

Defects, Diffusion, Deformation and Thermal Conductivity in the Lower Mantle and D''

John Brodholt

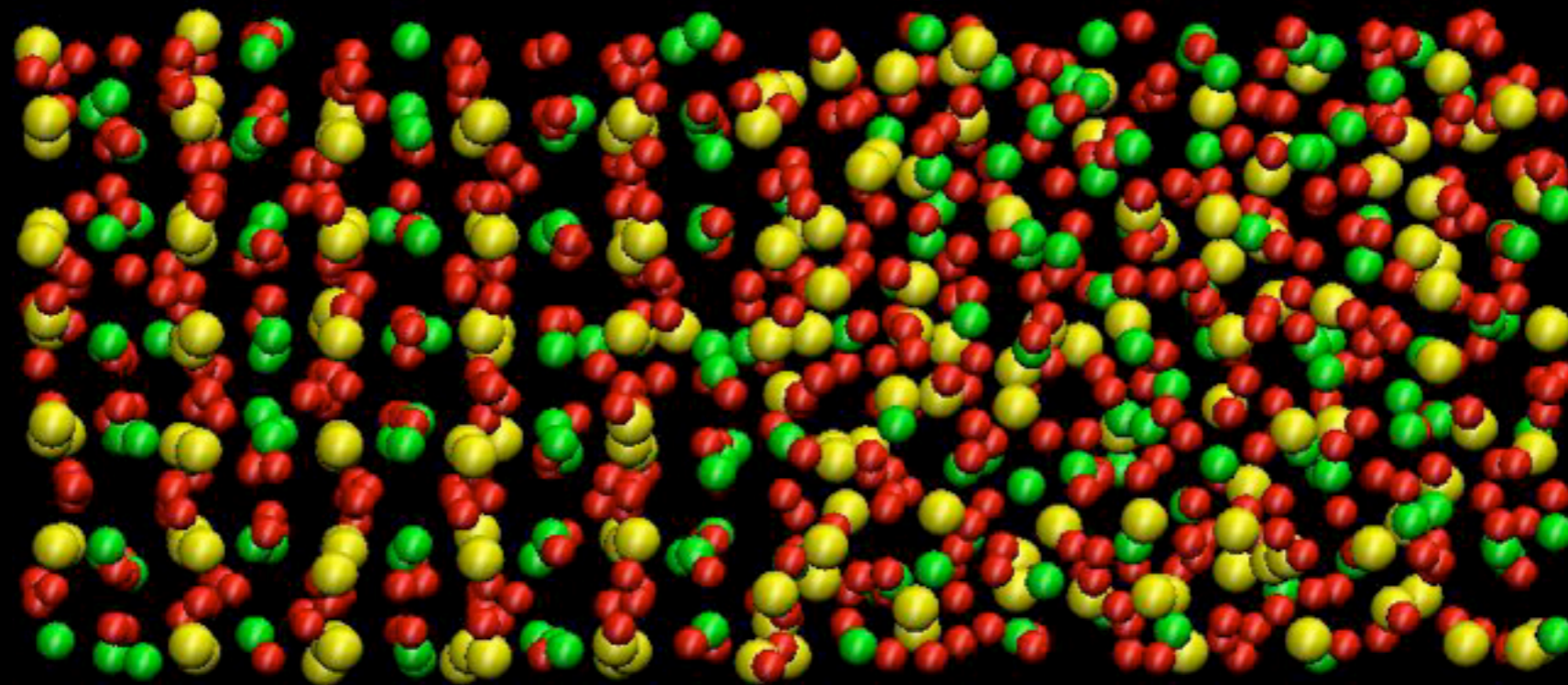
UCL

Thanks to: Michael Ammann, Simon Hunt,
James Wookey, Kai Wang, Andrew Walker and David Dobson

Defects and Diffusion is important:

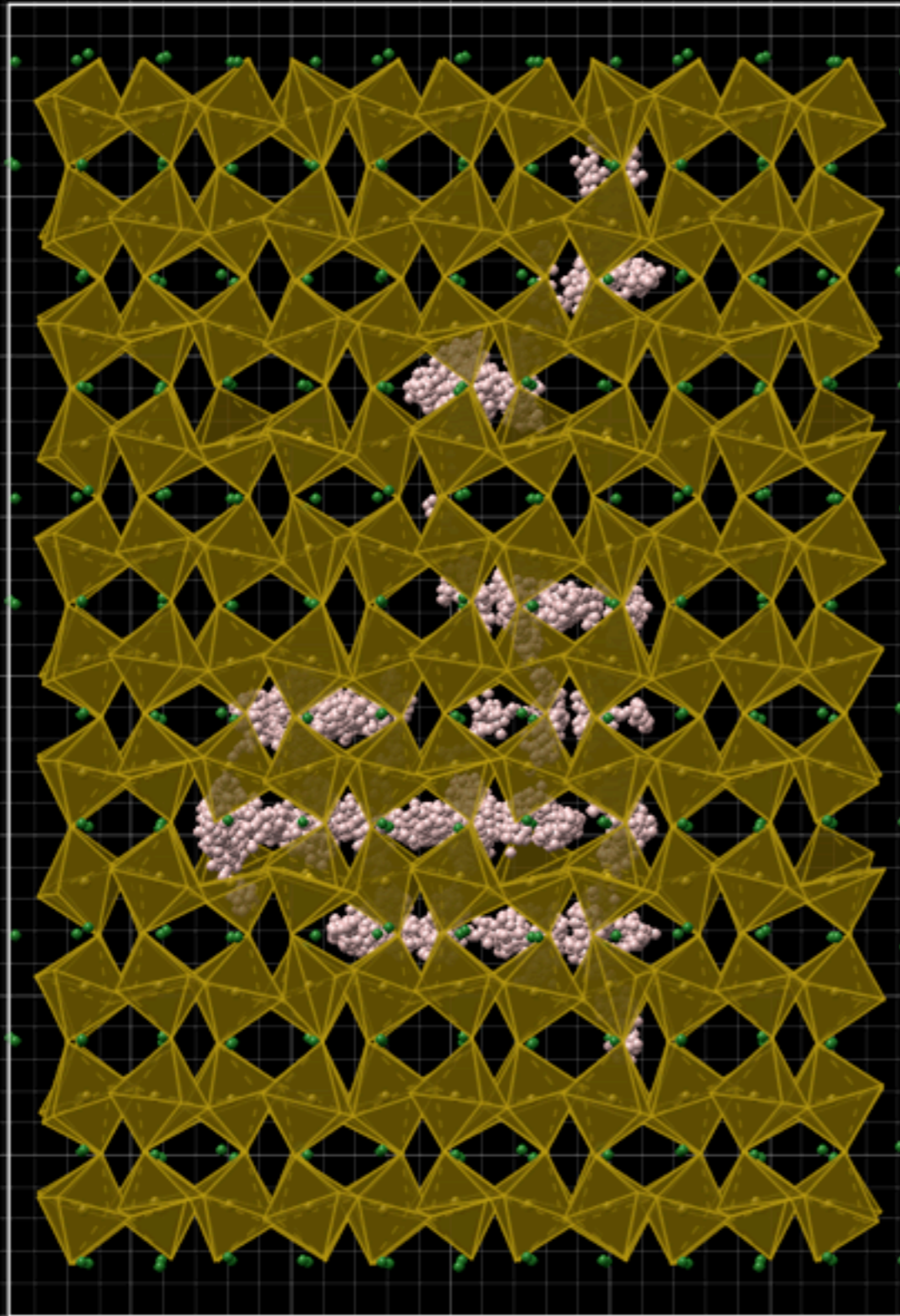
- Controls chemical exchange between crystalline, melt and fluid phases
- Degree of composition zoning in minerals
- Kinetics of phase transitions
- Rate at which minerals grow and their grain sizes
- Has a central role in controlling rheology
(Deformation)

- High P & T experiments on diffusion are hard.
- So our approach is to use a theoretical approach (ab initio or first principles).



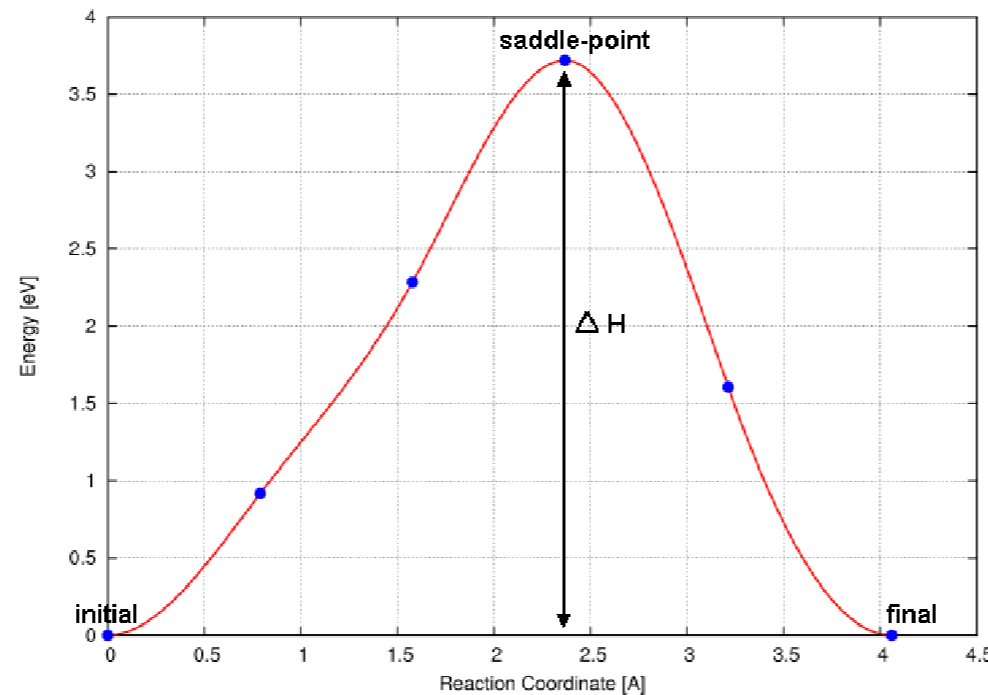
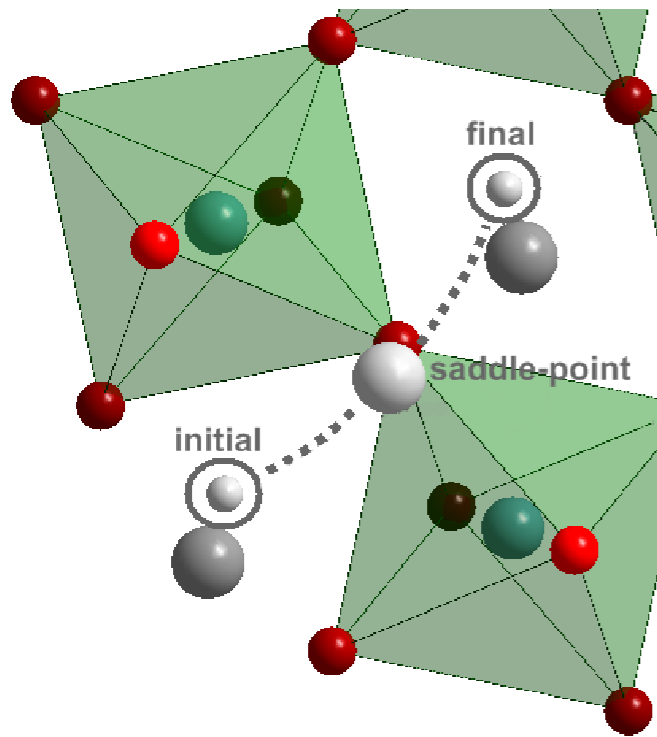
Forces are calculated either from “First principles” - quantum mechanics - or from “empirical potentials”.

He diffusion in
perovskite - the
exception.



He diffusion in perovskite - the exception.

Self Diffusion: $D_{sd} = AN_V \nu e^{\frac{\Delta S}{k}} e^{-\frac{\Delta H}{kT}}$



All energies and frequencies obtained from DFT calculations

Migration enthalpy: $\Delta H = E_{saddlepoint} - E_{initial}$

Attempt frequency and migration entropy (Vineyard Theory)

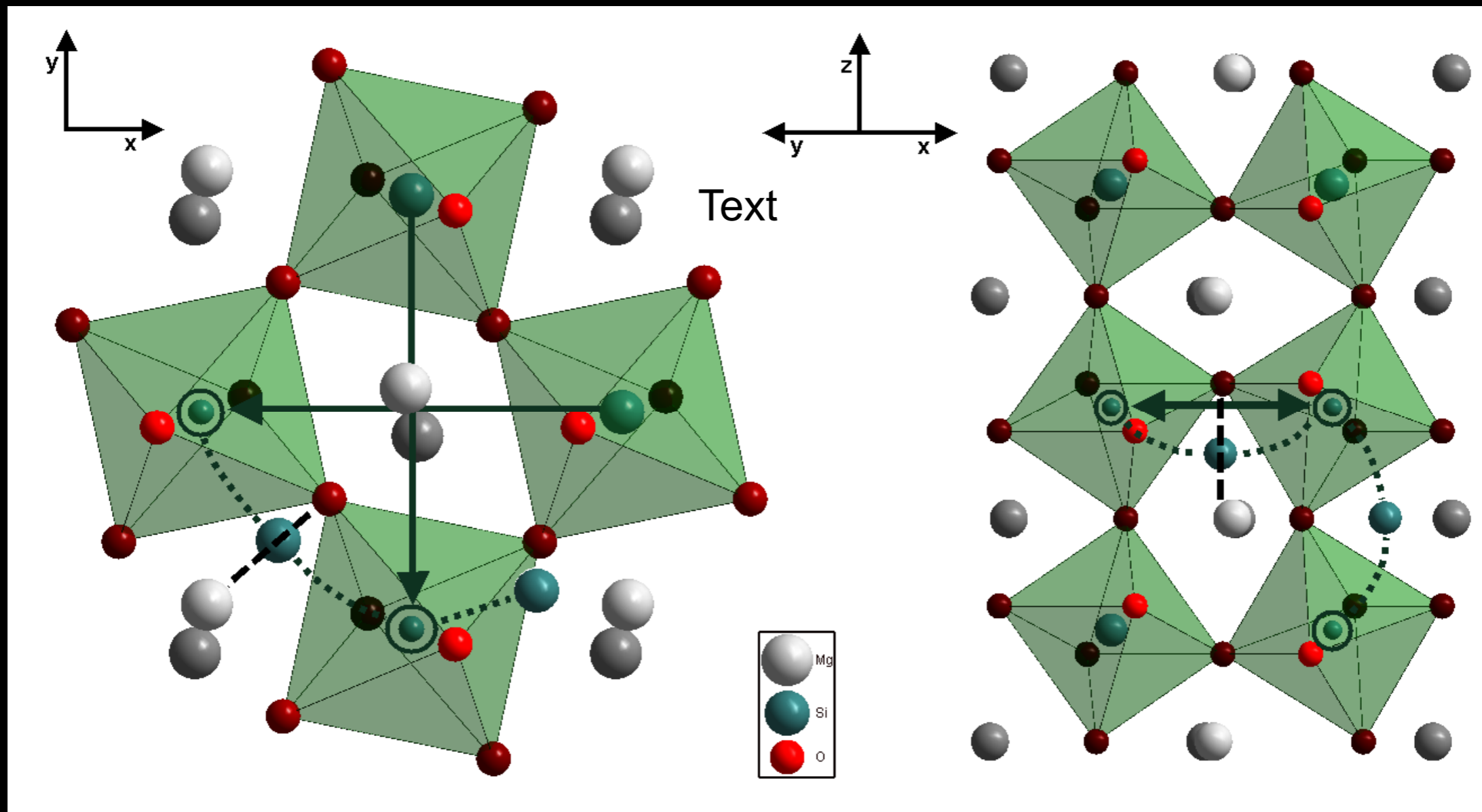
$$\nu e^{\frac{\Delta S}{k}} = \frac{\prod_{n=1}^N f_i}{\prod_{n=1}^{N-1} f'_i}$$

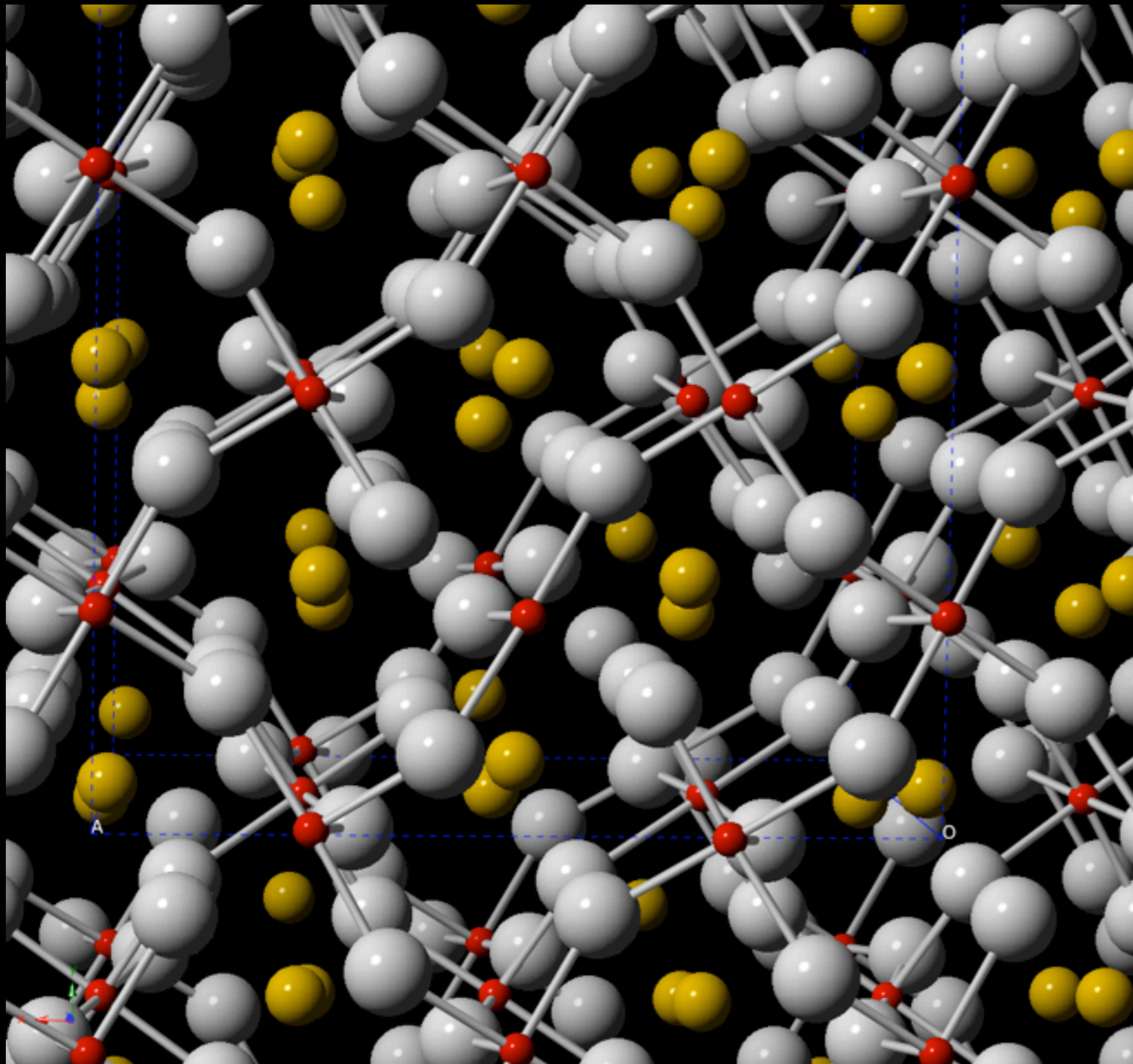
Theory: Wright & Price (1993) = 9.4 eV (empirical potentials)

Experiments: Yamazaki et al., (2000) = 3.6 eV

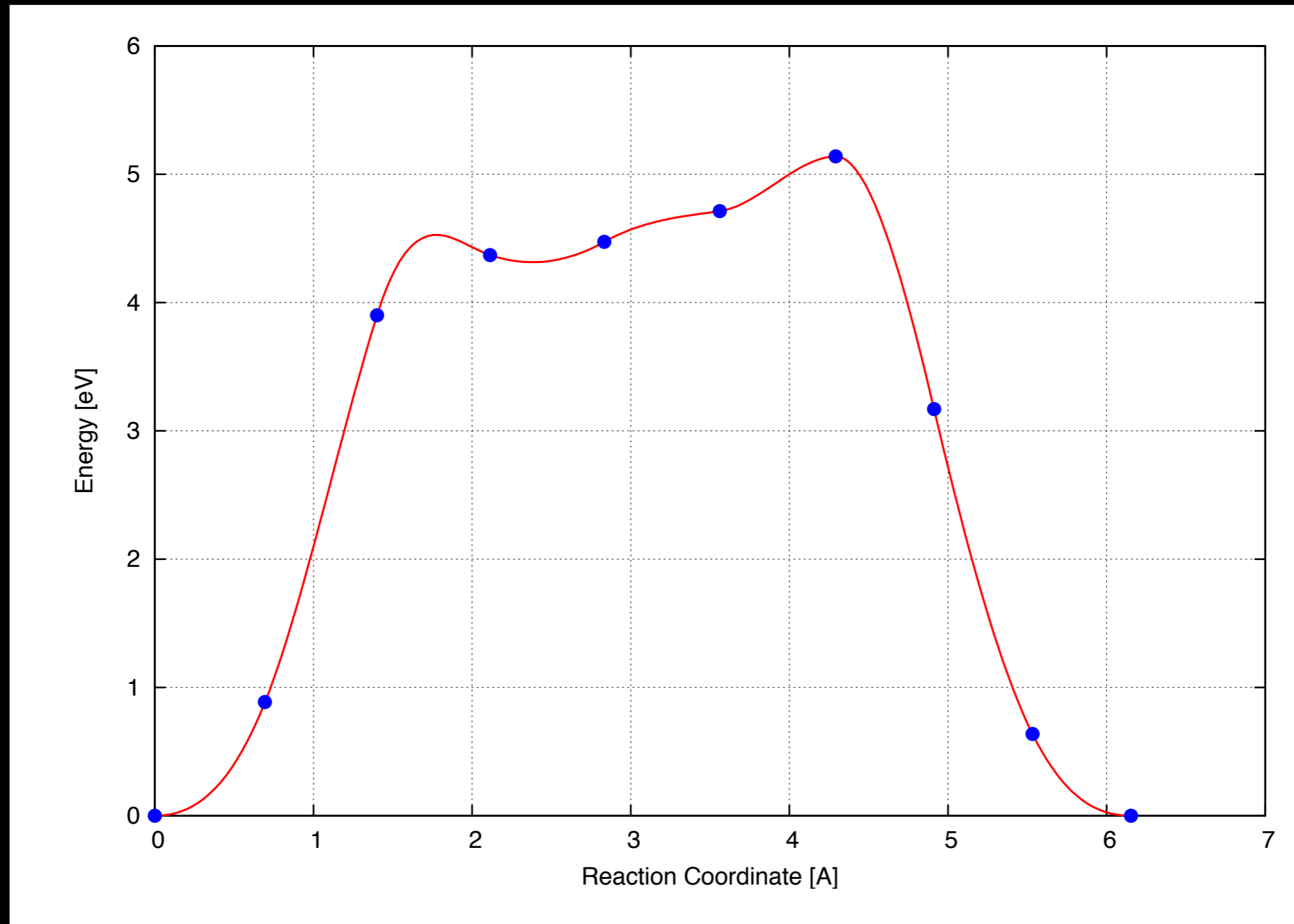
Theory: Karki and Khanduja (2007) = 9 eV (DFT)

Experiments: Dobson et al, (2008) = 3.7 eV

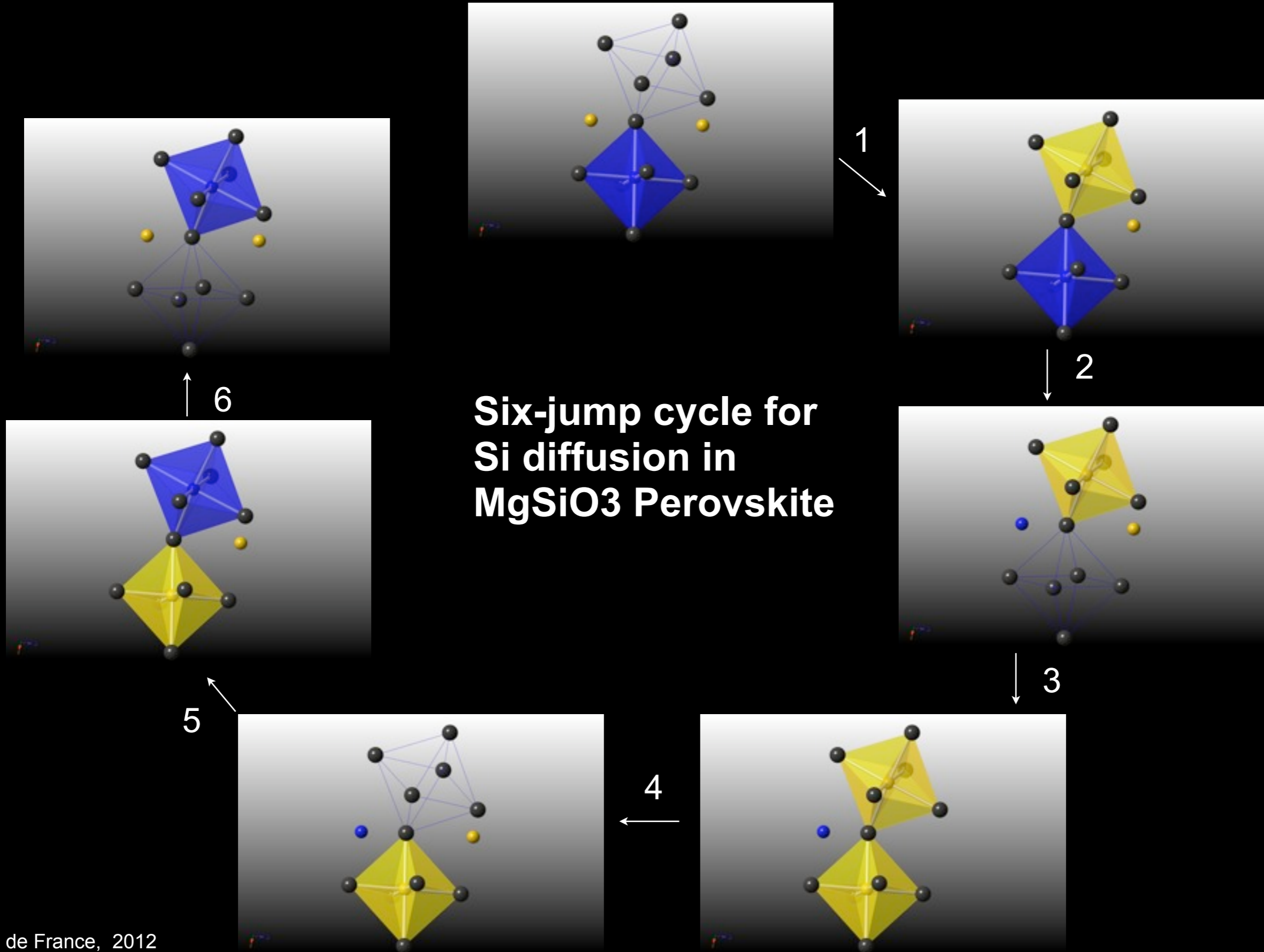




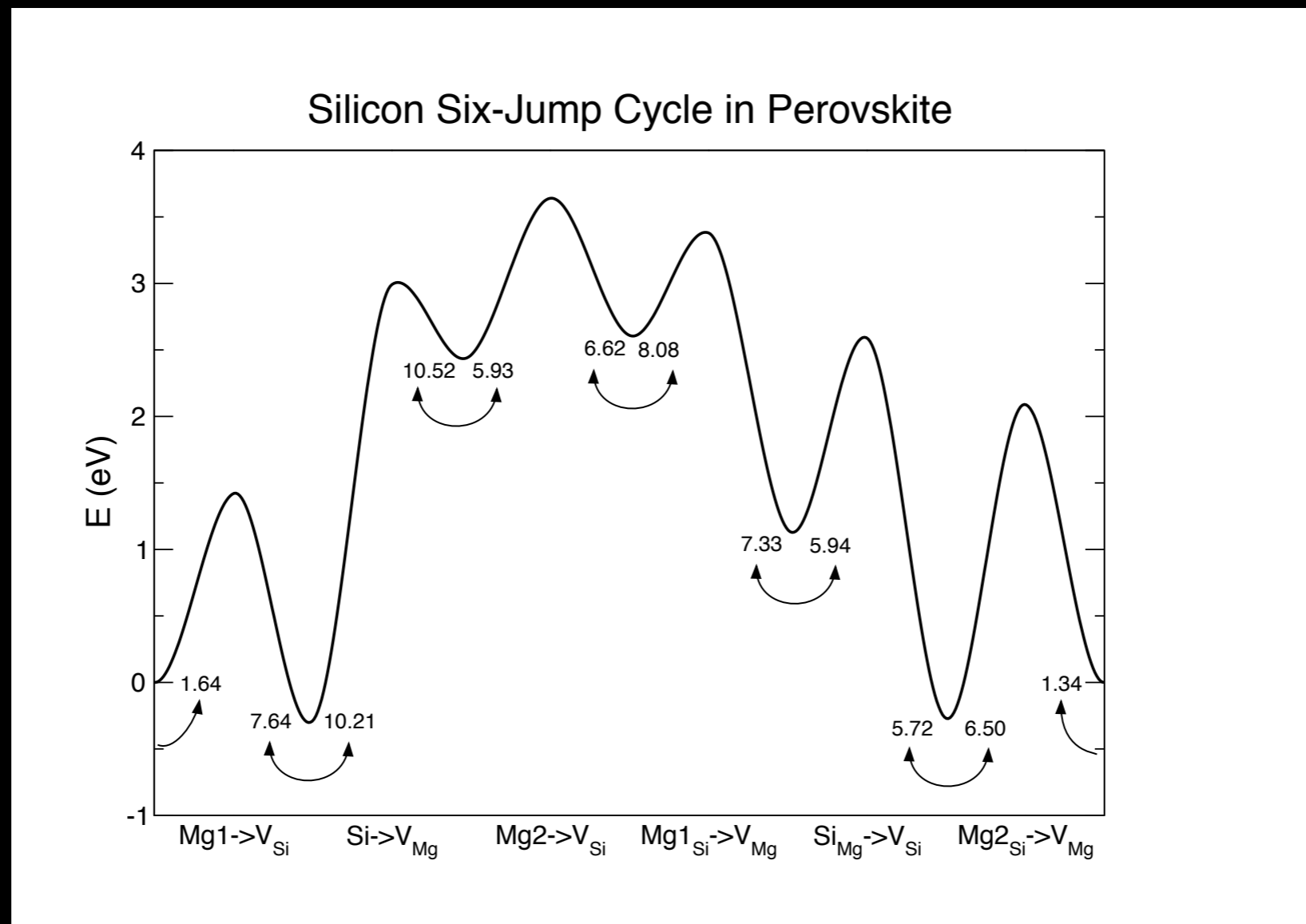
Si Diffusion in Perovskite



Agreement with experiment (3.6 eV) is better than previous estimates (9 eV) but still not great!

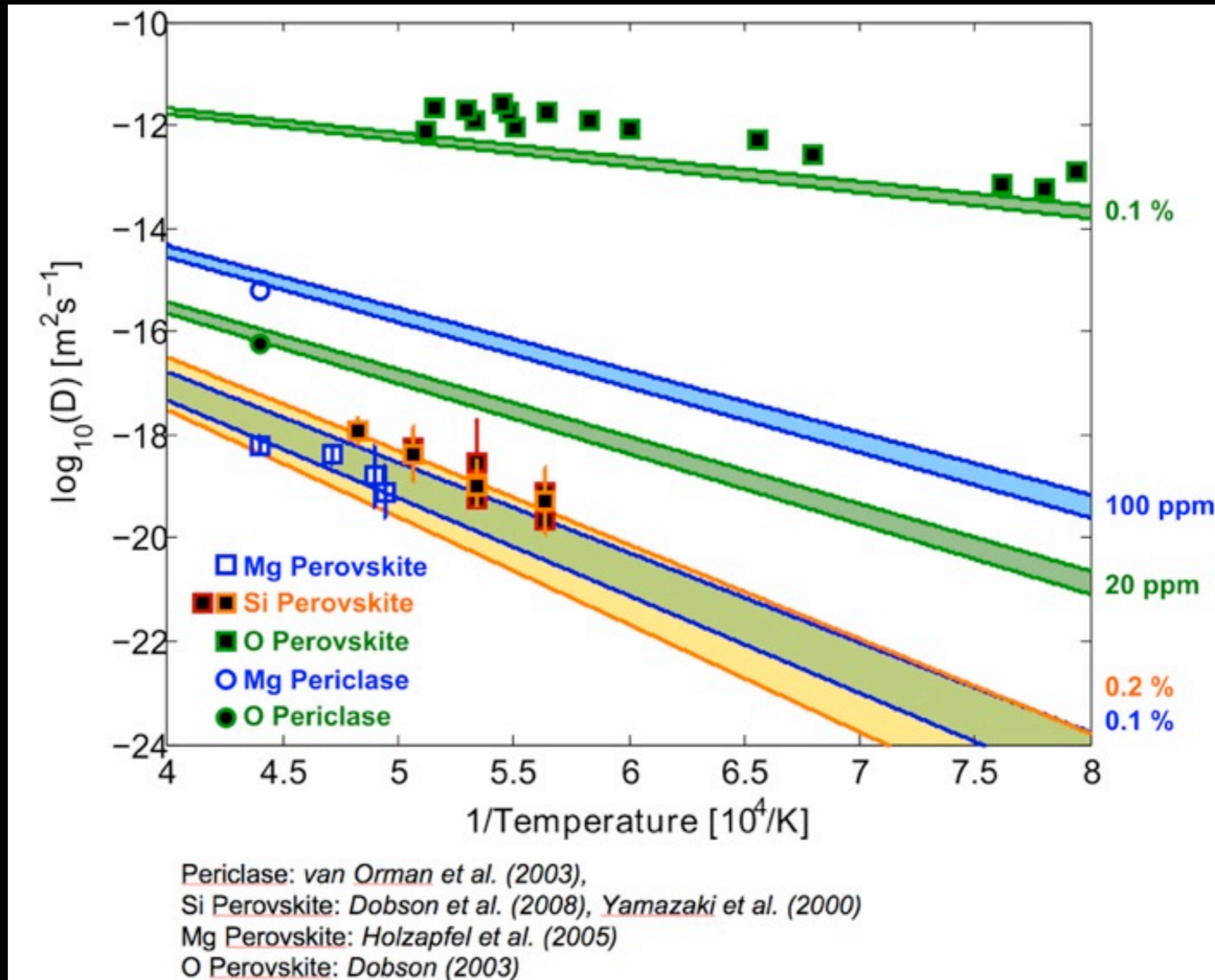


Si diffusion in perovskite does not seem to occur via a simple vacancy hopping mechanism

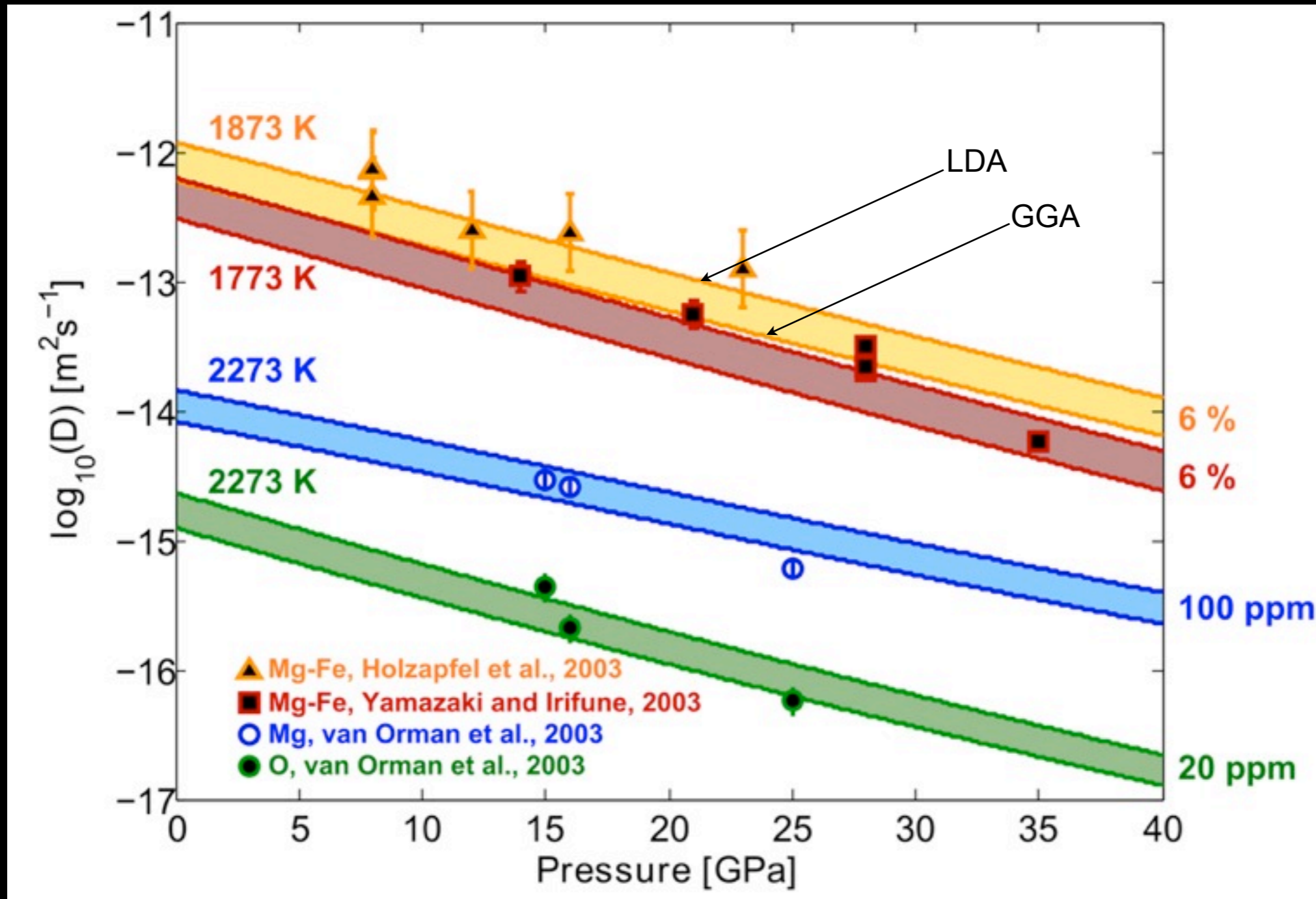


Apparent activation energy for the total cycle is 3.6 eV. This agrees well with 3.61 eV and 3.5 eV found by Dobson et al. (2008) and Yamazaki et al (2000).

PEROVSKITE DIFFUSION RATES

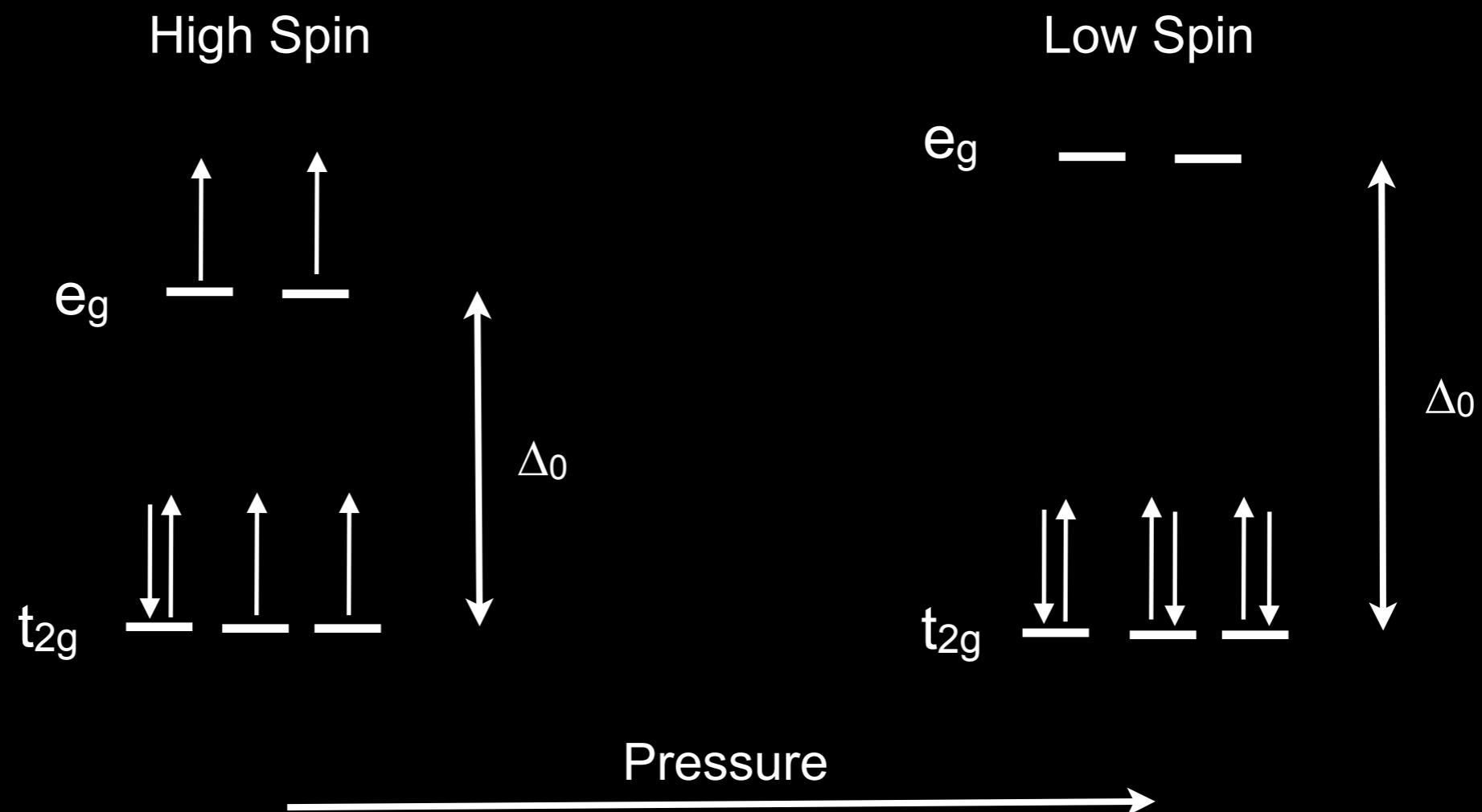


Absolute Diffusion Rates in MgO

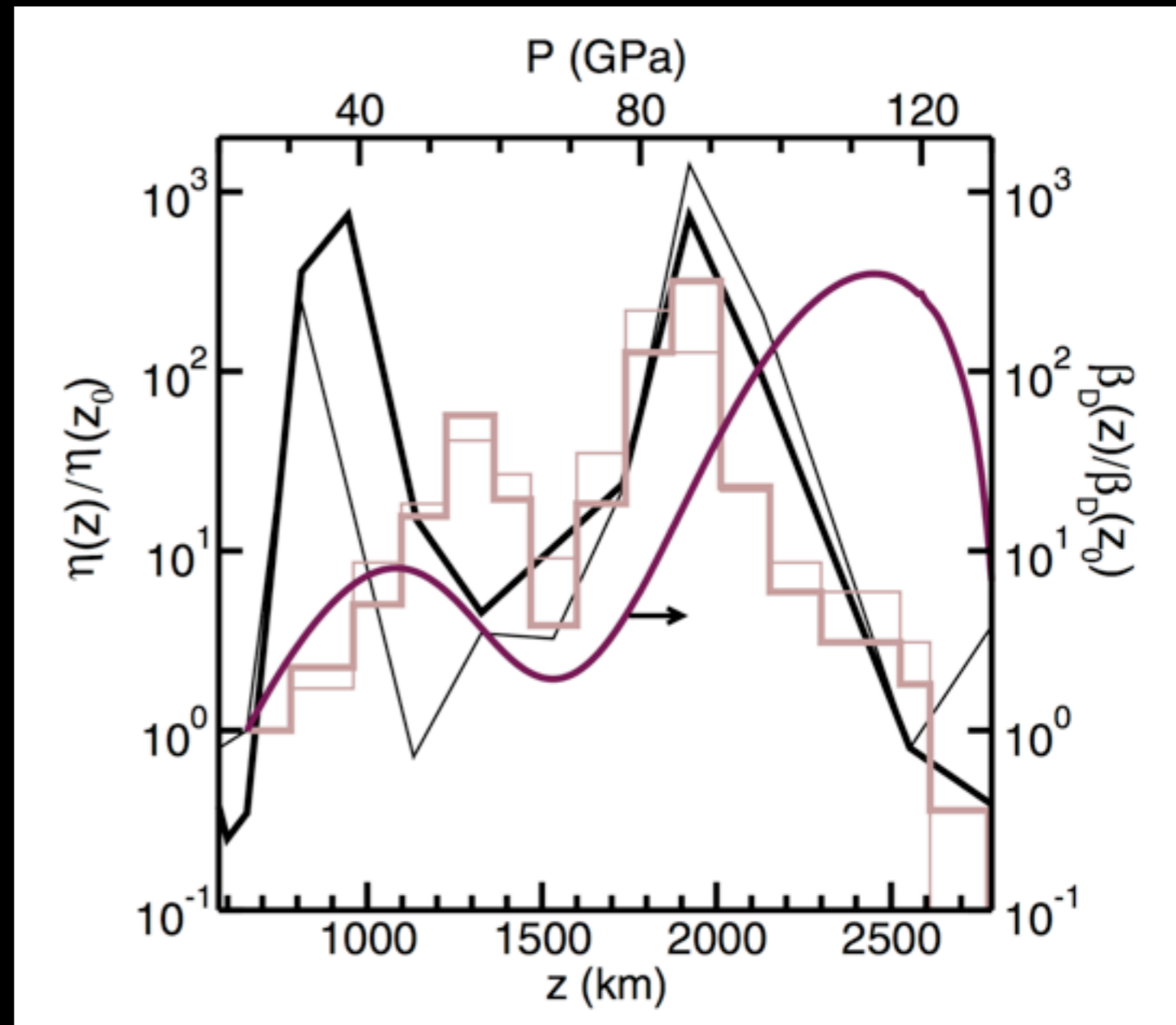
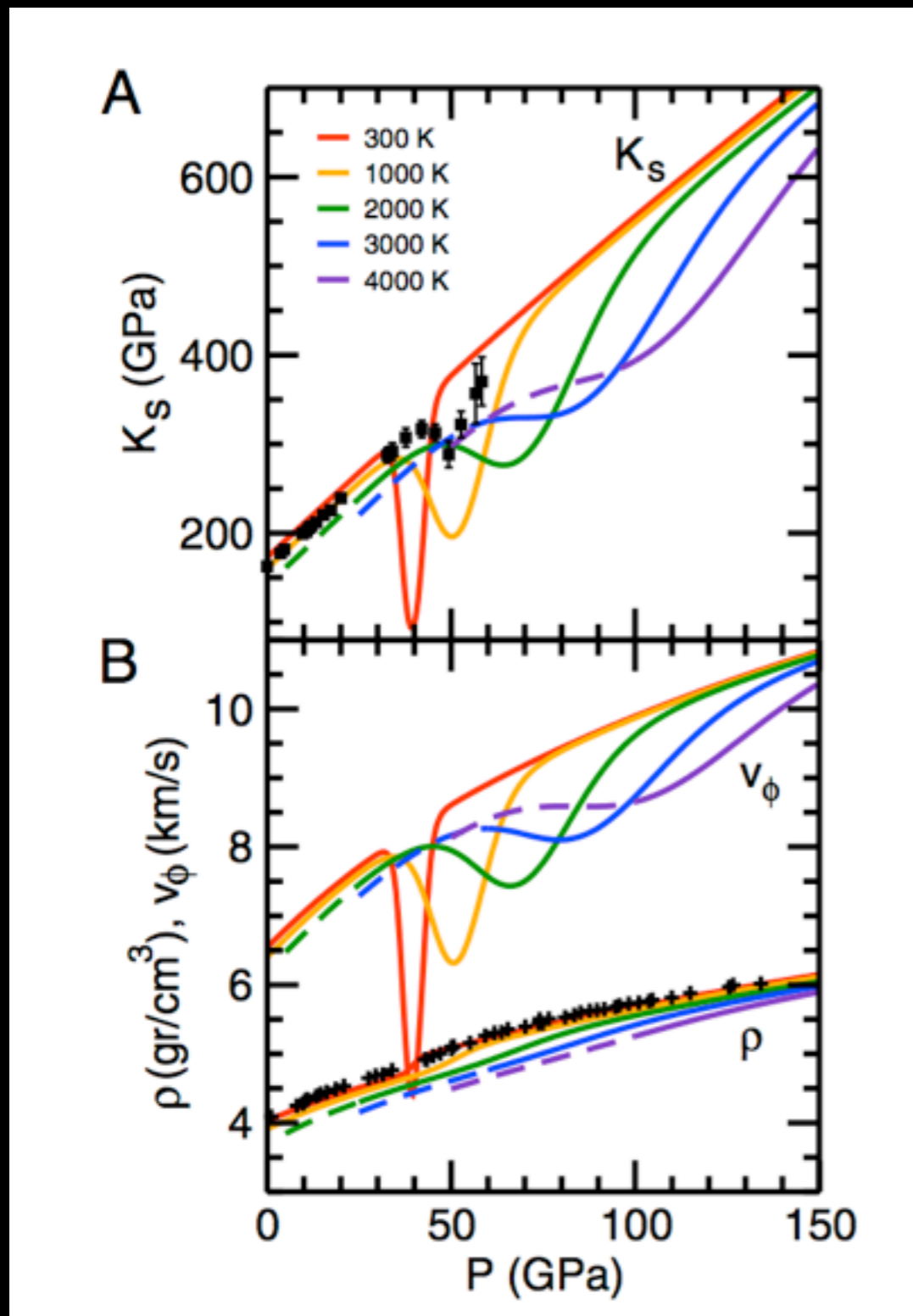


What about other components?
 Fe^{2+} , Fe^{3+} , Al^{3+} etc.

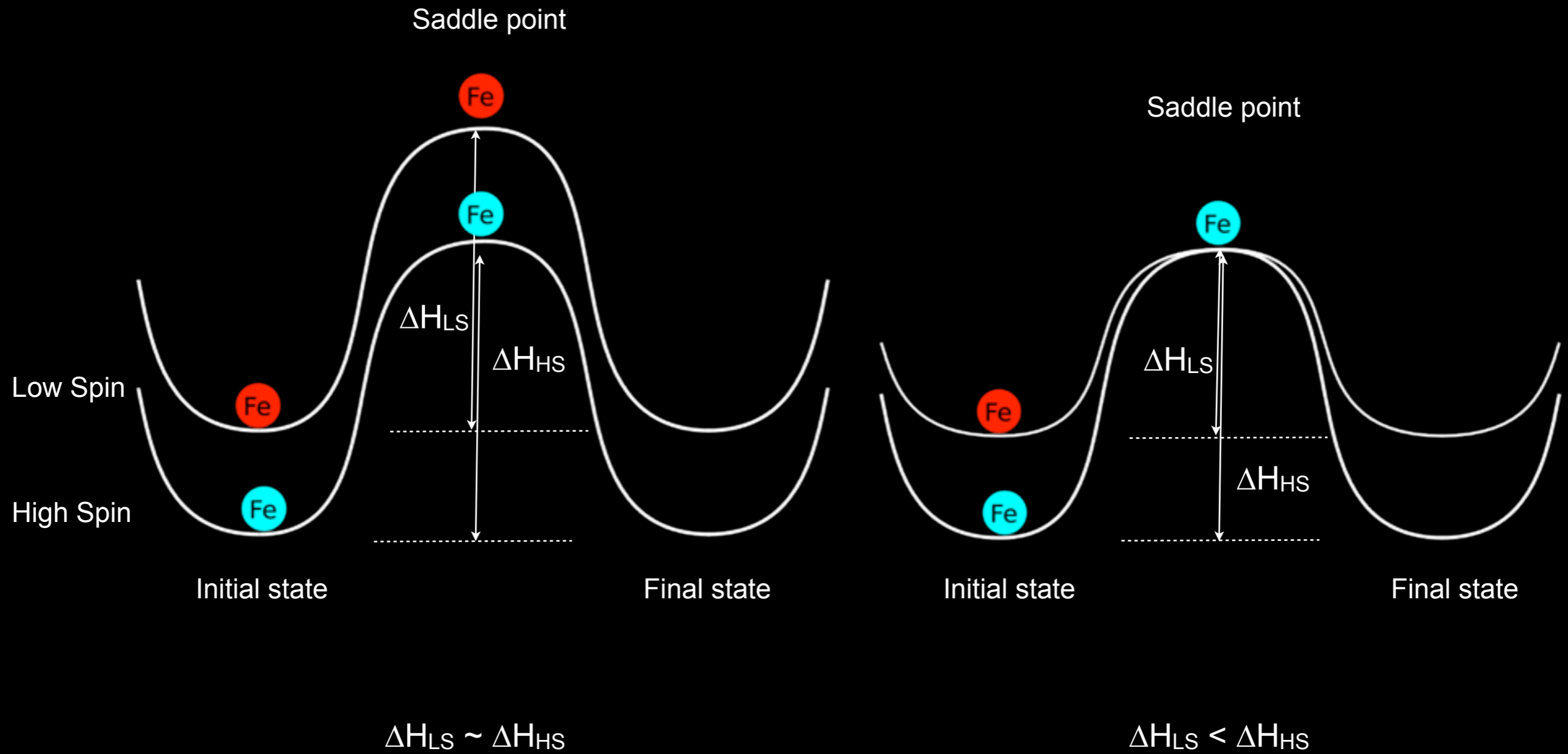
And in particular the effect of spin transition in Fe.



High - Low Spin Transition in Ferropericlase



Low Pressure

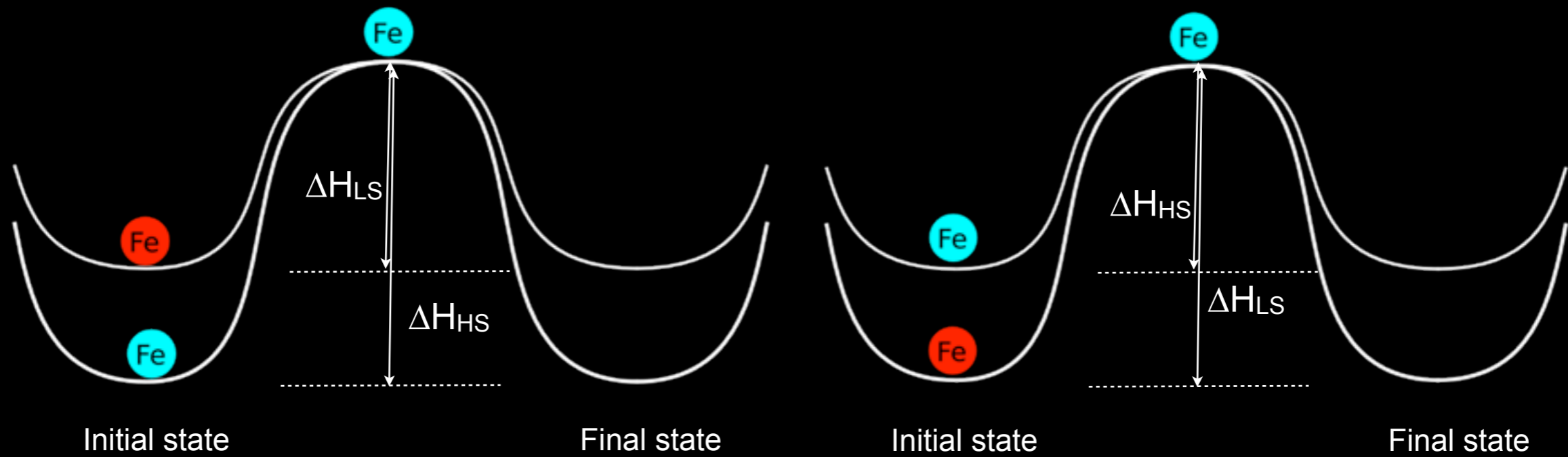


Low Pressure

High Pressure

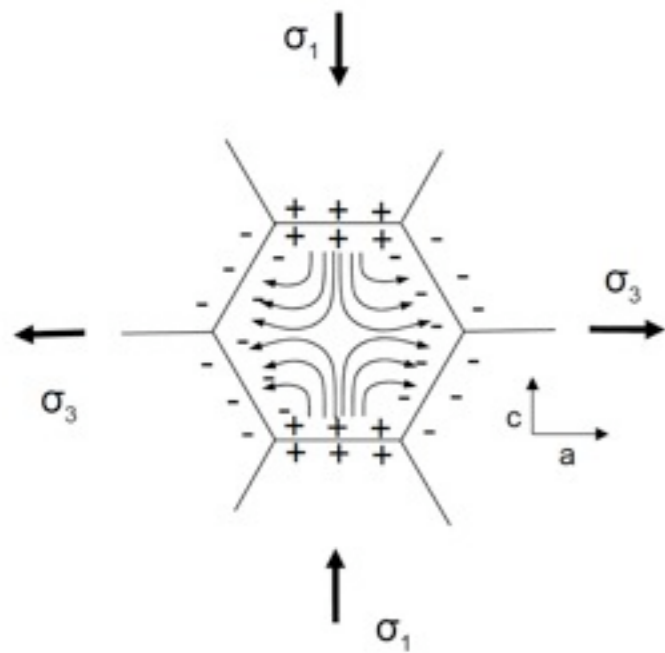
Saddle point

Saddle point



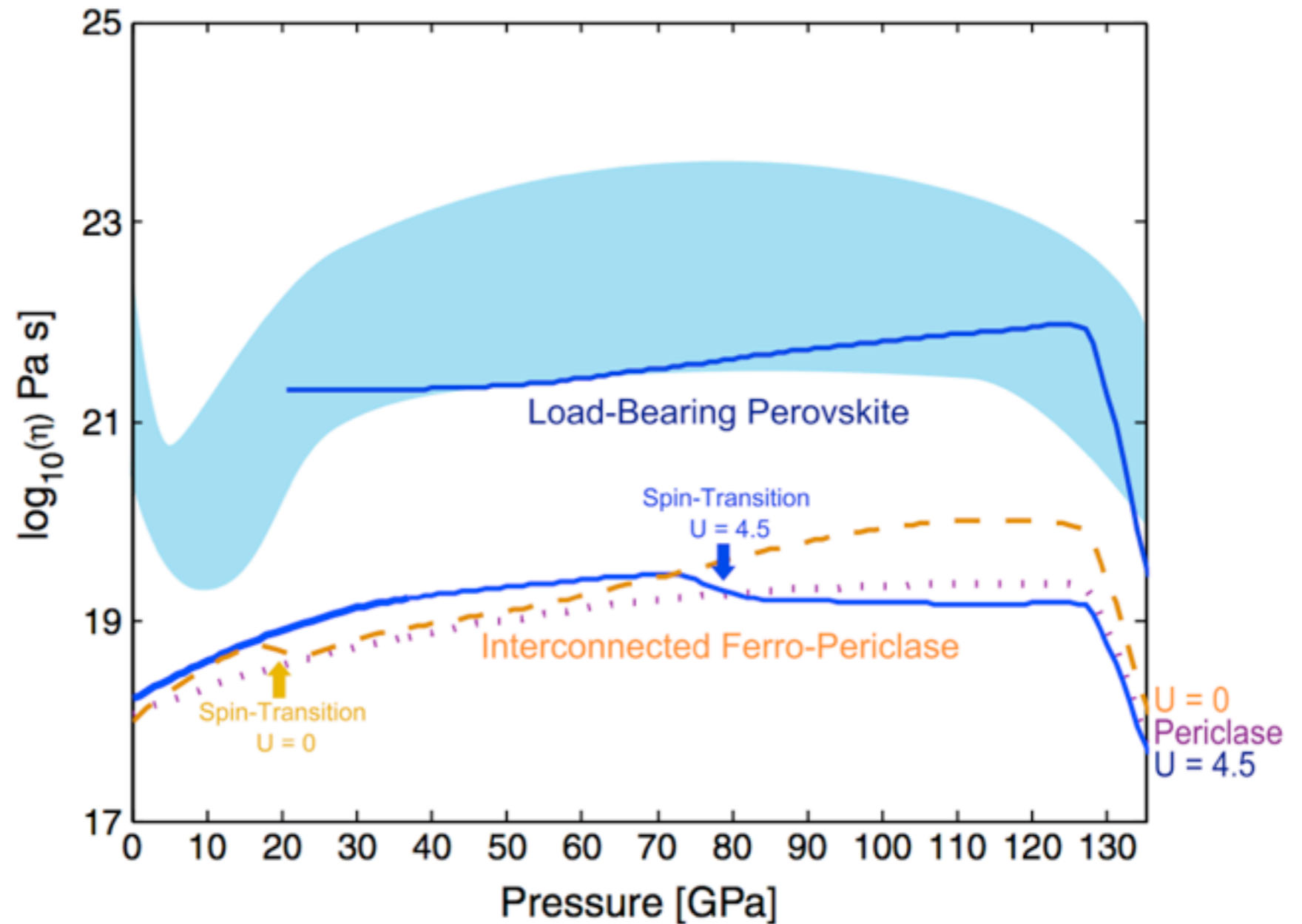
$$\Delta H_{LS} < \Delta H_{HS}$$

$$\Delta H_{LS} > \Delta H_{HS}$$

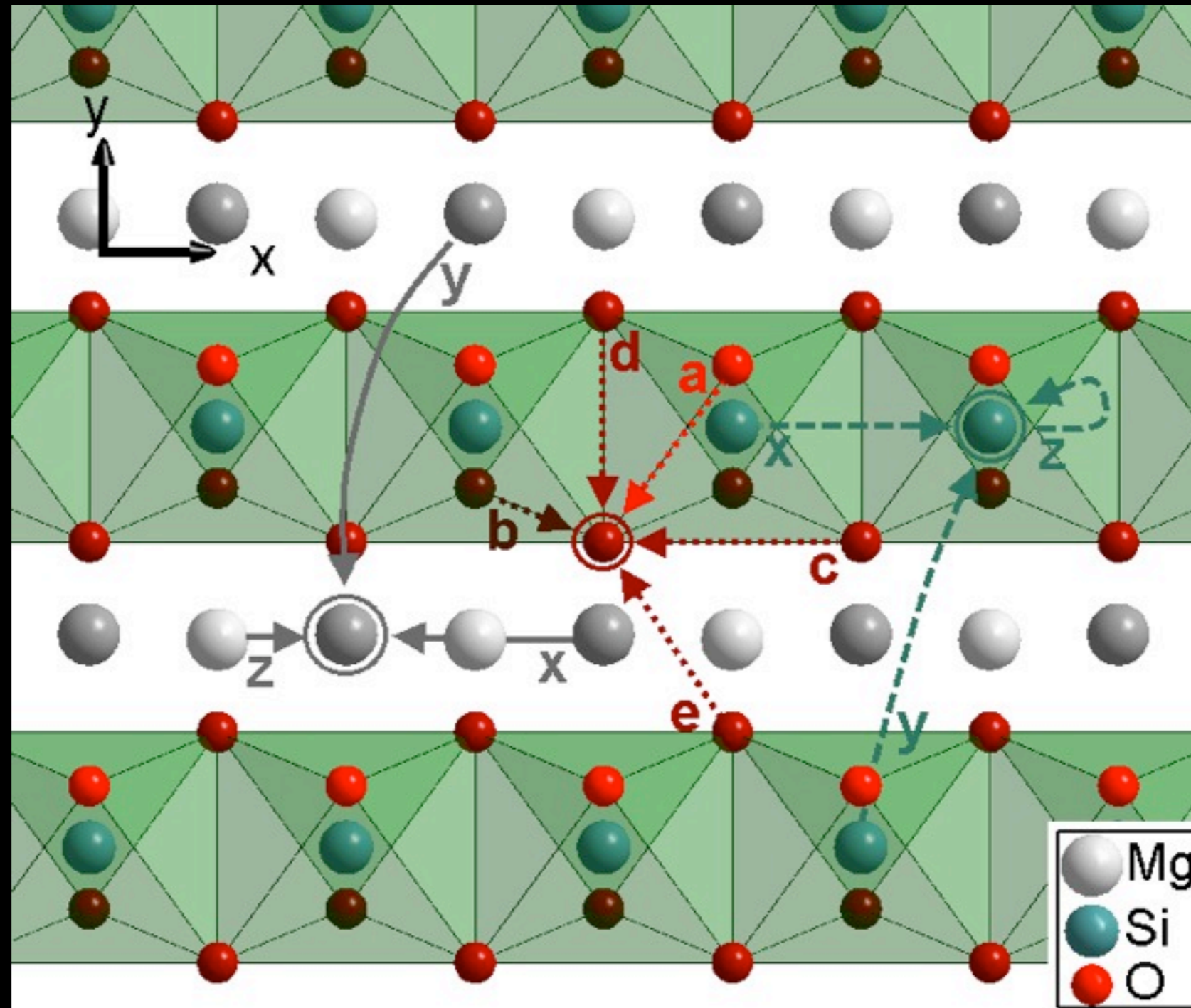


$$\eta = \frac{\sigma}{\dot{\epsilon}} = \frac{G^2 kT}{\alpha D \Omega}$$

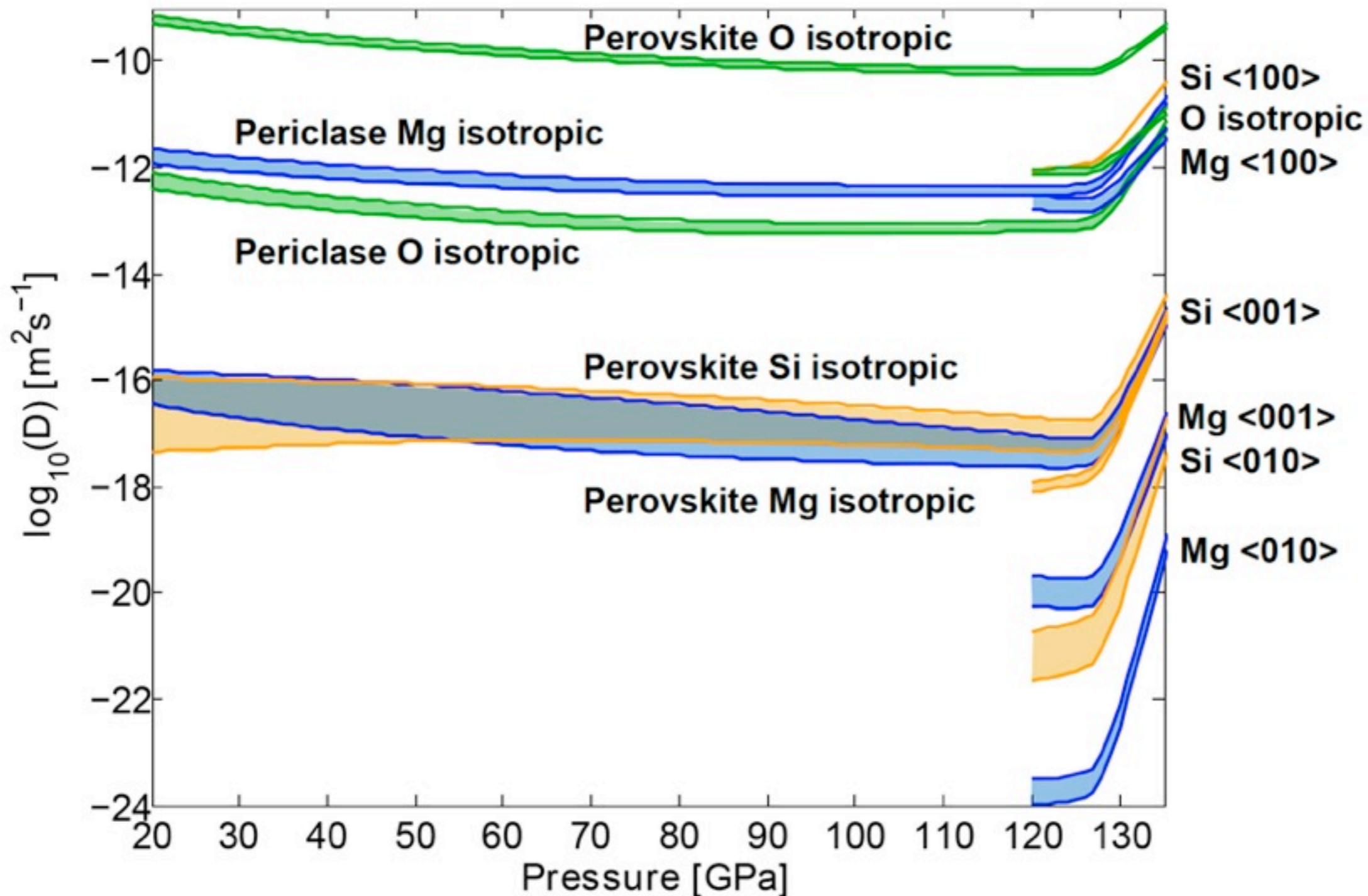
(Nabarro-Herring Creep)



Diffusion in post-perovskite

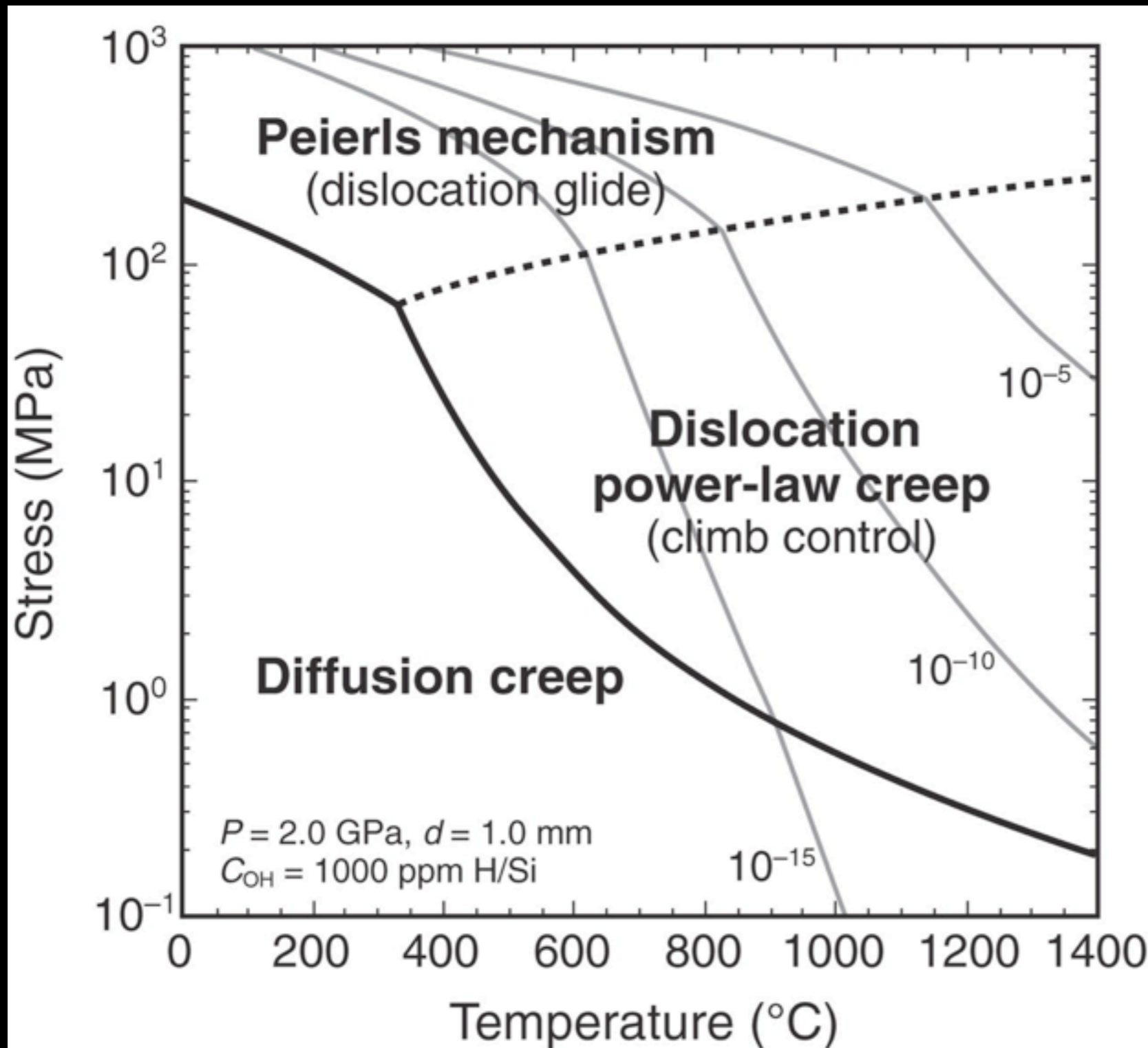


Diffusion in post-perovskite is extremely anisotropic.



Si and Mg diffusion in the $\langle 100 \rangle$ direction is very fast

Deformation Map for Olivine



Katayama, Geology (2012)

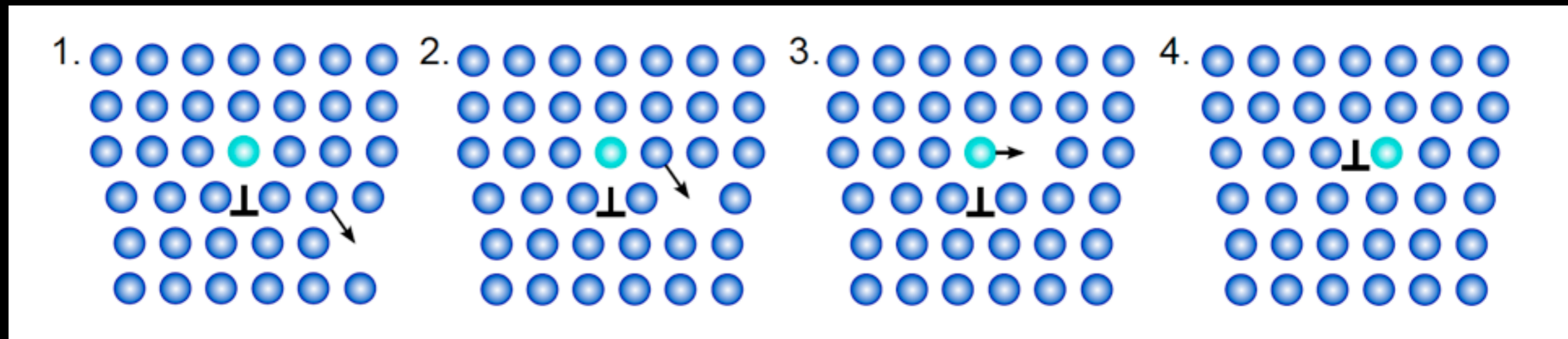
Don't know deformation map for PPV, but

If diffusion creep, PPV could be several orders of magnitude more viscous than PV.

If dislocation creep, PPV could be up to four orders of magnitude weaker than PV - if climb controlled.

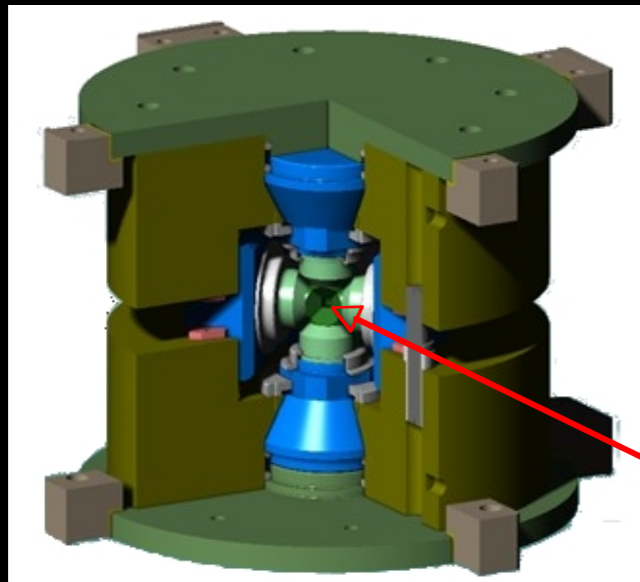
What about dislocation creep?

Dislocation creep is generally controlled by climb - and this is also a diffusion controlled mechanism.



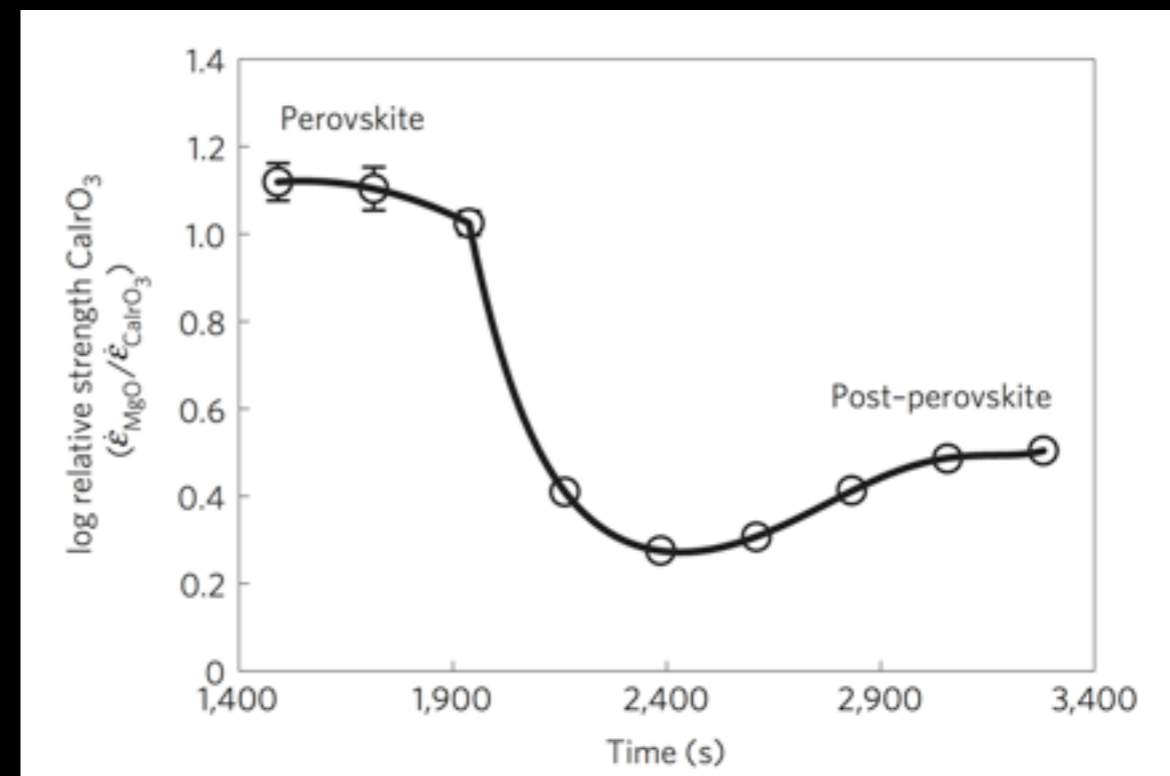
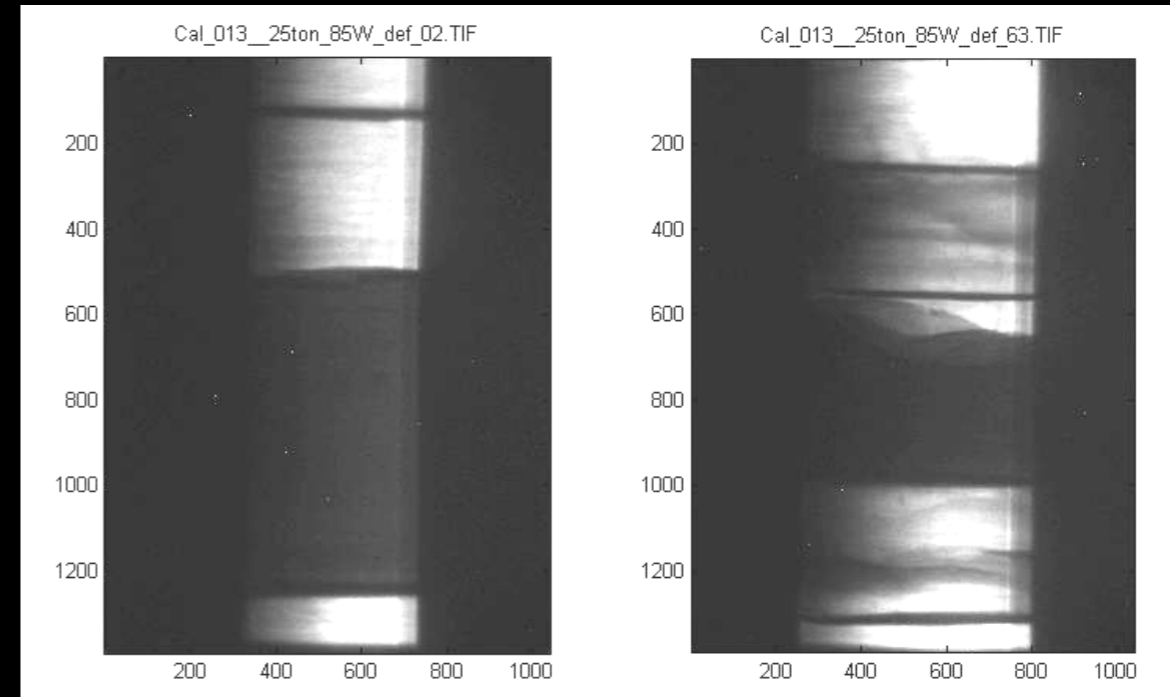
So the much faster diffusion of Si in the [100] direction in post-perovskite relative to perovskite should make it creep faster too - i.e., post-perovskite should be up to 4 orders of magnitude weaker than perovskite.

Experimental creep rates in transforming CaIrO_3

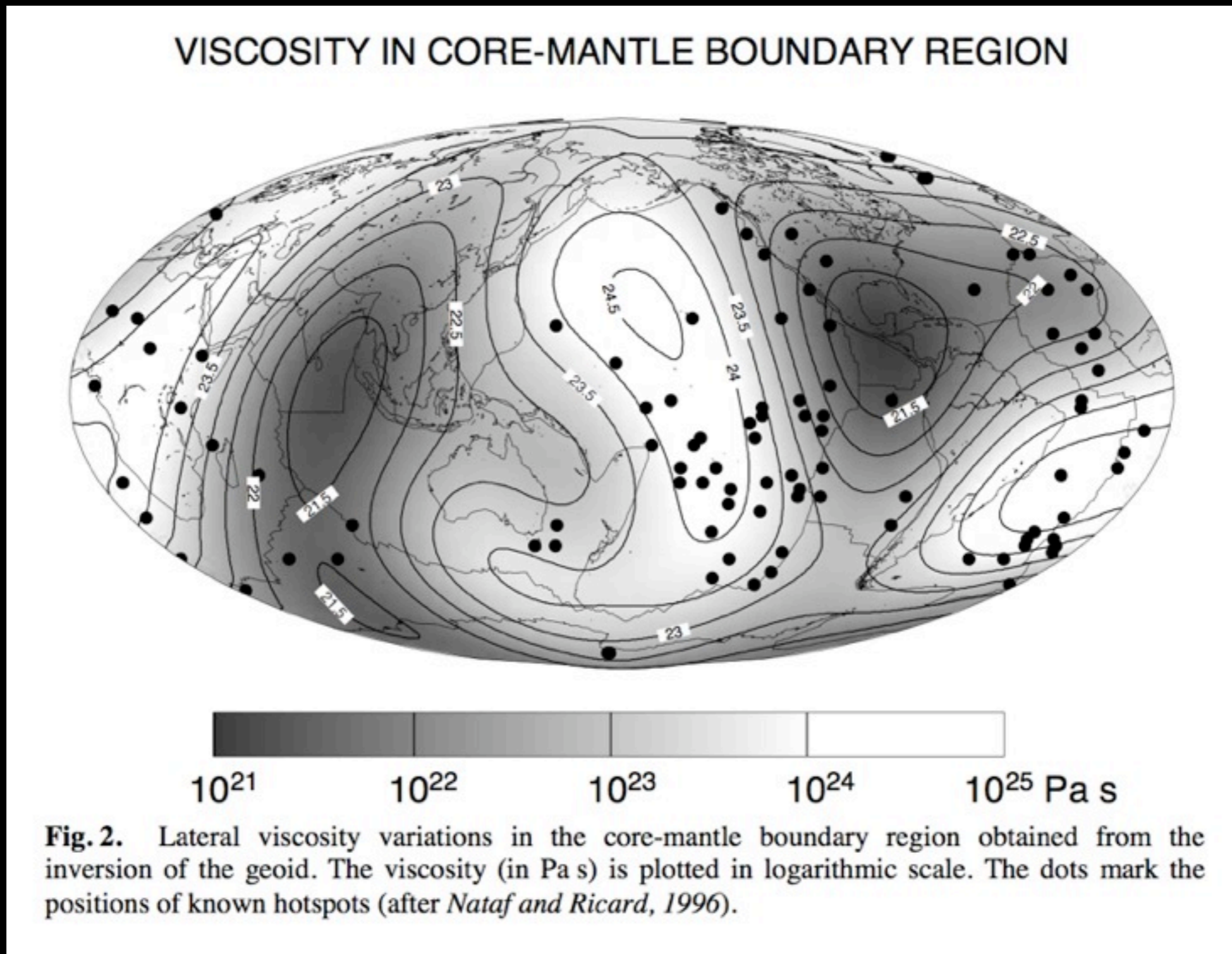


COMPRES beamline at the NSLS

Hunt et al, Nature Geoscience 2009

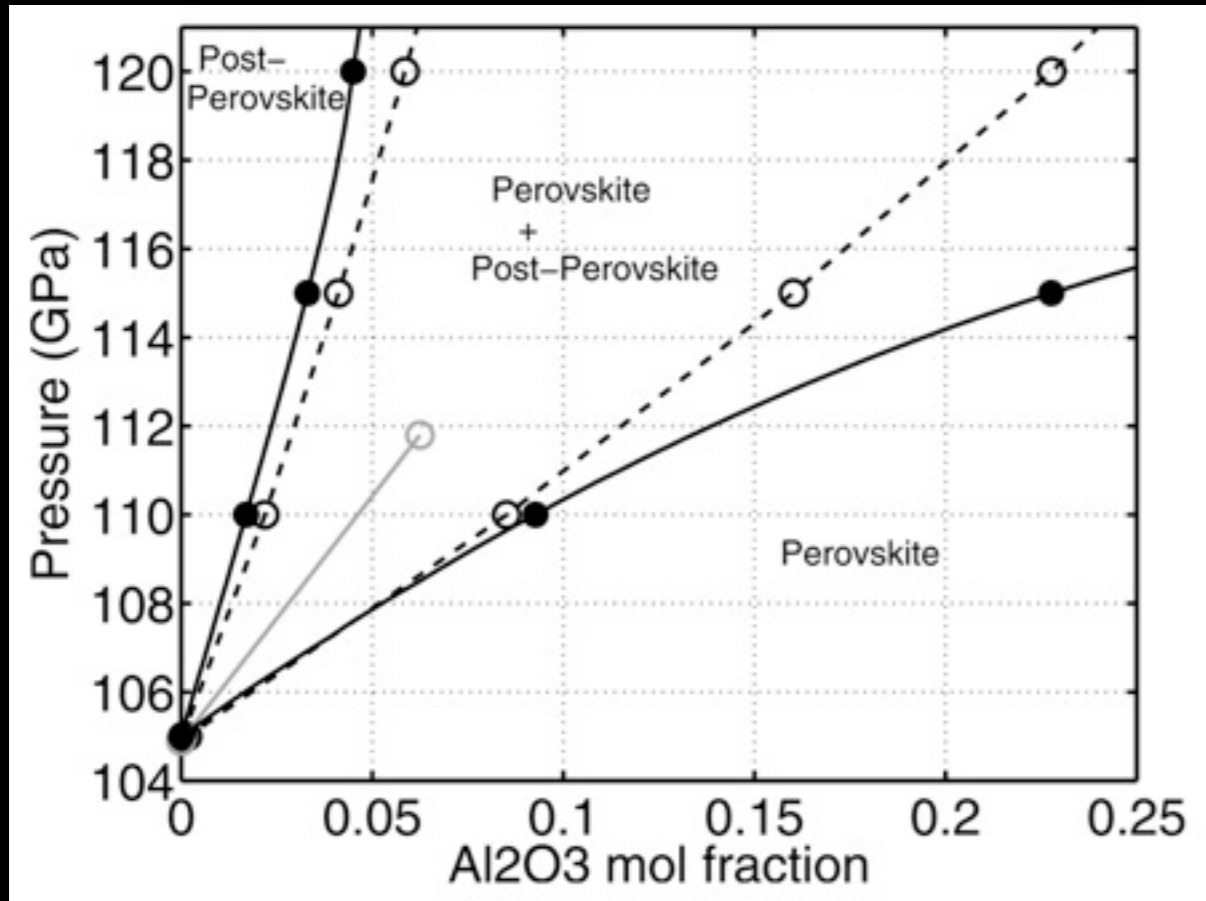


Other implications for weak post-perovskite:

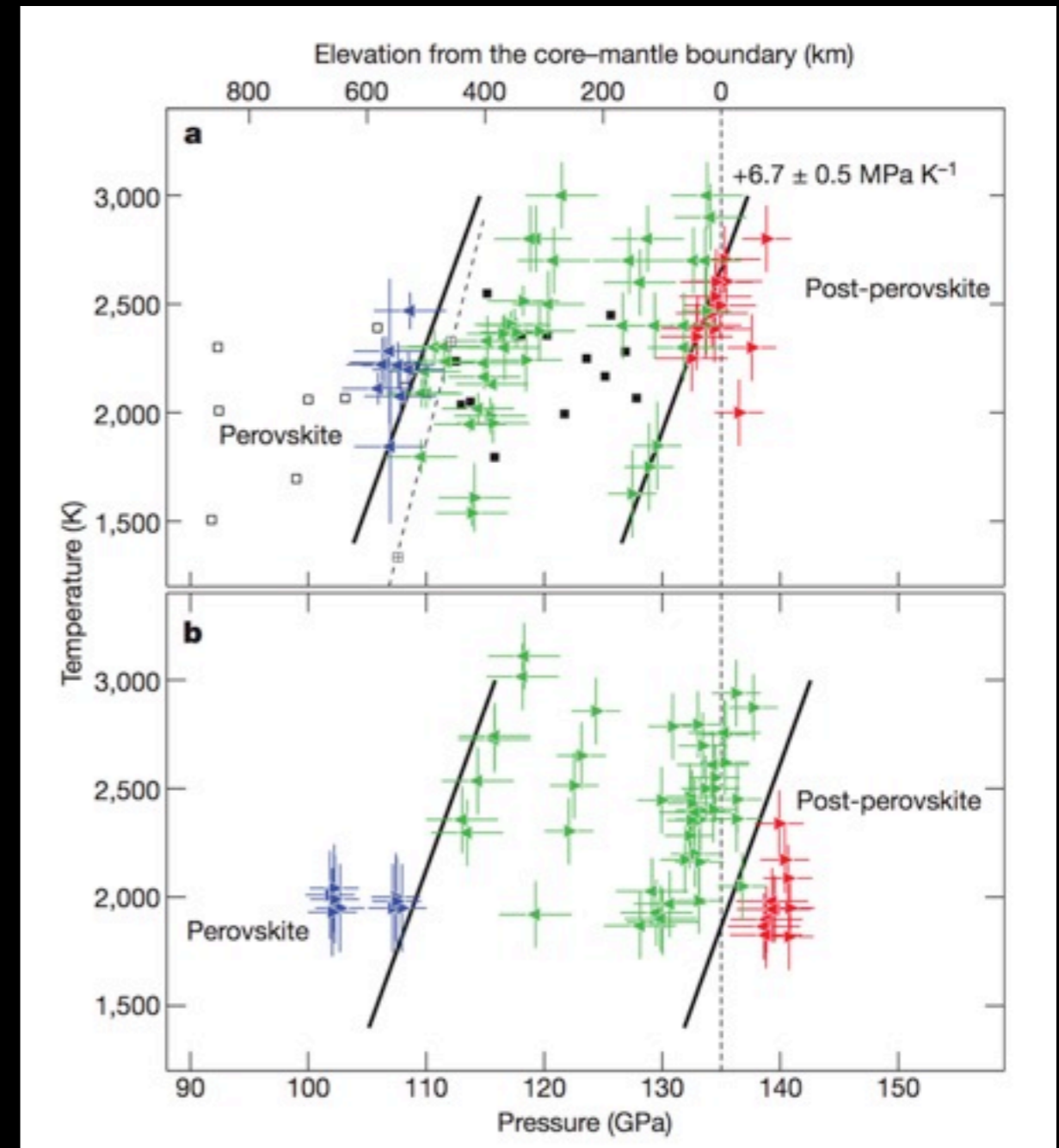


Cadek and Fleitout (2005)

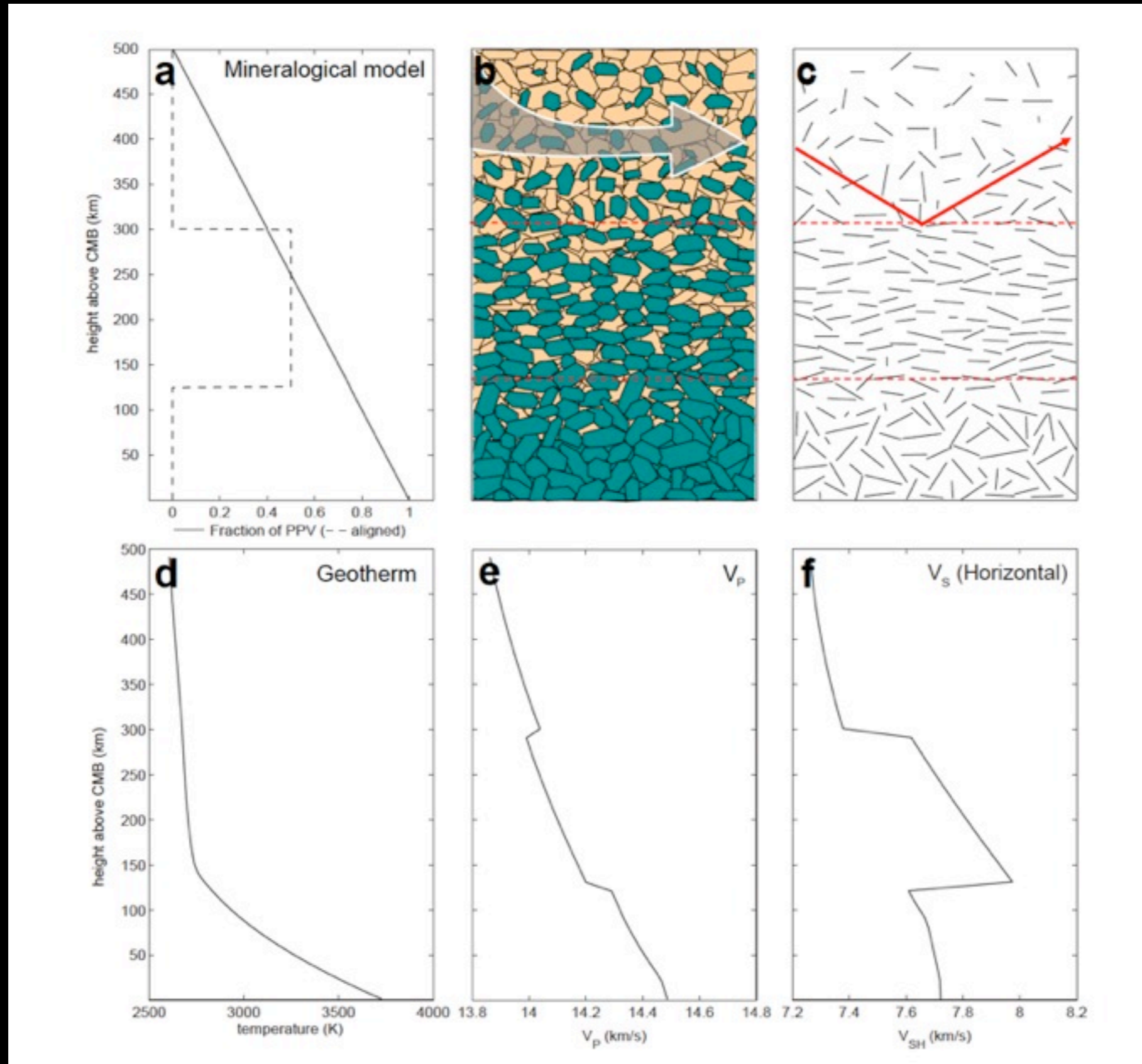
Is the phase transition too wide to be consistent with a sharp seismic refector?



Akber-Knutson et al, (2005)

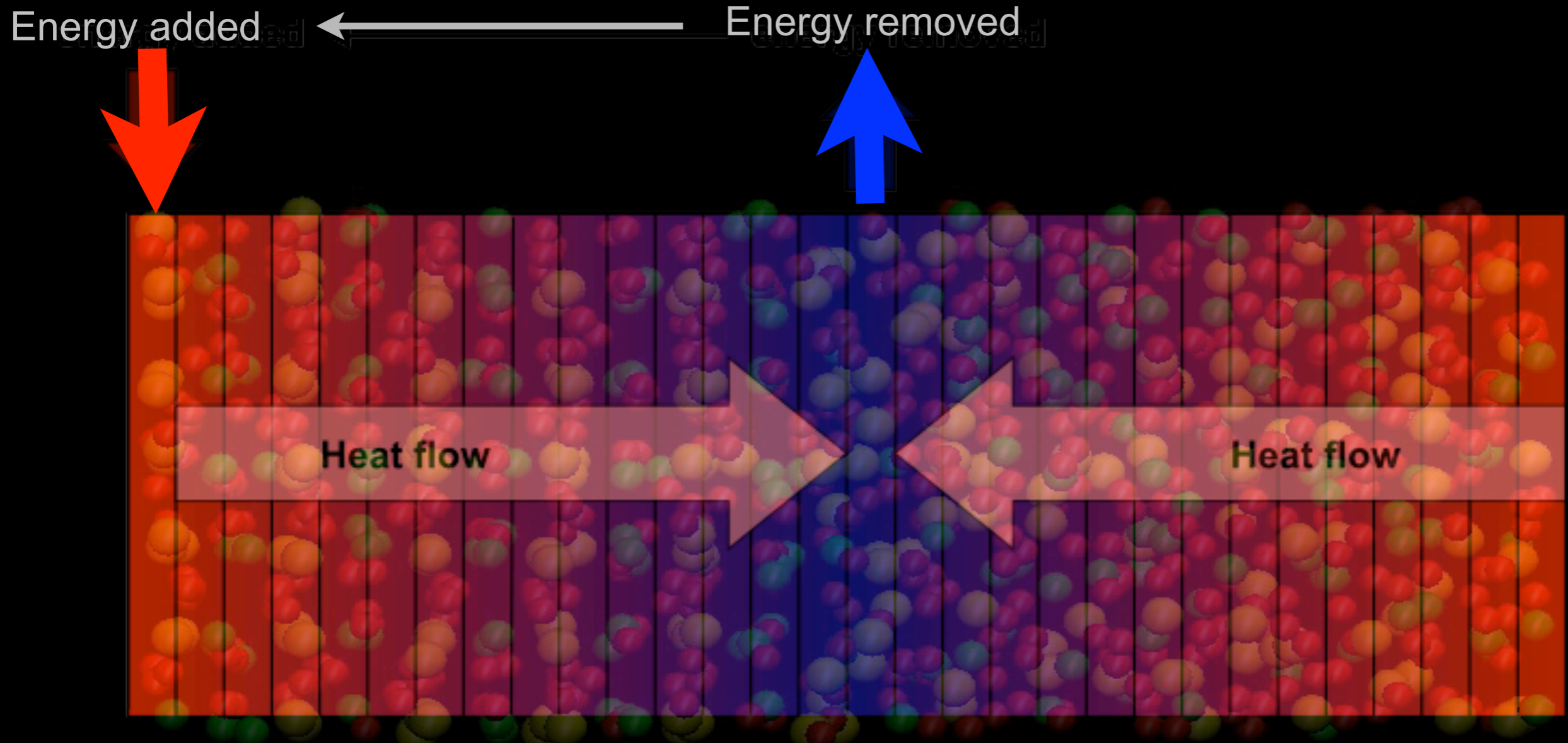


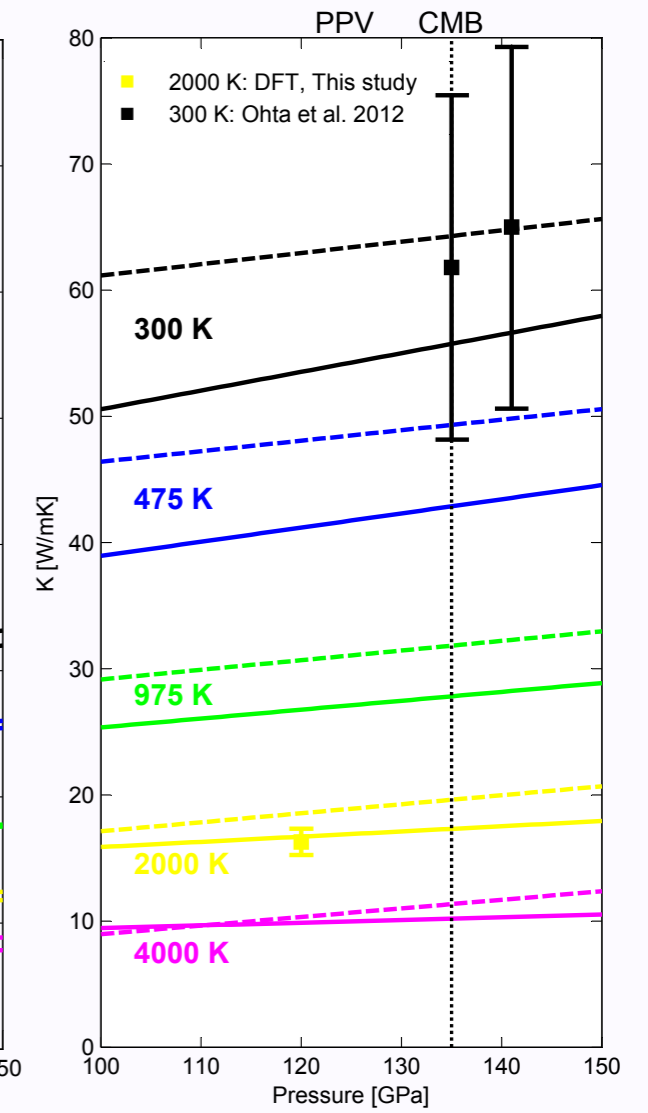
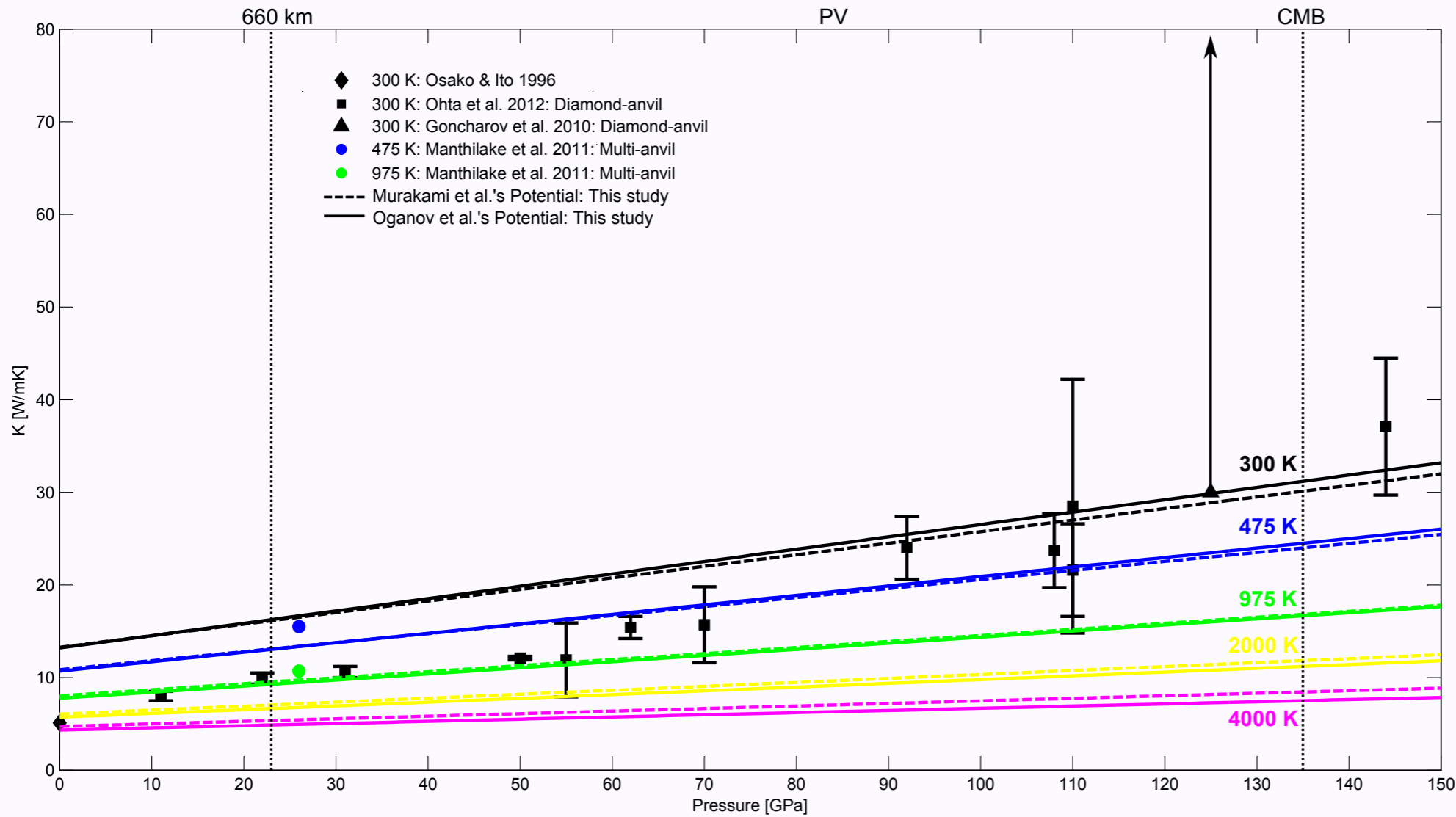
Catalli et al (2009)



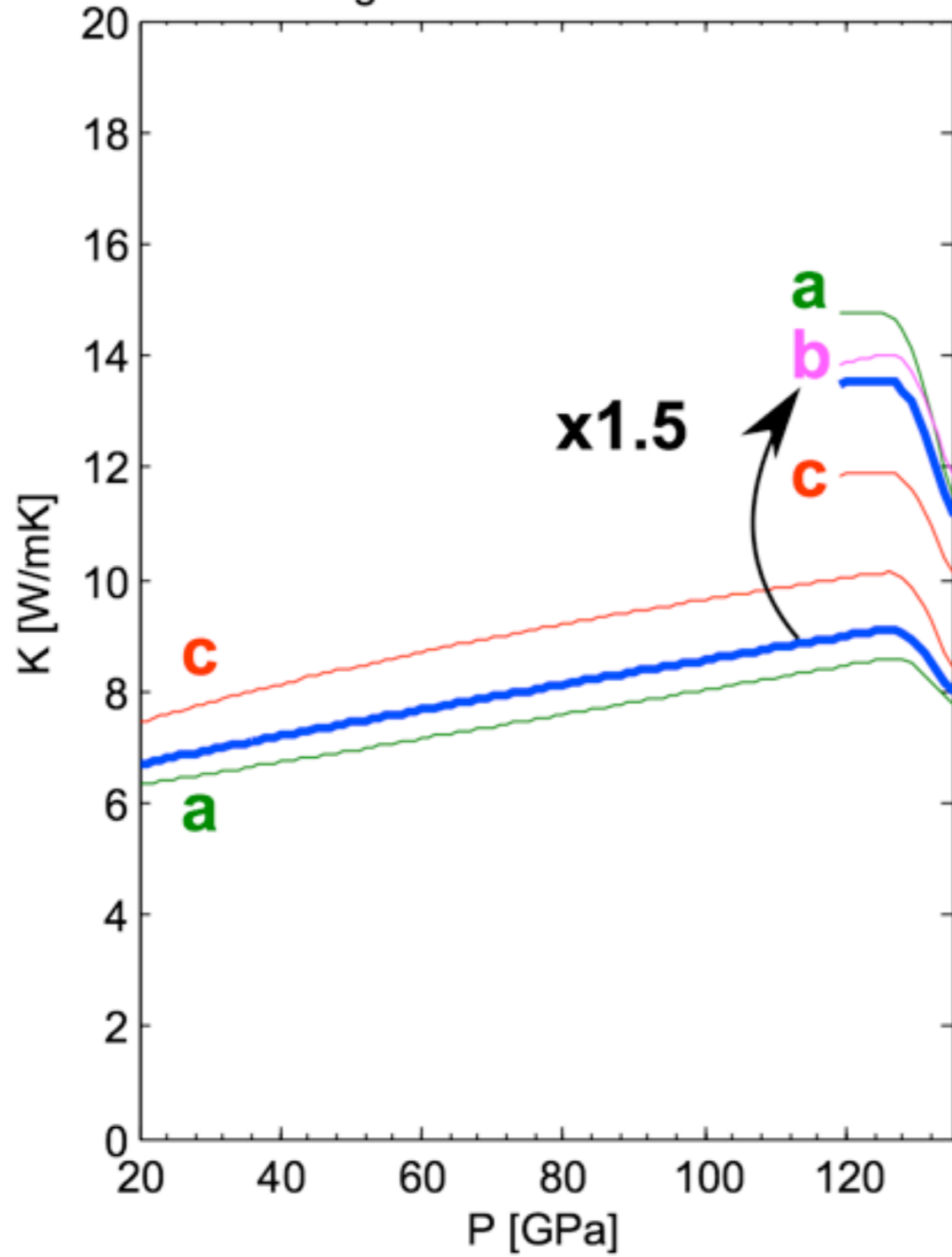
Ammann, Brodholt, Wookey and Dobson, Nature (2010)

Thermal conductivity (k)

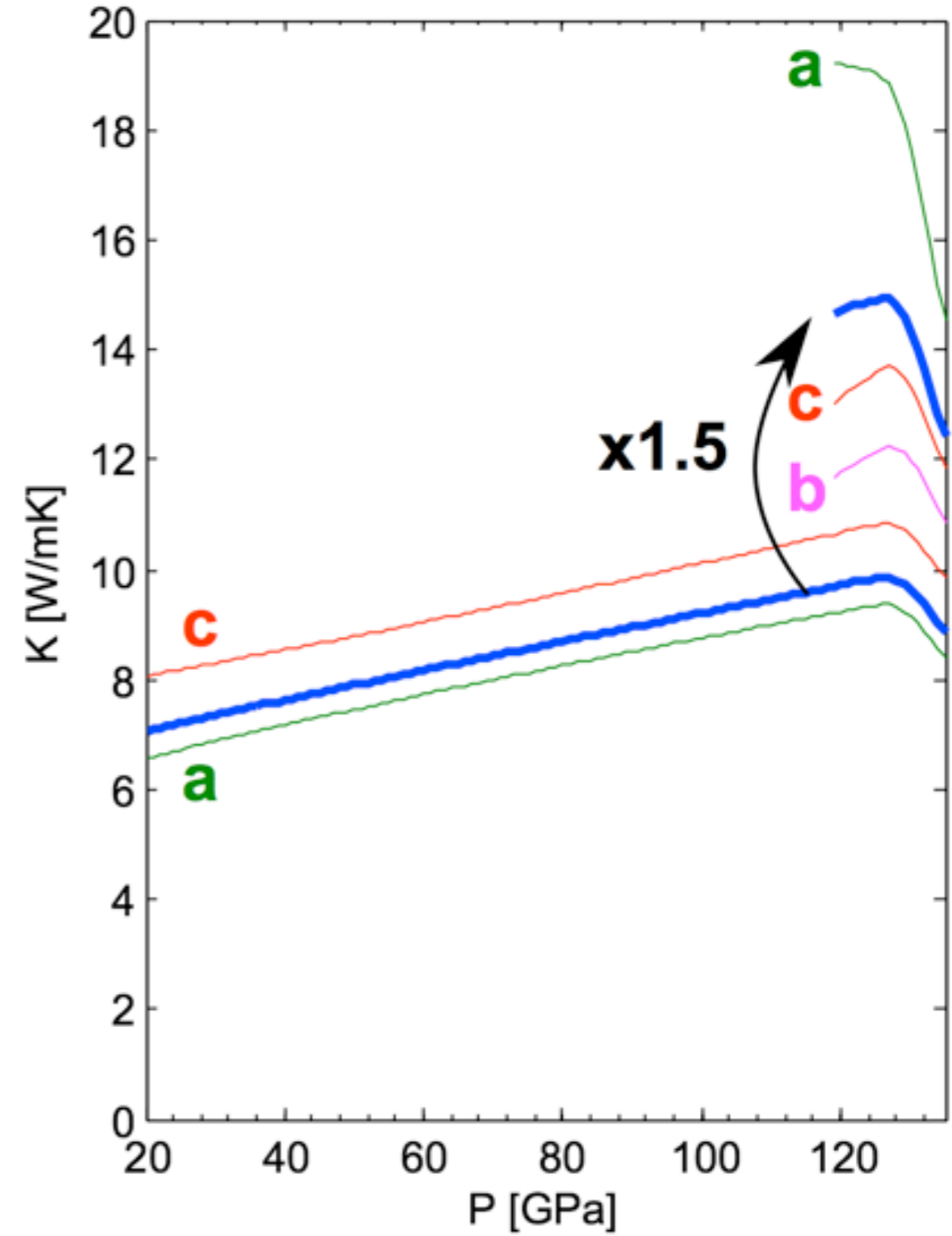




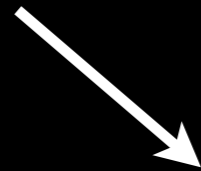
Oganov et al.'s Potential



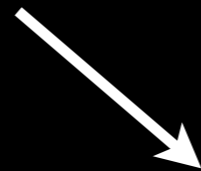
Murakami et al.'s Potential



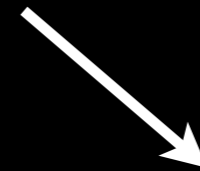
Flow model



Anisotropy (LPO)

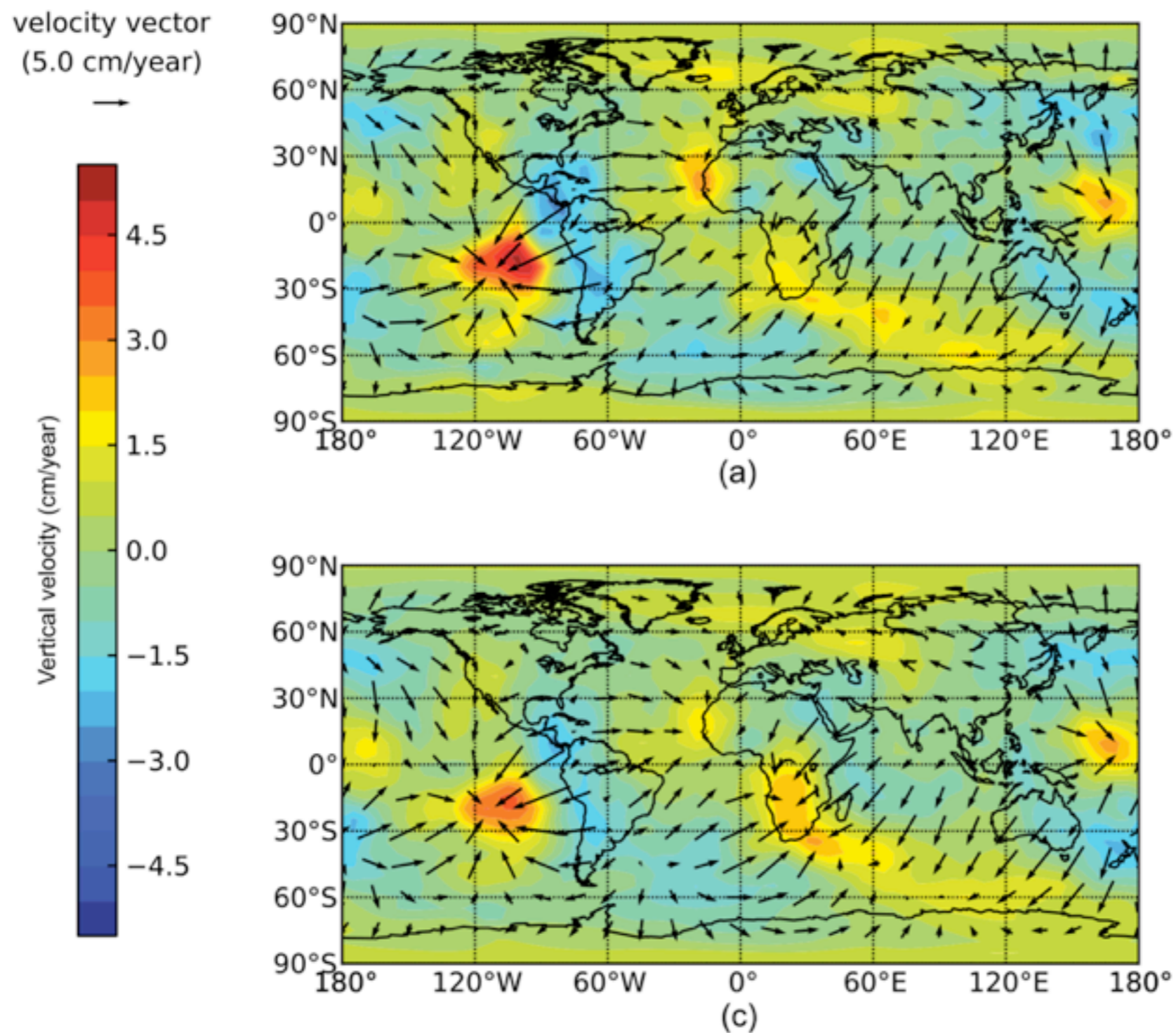


Conductive heat flow



Implications?

Flow Model



from Walker et al, 2011
and
Simmons et al
(2007,2008)

Anisotropy

Walker et al (2011)

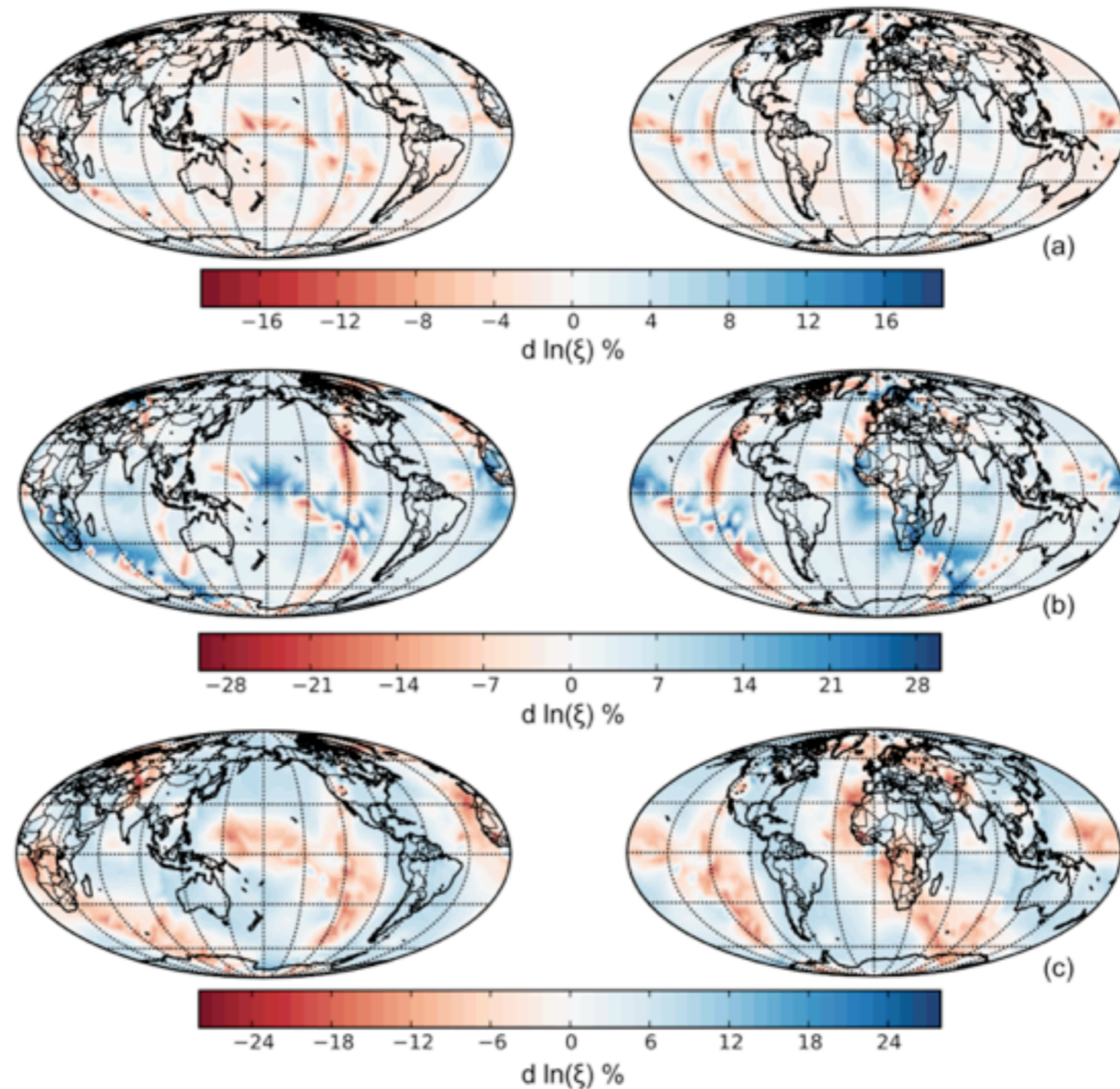
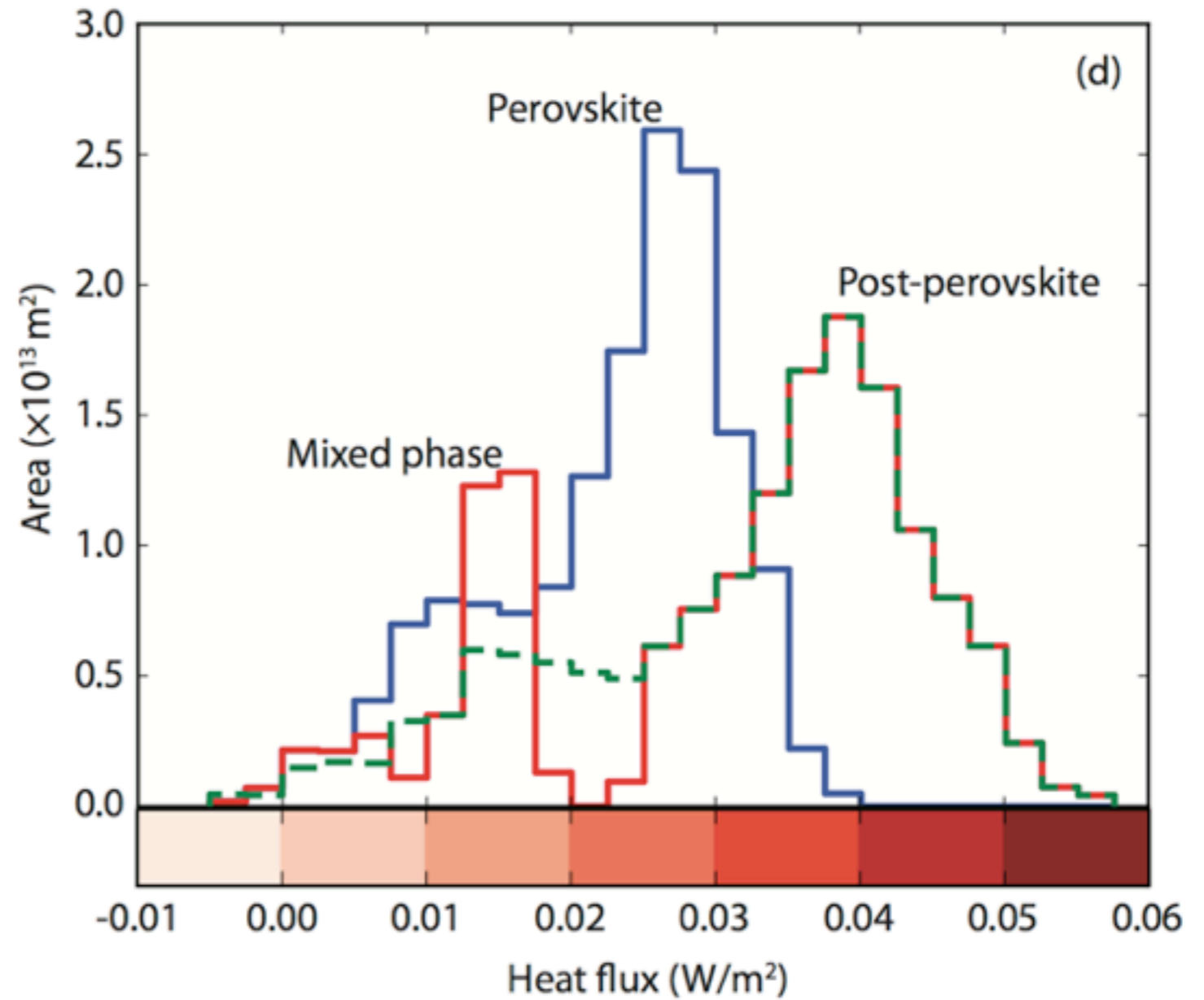
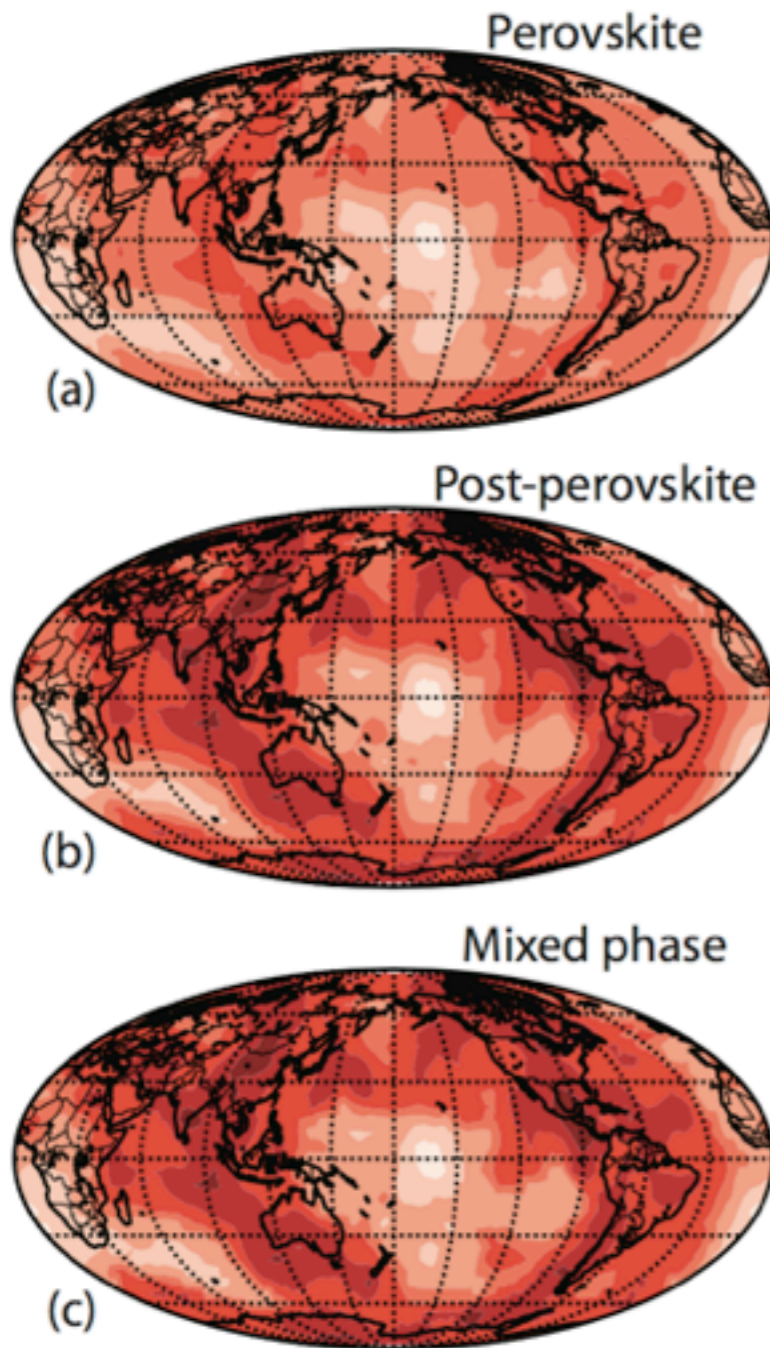


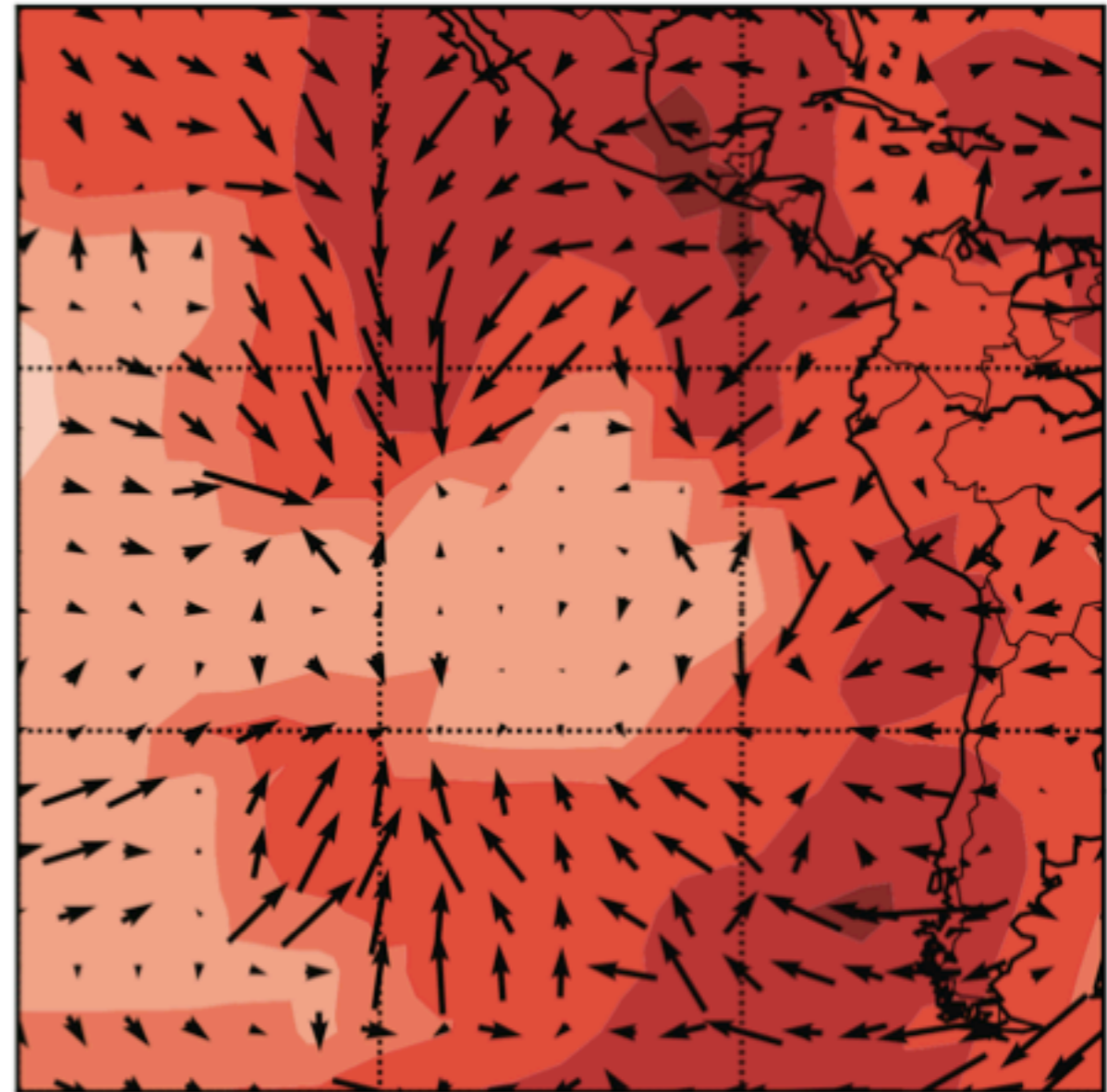
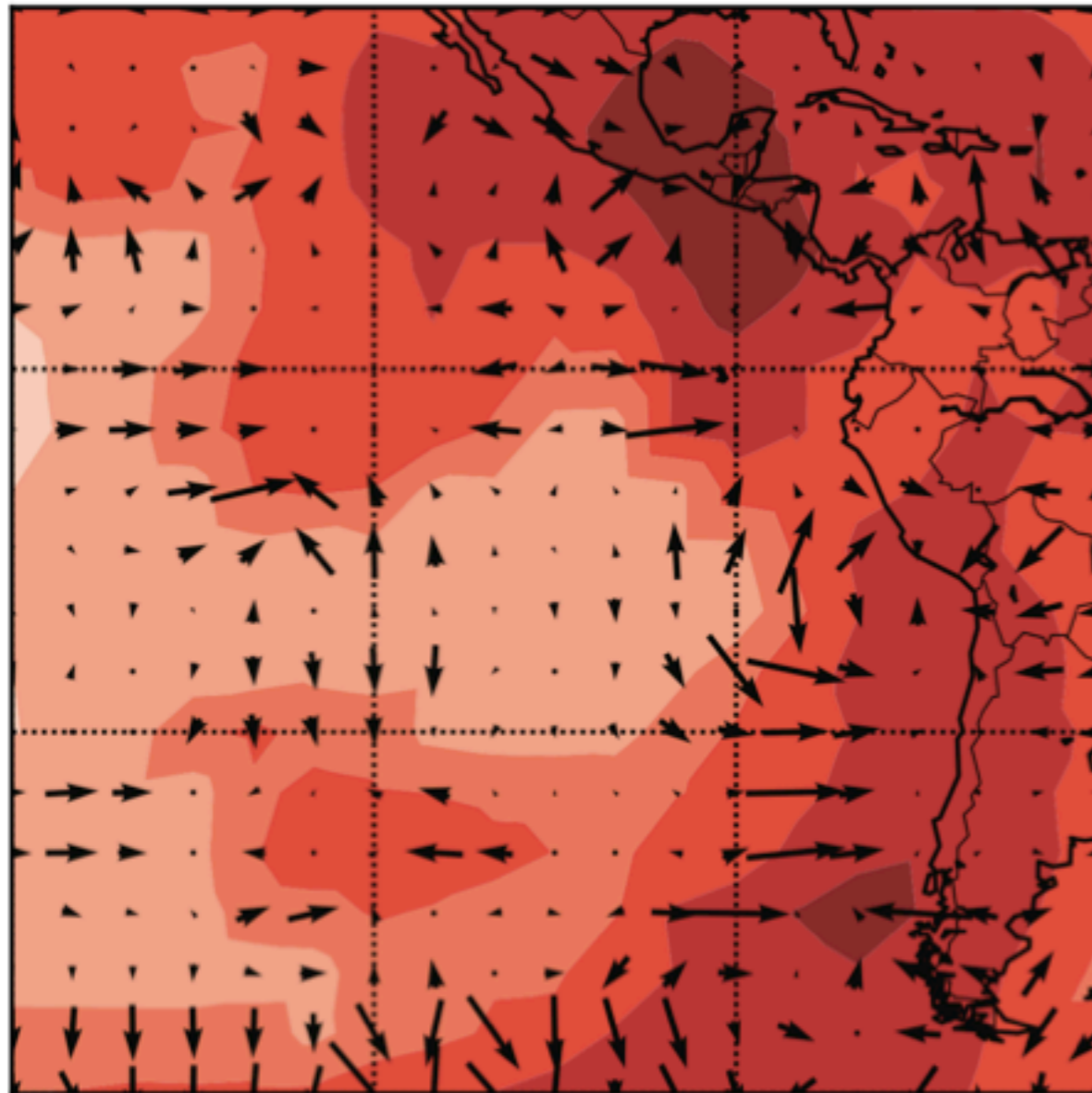
Figure 5. Calculated anisotropy, expressed as $\ln(\xi)$, for flow model TX2008.V1 75 km above the CMB with a temperature independent perovskite to post-perovskite phase transition 150 km above the CMB and three different single crystal plasticity models favoring dislocation motion on (010), (001) and (100). (a) TX2008.V1.P010; (b) TX2008.V1.P001; (c) TX2008.V1.P100.



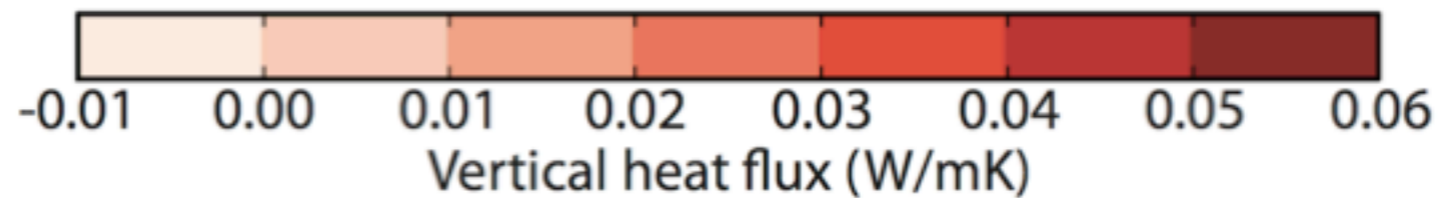
Total heat flow = 3.5, 4.9, 4.8 TW for three models

(a) Isotropic

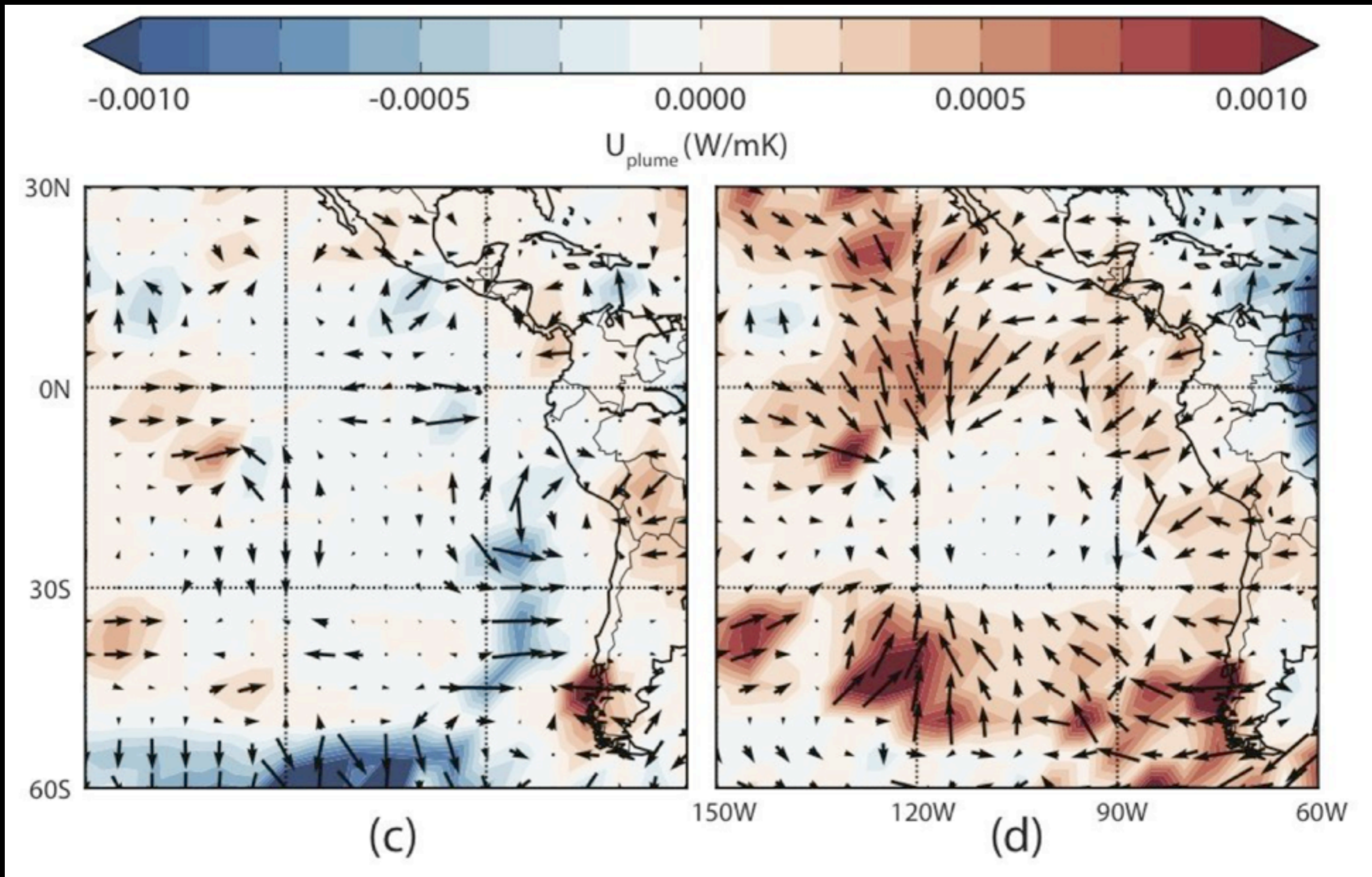
(b) Anisotropic



Horizontal heat flux
(0.005 W/mK)



Does this matter?



Conclusions

- Can use ab initio methods to calculate diffusion rates of minerals difficult to measure experimentally
- Ferro-periclase is much weaker than perovskite throughout the mantle. Spin transition slightly weakens it further.
- Lower mantle viscosity could be controlled by ferro-periclase in areas of high strains.
- Post-perovskite has very anisotropic diffusion rates and is probably much weaker than perovskite
- Post-perovskite has somewhat anisotropic conductivity, which may help stabilise plumes.