathbf{u} \in H_0(\operatorname{curl}, D); \mathbf{u} \neq 0; \\ \operatorname{curl} \operatorname{curl} \mathbf{u} - k^2 \mathbf{u} = 0 \quad \text{in } D. \end{cases}$$

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Existence of transmission eigenvalues:

- Colton-Kress (91') : Spherically, stratified medium: existence of infinitely many.
- Païvaïrinta-Sylvester ('08) : (for the scalar problem), Kirsch ('09) Cakoni-H. ('09) (generalization to the Maxwell problem): There exist eigenvalues for N sufficiently large (or sufficiently small).
- **③** Cakoni-H.-Gintides ('10) : existence of infinitely many without restrictions on N

$$\mathcal{A}_k(\mathbf{u},\mathbf{u}') - k^2 \mathcal{B}(\mathbf{u},\mathbf{u}') = 0$$
 for all $\mathbf{u}' \in \mathcal{U}_0(D)$,

Idea: Consider the eigenvalues of A_k with respect to the compact and non injective operator B, i.e. the $\lambda_j(k^2)$ such that $\ker(A_k - \lambda_j(k^2)B) \neq \{0\}$.

Theorem: The eigenvalues of $(\lambda_j(k^2))_{j\geq 1}$ are given by

$$\lambda_j(k^2) = \min_{W \subset \mathcal{U}_j} (\max_{u \in W \setminus \{0\}} (A_k u, u) / (Bu, u))$$

where U_j denotes the set of all *j*-dimensional subspaces W of $\mathcal{U}_0(D)$ such that $W \cap \ker(B) = \{0\}$.

The transmission eigenvalues are $k^2 = \tau$, solutions to

$$\lambda_j(\tau) = \tau$$
 for some $j \ge 1$.

Theorem: Assume that there exist two positive constants $\tau_0>0$ and $\tau_1>0$ such that

• $A_{ au_0} - au_0 B$ is positive on \mathcal{U}_0 ,

• $A_{\tau_1} - \tau_1 B$ is non positive on a *m*-dimensional subspace W_m of U. Then each of the equations $\lambda_j(\tau) = \tau$ for $j = 1, \ldots, m$, has at least one solution in $[\tau_0, \tau_1]$.

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Assume that $\mathcal{I}m(N) = 0$ a.e. in D and $||N(x)||_2 \ge \delta > 1$ for $x \in D$.

• $A_{\tau} - \tau B$ is positive for all $\tau < \tau_0 = \lambda(D) / \sup_D ||N||_2$, where $\lambda(D)$ is the first Dirichlet eigenvalue of $-\Delta$ on D. Theorem: Assume that there exist two positive constants $\tau_0>0$ and $\tau_1>0$ such that

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- $A_{\tau} \tau B$ is positive for all $\tau < \tau_0 = \lambda(D) / \sup_D ||N||_2$, where $\lambda(D)$ is the first Dirichlet eigenvalue of $-\Delta$ on D.
- Existence of τ_1 and W_m ?

Idea: use existence of eigenvalues for spherical geometries.

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Idea: use existence of eigenvalues for spherical geometries. Assume $\|N\|_\infty \geq n^* > 1$

 $A^*_{ au}$ same as $A_{ au}$ with N replaced with n_*

$$((A_{\tau} - \tau B)\mathbf{u}, \mathbf{u})_{\mathcal{U}_0} \le ((A_{\tau}^* - \tau B)\mathbf{u}, \mathbf{u})_{\mathcal{U}_0}$$

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 $\tau = k_1^*(r)^2$ first tr. eigen. associated with n_* for the sphere S of radius r. $\mathbf{u}_S = \mathbf{u}_1^*$ inside S, extended by 0 outside...

$$((A_{\tau} - \tau B)\mathbf{u}_S, \mathbf{u}_S)_{\mathcal{U}_0} \le 0$$



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$$\tau = k_1^*(r)^2 \text{ first tr. eigen. associated with } n_* \text{ for the sphere } S \text{ of radius } r.$$

$$\mathbf{u} \in \text{Span}\{\mathbf{u}_{S_j} ; j = 1, \dots, N(r)\}$$

$$((A_{\tau} - \tau B)\mathbf{u}, \mathbf{u})_{\mathcal{U}_0} \leq 0$$

$$\implies \text{Existence of } N(r) \text{ tr. eigenvalues } !$$

But: $N(r) \to \infty$ as $r \to 0$.

First application: a priori lower bound on material coefficients

• If $||N(x)||_2 \ge \delta > 1$ for all $x \in D$ then for any positive transmission eigenvalue k we have

$$\sup_{D} \|N(x)\| \ge \lambda(D)/k^2$$

where $\lambda(D)$ denotes the first Dirichlet eigenvalue of $-\Delta$ on D.

• $N \rightarrow k_1(N, D)$ "is decreasing":

$$n_* \le N \le n^* \Longleftrightarrow k_1(n^*, D) \le k_1(N, D) \le k_1(n_*, D)$$

• $D \rightarrow k_1(N, D)$ "is decreasing".

How to compute the transmission eigenvalues?

Without knowing N!

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Idea: use the failure of the LSM for those values: the existence of nearby solution

$$(F g_{\mathbf{z}})(x) \simeq \nu(x) \times \mathbb{G}(x; \mathbf{z}) q(=: E_{e}(\cdot; \mathbf{z}, q))$$

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with bounded "Herglotz norm" is no longer guaranteed if k is a transmission eigenvalue.

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with bounded "Herglotz norm" is no longer guaranteed if k is a transmission eigenvalue.

Algorithm: We fix a point z in D and evaluate

$$k \mapsto \|\boldsymbol{g}_{\alpha}(\cdot; \boldsymbol{z}, q)\|$$

where

$$(\alpha + F^*F) \ g_{\alpha}(\cdot; z, q) = F^* \left(E_e(\cdot; z, q) \right)$$

We expect to observe peaks at the irregular frequencies...

How to compute the transmission eigenvalues

Justification: Let F_{δ} be the operator corresponding to noisy measurements

 $\|F_{\delta} - F\| \le \delta$

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$$\|F_{\delta} - F\| \le \delta$$

$$(\alpha(\delta) + F_{\delta}^*F_{\delta}) \ g_{\delta}(\cdot; z, q) = F_{\delta}^* \left(E_e(\cdot; z, q) \right)$$

Assumptions:

• F_{δ} is a family of uniformly compact operators with dense range • $\sqrt{\alpha(\delta)}/\delta \to 0$ as $\delta \to 0$

Theorem: Let k be a transmission eigenvalue, then for a.e. $z \in D$

$$\lim_{\delta \to 0} \|\mathcal{H}g_{\delta}(\cdot; z, q)\|_{L^2(D)} = \infty$$

Main "formal idea of the proof": $E_e(\cdot; z, q)$ cannot be "orthogonal" to transmission eigenvectors for all z in an open set of D.

Validation in the spherical case



Blue line: $||g(\cdot, z)||$ in terms of the wave number k

Dark line: Computed eigenvalues using numerical solutions to ITP Circles: Computed eigenvalues using speration of variables **Idea:** Use eigenfrequencies to obtain a qualitative information on the presence of faults inside D

Advantage: We do not need to know N or to solve any direct problem Limitation: Only a qualitative information can be obtained

First investigations: test the presence of cavities inside D (application: cavities inside composite materials)

N=1 in $D_0 \subset D$

 \implies difficulty in the mathematical analysis:

N-1 is not invertible a.e. in D!



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ITP with presence of cavities Cossonnière-H.

with $\mathcal{U}_0(D) := \{ \mathbf{u} \in H_0(\operatorname{curl}, D); \operatorname{curl} \mathbf{u} \in H_0(\operatorname{curl}, D) \}$ and

 $\mathcal{U}(D, D_0, k^2) := \{ \mathbf{u} \in \mathcal{U}_0(D) \text{ such that } \operatorname{curlcurl} \mathbf{u} - k^2 \mathbf{u} = 0 \text{ in } D_0 \}$

Assume that k^2 is not a *Dirichlet and Neumann* eigenvalue of $-\Delta$ in D_0 and that N and $(N-1)^{-1}$ are positive definite in $D \setminus D_0$.

Then, solving ITP is equivalent to solve for $\mathbf{u} \in \mathcal{U}(D,D_0,k^2)$ such that

$$\boxed{\mathcal{A}_k(\mathbf{u},\mathbf{u}')-k^2\mathcal{B}(\mathbf{u},\mathbf{u}')=0} \ \text{ for all } \mathbf{u}'\in\mathcal{U}(D,D_0,k^2),$$

 $\begin{aligned} \mathcal{A}_k(\mathbf{u},\mathbf{u}') &:= \left((N-1)^{-1} (\operatorname{curlcurl} \mathbf{u} - k^2 \mathbf{u}), \, (\operatorname{curlcurl} \mathbf{u}' - k^2 \mathbf{u}') \right)_{D \setminus D_0} \\ &+ k^4 \, (\mathbf{u},\,\mathbf{u}')_D \end{aligned}$

 $\mathcal{B}(\mathbf{u},\mathbf{u}') \quad := \quad (\operatorname{curl}\mathbf{u},\operatorname{curl}\mathbf{u}')_D$

ITP with presence of cavities: existence of transmission eigenvalues

Difficulty: The variational space

 $\mathcal{U}(D, D_0, k^2) := \{ \mathbf{u} \in \mathcal{U}_0(D) \text{ such that } \operatorname{curlcurl} \mathbf{u} - k^2 \mathbf{u} = 0 \text{ in } D_0 \}$

depends on k...

Theorem: $k \mapsto P_k$, where $P_k : \mathcal{U}_0(D) \to \mathcal{U}(D, D_0, k^2) \subset \mathcal{U}_0(D)$ denotes the orthogonal projection operator, is continuous.

 \Longrightarrow Similar analysis as in the case without cavities carry over by considering

$$I + k^2 P_k^* A_k^{-1/2} B_k A_k^{-1/2} P_k : \mathcal{U}_0(D) \to \mathcal{U}_0(D)$$

 \implies Existence of infinitely many transmission eigenvalues.

Countability of transmission eigenvalues

Countability of transmission eigenvalues cannot be obtained this way since $k \mapsto P_k$ cannot be analytic.

We substitute to P_k a projection like, but analytic operator: $\tilde{P}_k: \mathcal{U}_0(D) \to \mathcal{U}(D, D_0, k^2) \subset \mathcal{U}_0(D)$

$$\tilde{P}_k \mathbf{u} := \mathbf{u} - \chi \theta_k \mathbf{u}$$

 χ is a C^{∞} cutoff function = 1 in D_0 and 0 outside D. $\theta_k \mathbf{u} := \int_{D_0} \Im \mathbb{G}(x, y) (\operatorname{curlcurl} \mathbf{u} - k^2 \mathbf{u})(y) \, dy$ **Theorem:** If k^2 is a transmission eigenvalue then there exists non trivial solution $\mathbf{u} \in \mathcal{U}_0(D)$ to

$$\mathcal{A}_k(\tilde{P}_k\mathbf{u},\overline{\tilde{P}_k\overline{\mathbf{u}'}}) + \alpha(\theta_k\mathbf{u},\overline{\theta_k\overline{\mathbf{u}'}})_{\mathcal{U}} - k^2\mathcal{B}(\tilde{P}_k\mathbf{u},\overline{\tilde{P}_k\overline{\mathbf{u}'}}) = 0$$

for all $\mathbf{u}' \in \mathcal{U}_0(D)$.

ITP with presence of cavities

- The set of transmission eigenvalues is discrete and $+\infty$ is the only point of accumulation.
- If $D_0 \subset D'_0$ and $N_1 \ge N_2 > 1$ then $k_1(D, D'_0, N_1) > k_1(D, D_0, N_2)$.

• Transmission eigenvalues satisfy

$$k^{2} \geq k_{*}^{2}(D, D_{0}) = \frac{\lambda(D, D_{0})\mu(D, D_{0})}{\lambda(D, D_{0}) + \sup |N| \mu(D, D_{0})}$$
$$\lambda(D, D_{0}) = \inf_{k \geq 0} \left(\inf_{\psi \in \mathcal{U}_{0}(D, D_{0}, k)^{*}} \|\nabla\psi\|_{L^{2}(D)}^{2} / \|\psi\|_{L^{2}(D\setminus\overline{D}_{0})}^{2} \right) \cdot$$
$$\mu(D, D_{0}) = \inf_{k \geq 0} \left(\inf_{\psi \in \mathcal{U}_{0}(D, D_{0}, k)^{*}} \|\nabla\psi\|_{L^{2}(D)}^{2} / \|\psi\|_{L^{2}(D_{0})}^{2} \right)$$

• If k^2 is a transmission eigenvalue then for a.e. $z \in D$, $\lim_{\delta \to 0} \|\mathcal{H}g_{\delta}\|_{L^2(D)} = \infty.$

Sphere of radius 1, index n = 4

Sphere of radius 1, index n = 4



Sphere of radius 1, index n = 4



Sphere of radius 1, index n = 4



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Some open problems and perspectives

- Improve a-priori estimates on physical properties in terms of transmission eigenfrequencies.
- Extend applications to non-destructive testing of other types of faults and materials (dispersive).
- The case of N I changing sign (photo-acoustic tomography).
- Stable identification from limited aperture data.
- What about the structure of transmission eigenfunctions... and can we say more about physical interpretations of transmission eigenfrequencies?

• Inverse spectral problem?

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Thank you!