



COLLÈGE  
DE FRANCE  
1530



CIFAR  
CANADIAN INSTITUTE  
for ADVANCED RESEARCH

# Lecture 2: Doped Mott insulators Strongly correlated superconductivity and its normal phase

André-Marie Tremblay

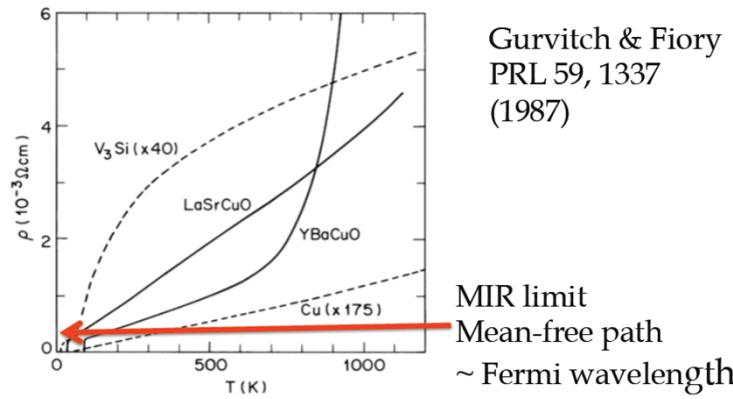
 UNIVERSITÉ DE  
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Collège de France, 16 mars 2015  
17h00 à 18h30



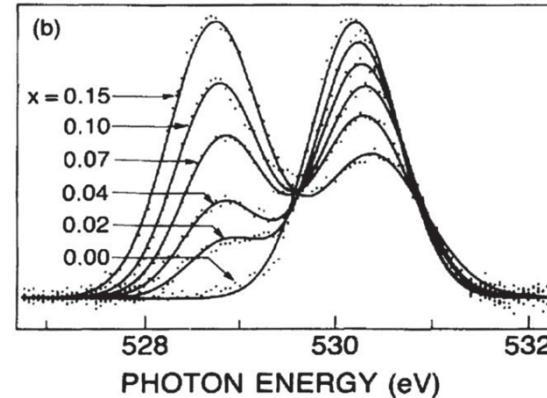
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# Last time

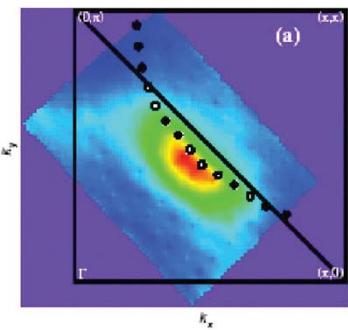


Gurvitch & Fiory  
PRL 59, 1337  
(1987)

MIR limit  
Mean-free path  
~ Fermi wavelength



Hole-doped, 10%



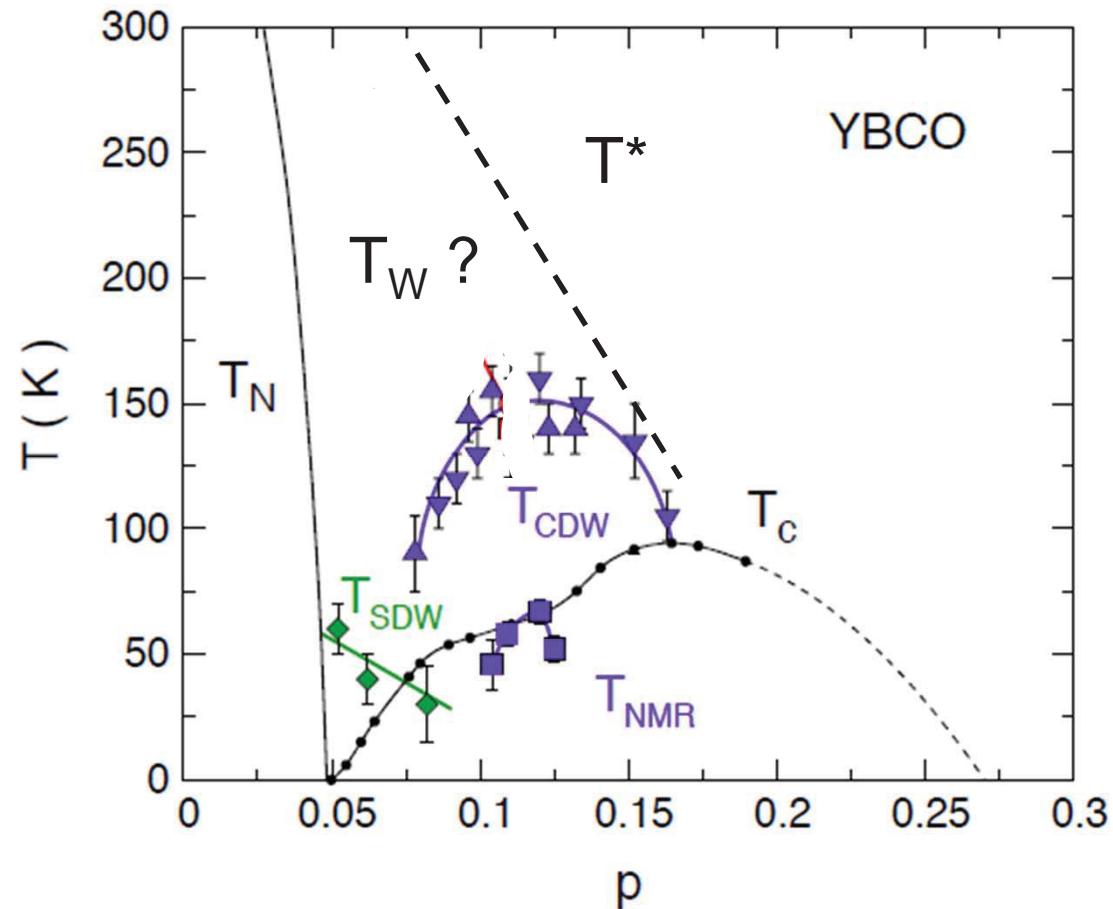
F. Ronning et al. Jan. 2002,  $\text{Ca}_{2-x}\text{Na}_x\text{CuO}_2\text{Cl}_2$

$$H = \sum_{ij\sigma} (t_{ij} - \delta_{ij}\mu) c_{i\sigma}^\dagger c_{j\sigma} + U \sum_i n_{i\uparrow} n_{i\downarrow}$$



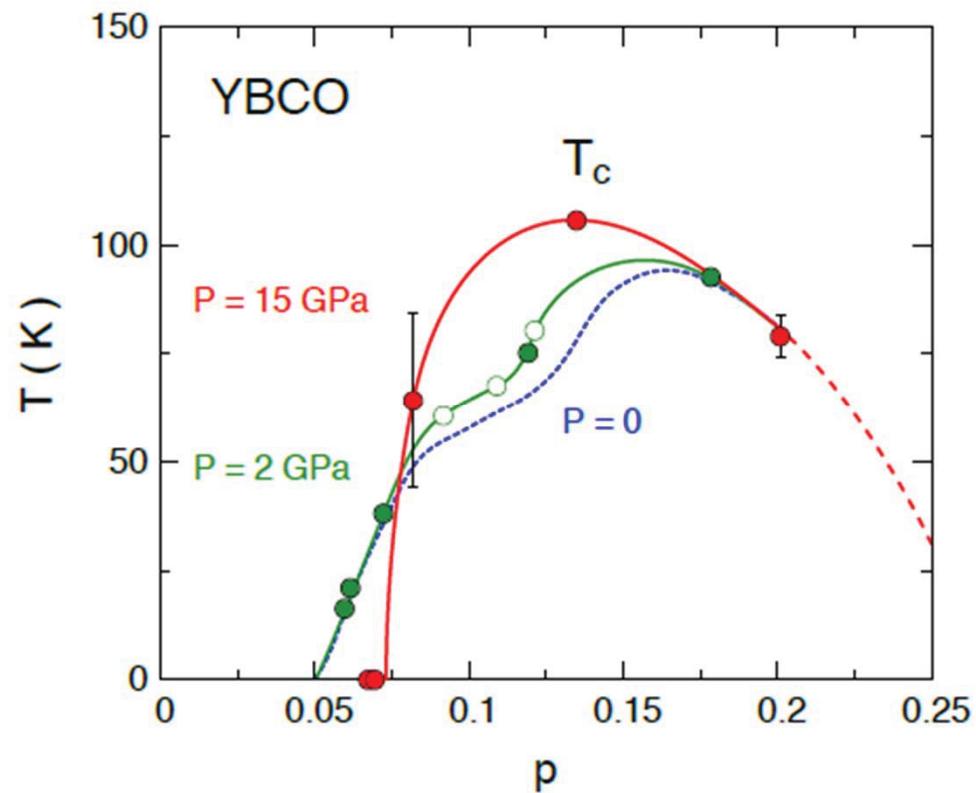
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# Phase diagram for hole-doped cuprates



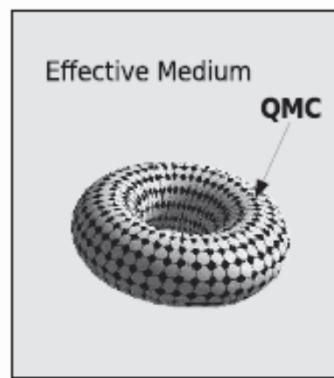
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# Getting rid of the CDW



Cyr-Choinière et al, arxiv1503.02033

# *2d Hubbard: Quantum cluster method*

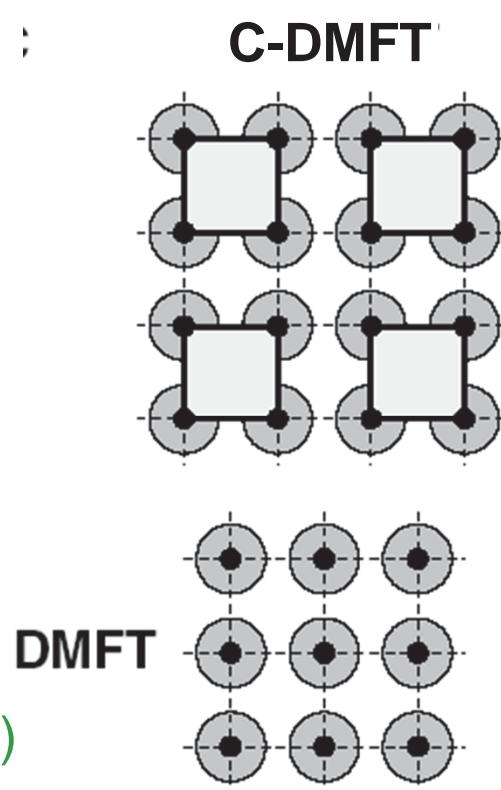


**DCA**

Hettler ... Jarrell ... Krishnamurty PRB **58** (1998)

Kotliar et al. PRL **87** (2001)

M. Potthoff et al. PRL **91**, 206402 (2003).



**REVIEWS**

Maier, Jarrell et al., RMP. (2005)

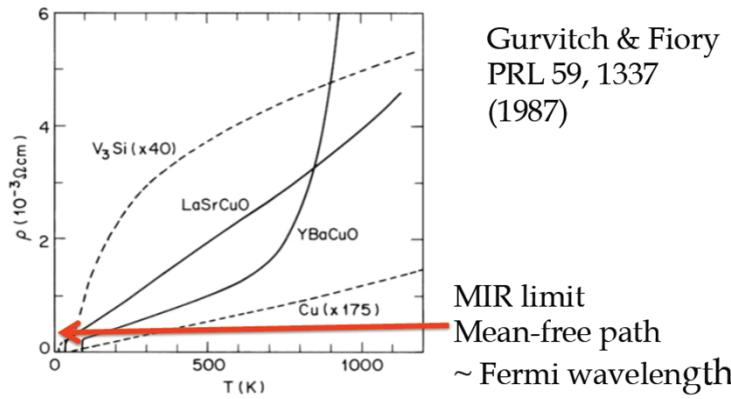
Kotliar et al. RMP (2006)

AMST et al. LTP (2006)



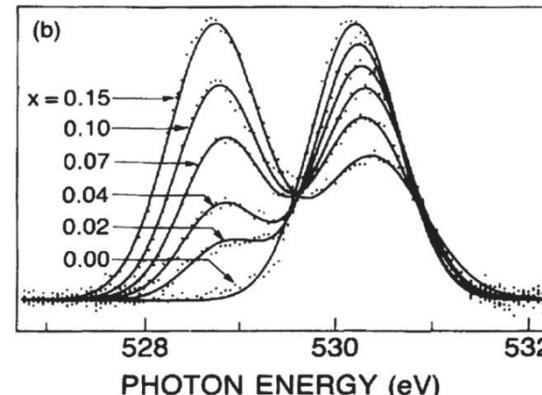
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# Last time

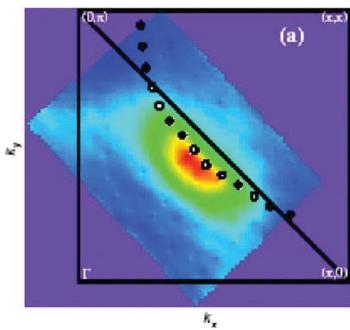


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MIR limit  
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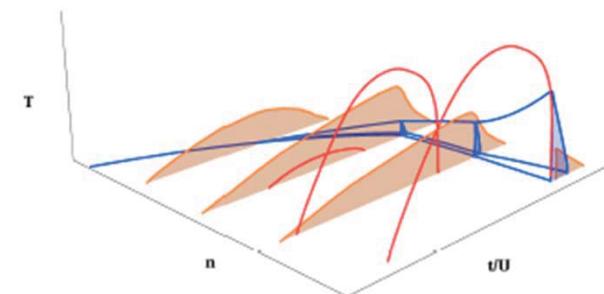
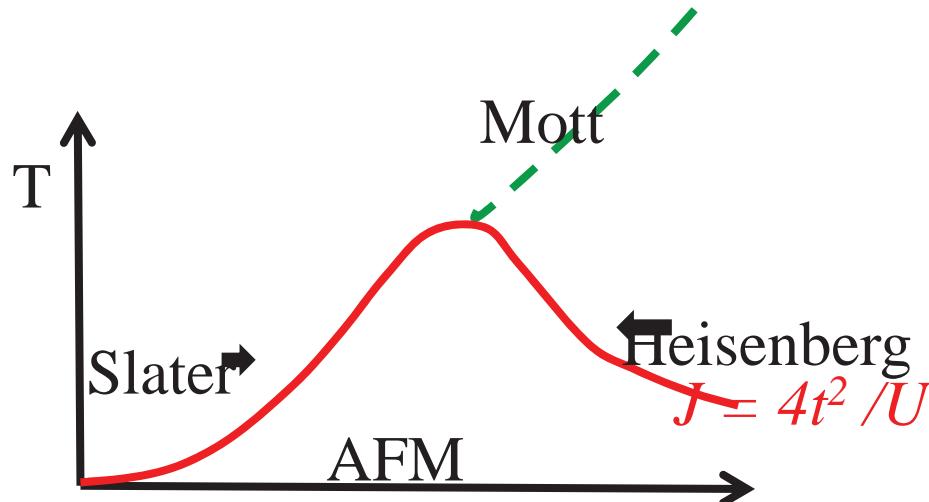


Hole-doped, 10%

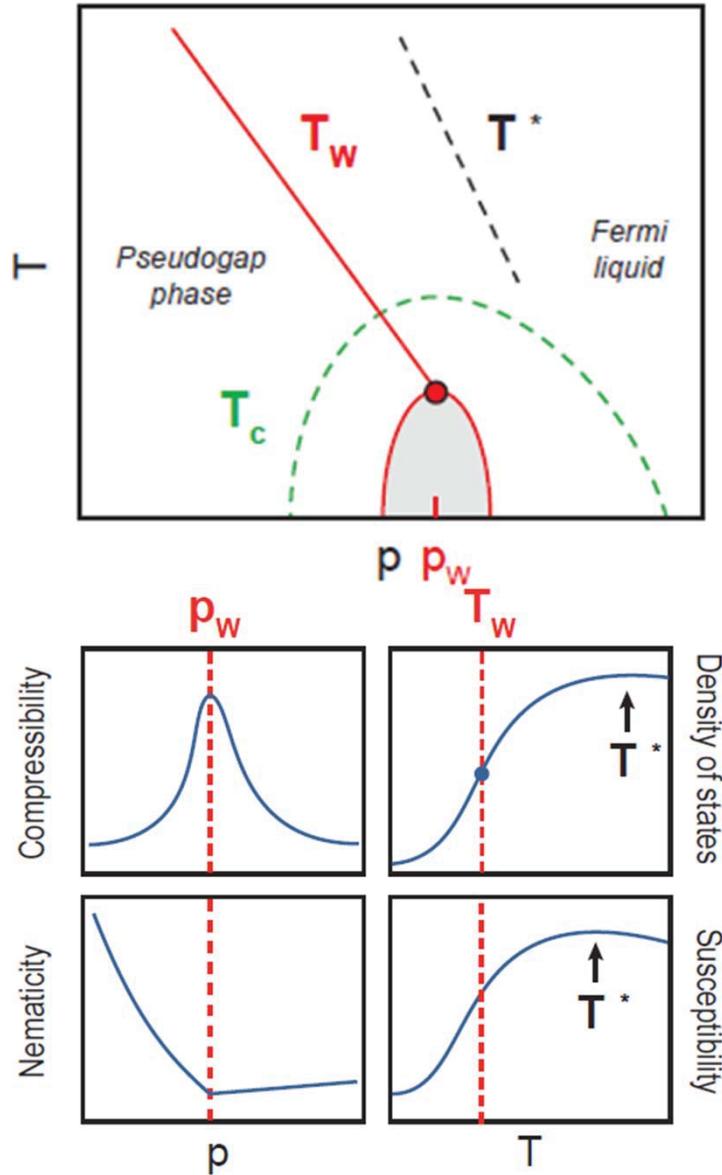


F. Ronning et al. Jan. 2002,  $\text{Ca}_{2-x}\text{Na}_x\text{CuO}_2\text{Cl}_2$

$$H = \sum_{ij\sigma} (t_{ij} - \delta_{ij}\mu) c_{i\sigma}^\dagger c_{j\sigma} + U \sum_i n_{i\uparrow} n_{i\downarrow}$$



# CDMFT: Emergent first-order transition



- Is the pseudogap (PG) a crossover or a phase transition ?
- Relation between CDW and the PG ?
- Why CDW peaked at 12% doping ?
- Origin of nematicity ?
- Why a dome of SC ?
- Why superconducting ?
- Does a one-band model capture the key physics ?
- AFM QCP important?
- Lessons from other SC?



# Today

- « Normal » state of cuprates
  - Signatures of Mott physics away from  $n=1$
- Superconductivity
  - What is special about strongly correlated SC
  - Origin

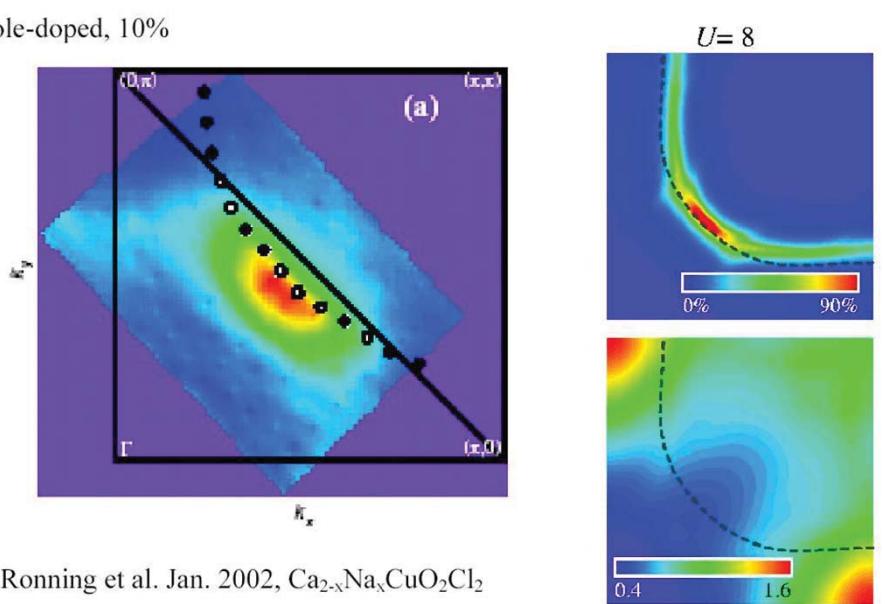


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# ARPES pseudogap

# Strong correlation pseudogap ( $U > 8t$ )

- Different from Mott gap that is local (all  $k$ ) not tied to  $\omega=0$ .
- Pseudogap close to  $\omega=0$  and only in regions nearly connected by  $(\pi,\pi)$ . (e and h),
- Pseudogap is independent of cluster shape (and size) in CPT.
- Not caused by AFM LRO
  - No LRO, few lattice spacings.
  - Not very sensitive to  $t'$
  - Scales like  $t$ .



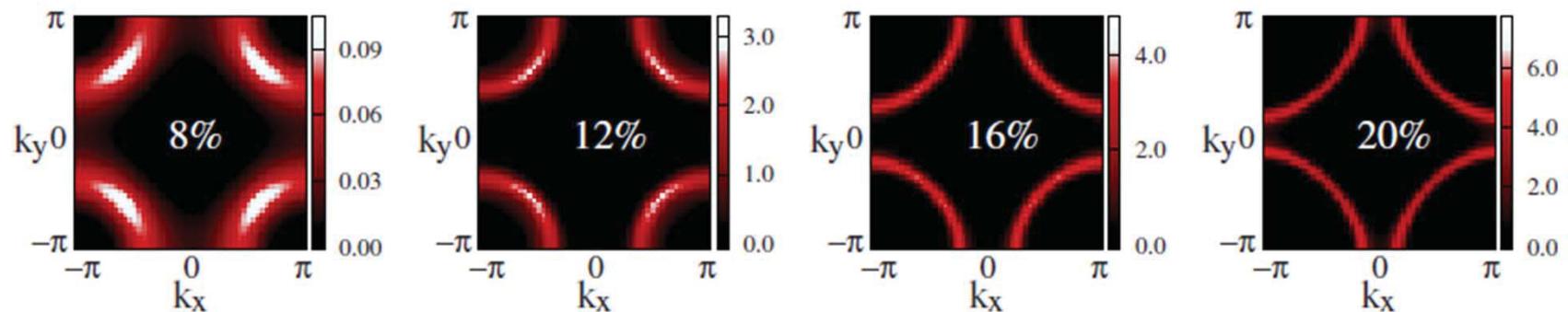
F. Ronning et al. Jan. 2002,  $\text{Ca}_{2-x}\text{Na}_x\text{CuO}_2\text{Cl}_2$

Sénéchal, AMT, PRL 92, 126401 (2004).



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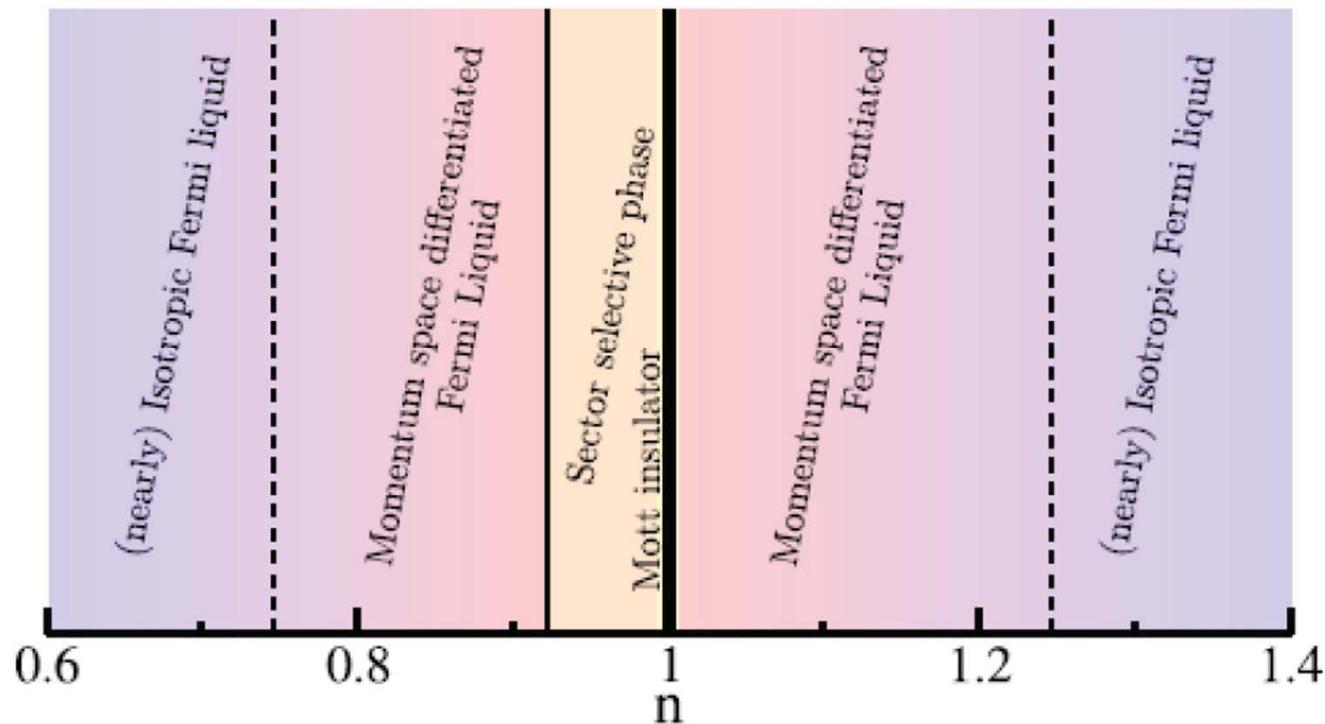
# Can be seen with 2 site DCA



Michel Ferrero, P. S. Cornaglia, L. De Leo, O. Parcollet, G. Kotliar, A. Georges  
PRB **80**, 064501 (2009)

Seen by all groups and DCA, CDMFT

# Momentum dependence of $\Sigma$



Gull, Werner, Millis, (2009)



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# Mott transition at $n = 1$

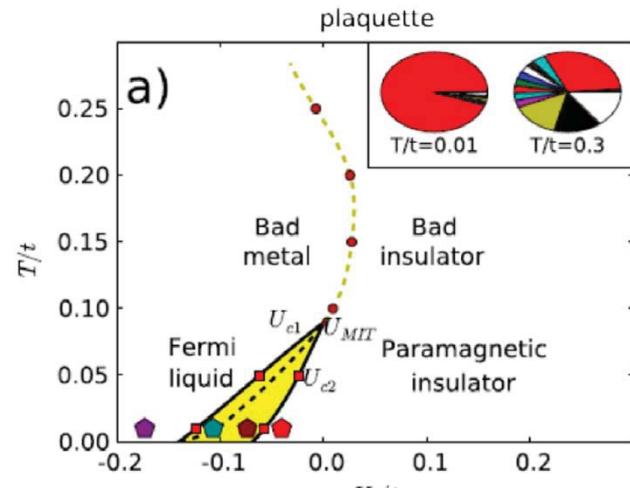
# Interaction-induced Mott transition, $n = 1$

4 sites (CDMFT-CTQMC)

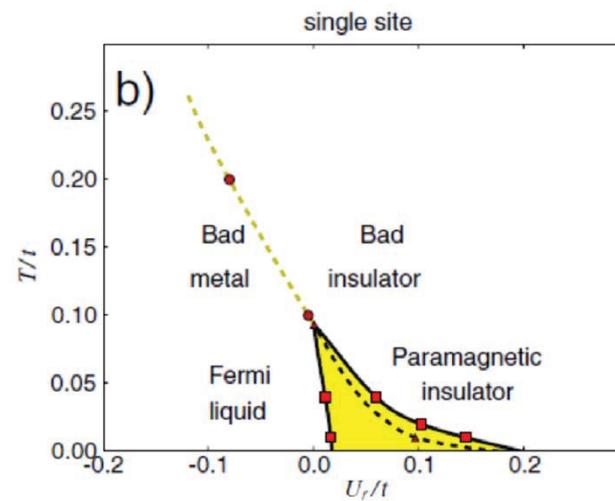
Georges, A., and W. Krauth,  
PRB **48**, 7167 (1993)

Single site

$$U_r = (U - U_{MIT})/U_{MIT}$$



$$U_{MIT} = 6.05$$

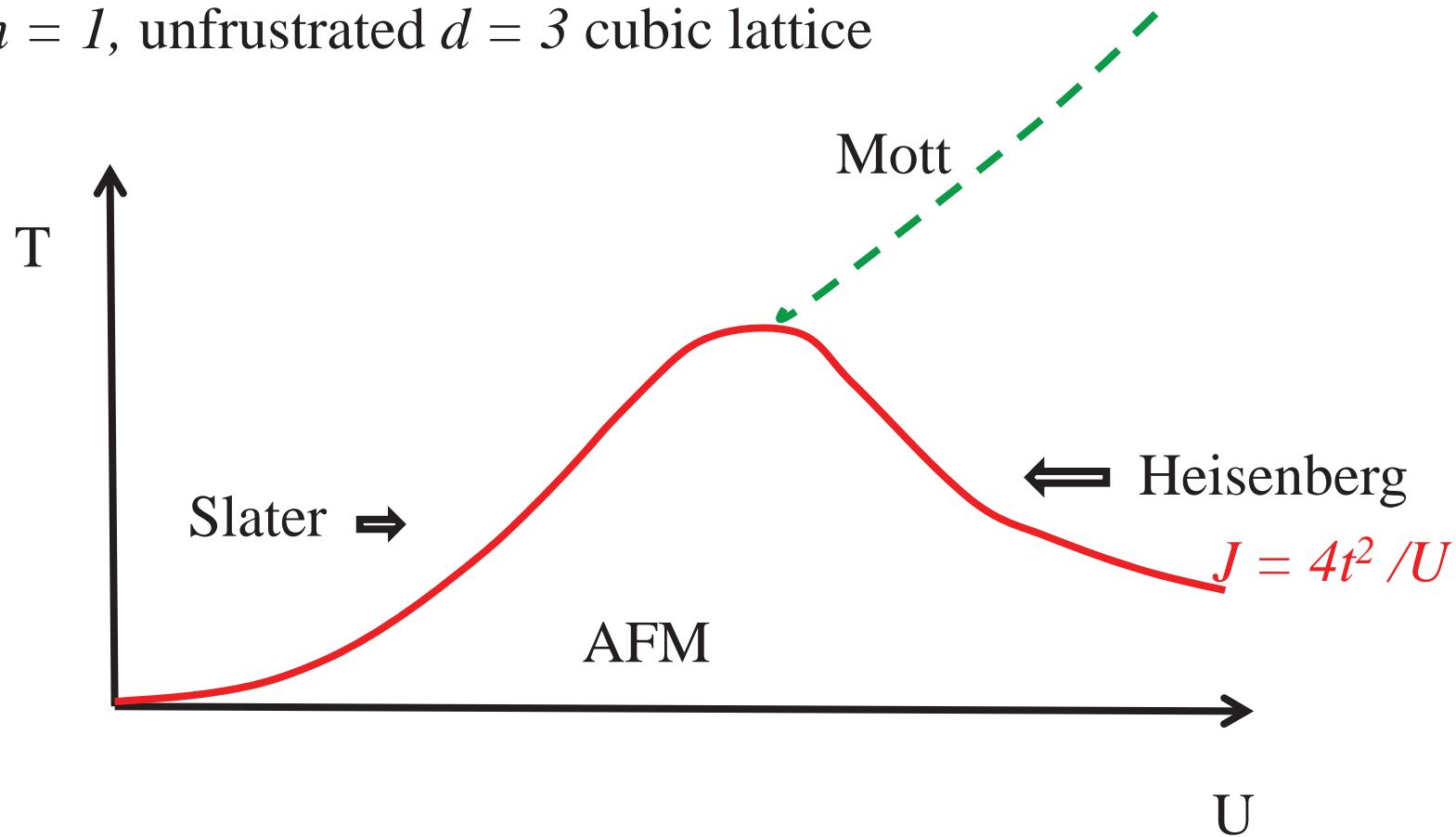


$$U_{MIT} \sim 12$$

H. Park, K. Haule, and G. Kotliar PRL **101**, 186403 (2008)  
Balzer, Kyung, Sénecal, Tremblay, Potthof EPL, **85** (2009) 17002

# Local moment and Mott transition

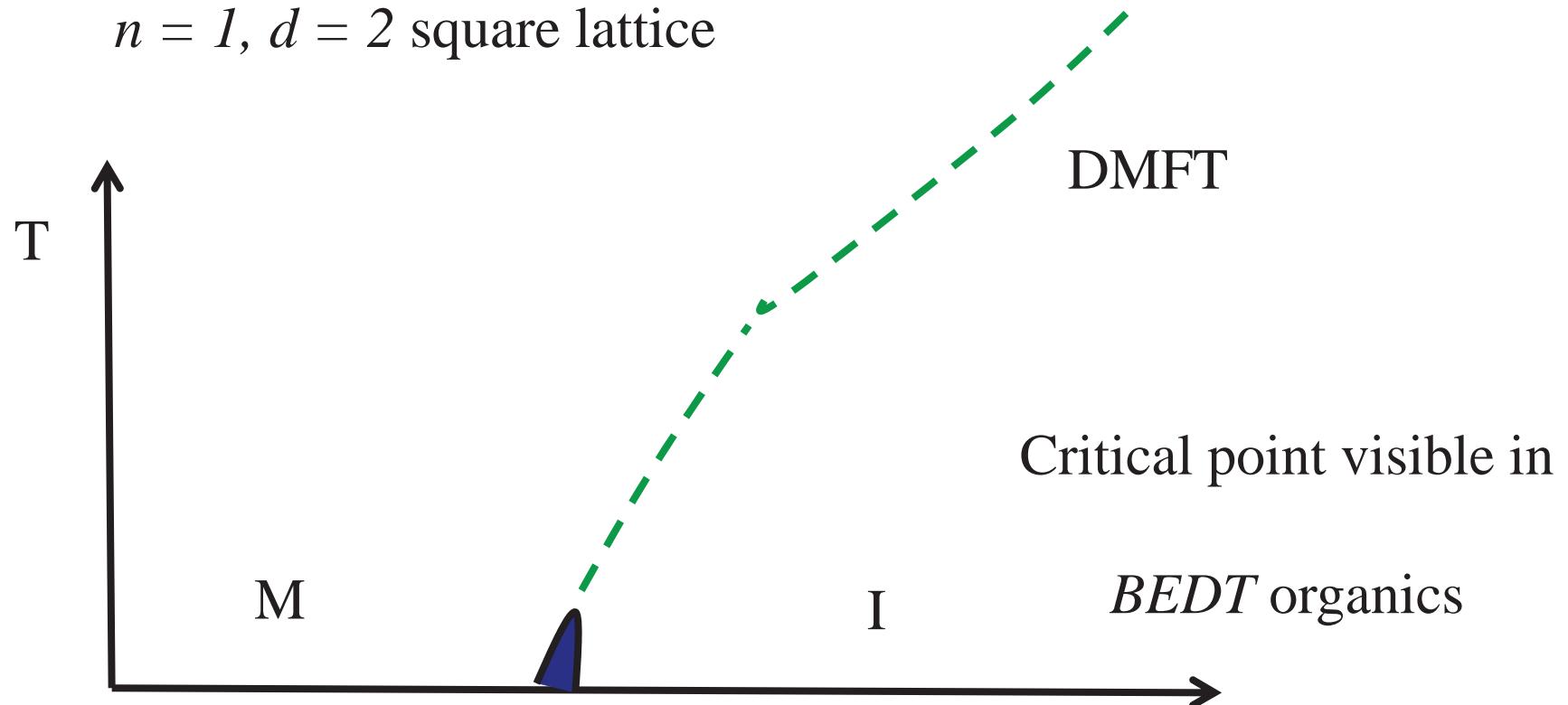
$n = 1$ , unfrustrated  $d = 3$  cubic lattice



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# Local moment and Mott transition

$n = 1, d = 2$  square lattice



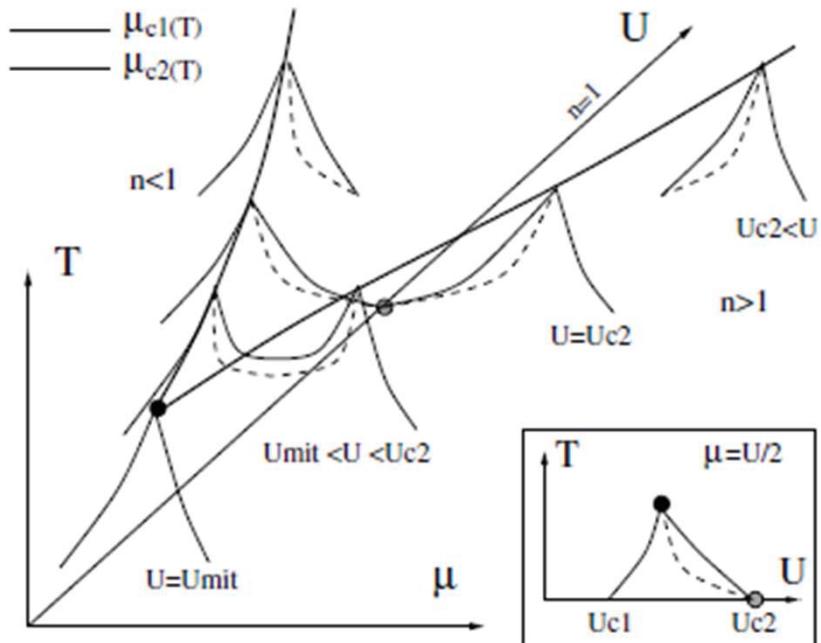
Understanding finite temperature phase from a *mean-field theory* down to  $T = 0$



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# Doped Mott insulator

# Compressibility divergence at Mott and coexistence (single-site DMFT)



G. Kotliar, S. Murthy, and M. J. Rozenberg, Phys. Rev. Lett. **89**, 046401 (2002).

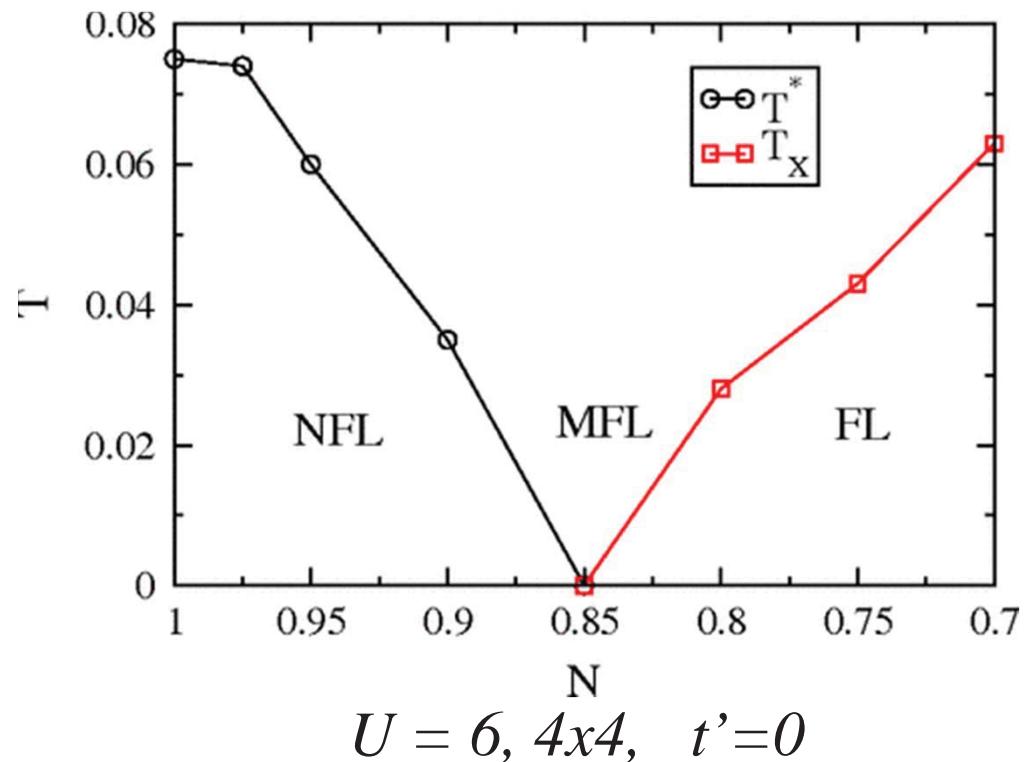
S. Murthy, Rutgers thesis 2004

K. Frikach, M. Poirier, et al.  
PRB **61**, R6491 (2000).  
S. R. Hassan, A. Georges,  
and H. R. Krishnamurthy  
PRL **94**, 036402 (2005)

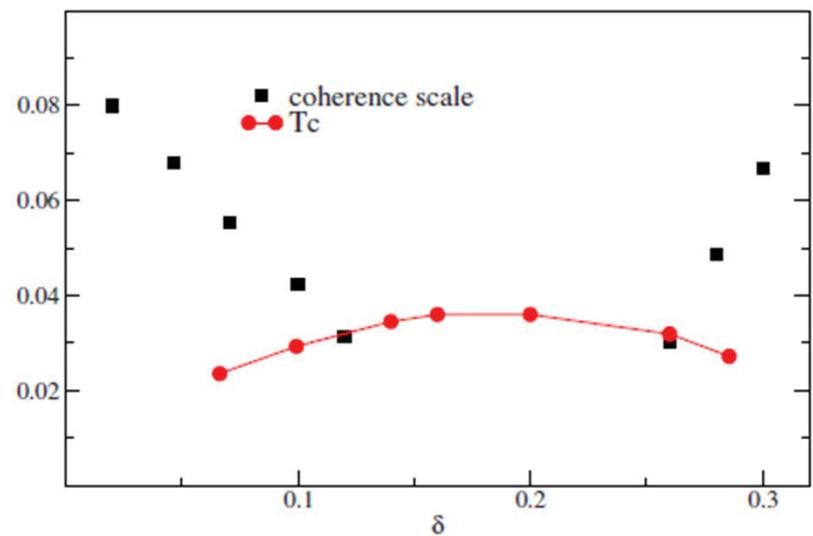
# Anomalous metallic state near half-filling (examples)

- Pseudogap
  - B. Kyung et al., PRB 73, 165114 (2006).
  - N. S. Vidhyadhiraja et al., PRL 102, 206407 (2009).
  - A. Liebsch and N.-H. Tong, PrB 80, 165126 (2009).
- Momentum selective transition
  - P. Werner et al., PRB 80, 045120 (2009).
  - M. Ferrero et al., EPL 85, 57 009 (2009).
- Competition between Kondo and J
  - K. Haule and G. Kotliar, Phys. Rev. B 76, 104509 (2007).
  - M. Ferrero et al., Europhys. Lett. 85, 57 009 (2009).
  - K. Haule and G. Kotliar, Phys. Rev. B 76, 092503 (2007).

# Previous cluster results at finite doping



N. S. Vidhyadhiraja, A. Macridin, C. Sen,  
M. Jarrell, and M. Ma,  
Phys. Rev. Lett. **102**, 206407 (2009).



$J = 0.3 (U \sim 13)$   
EDCA-NCA 2x2

K. Haule and G. Kotliar, Phys. Rev. B **76**, 092503 (2007)

# Previous results

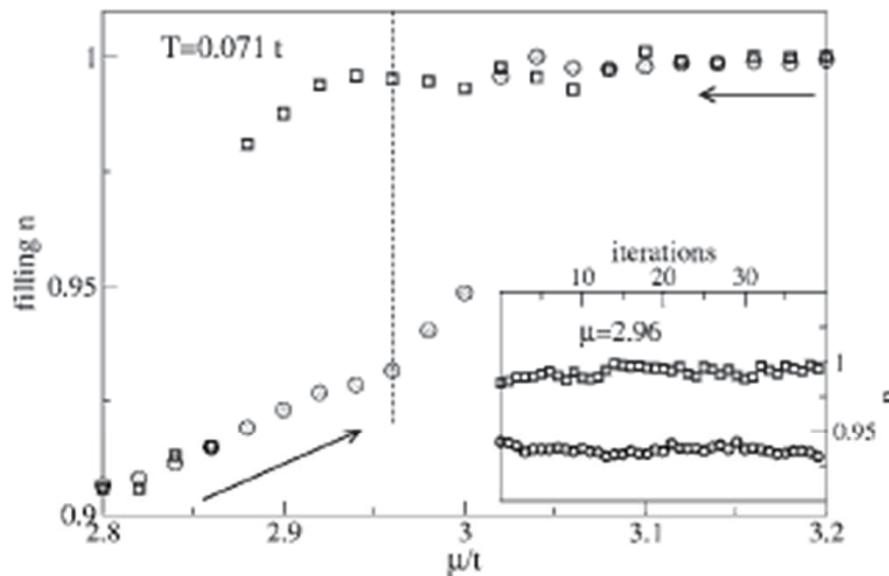
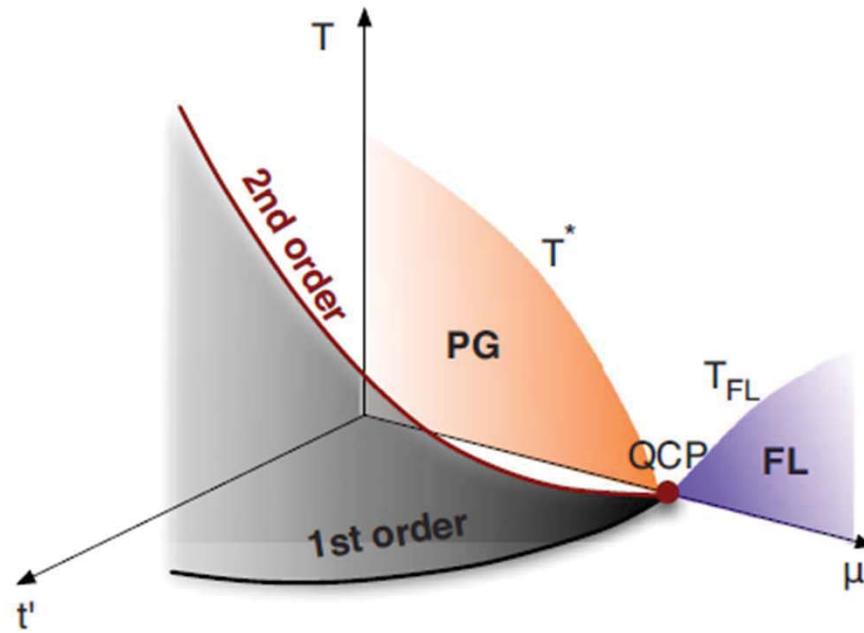


FIG. 2.  $N_c=8$  results. Filling  $n$  versus chemical potential below  $T_c$ , at  $T=0.071t$ . Two solutions describing a hysteresis are found: one incompressible with  $n \approx 1$  (squares) and a doped one (circles). Inset: stability of the two solutions versus DCA iterations when  $\mu = 2.96t$  (middle of the hysteresis, corresponding to the dotted line in the main figure).

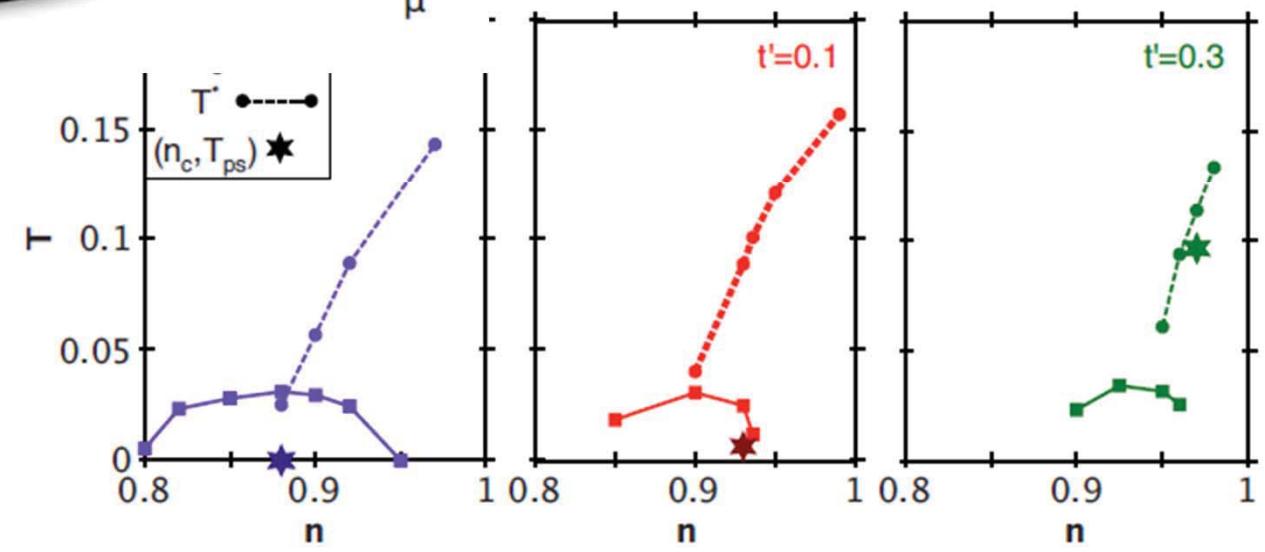
A. Macridin, M. Jarrell, and T. Maier,  
Phys. Rev. B **74**, 085104 (2006)

# Phase separation on electron-doped side

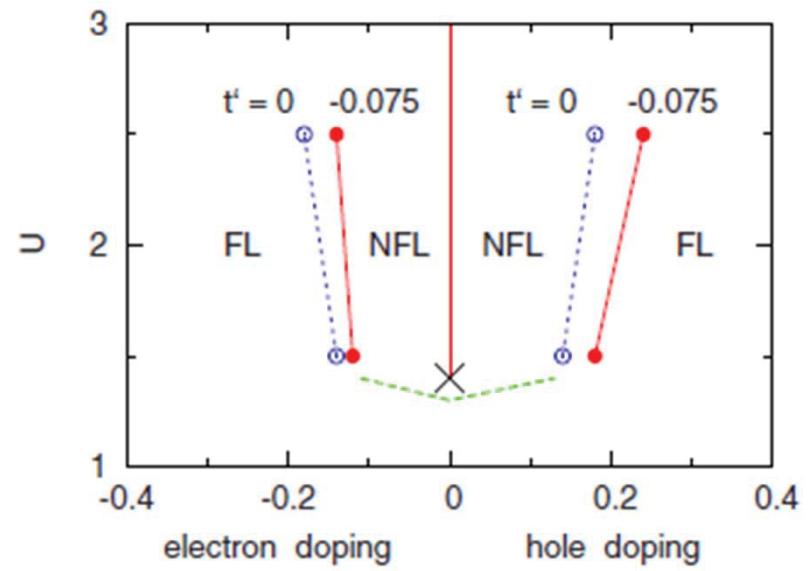
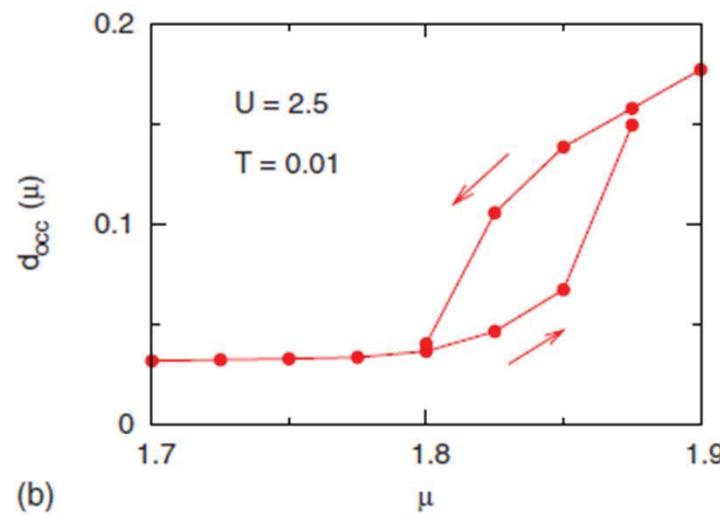


$U=8, N_c = 8, DCA$

E. Khatami,  
K. Mikelsons,  
D. Galanakis,  
A. Macridin,  
J. Moreno,  
R. T. Scalettar, and  
M. Jarrell  
PRB 81, 201101(R)  
2010



# Crossovers and transition

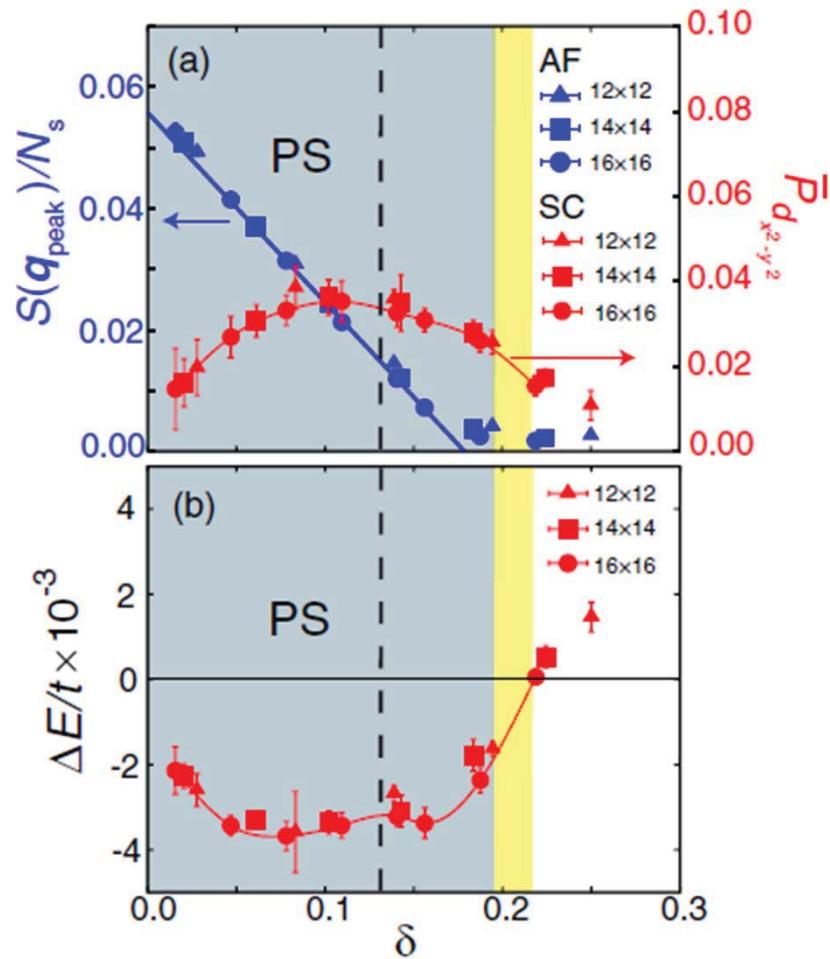


A. Liebsch, N.H. Tong, PRB **80**, 165126 (2009)



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# Variational Monte Carlo



T. Misawa M. Imada PRB **90**, 115137 (2014)



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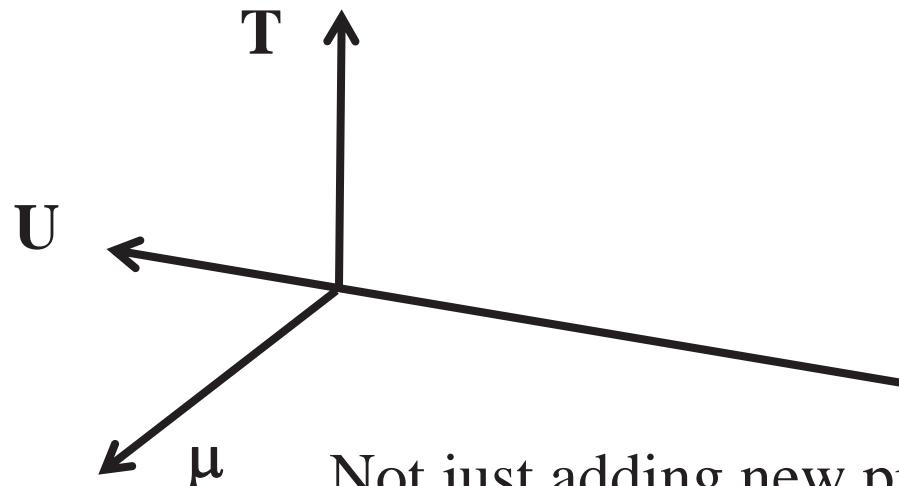


Giovanni Sordi

G. Sordi, K. Haule, A.-M.S.T  
PRL, **104**, 226402 (2010)  
and

Phys. Rev. B, **84**, 075161 (2011)

## Doping-induced Mott transition ( $t'=0$ )



Not just adding new piece:

Lesson from DMFT, first order transition + critical  
point governs finite  $T$  phase diagram

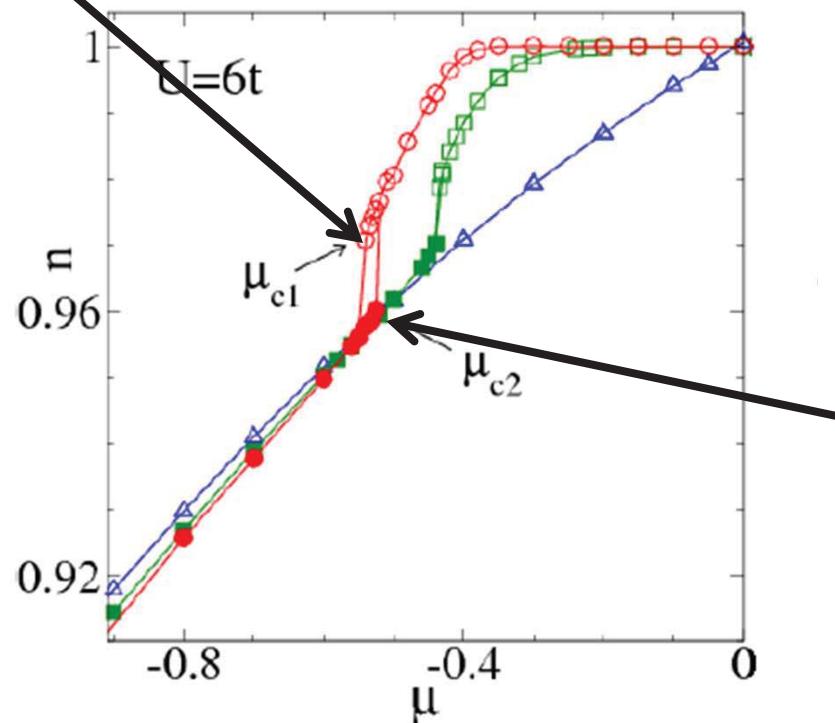


Kristjan Haule

# First order transition at finite doping

Spinodals

$t' = 0$



$n(\mu)$  for several temperatures:  
 $T/t = 1/10, 1/25, 1/50$

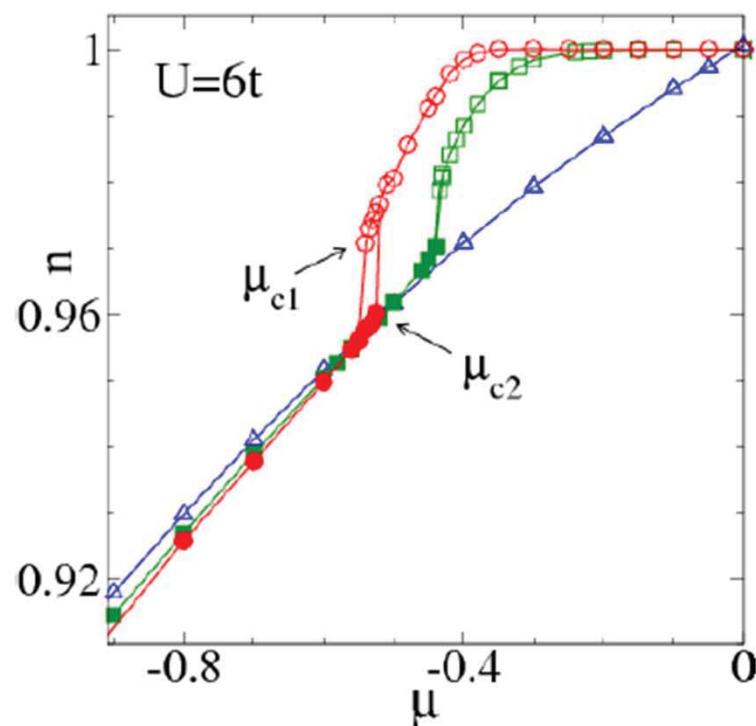
Sordi et al. PRL 2010, PRB 2011



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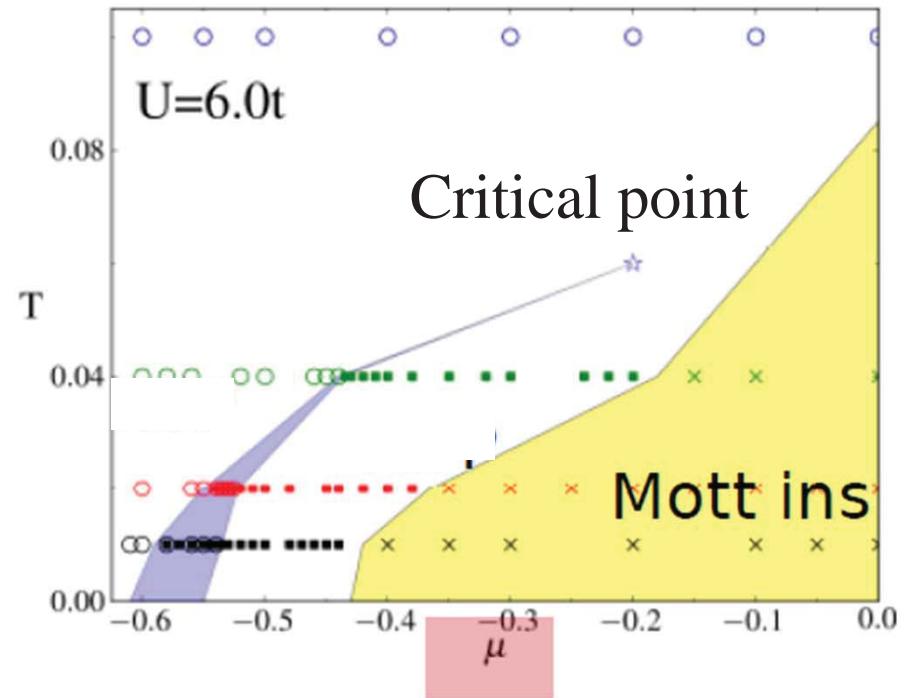
# First order transition at finite doping

$$t' = 0$$



$n(\mu)$  for several temperatures:  
 $T/t = 1/10, 1/25, 1/50$

Sordi et al. PRL 2010, PRB 2011

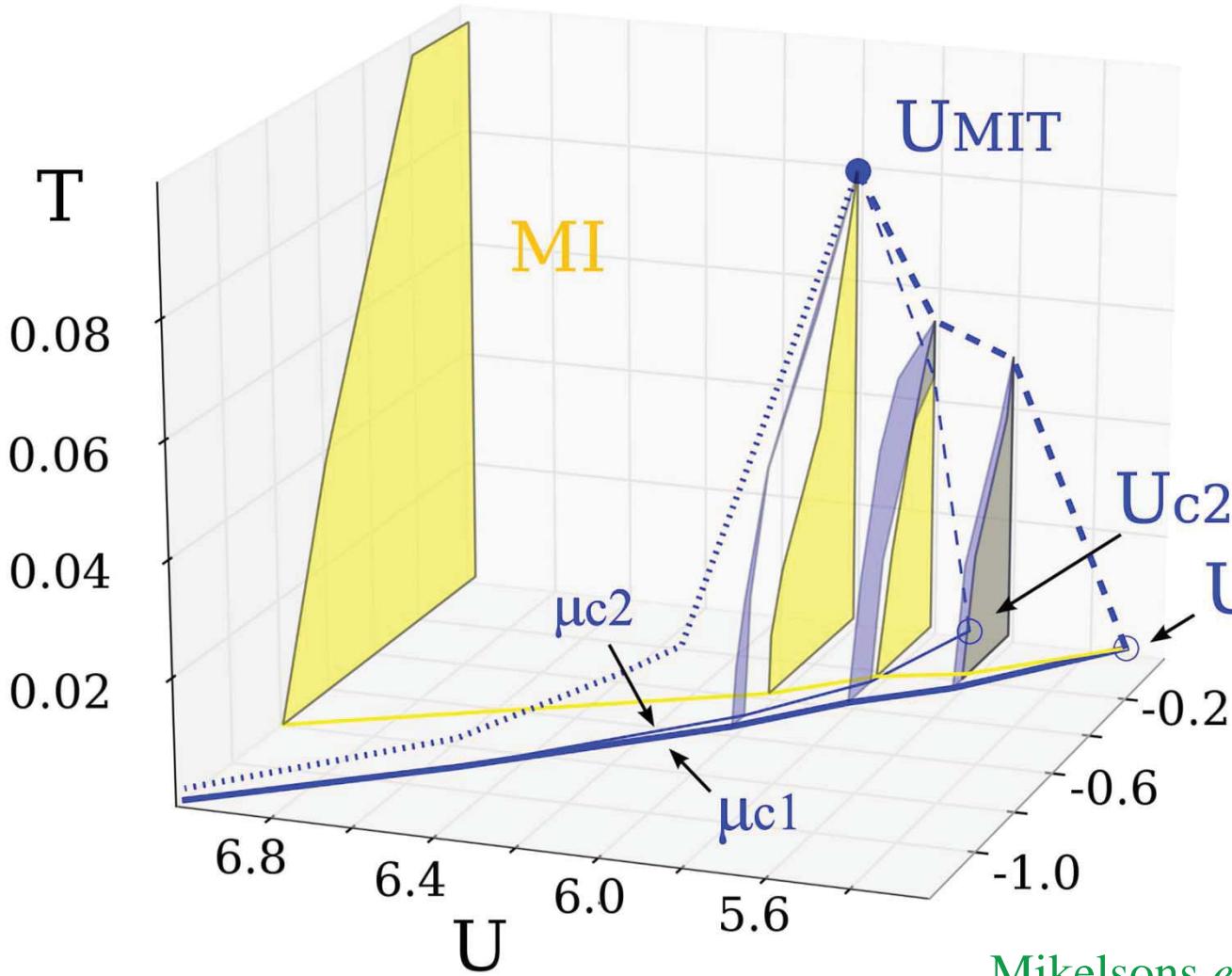


Hysteretic behavior:  
fingerprint first order  
transition!



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# Overall phase diagram



Clausius-Clapeyron

$$\left( \frac{dT_c}{d\mu_c} \right)_U = \frac{(n_1 - n_2)}{(S_2 - S_1)}$$

$$\left( \frac{dU_c}{d\mu_c} \right)_T = \frac{(n_1 - n_2)}{(D_1 - D_2)}$$

Underdoped phase  
lower  $D$  and  $S$   
 $U_{c1}$  More compressible

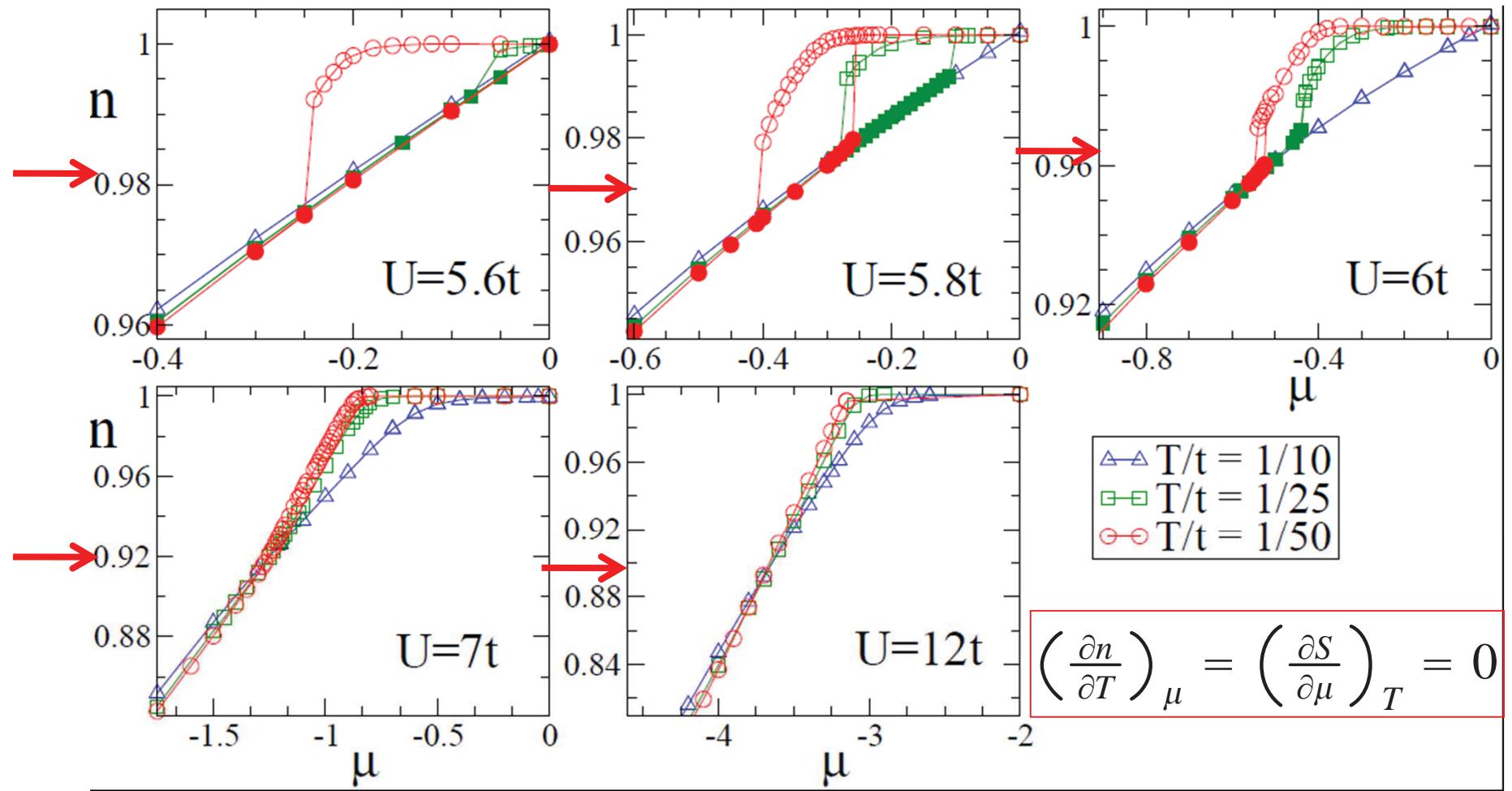
$$\left( \frac{\partial S}{\partial n} \right)_T = 0$$

Mikelsons *et al.* PRB **80**, (2009)



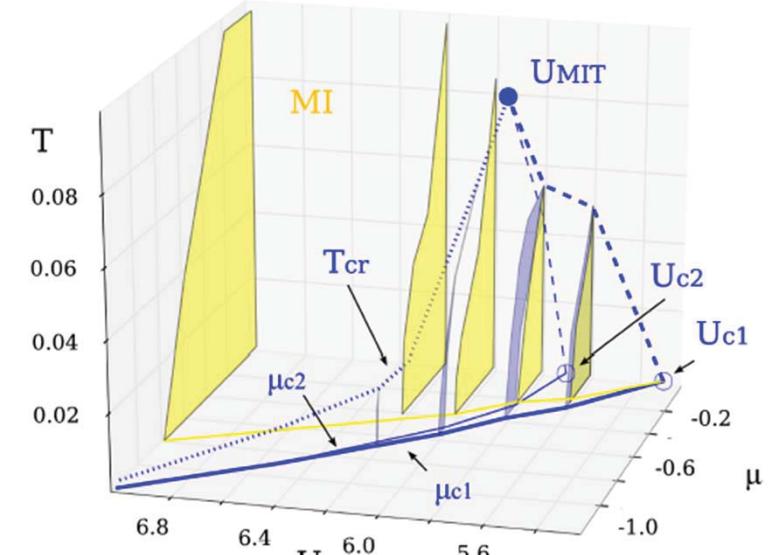
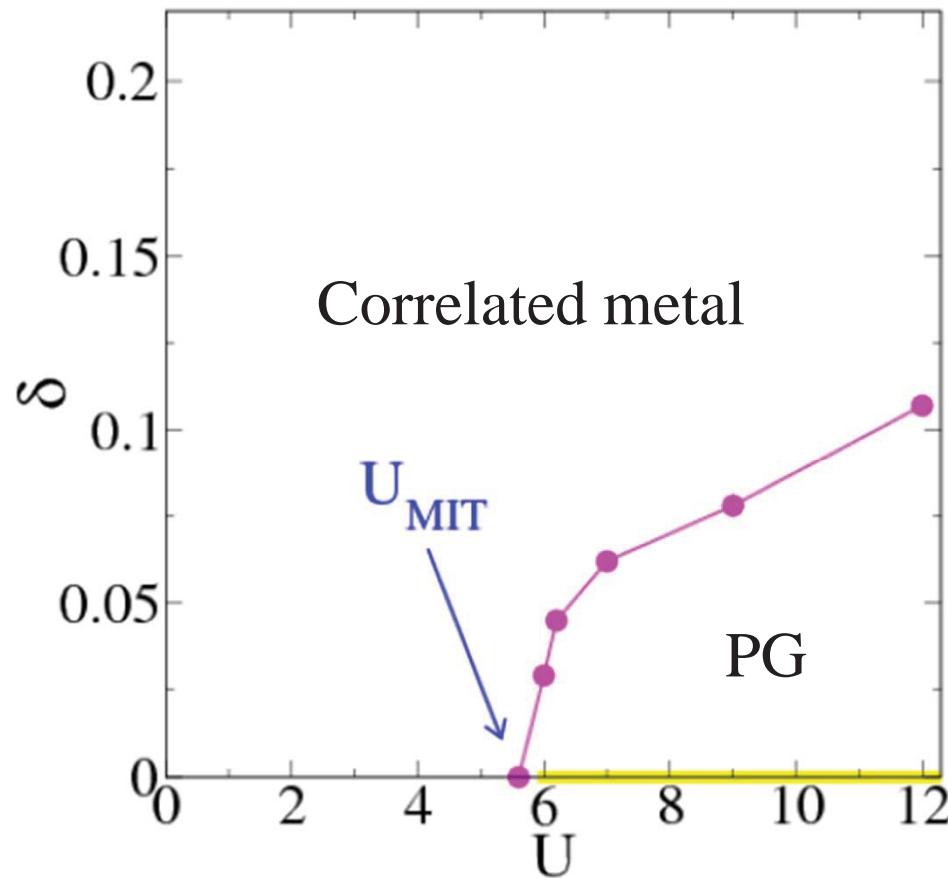
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# Critical doping as a function of $U$ increases



# A finite-doping first order transition, linked to Mott transition up to optimal doping

Doping dependence of critical point as a function of  $U$

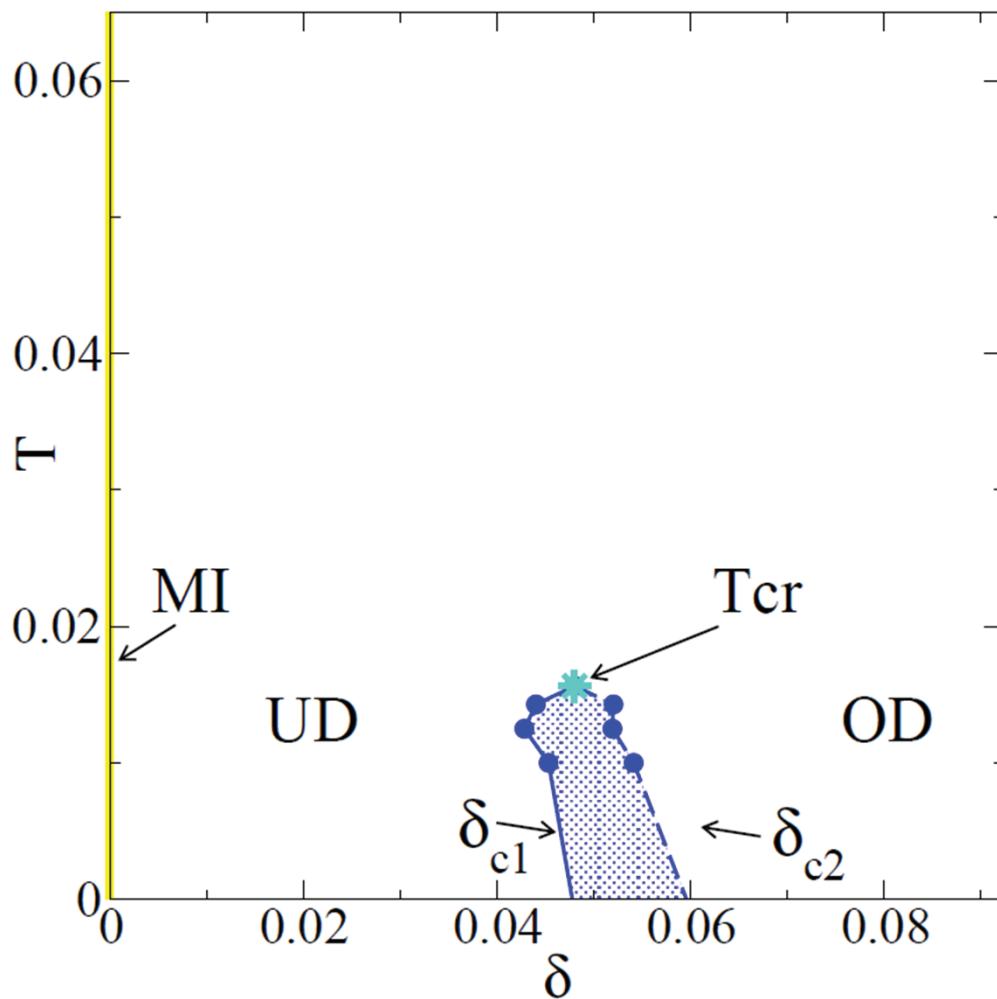


Sordi et al. PRL 2010, PRB 2011



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# Characterisation of the phases ( $U=6.2t$ )



$U > U_{MIT}$ :

1. Mott insulator (MI)
2. Underdoped phase (UD):  
 $\delta < \delta_c$
3. Overdoped phase (OD):  
 $\delta > \delta_c$
4. Coexistence/forbidden region

Here “optimal doping”  $\delta_c$  = doping at which the 1st order transition occurs

How does the UD phase differ from the OD phase?



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Giovanni Sordi



Patrick Sémon



Kristjan Haule

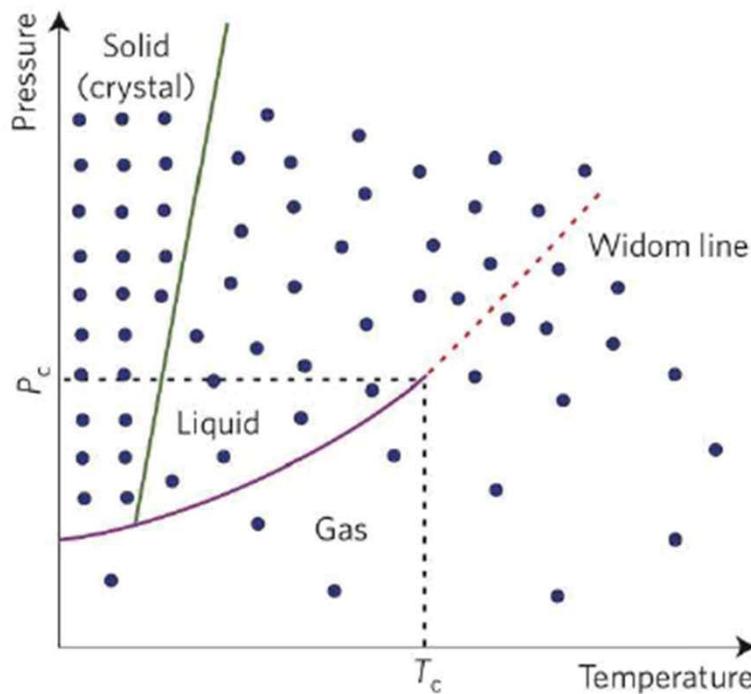
## The Widom line ( $t' = 0$ )

G. Sordi, *et al.* Scientific Reports 2, 547 (2012)



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# What is the Widom line?



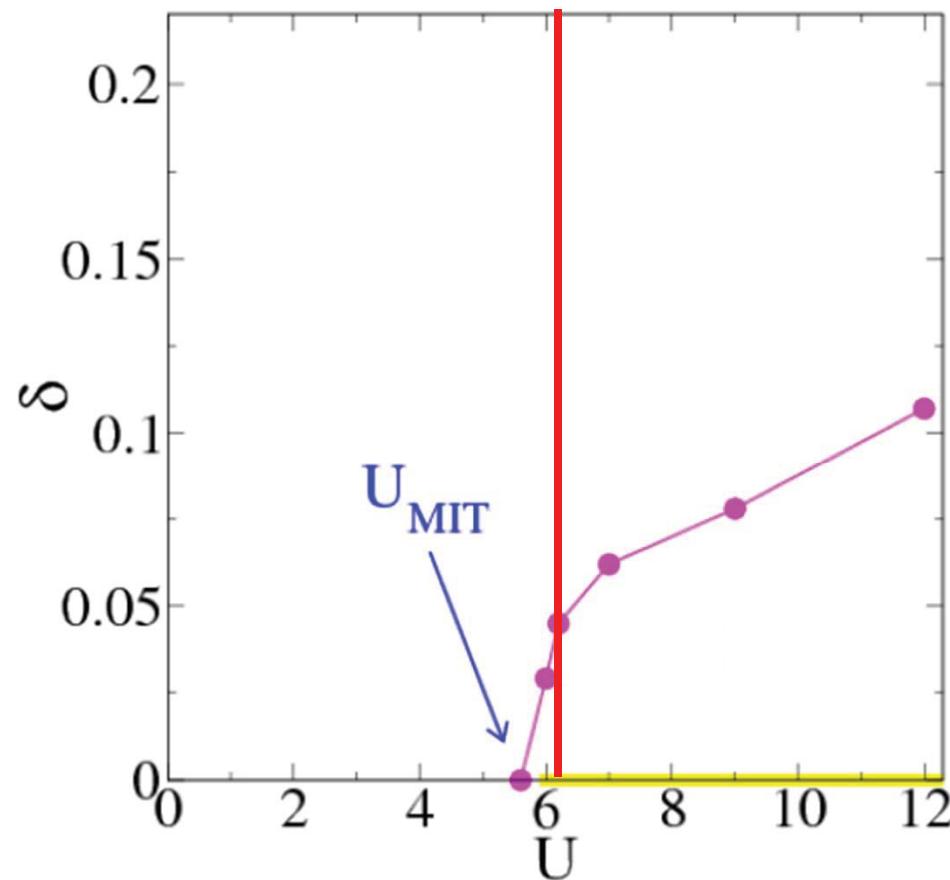
McMillan and Stanley, Nat Phys 2010

- ▶ it is the continuation of the coexistence line in the supercritical region
- ▶ line where the **maxima of different response functions** touch each other asymptotically as  $T \rightarrow T_p$
- ▶ liquid-gas transition in water: max in isobaric heat capacity  $C_p$ , isothermal compressibility, isobaric heat expansion, etc

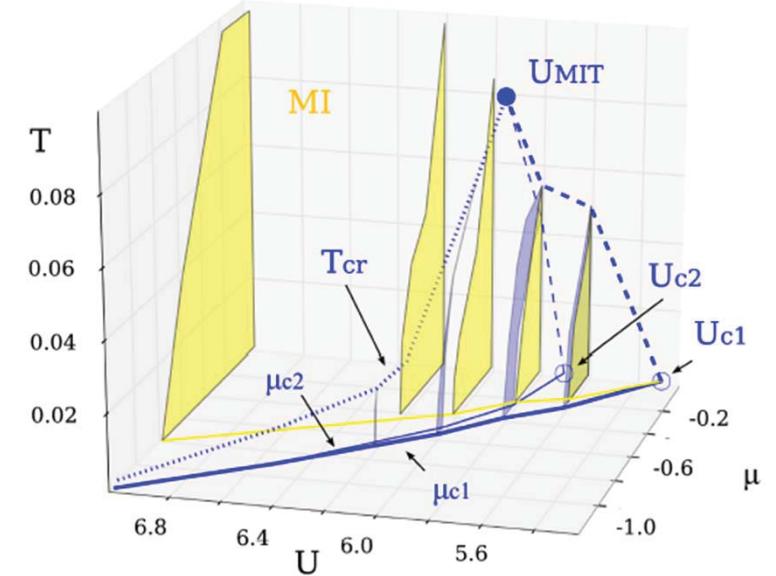
- ▶ **DYNAMIC crossover arises from crossing the Widom line!**  
water: Xu et al, PNAS 2005,  
Simeoni et al Nat Phys 2010

# Link to Mott transition up to optimal doping

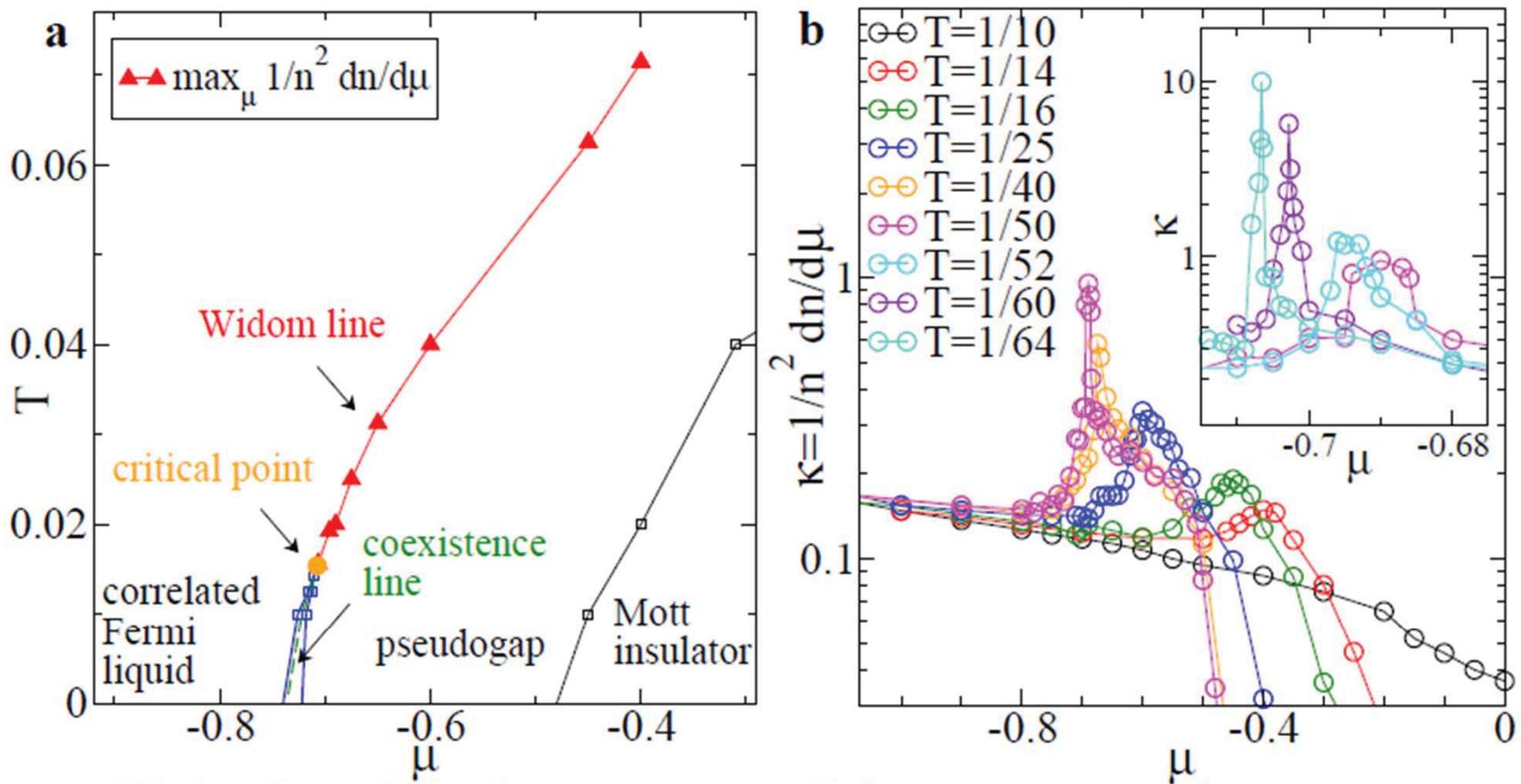
Doping dependence of critical point as a function of  $U$



Smaller  $D$  and  $S$



# Pseudogap $T^*$ along the Widom line



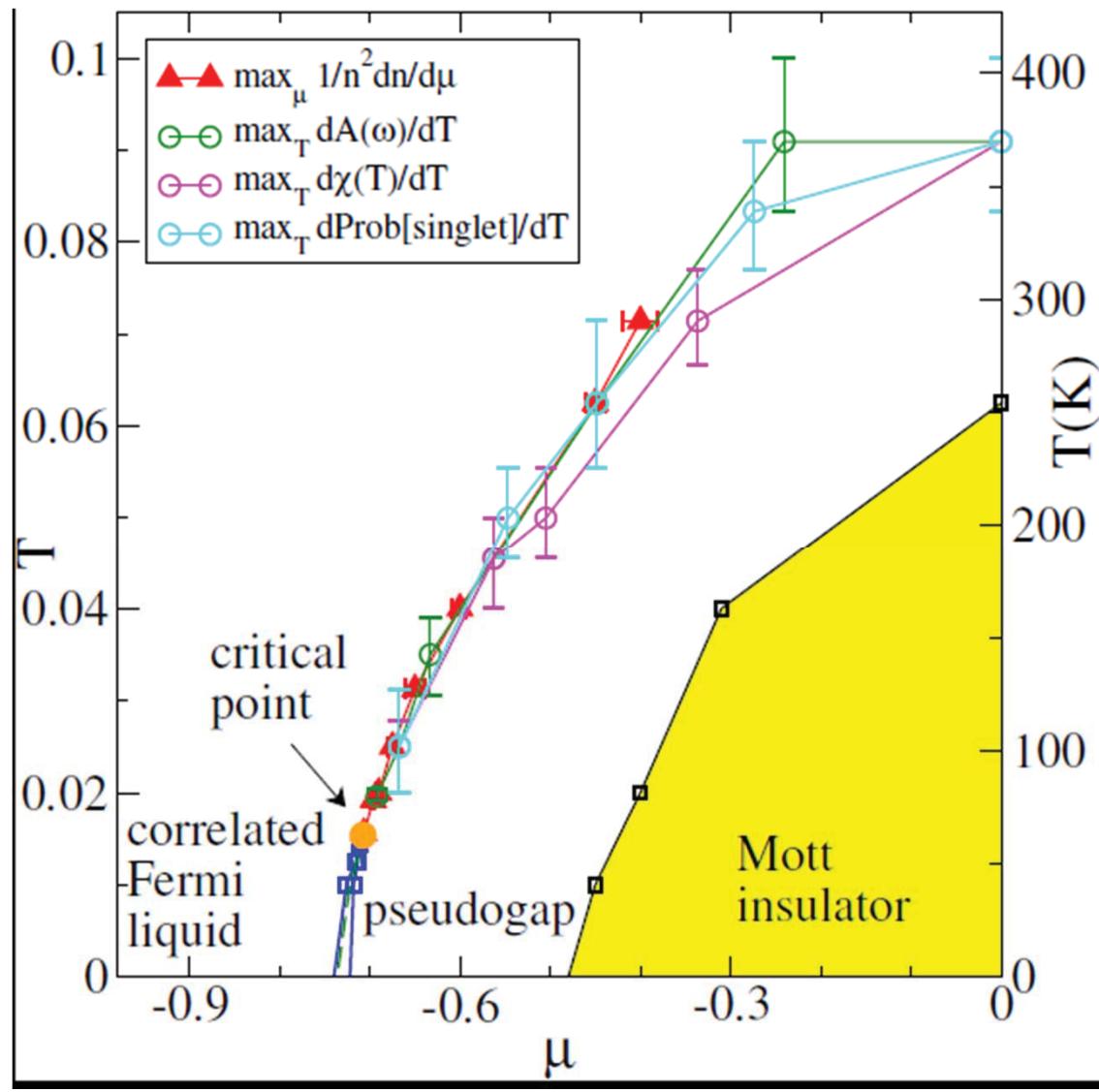
Widom line: defined from maxima of charge compressibility

$$\kappa = 1/n^2(dn/d\mu)_T$$

divergence of  $\kappa$  at the (classical) critical point!



# Rapid change also in dynamical quantities



Compare a few results for cuprates

Caveats:

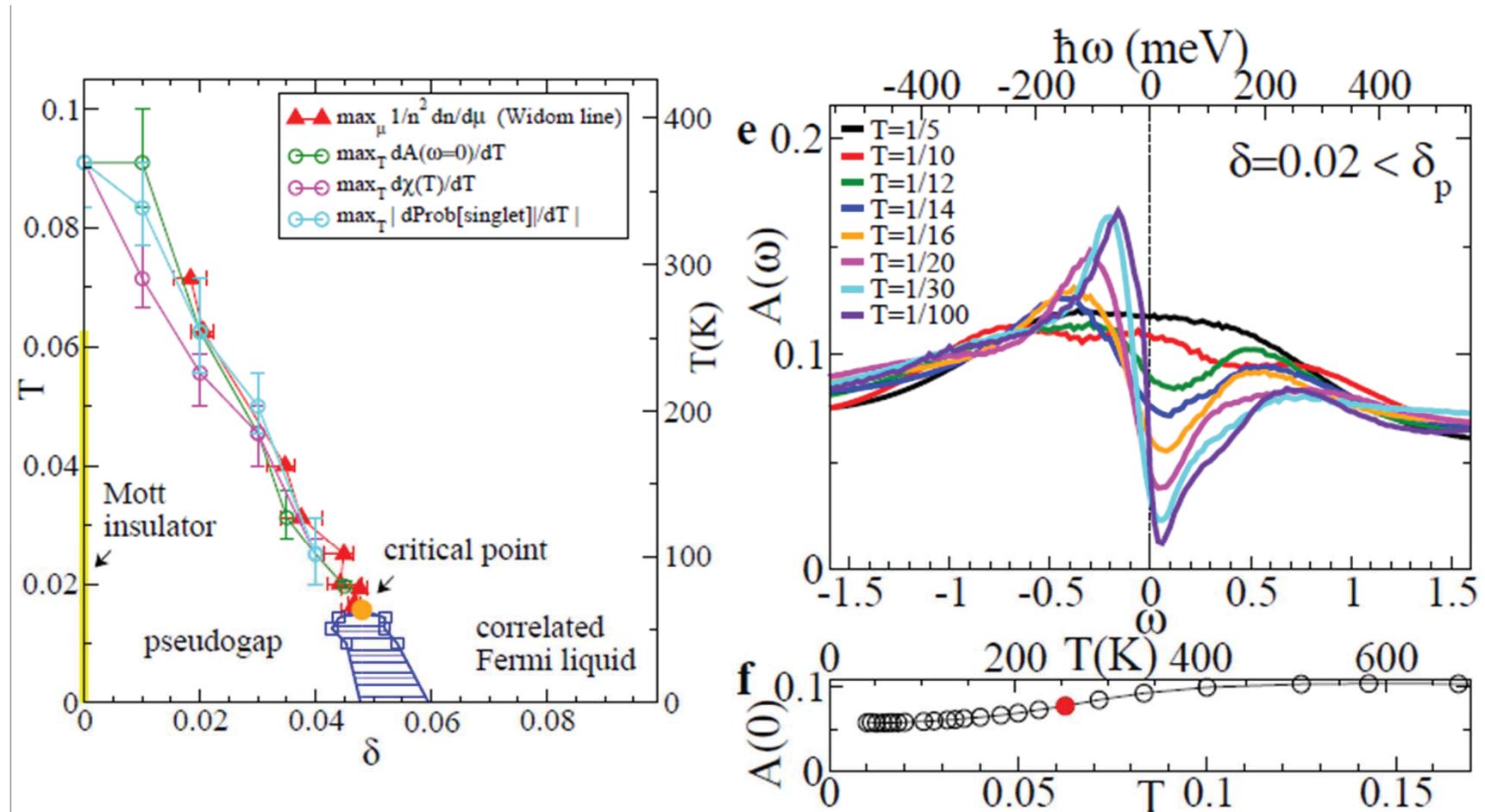
$U$  not large enough

$t' = 0$

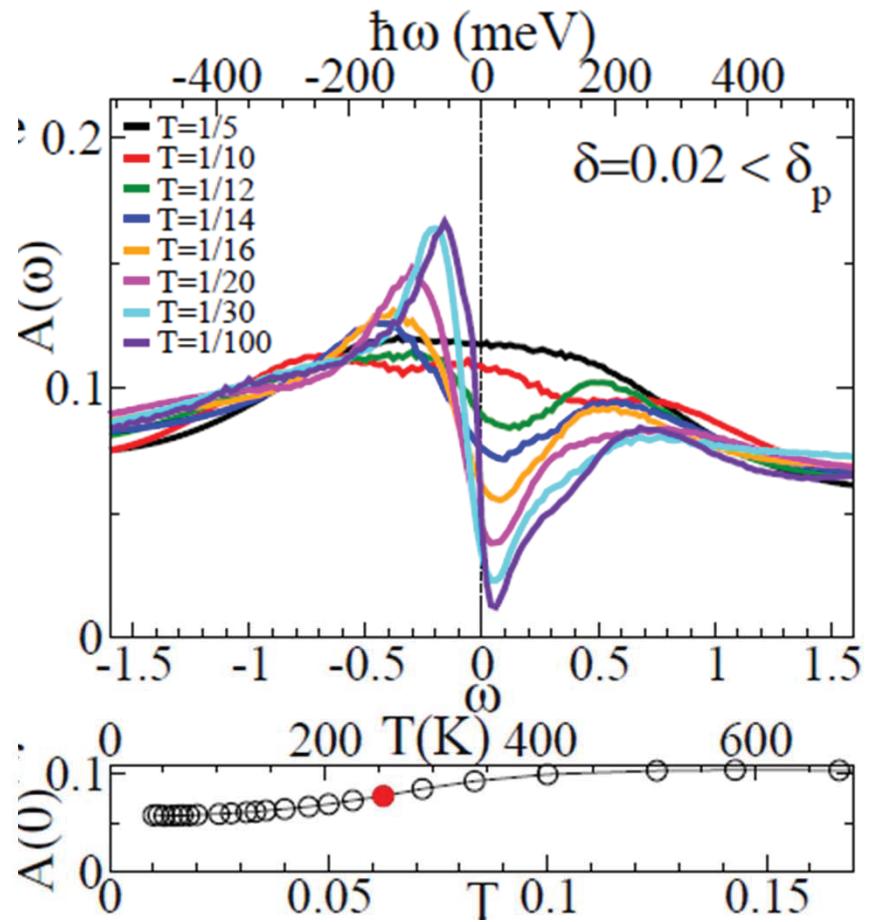
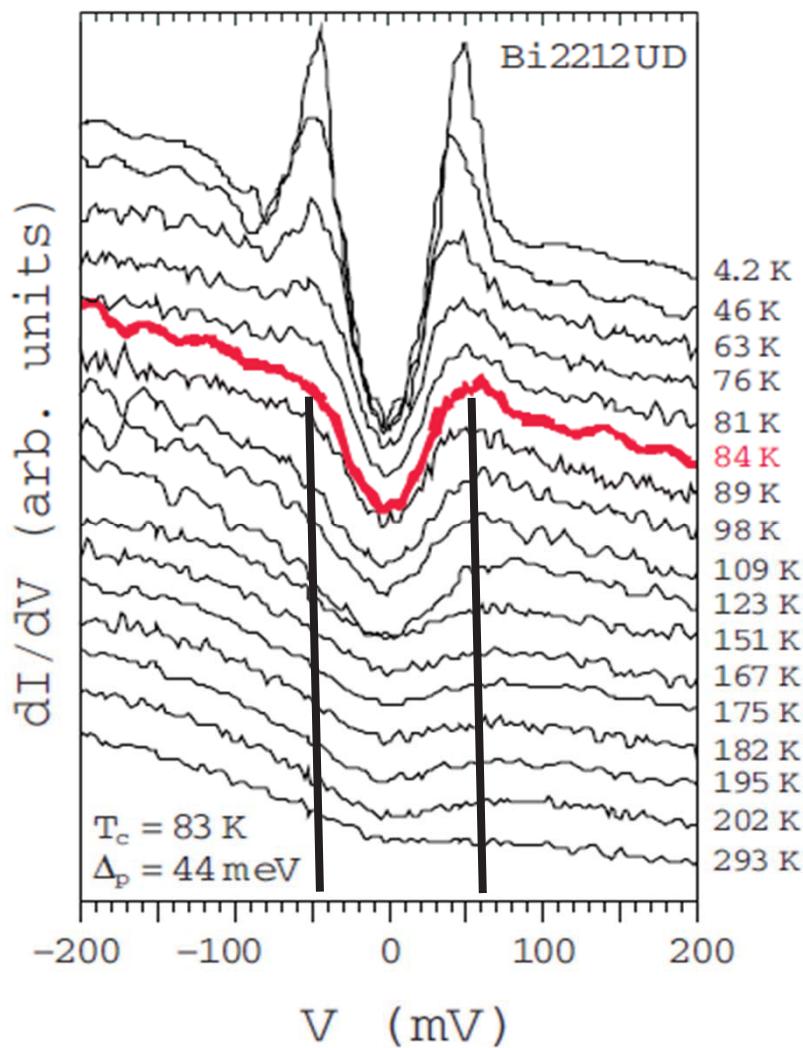


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# Density of states



# Density of states



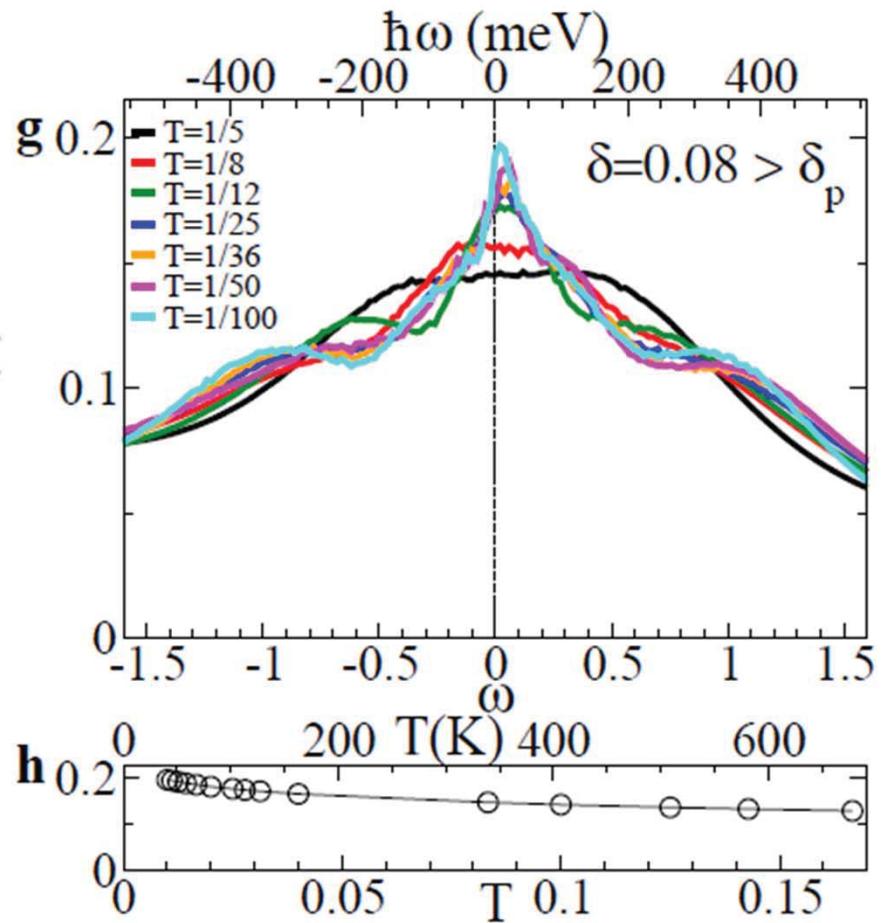
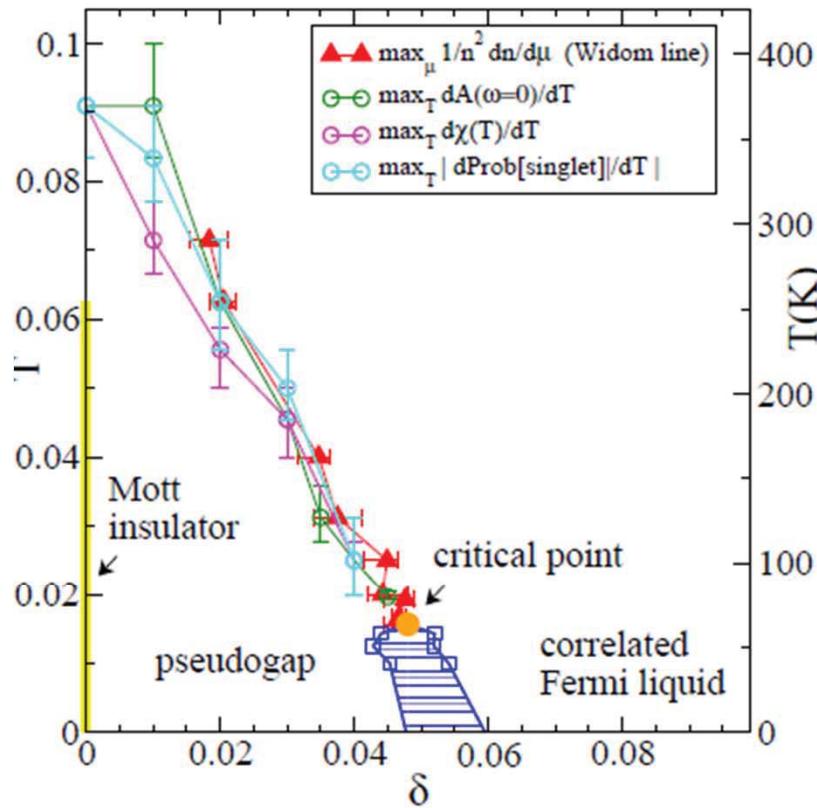
h. Renner, B. Revaz, J.-Y. Genoud,  
K. Kadowaki, and Ø. Fischer

PRL 80, 149 (1998)



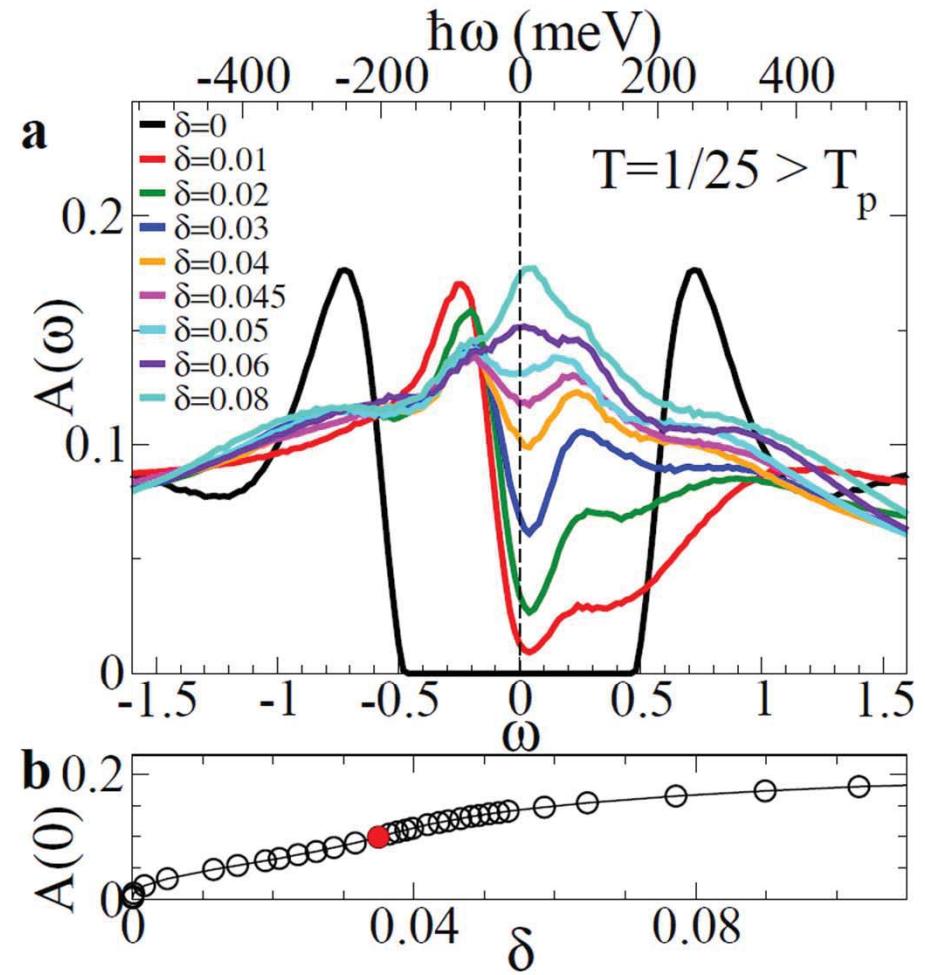
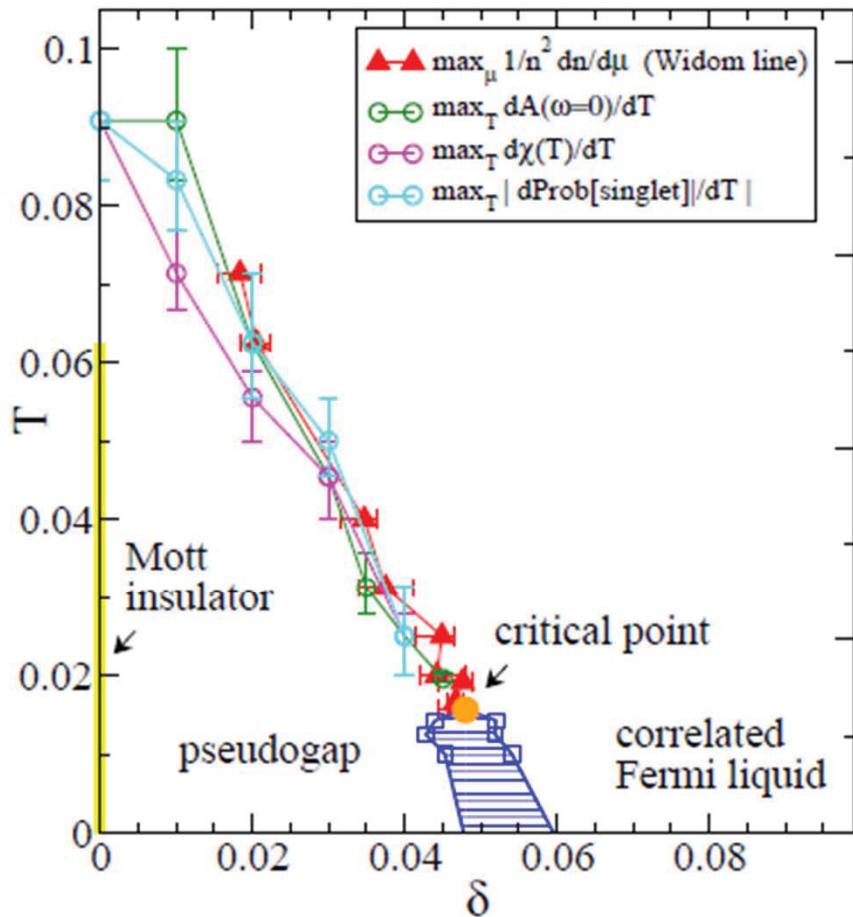
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# Density of states

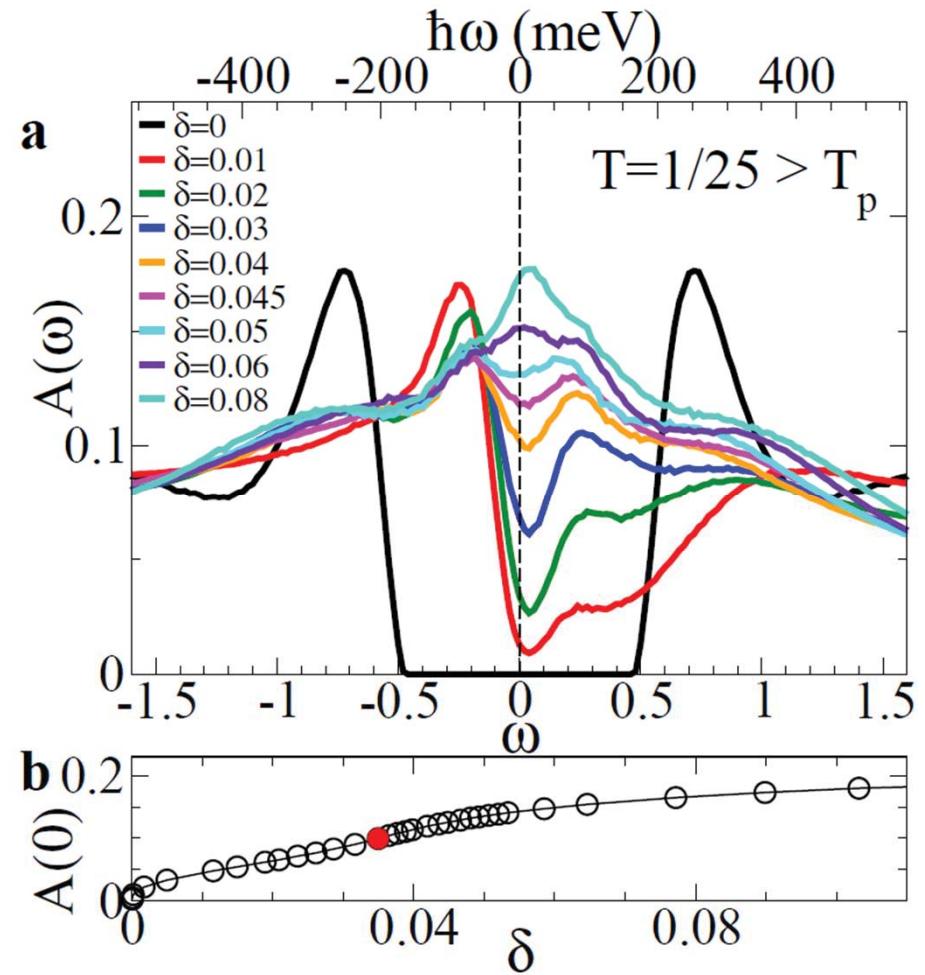
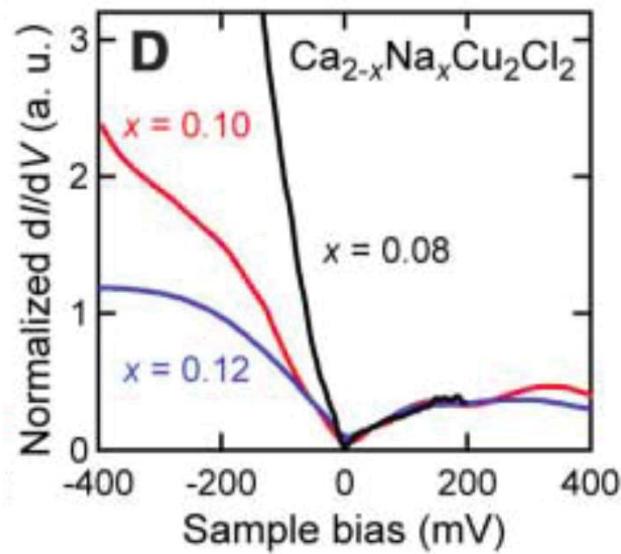


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# Density of states

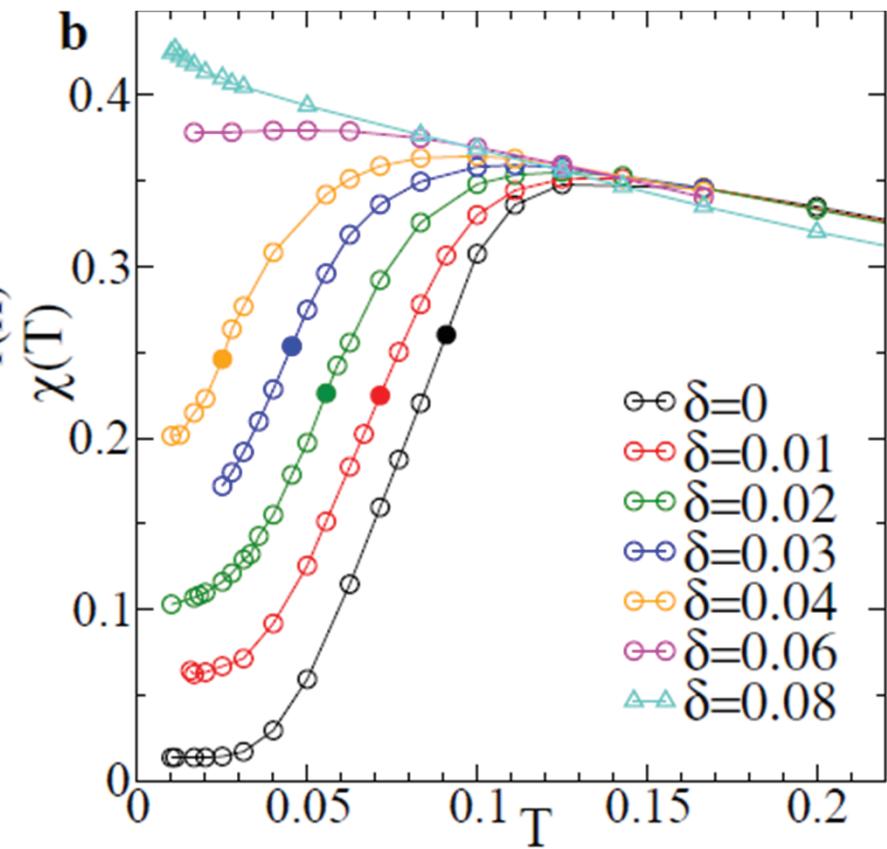
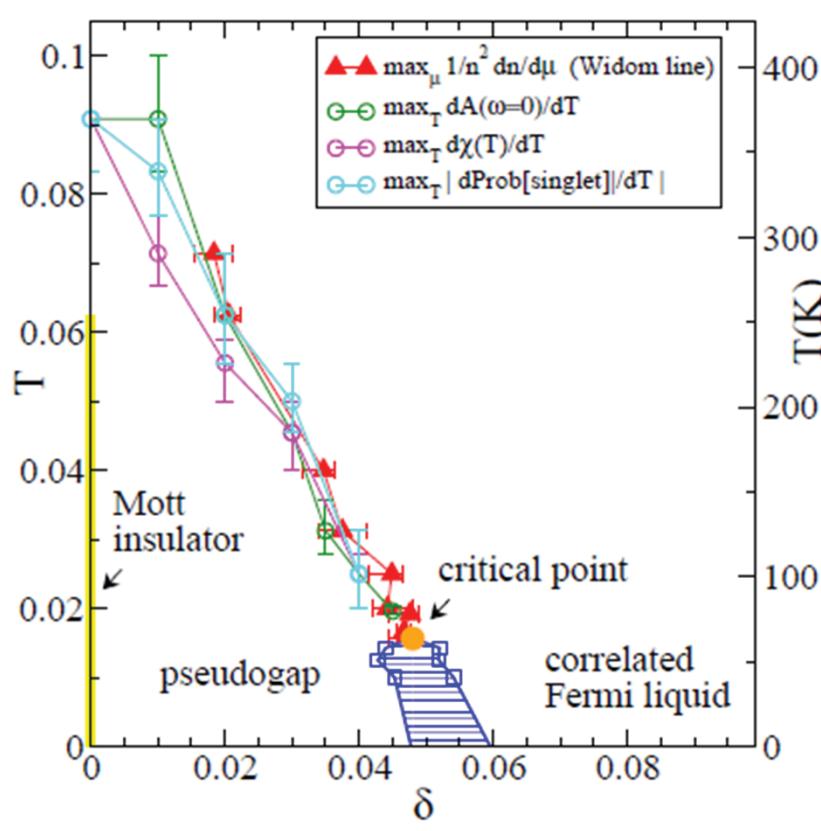


# Density of states

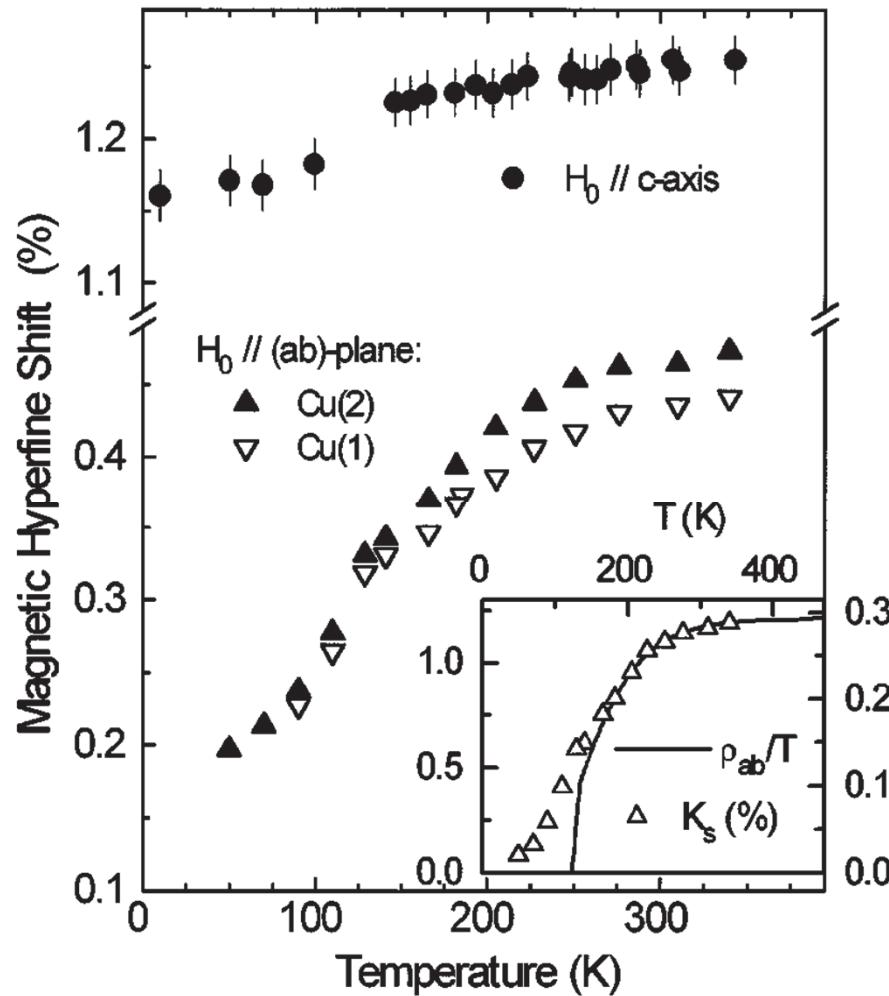


Khosaka et al. *Science* **315**, 1380 (2007);

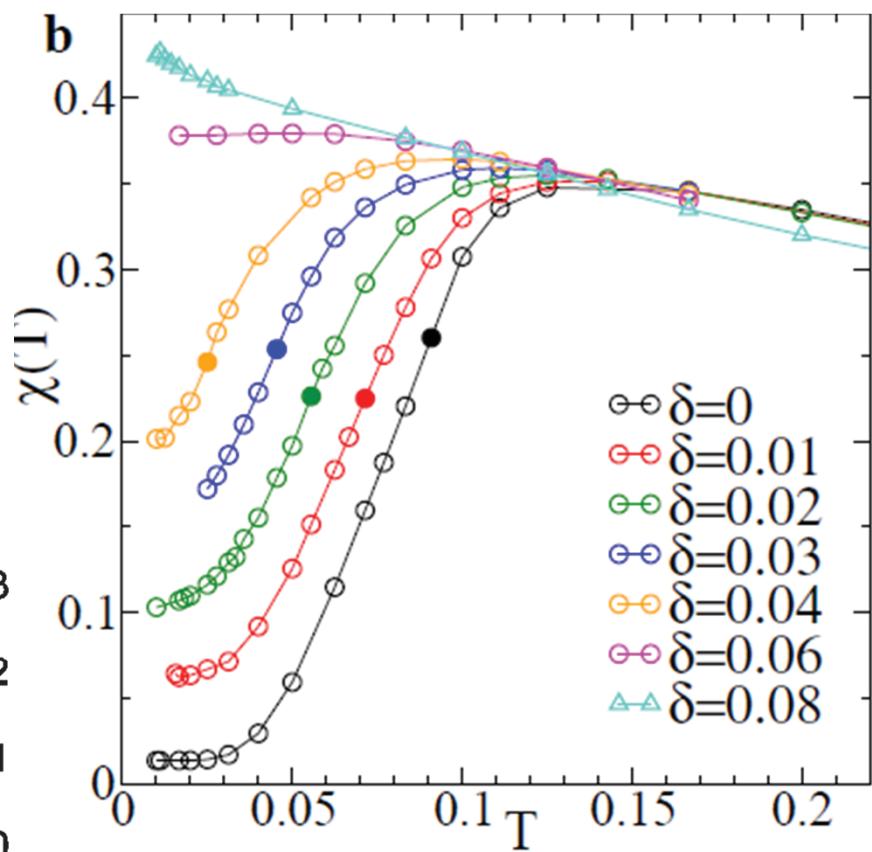
# Spin susceptibility



# Spin susceptibility



Underdoped Hg1223  
Julien et al. PRL 76, 4238 (1996)



# What is the minimal model?

H. Alloul arXiv:1302.3473  
C.R. Académie des Sciences, (2014)

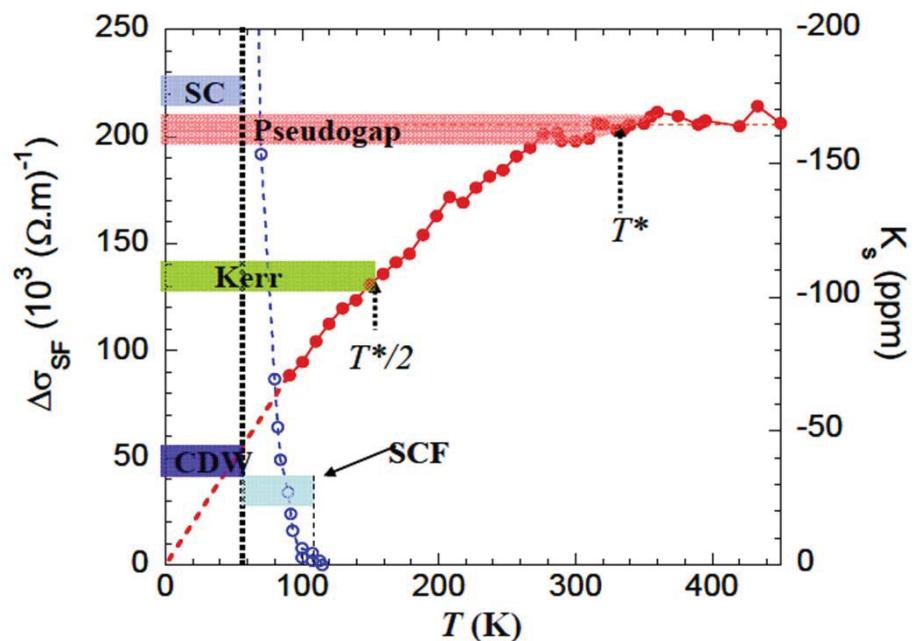


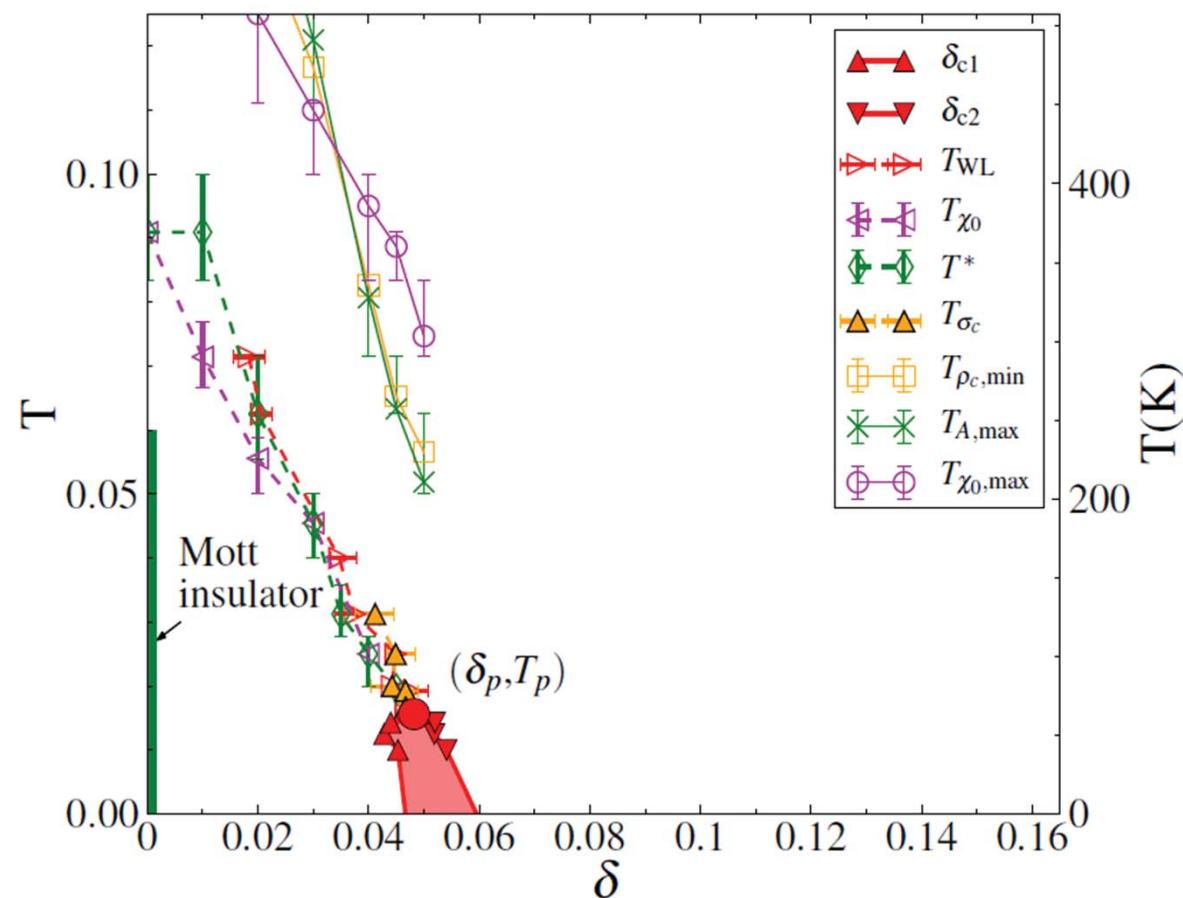
Fig 1 Spin contribution  $K_s$  to the  $^{89}\text{Y}$  NMR Knight shift [11] for  $\text{YBCO}_{6.6}$  permit to define the PG onset  $T^*$ . Here  $K_s$  is reduced by a factor two at  $T \sim T^*/2$ . The sharp drop of the SC fluctuation conductivity (SCF) is illustrated (left scale) [23]. We report as well the range over which a Kerr signal is detected [28], and that for which a CDW is evidenced in high fields from NMR quadrupole effects [33] and ultrasound velocity data [30]. (See text).



Giovanni Sordi



Patrick Sémon



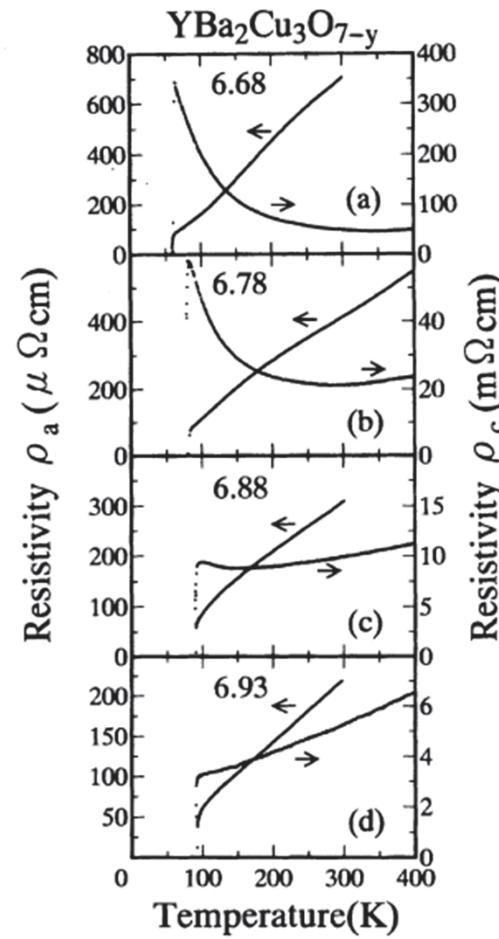
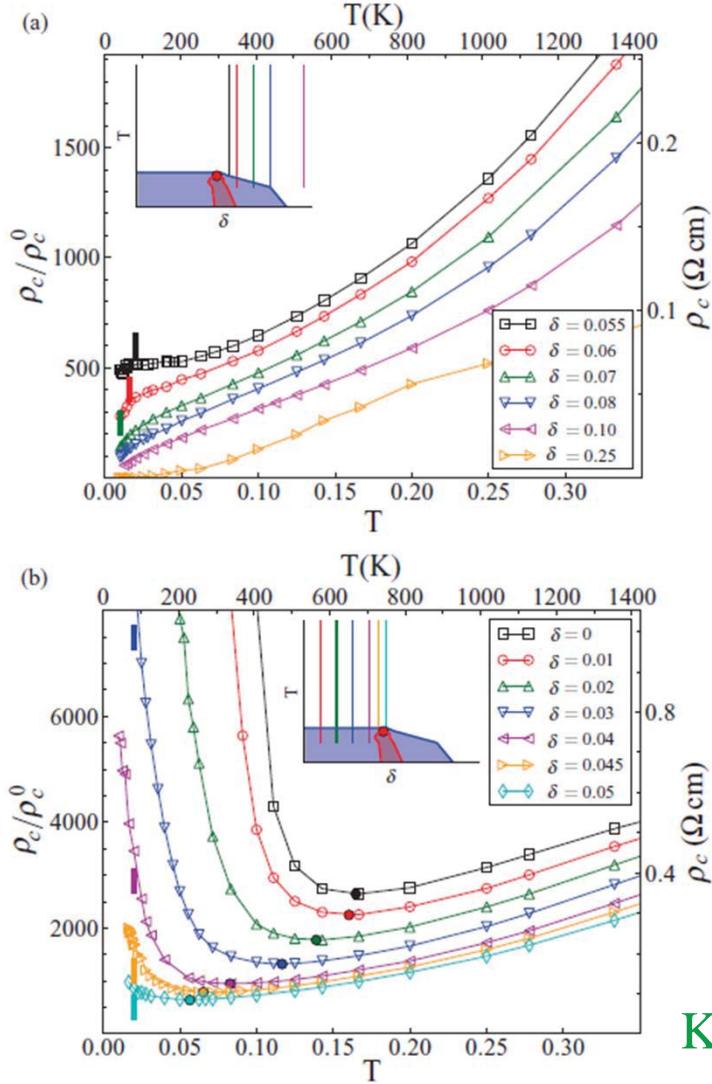
G. Sordi et al. Phys. Rev. Lett. 108, 216401/1-6 (2012)

P. Sémon, G. Sordi, A.-M.S.T., Phys. Rev. B **89**, 165113/1-6 (2014)



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# C-axis resistivity

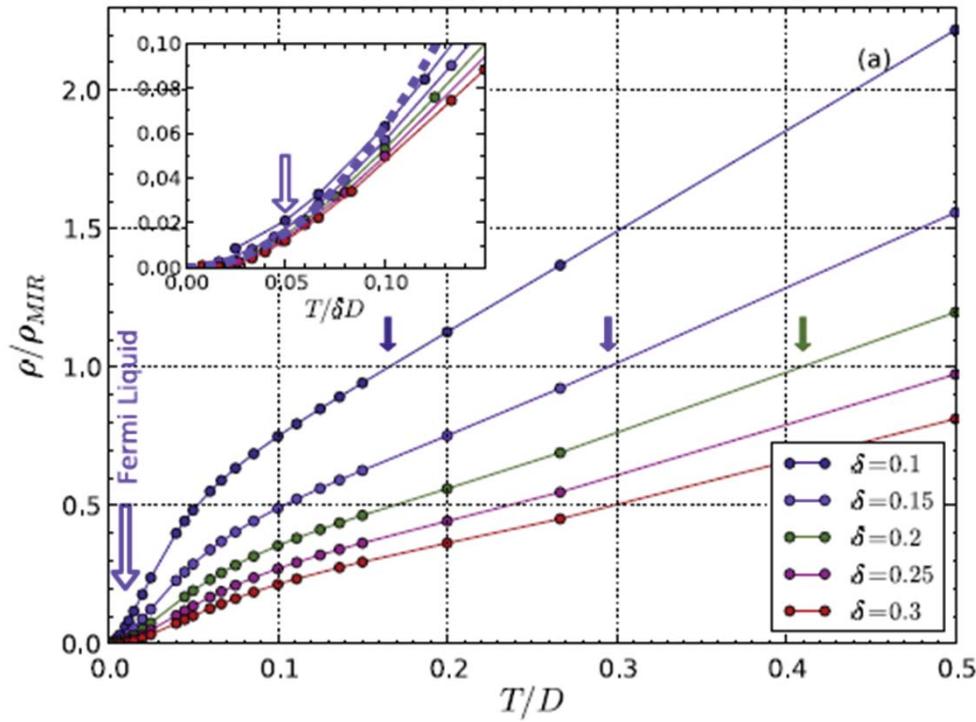


K. Takenaka, K. Mizuhashi, H. Takagi, and S. Uchida,  
Phys. Rev.B 50, 6534 (1994).



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# Mott-Ioffe-Regel limit



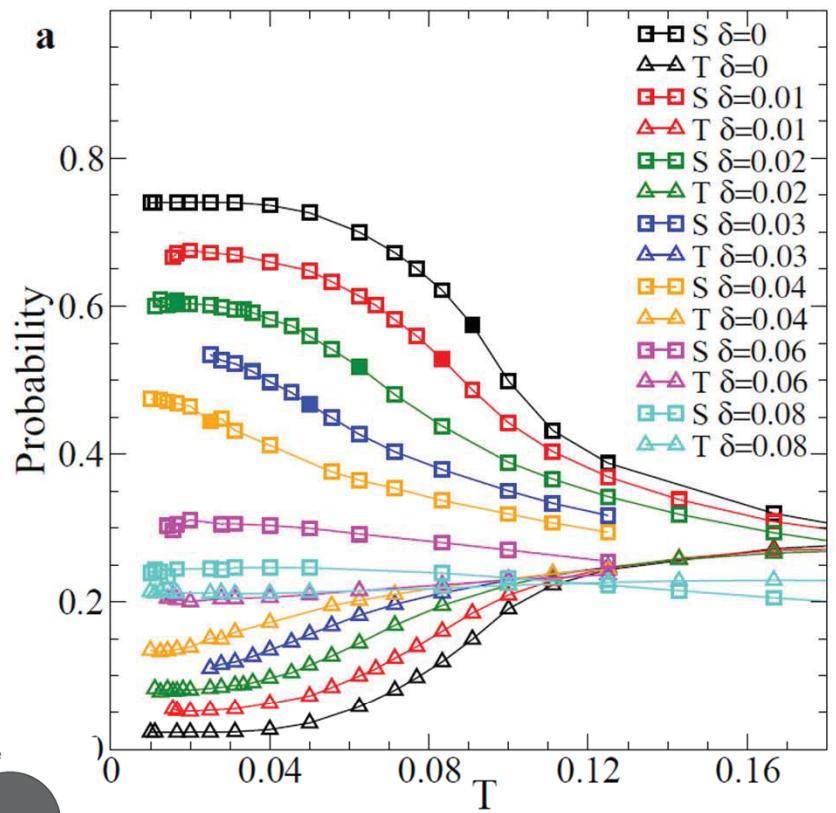
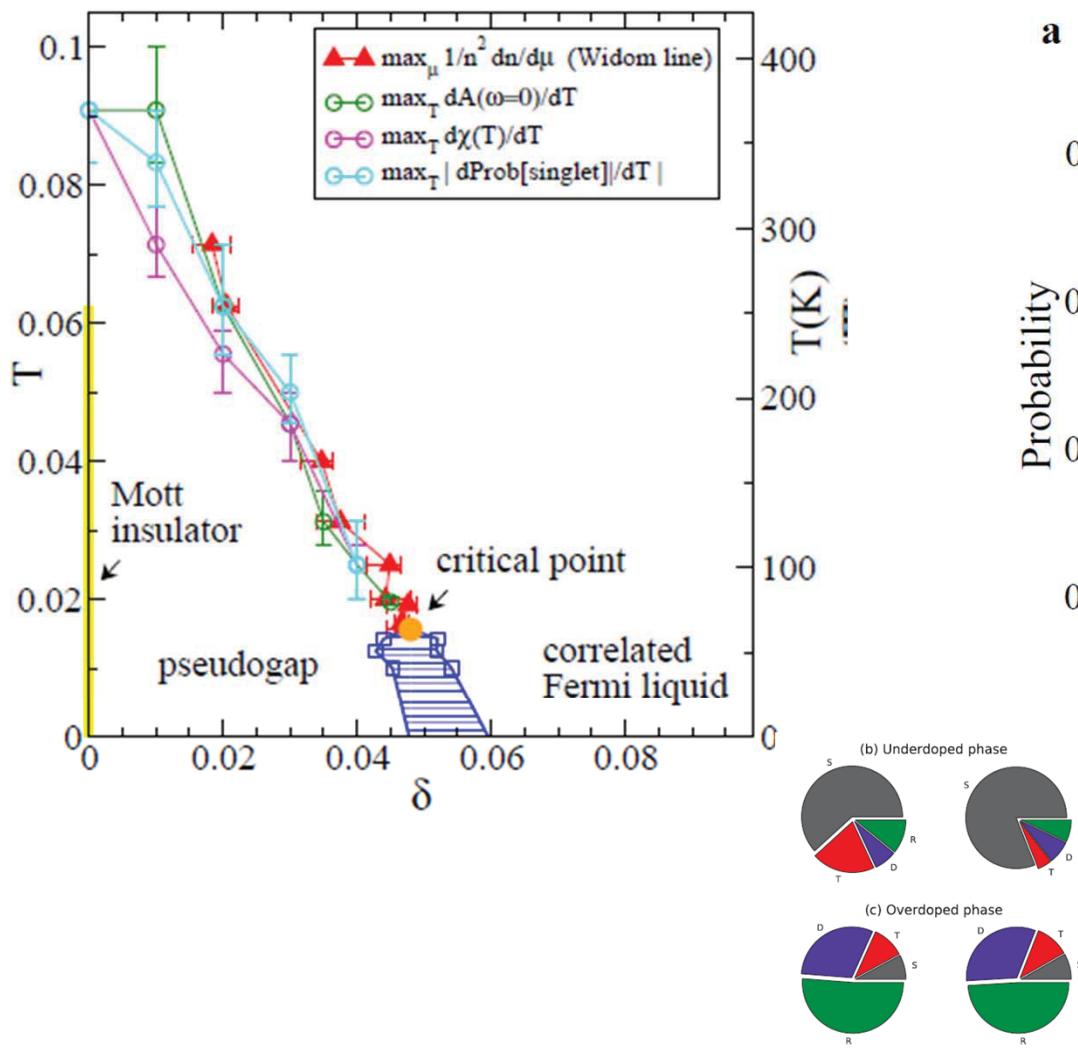
X. Deng, J.j Mravlje, R. Zitko, M. Ferrero, G. Kotliar, and A. Georges  
PRL 110, 086401 (2013)

# Physics



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# Plaquette eigenstates

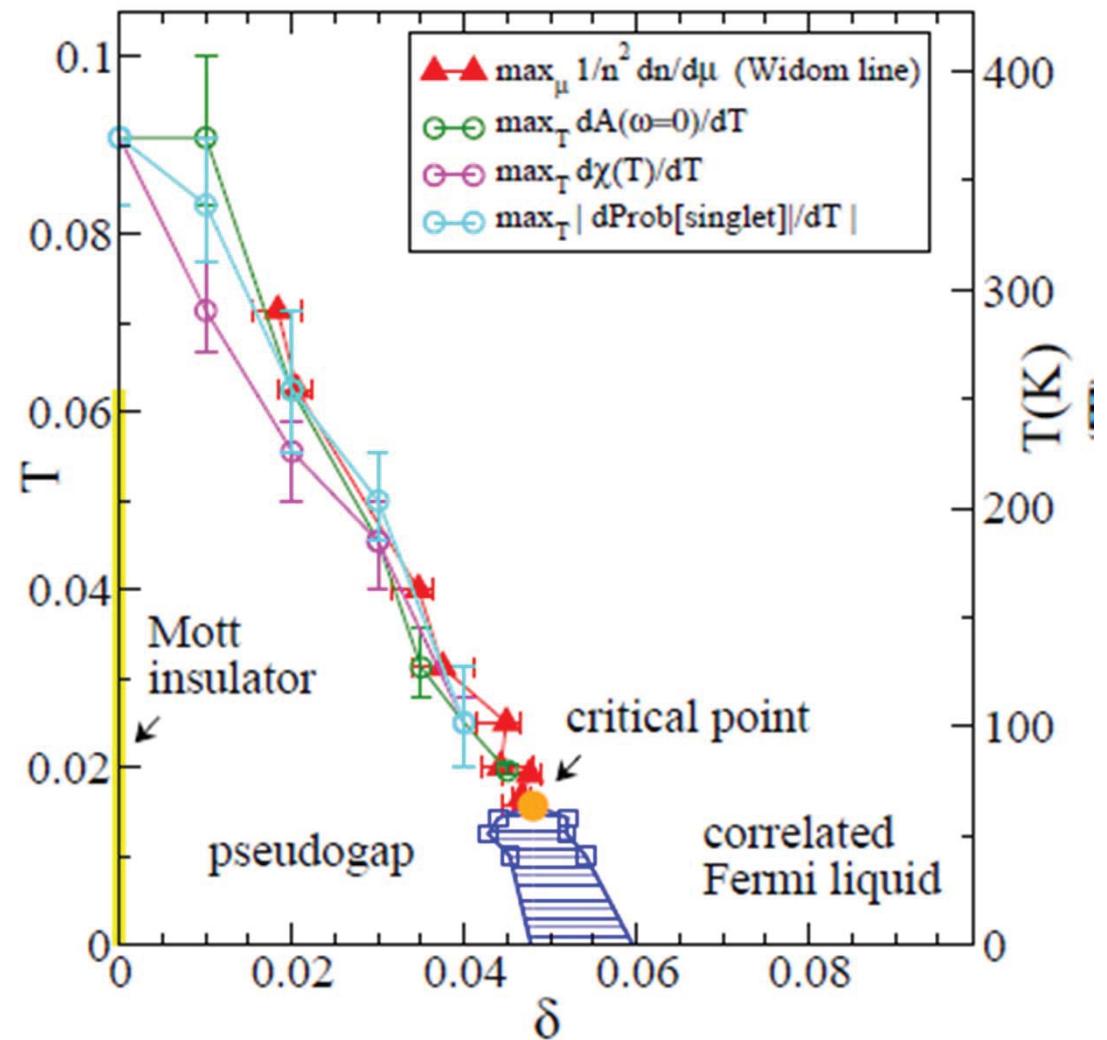


$$T = 1/10, 1/50$$

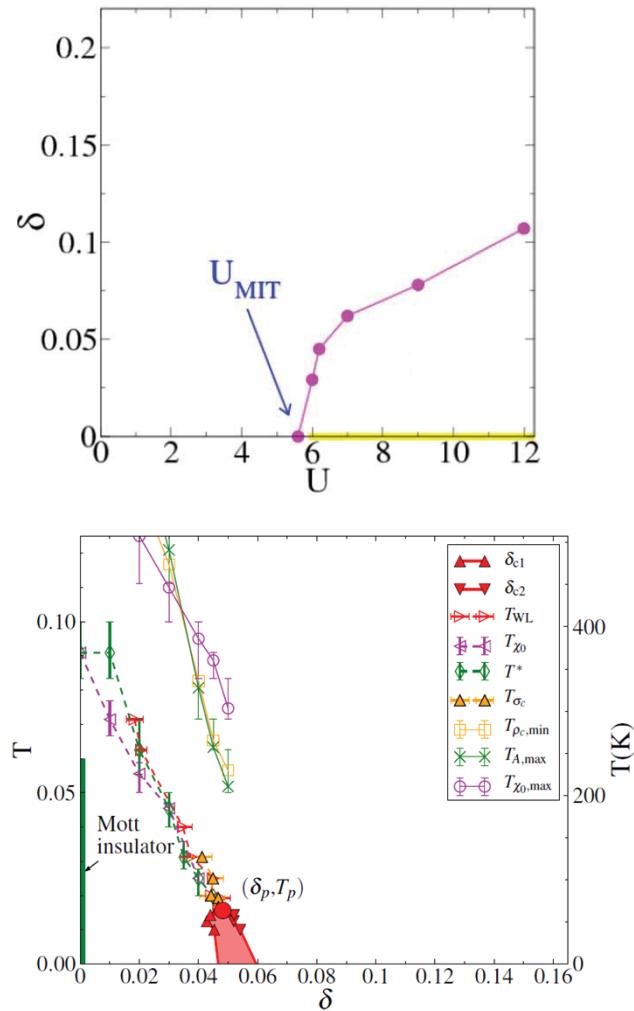


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# Pseudogap along the Widom line $T_W$

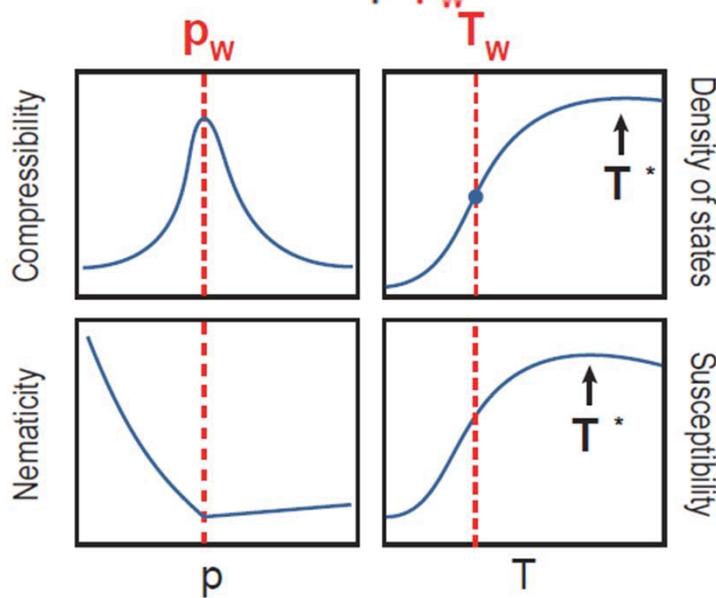
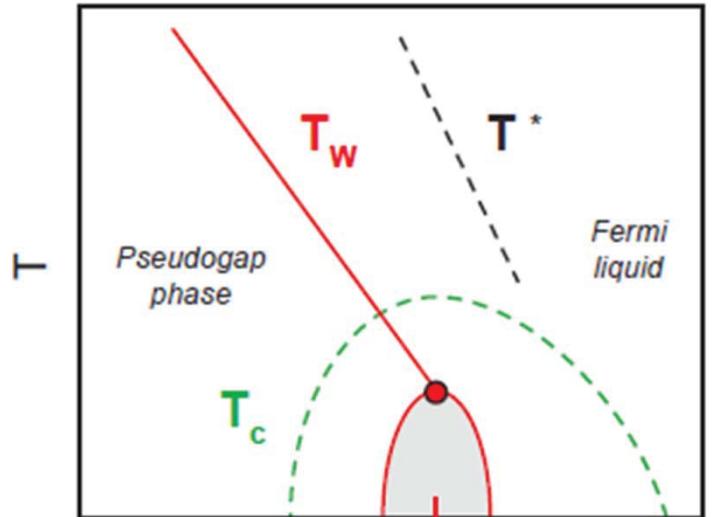


# Summary: normal state



- Signatures of Mott physics extend way beyond half-filling
- Pseudogap is a phase
- Pseudogap  $T^*$  controlled by a Widom line and its precursor
- High compressibility (stripes?)

# Organizing principle



- Is the pseudogap (PG) a crossover or a phase transition ?
- Relation between CDW and the PG ?
- Why CDW peaked at 12% doping ?
- Origin of nematicity ?
- Why superconducting ?
- Why a dome of SC ?
- Does a one-band model capture the key physics ?
- AFM QCP important?
- Lessons from other SC?



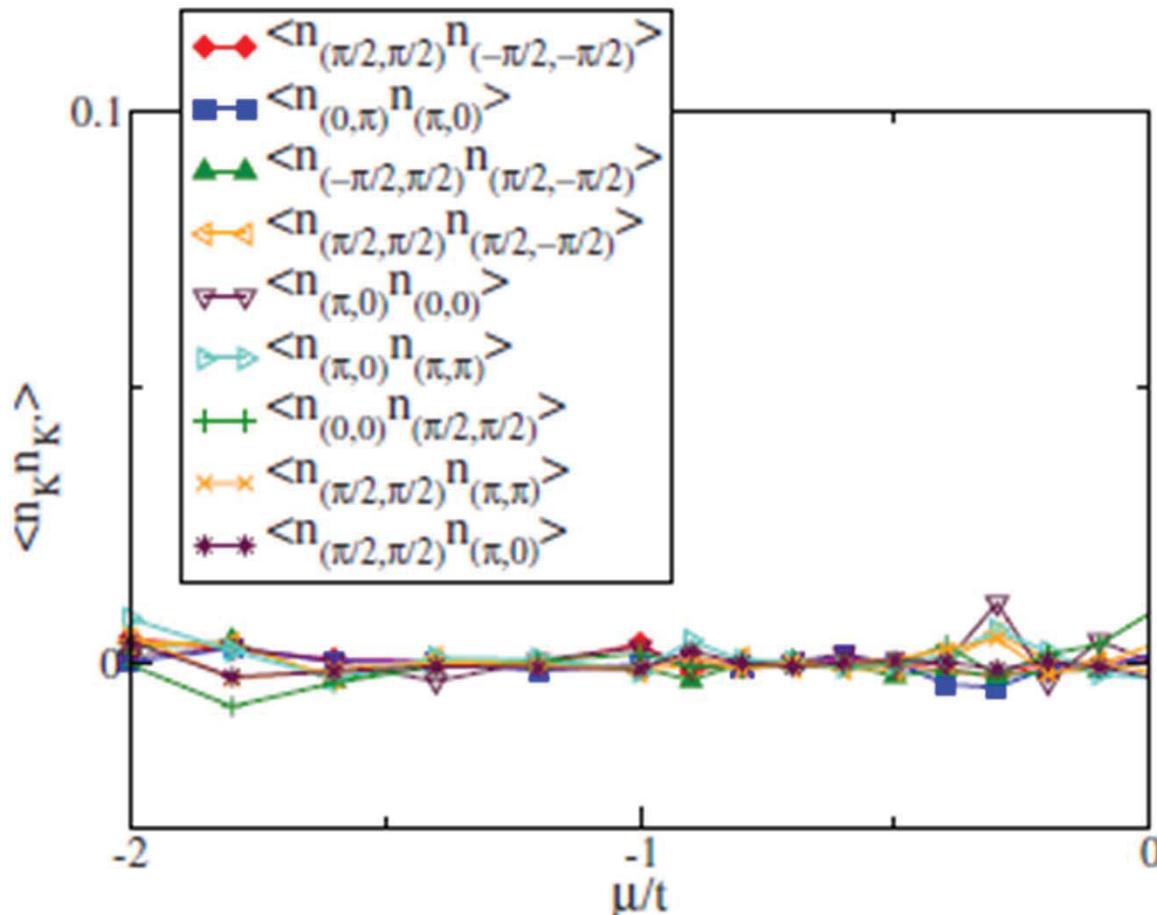
# Anisotropy (nematicity)

Normal state and large anisotropy  
in an *orthorhombic* crystal



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# No spontaneous tendency to nematicity in tetragonal crystal



E.Gull, O. Parcollet, P. Werner, and A.J. Millis  
PRB **80**, 245102 (2009)



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# Underdoped metal very sensitive to anisotropy

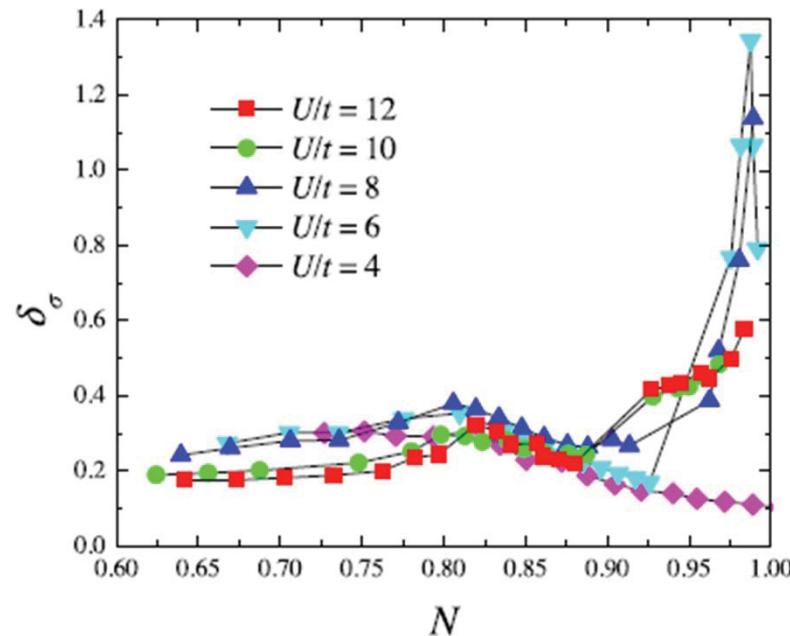
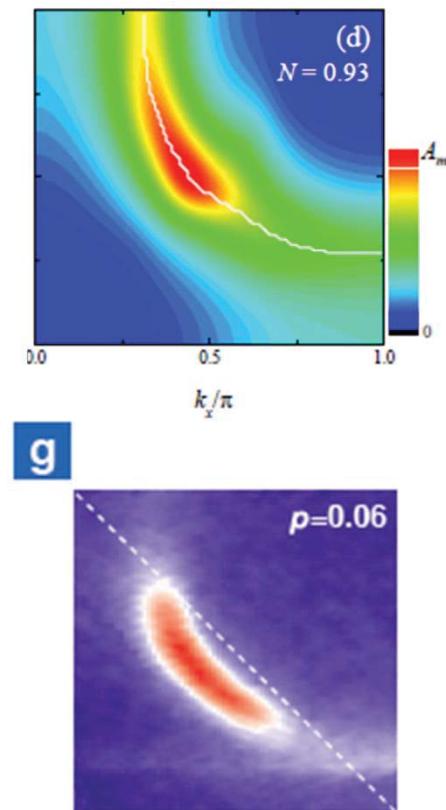


FIG. 3: (Color online) Anisotropy in the CDMFT conductivity  $\delta\sigma = 2 [\sigma_x(0) - \sigma_y(0)] / [\sigma_x(0) + \sigma_y(0)]$  as a function of filling  $N$  for various values of  $U$  and  $\eta = 0.1$ ,  $\delta_0 = 0.04$ .



Satoshi Okamoto



David Sénéchal



Okamoto, Sénéchal, Civelli, AMST  
Phys. Rev. B 82, 180511R 2010

D. Fournier *et al.* Nature Physics ( Marcello Civelli )

# At finite temperature anisotropy in Z

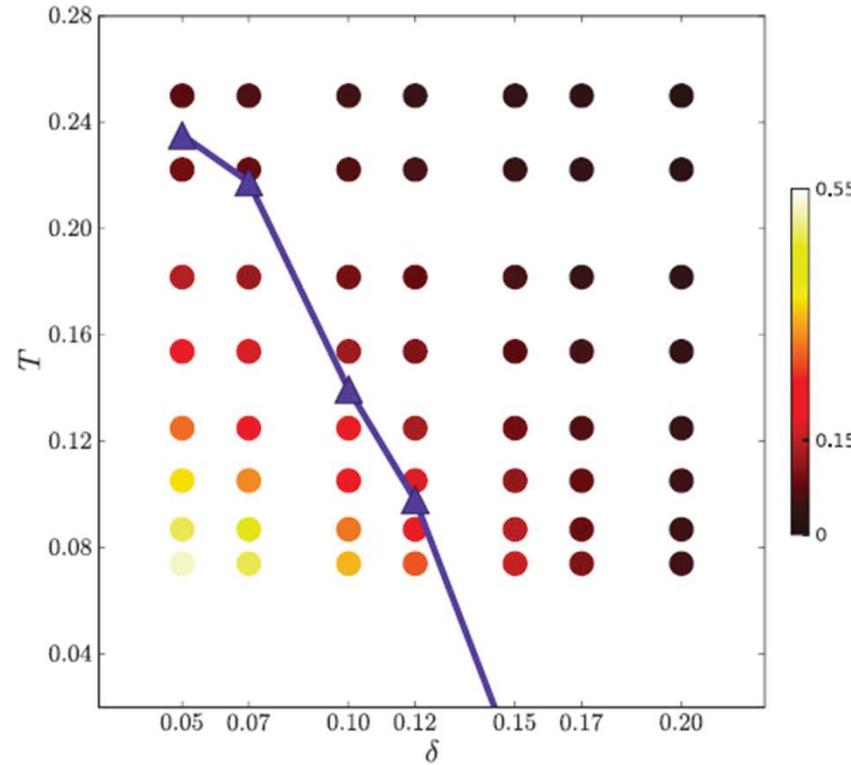
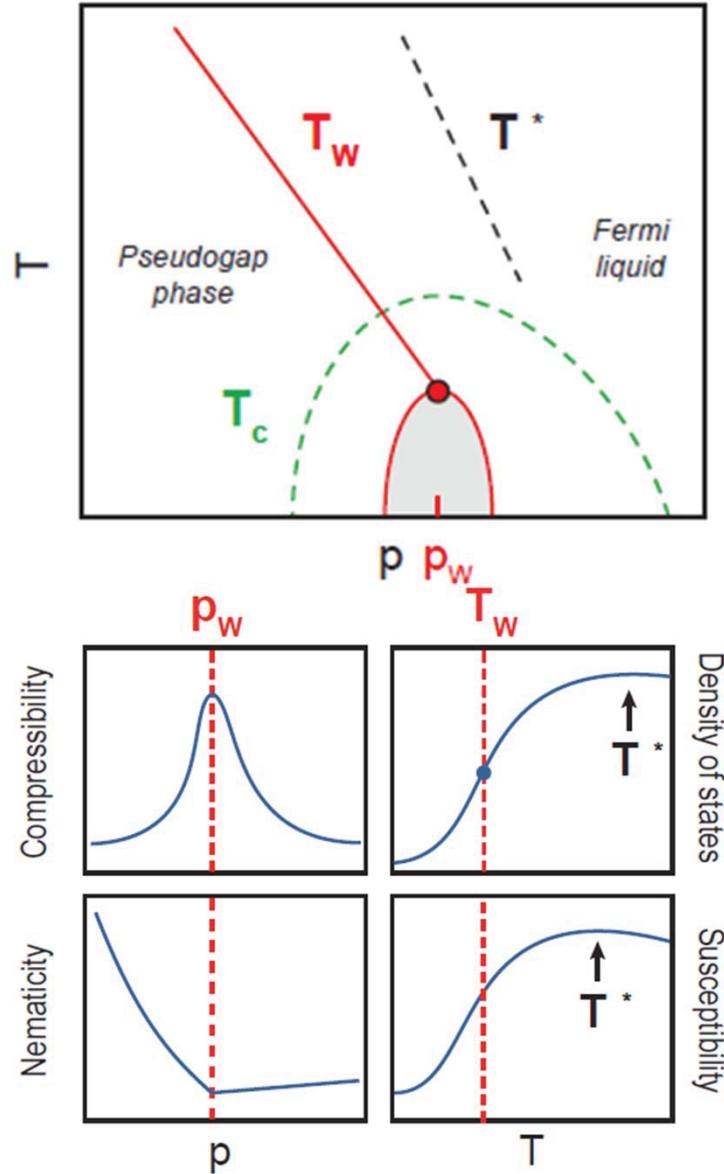


FIG. 3. (Color online) Color map of the anisotropic ratio of the quasiparticle weight  $\sigma_Z$  over the temperature-doping plane, for  $U = 6t$ . The solid blue curve indicates the pseudogap temperature  $T^*(\delta)$  which is obtained as the temperature at which the uniform magnetic susceptibility  $\chi_m[q = (0,0), T]$  has a maximum.

$U = 6t$ , DCA, 4x4

Su, Maier, PRB **84**, 220506(R) (2011)

# An emergent phenomenon in CDMFT



- Is the pseudogap (PG) a crossover or a phase transition ?
- Relation between CDW and the PG ?
- Why CDW peaked at 12% doping ?
- Origin of nematicity ?
- **Why superconducting ?**
- Why a dome of SC ?
- Does a one-band model capture the key physics ?
- AFM QCP important?
- Lessons from other SC?



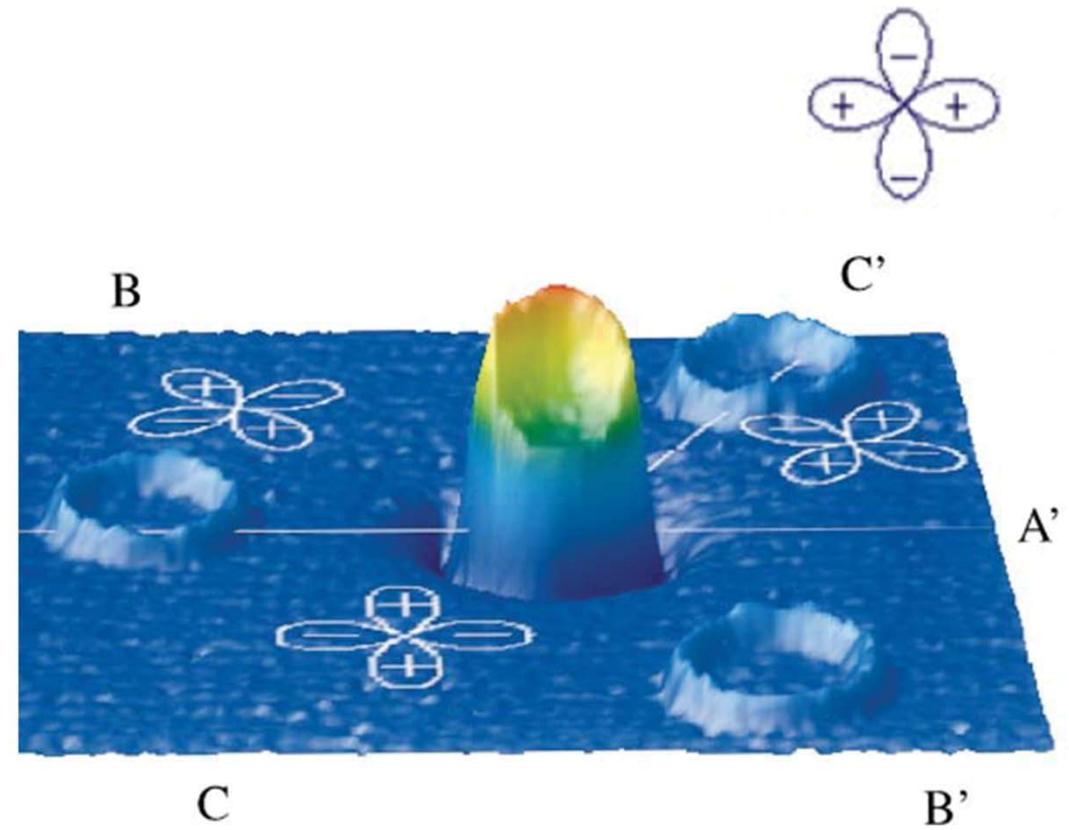
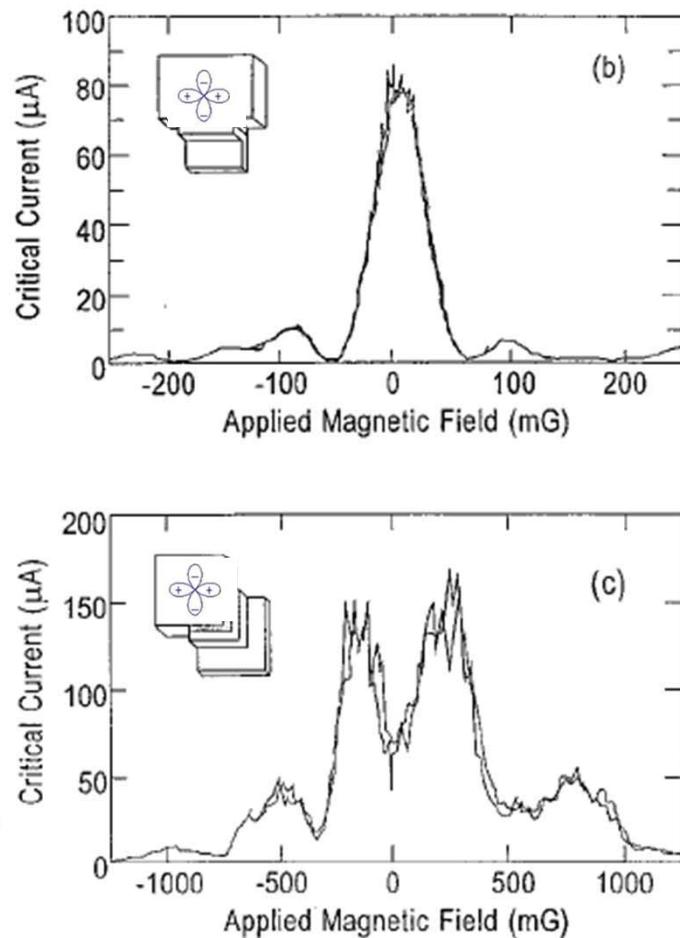
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# d-wave superconductivity



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# High T<sub>c</sub> are d-wave (interference)



Tsuei Kirtley, Rev. Mod. Phys. 2000

Wollman et al. PRL 1993



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# *d*-wave superconductivity

- Weak coupling

- C. J. Halboth and W. Metzner, Phys. Rev. Lett. 85, 5162 (2000).
- B. Kyung, J.-S. Landry, and A. M. S. Tremblay, Phys. Rev. B 68, 174502 (2003).
- C. Bourbonnais and A. Sedeki, Physical Review B 80, 085105 (2009).
- D. J. Scalapino, Physica C: Superconductivity 470, Supplement 1, S1 (2010), ISSN 0921-4534, proceedings of the 9th International Conference on Materials and Mechanisms of Superconductivity.

- Renormalized Mean-Field Theory

- P. W. Anderson, P. A. Lee, M. Randeria, T. M. Rice, N. Trivedi, and F. C. Zhang, Journal of Physics: Condensed Matter 16, R755 (2004).
- K.-Y. Yang, T. M. Rice, and F.-C. Zhang, Phys. Rev. B 73, 174501 (2006).

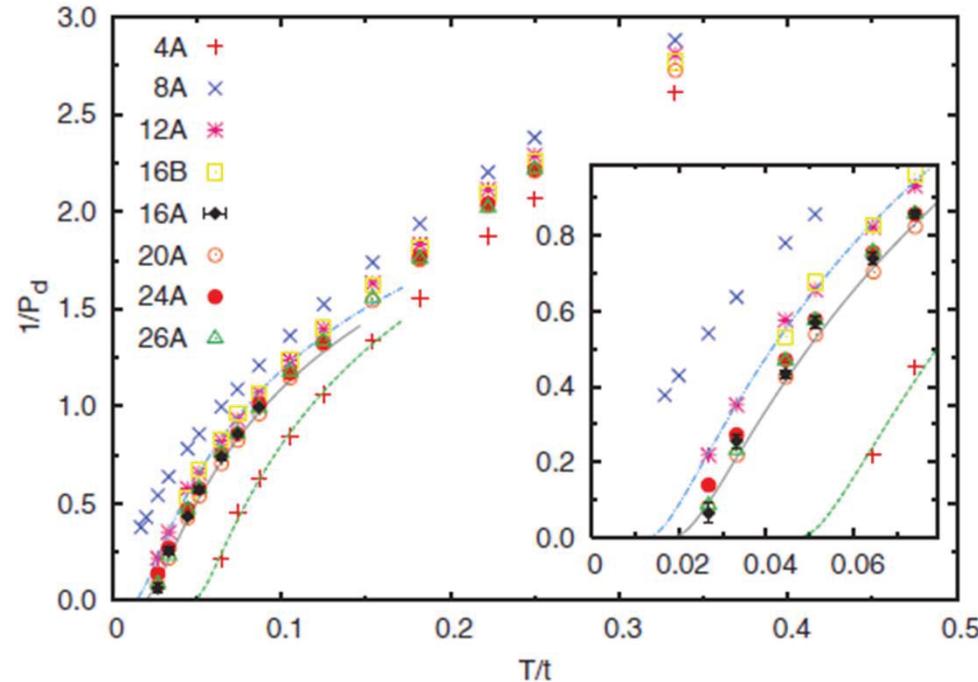
- Slave particles

- P. A. Lee, N. Nagaosa, and X.-G. Wen, Rev. Mod. Phys. 78, 17 (2006).
- M. Imada, Y. Yamaji, S. Sakai, and Y. Motome, Annalen der Physik 523, 629 (2011)

- Variational approaches

- T. Giamarchi and C. Lhuillier, Phys. Rev. B 43, 12943 (1991).
- A. Paramekanti, M. Randeria, and N. Trivedi, Phys. Rev. B 70, 054504 (2004).

# Divergence of d-wave: finite size study



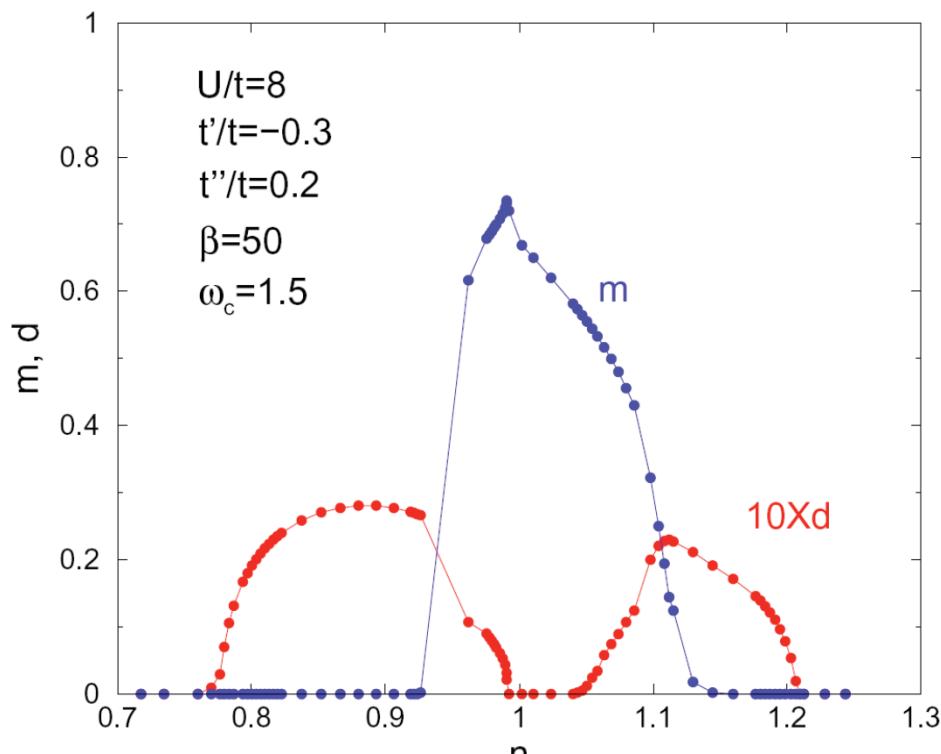
DCA,  $U=4$

T. A. Maier, M. Jarrell, T. C. Schulthess,  
P. R. C. Kent, and J. B. White PRL **95**, 237001 (2005)



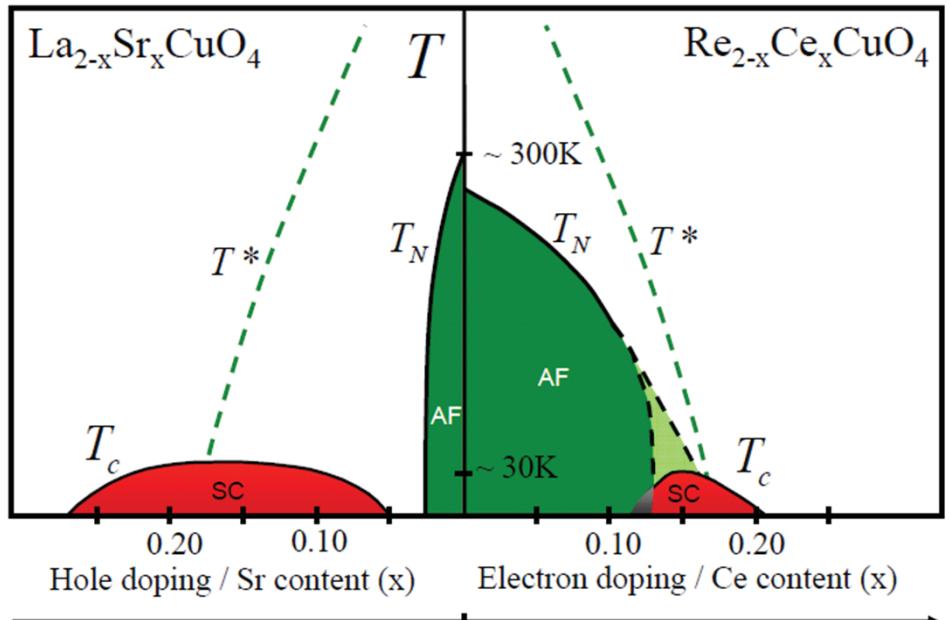
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# CDMFT global phase diagram



Kancharla, Kyung, Civelli,  
Sénéchal, Kotliar AMST

Phys. Rev. B (2008)  
AND Capone, Kotliar PRL (2006)



Armitage, Fournier, Greene, RMP (2009)





Giovanni Sordi



Patrick Sémon



Kristjan Haule

# Finite $T$ phase diagram Superconductivity $t' = 0$

Sordi et al. PRL **108**, 216401 (2012)



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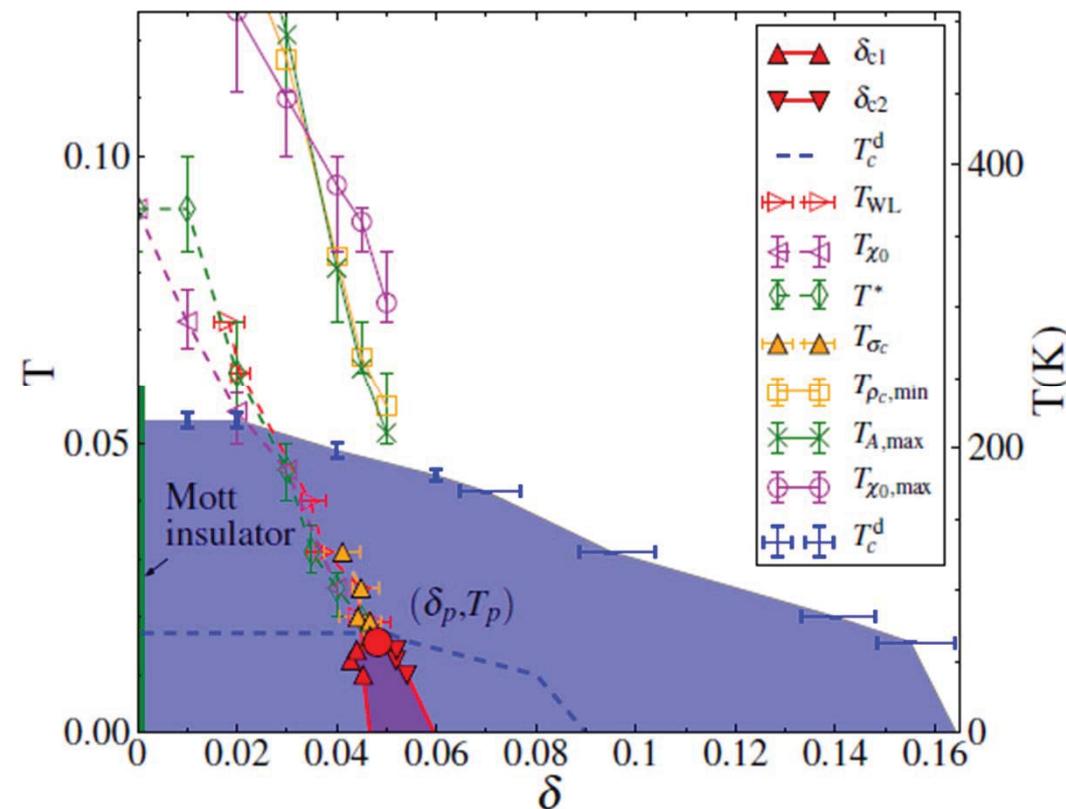


Giovanni Sordi



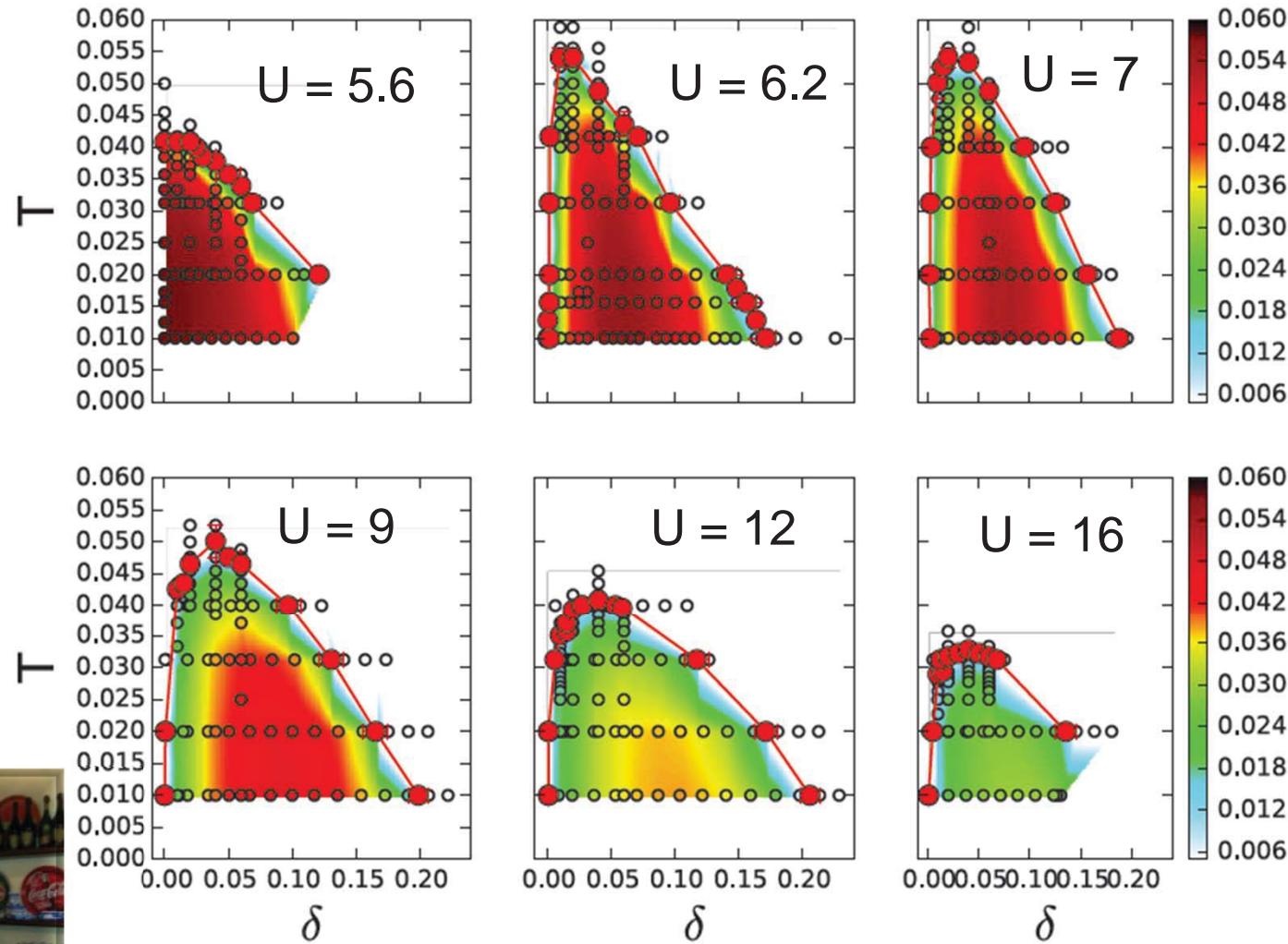
Patrick Sémon

# Phase diagram for $U = 6.2 t$



G. Sordi et al. Phys. Rev. Lett. 108, 216401/1-6 (2012)  
P. Sémon, G. Sordi, A.-M.S.T., Phys. Rev. B **89**, 165113/1-6 (2014)

# Order parameter (color) and $T_c$



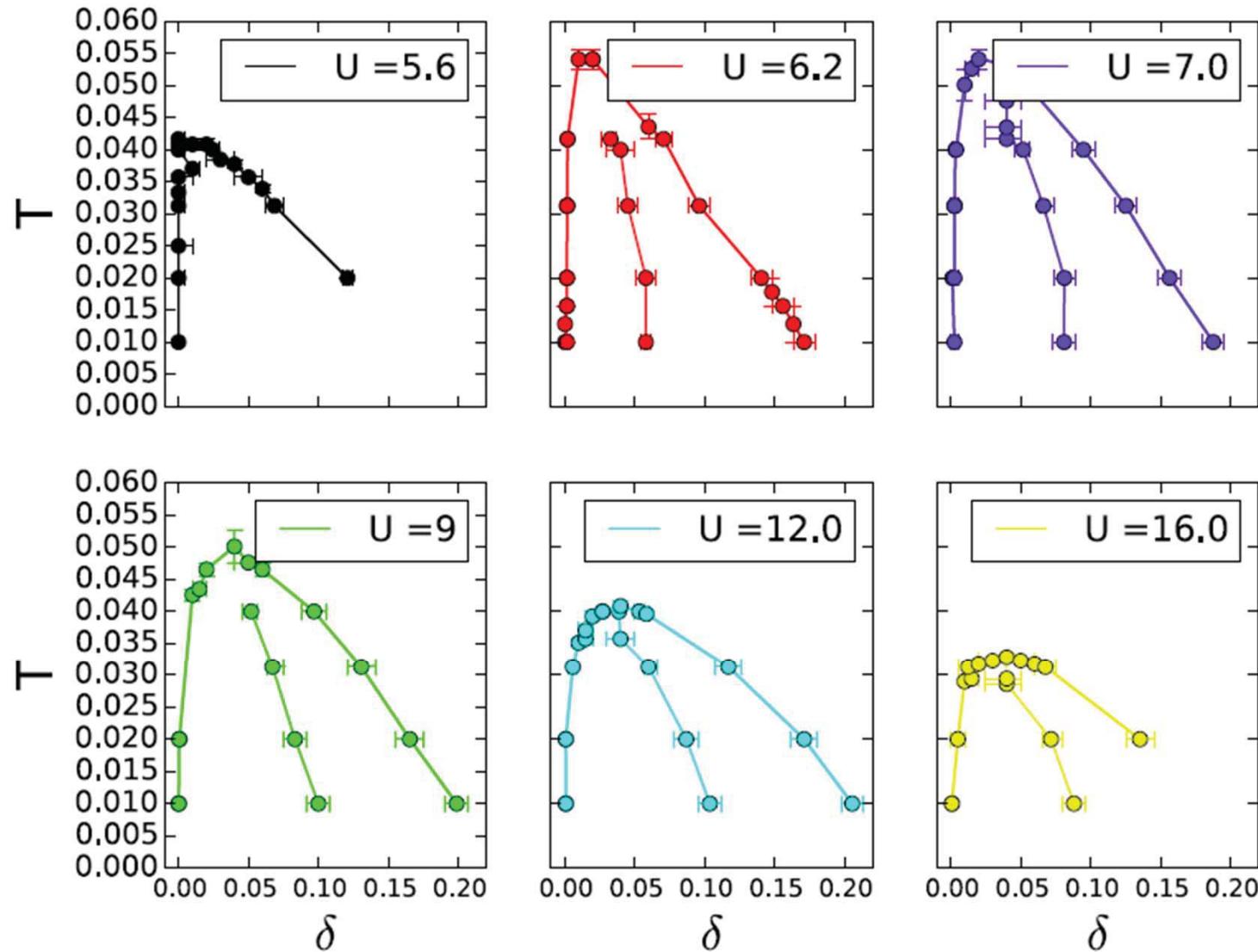
Lorenzo Fratino

L. Fratino, G. Sordi (unpublished)



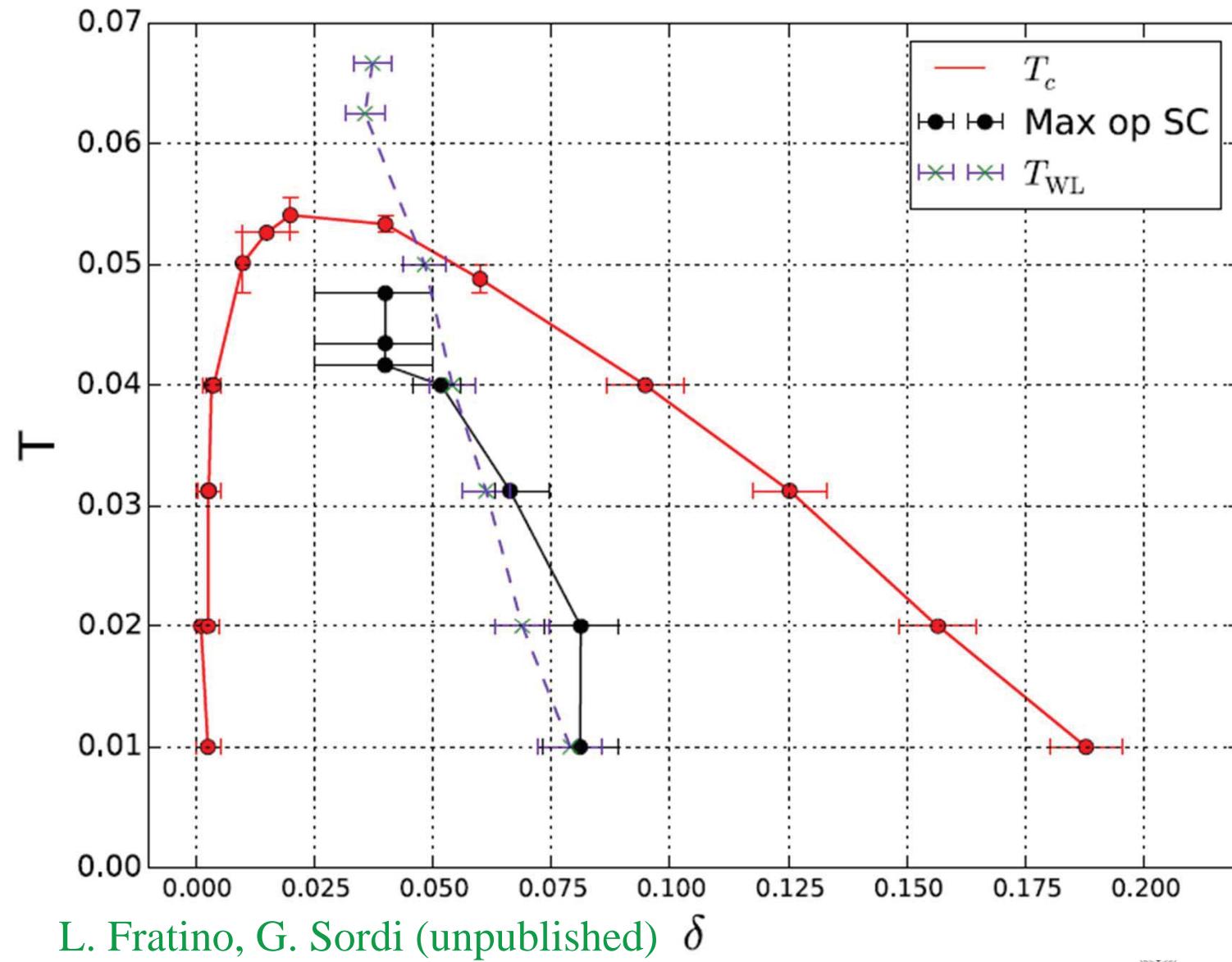
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# $T_c$ vs $T_{\max}$ order parameter



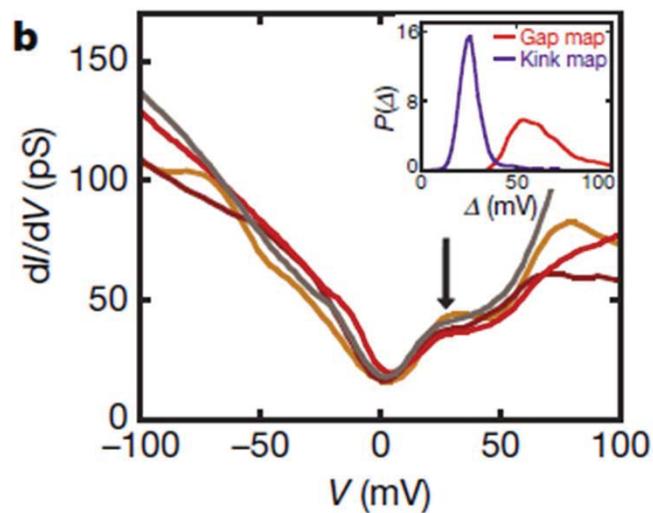
L. Fratino, G. Sordi (unpublished)

# $U=7$ , $T_W$ vs $T_c$ vs $T_{\max}$ order parameter



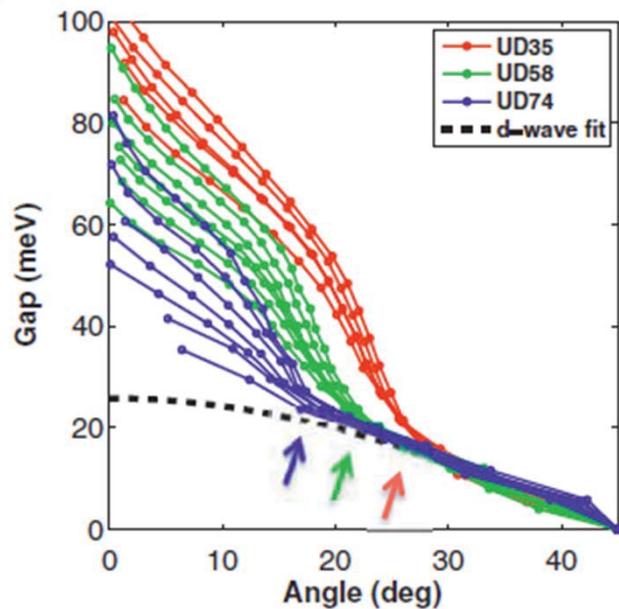
# Meaning of $T_c$

- Local pair formation



K. K. Gomes, A. N. Pasupathy, A. Pushp,  
S. Ono, Y. Ando, and A. Yazdani,  
Nature **447**, 569 (2007)

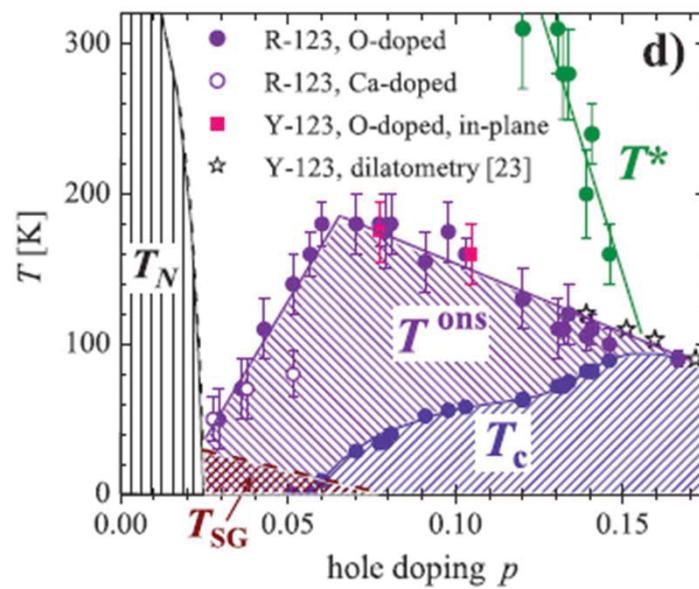
# Meaning of $T_c^d$ : Local pair formation



A. Pushp, Parker, ... A. Yazdani,  
Science **364**, 1689 (2009)

However, our measurements demonstrate that the nodal gap does not change with reduced doping. The pairing strength does not get weaker or stronger as the Mott insulator is approached; rather, it saturates.

# Fluctuating region

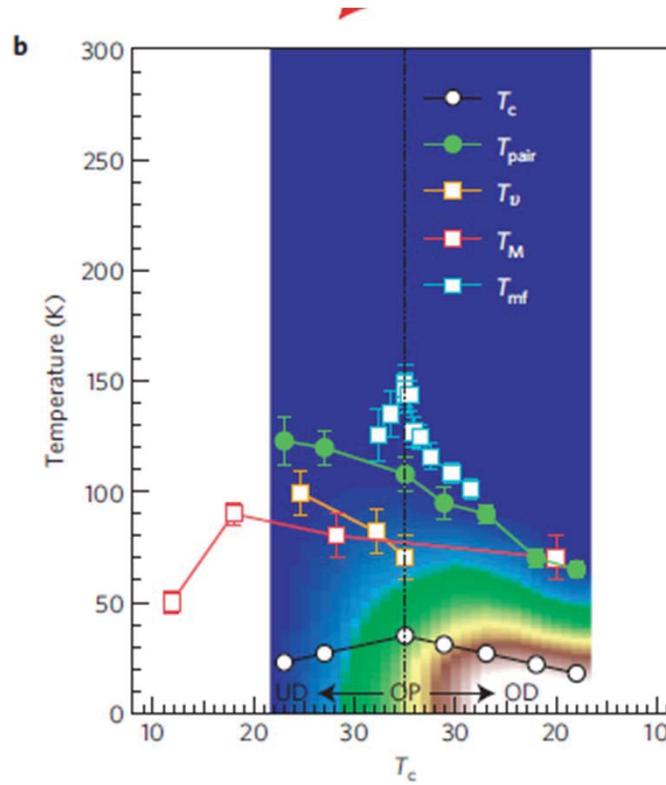


Infrared response

Dubroka et al. PRL 106, 047006 (2011)

# T<sub>pair</sub>

ARPES  
Bi2212

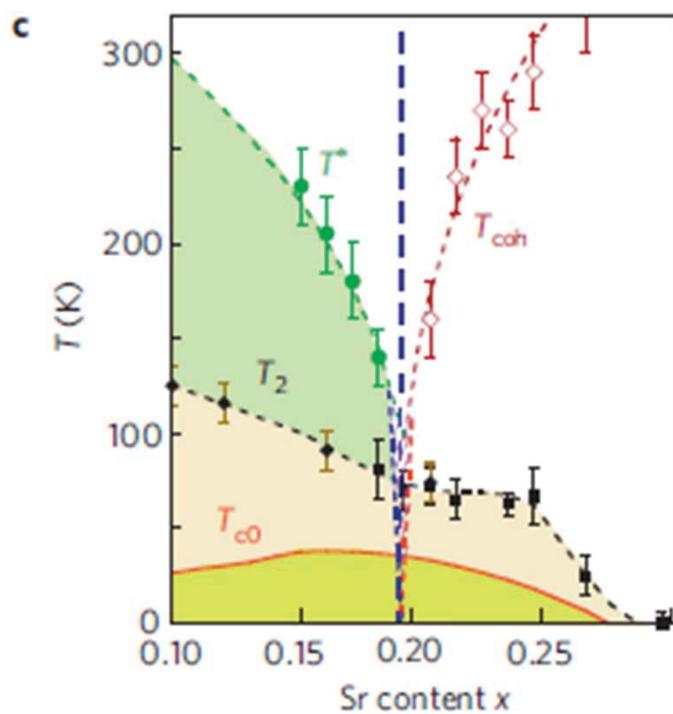


Kondo, Takeshi, et al. Kaminski Nature  
Physics 2011, 7, 21-25



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$T_2$



Magnetoresistance, LSCO  
Fluctuating vortices

Patrick M. Rourke, et al. Hussey Nature Physics 7, 455–458 (2011)

# Giant proximity effect

$$\begin{aligned}T_c &= 32 \text{ K} \\T_c &< 5 \text{ K}\end{aligned}$$

Morenzoni et al.,  
Nature Comms. 2 (2011)

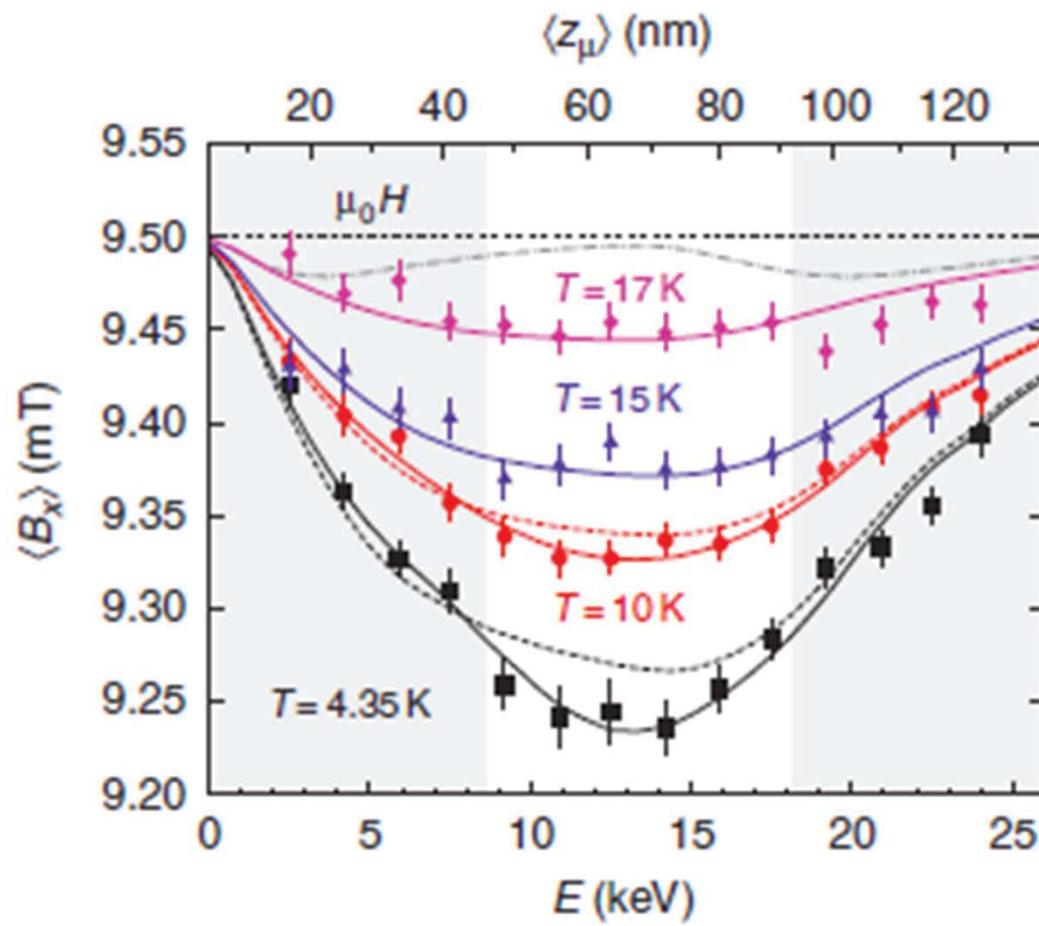


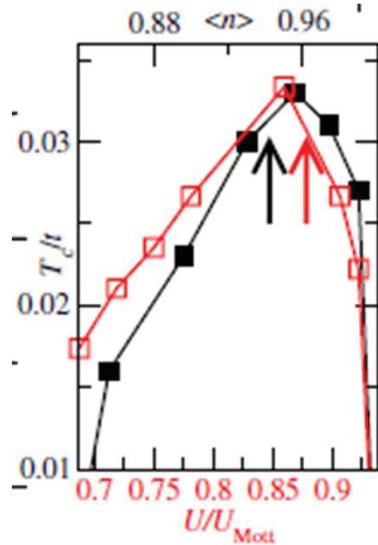
Figure 6 | Depth profile of the local field at different temperatures. The

# Actual $T_c$ in underdoped

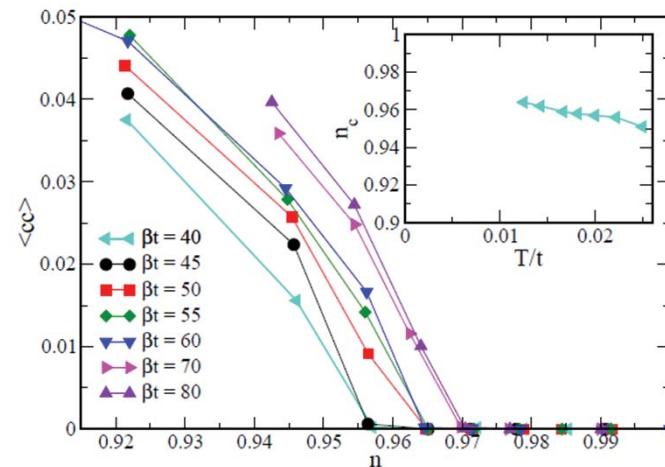
- Quantum and classical phase fluctuations
  - V. J. Emery and S. A. Kivelson, Phys. Rev. Lett. **74**, 3253 (1995).
  - V. J. Emery and S. A. Kivelson, Nature **374**, 474 (1995).
  - D. Podolsky, S. Raghu, and A. Vishwanath, Phys. Rev. Lett. **99**, 117004 (2007).
  - Z. Tesanovic, Nat Phys **4**, 408 (2008).
- Magnitude fluctuations
  - I. Ussishkin, S. L. Sondhi, and D. A. Huse, Phys. Rev. Lett. **89**, 287001 (2002).
- Competing order
  - E. Fradkin, S. A. Kivelson, M. J. Lawler, J. P. Eisenstein, and A. P. Mackenzie, Annual Review of Condensed Matter Physics **1**, 153 (2010).
- Disorder
  - F. Rullier-Albenque, H. Alloul, F. Balakirev, and C. Proust, EPL (Europhysics Letters) **81**, 37008 (2008).
  - H. Alloul, J. Bobro, M. Gabay, and P. J. Hirschfeld, Rev. Mod. Phys. **81**, 45 (2009).

# Larger clusters

- In 2x2  $T_c$  vanishes extremely close to half-filling. In larger cluster, earlier.
- Local pairs in underdoped (2x2)



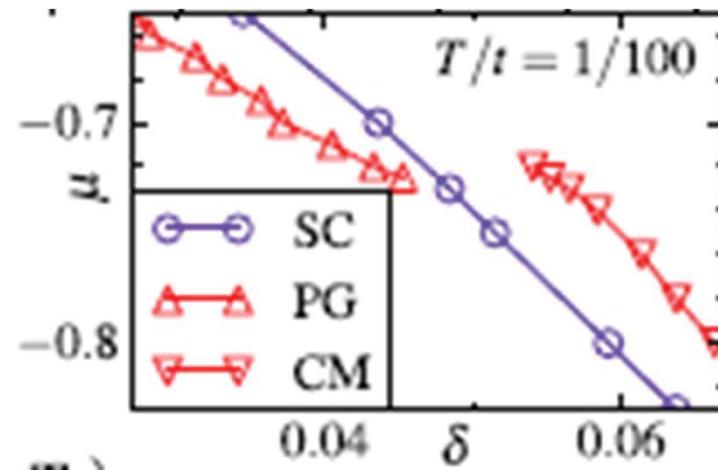
8 site DCA,  $U=6t$



8 site DCA,  $U=6.5t$

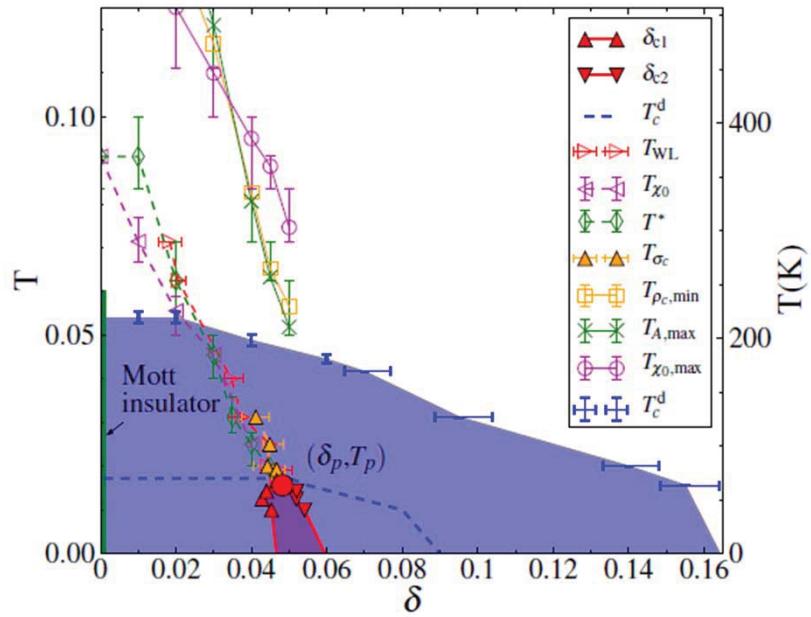
Gull Parcollet Millis,  
PRL 110, 216405 (2013)

# Fate of the first order transition in SC state



G. Sordi et al. Phys. Rev. Lett. 108, 216401/1-6 (2012)

# Summary

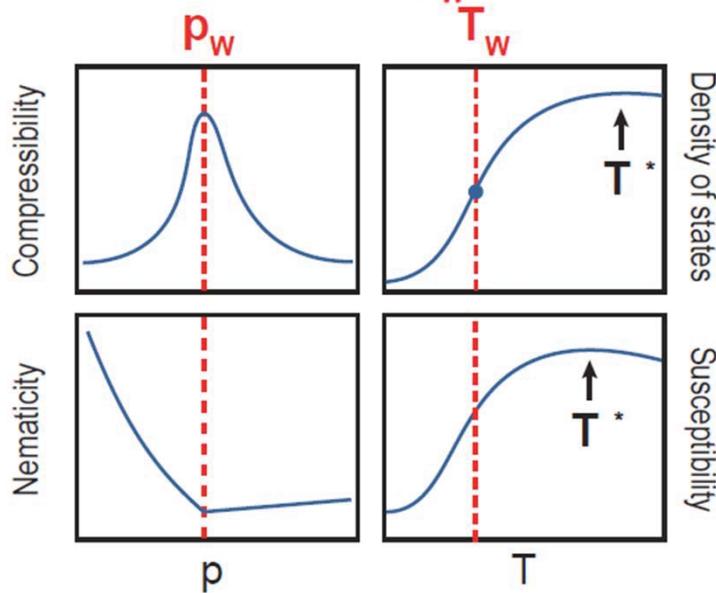
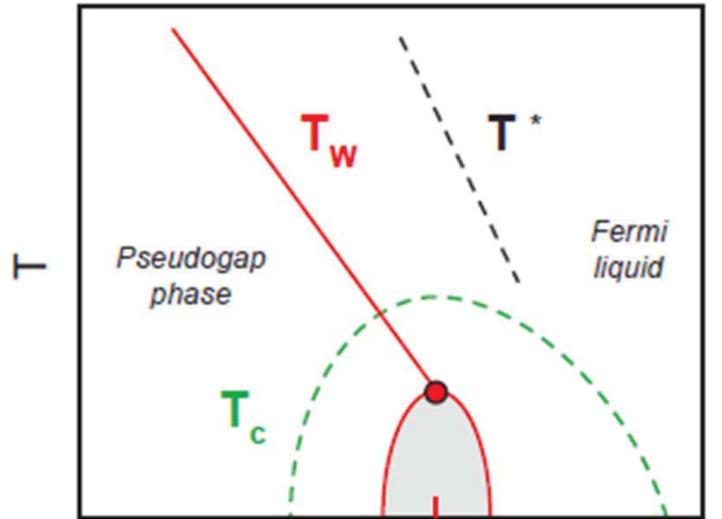


- Below the dome, not QCP (but Mott)
- Maximum near Widom line
- $T^*$  different from  $T_c^d$
- First-order transition destroyed (but traces in the dynamics)
- Actual  $T_c$  in underdoped
  - Competing order
  - Long wavelength fluctuations (see O.P.)



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# Organizing principle



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- Origin of nematicity ?
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- Does a one-band model capture the key physics ?
- AFM QCP important?
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Bio break

