

# Interplay between SDW and PDW orders in the cuprates

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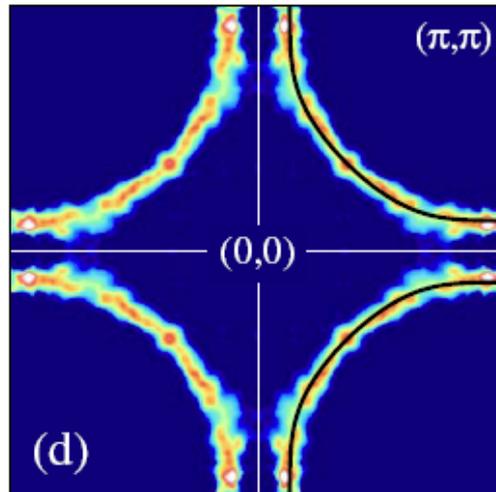
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Original motivation: to verify whether both superconductivity and charge order in underdoped, yet still metallic cuprates, come from the same spin-fluctuation exchange



It ended up with: a generic analysis of charge orders in systems with cuprate Fermi surface



# Our first interest was to analyze CDW order with $(Q,0)/(0,Q)$

following Metlitski and Sachdev  
Efetov, Pepin and Meier on  $(Q,Q)$  order  
La Placa and Sachdev on a generic CDW

- This order does develop, despite anti-nesting of fermions in blue region
- $\Delta_k$  changes sign between red and blue regions, but magnitudes are not equal.

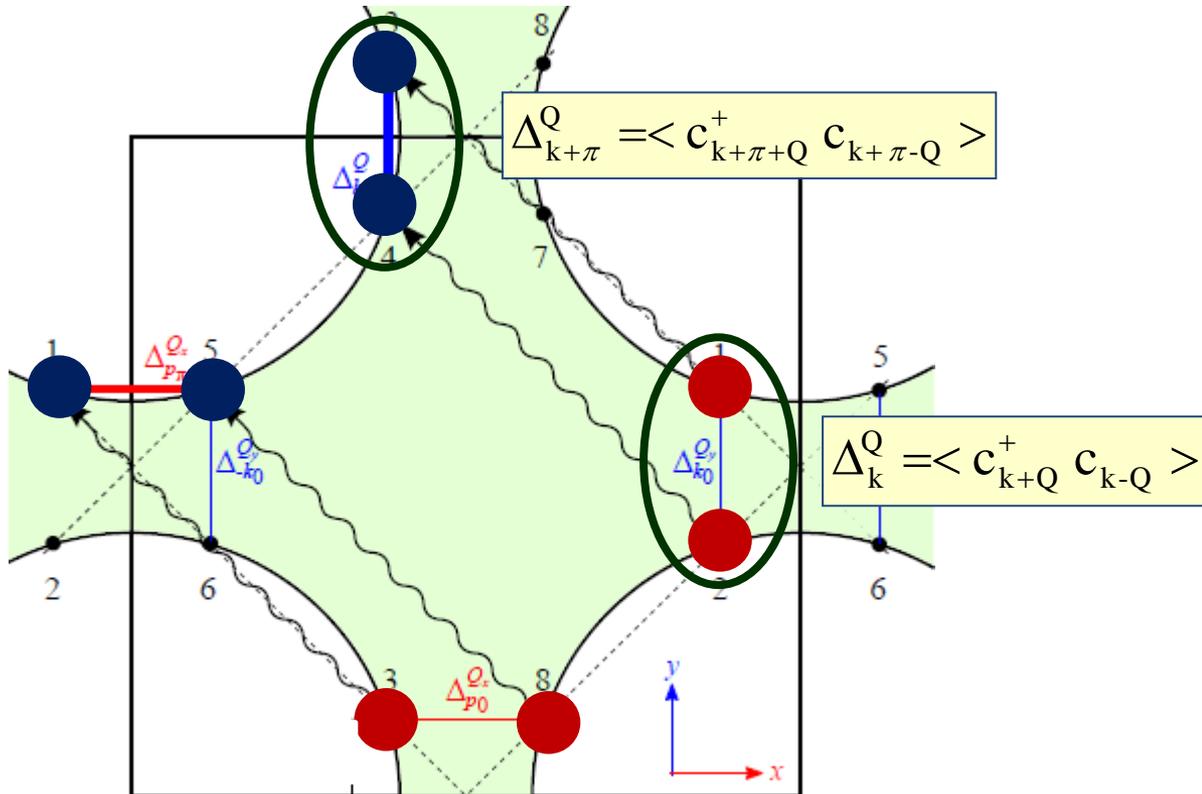
$$\text{sign}(\Delta_{k+\pi}^Q) = -\text{sign}(\Delta_k^Q)$$

$$|\Delta_{k+\pi}^Q| \neq |\Delta_k^Q|$$

The outcome:

$$\Delta_k^Q = A + \bar{A} (\cos k_x + \cos k_y) + B (\cos k_x - \cos k_y) + \dots$$

d-wave component is larger because interaction is repulsive and is peaked at large momentum transfer  $((\pi, \pi)$  in our case)



d-wave (bond order) and s-wave (CDW) are both present, d-wave is larger

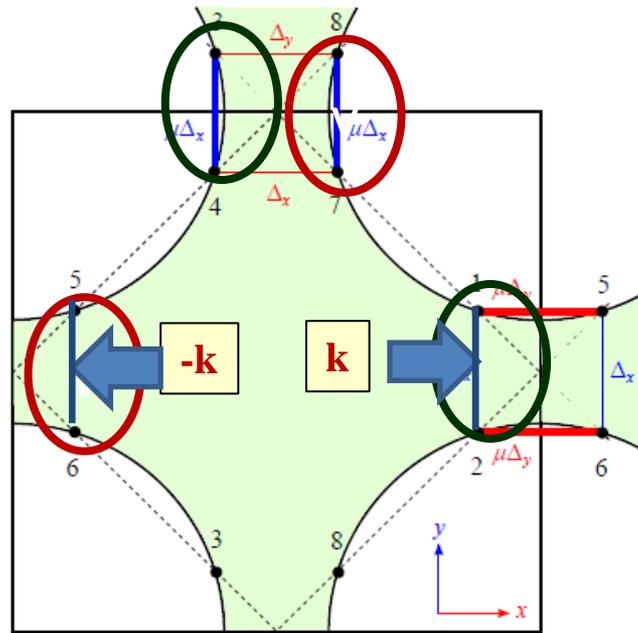
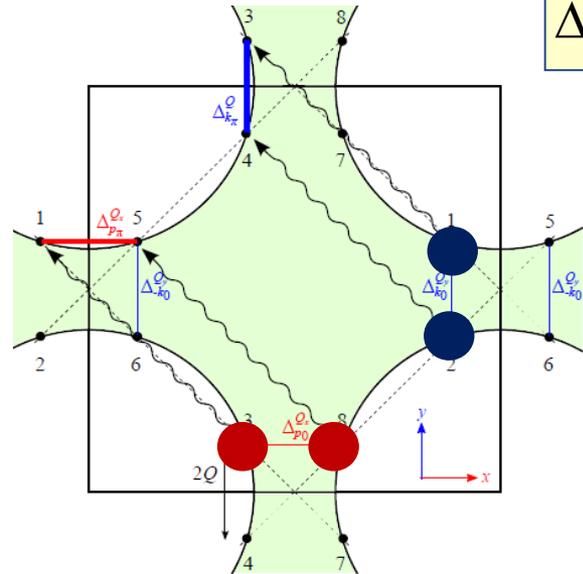
Davis, Sachdev

# There is more: $(Q,0)/(0,Q)$ order can break several symmetries

Expected: incommensurate CDW breaks translational  $U(1)$ :

$$\Delta_k^Q = \langle c_{k+Q}^+ c_{k-Q} \rangle = (\Delta_k^Q)^* = |\Delta_k^Q| e^{i\varphi}$$

It may also break  $C_4$  lattice rotational symmetry by choosing  $(Q,0)$  or  $(0,Q)$



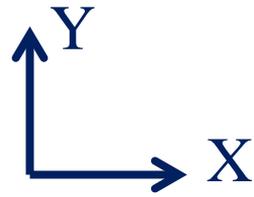
If the phases of  $\Delta_k^Q$  and  $\Delta_{-k}^Q$  differ by  $\pm \pi/2$ ,  $T$  symmetry is broken and the order has incommensurate density modulation and incommensurate current

$$\widehat{T} \Delta_k^Q = c_{-k+Q}^+ c_{-k-Q} = \Delta_{-k}^Q$$

Less expected: CDW can also break time-reversal symmetry

Real space picture for stripe CDW with

$$\Delta_{-k}^Q = i \Delta_k^Q$$



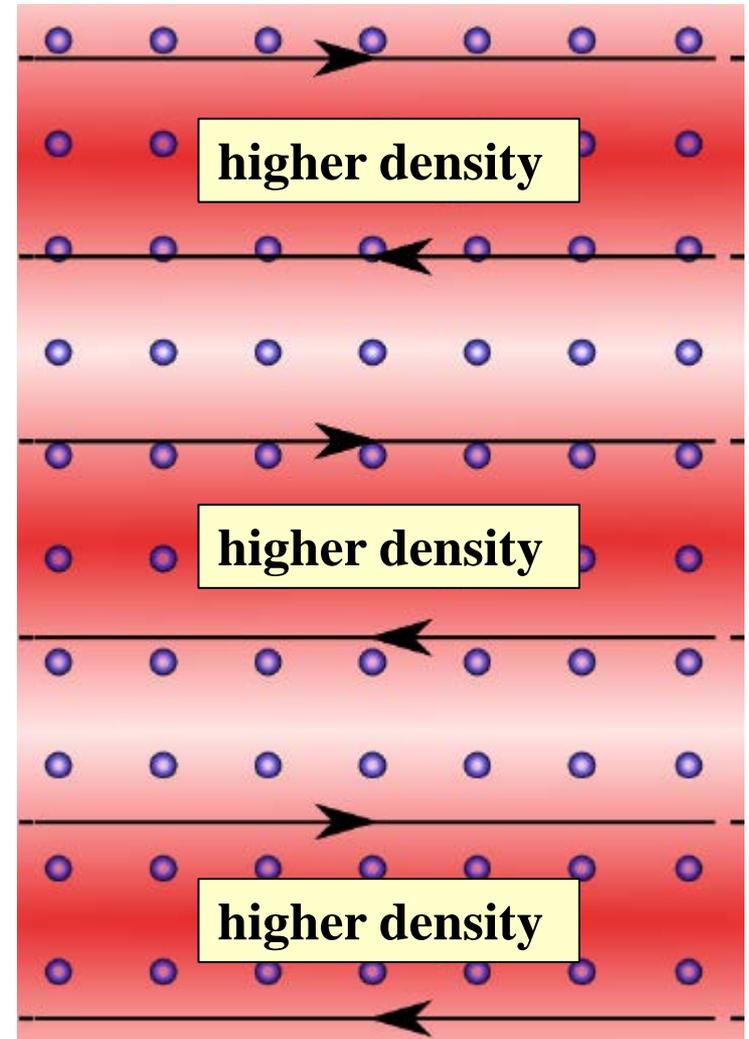
$$\delta\rho = \Delta_k^Q + \Delta_{-k}^Q \propto \cos Qy,$$

$$j_x = \Delta_k^Q - \Delta_{-k}^Q \propto \sin Qy$$

$$Q = (0, Q)$$

Such an order also breaks  
X and Y mirror symmetries  
and give rise to Kerr effect

It generates incommensurate  
magnetic field but no bulk  
homogeneous magnetic field





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A possibility that CDW order breaks additional discrete symmetries is quite encouraging as there is exp evidence for discrete symmetry breaking near the onset of CDW

- NMR, structural, X-ray, STM, ARPES measurements – incommensurate static charge order  $U(1)$

Julien, Proust, Le Tacon, Chang, Davis, Damascelli

- resistivity anisotropy measurements – lattice rotational  $C_4$  symmetry is broken down to  $C_2$   $Z_2$

Taillefer, Kivelson, Tremblay

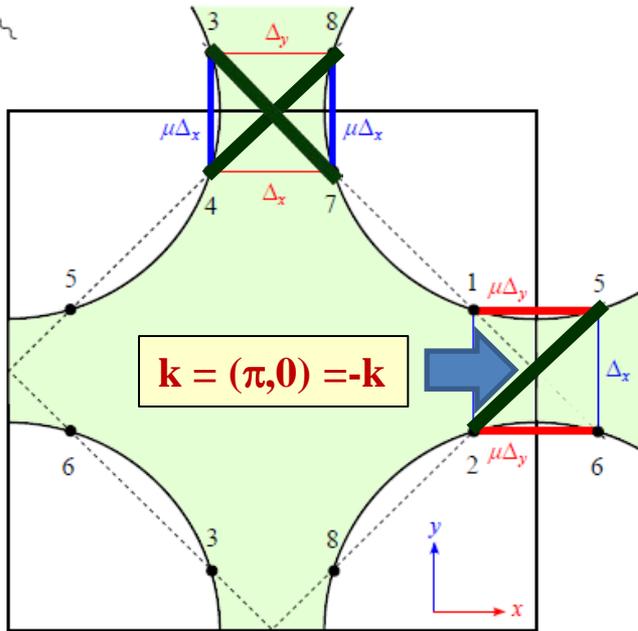
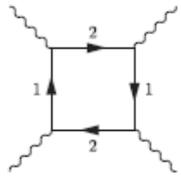
- Kerr effect and neutron scattering measurements  $Z_2$   
time-reversal symmetry is broken

Kapitulnik, Bourges

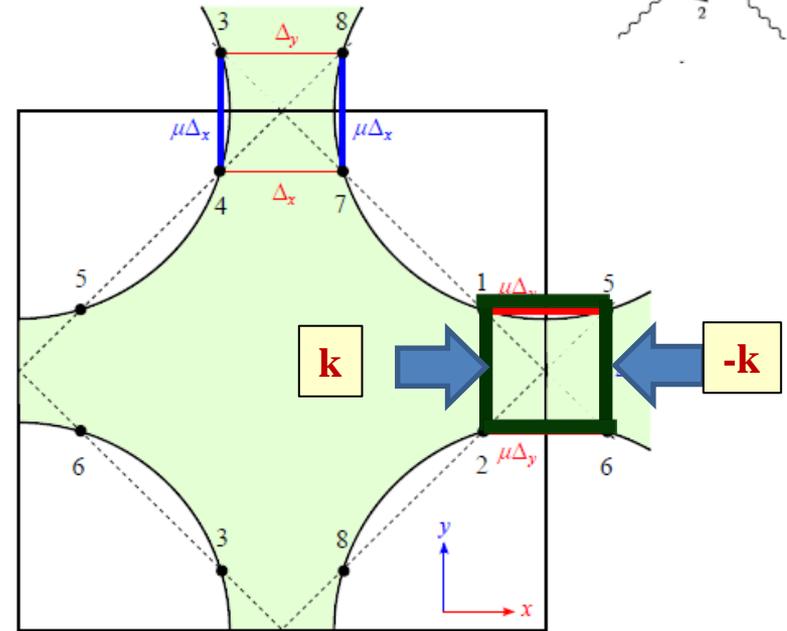
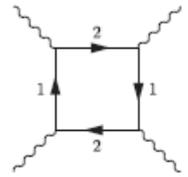
# No discrete symmetry breaking for bond order with (Q,Q)

Orders with (Q,Q) and (-Q,Q) do not couple to each other, hence no  $C_4$  symmetry breaking between the two

(Q,Q) and (0,Q) orders do couple to each other, and the coupling, if strong enough, leads to  $C_4$  breaking (stripes)



No T-breaking

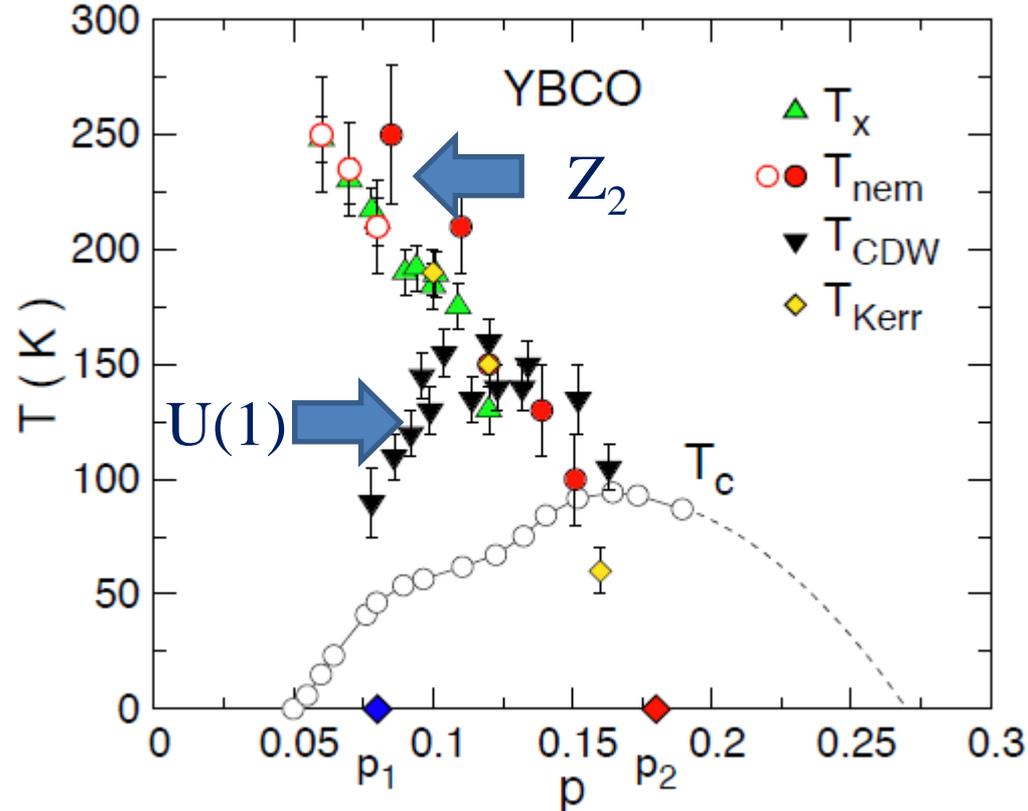


Allows for T-breaking

A generic reasoning: if the ordered state breaks both continuous and discrete symmetries ( $U(1)$  and  $Z_2 * Z_2$  for CDW), discrete symmetries generally get broken at higher  $T$  and at intermediate  $T$  the system has a “nematic” order

Kivelson

At least at some dopings discrete symmetry breaking has been observed at higher temperatures than static CDW



# Let's see what theory says about discrete symmetry breakings

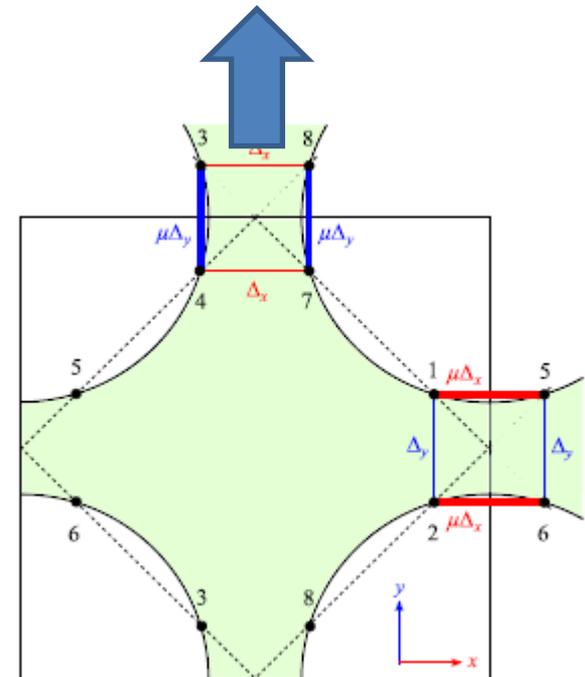
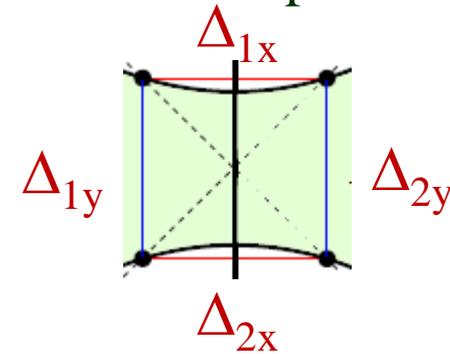
- Ginzburg-Landau analysis for 4-component order parameter

$$\Delta_{1x}, \Delta_{2x}, \Delta_{1y}, \Delta_{2y}$$

“Hot spot” CDW model

(low-energy fermions separated by Q)

Parameters: coupling  $g$  and energy cutoff  $\Lambda$ ,  
theory is under control when  $g < \Lambda$



- We did obtain stripe phase,  
with  $\Delta_{1x} = \pm i \Delta_{2x}$   
but ONLY when  $g \sim \Lambda$



- When  $g < \Lambda$ , we did find  
a checkerboard state



The  
results  
of the  
GL  
analysis

And this forced us to realize that the analysis is incomplete

## Back to Catherine's talk

d-wave superconductivity and  
(Q,Q) charge order are degenerate “twins”

Metlitski and Sachdev  
Efetov, Pepin and Meier

Does CDW has a “twin”?

Yes, its twin is superconductivity  
with a finite momentum of a pair (PDW)

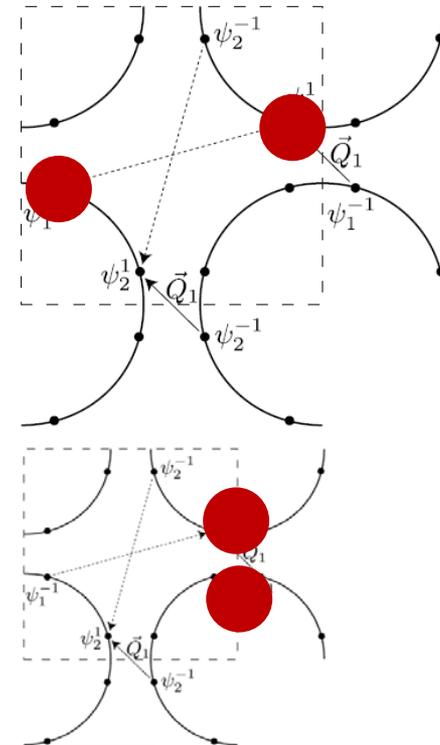
Kivelson and Fradkin  
P.A. Lee, D. Agreberg....

$$\Delta_{\mathbf{k}}^Q = \langle c_{\mathbf{k}+\mathbf{Q}}^+ c_{\mathbf{k}-\mathbf{Q}} \rangle$$

CDW

$$\bar{\Delta}_{\mathbf{k}}^Q = \langle c_{\mathbf{k}+\mathbf{Q}}^+ c_{\mathbf{k}-\mathbf{Q}}^+ \rangle$$

PDW

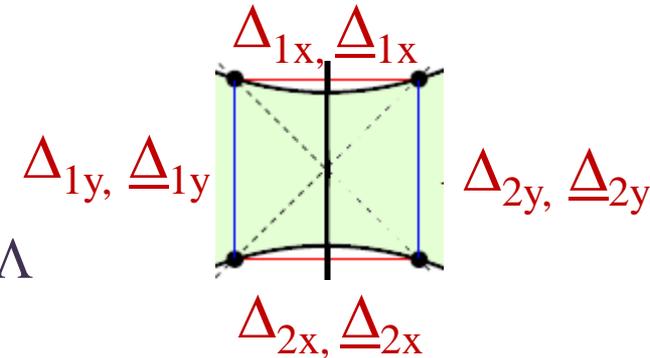


- Ginzburg-Landau analysis for 8-component order parameter

CDW:  $\Delta_{1x}, \Delta_{2x}, \Delta_{1y}, \Delta_{2y}$     PDW:  $\underline{\Delta}_{1x}, \underline{\Delta}_{2x}, \underline{\Delta}_{1y}, \underline{\Delta}_{2y}$

Same hot spot model

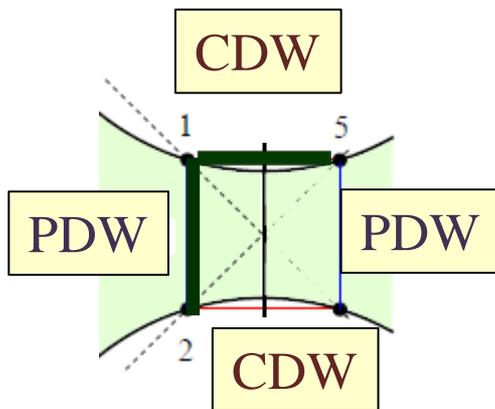
(low-energy fermions separated by Q)



Same parameters: coupling  $g$  and energy cutoff  $\Lambda$

Let's set  $g < \Lambda$  and play "by the rules"

We still get a checkerboard state, but a rather specific one:  
CDW along one direction, PDW along the other



CDW between 1 and 5:  $1 = (\pi-k, k)$ ,  $5 = (\pi+k, k)$ ,  
transferred momentum  $1-5 = (-2k, 0)$

PDW between 1 and 2:  $1 = (\pi-k, k)$ ,  $2 = (\pi-k, -k)$ ,  
total momentum  $1+2 = (-2k, 0)$

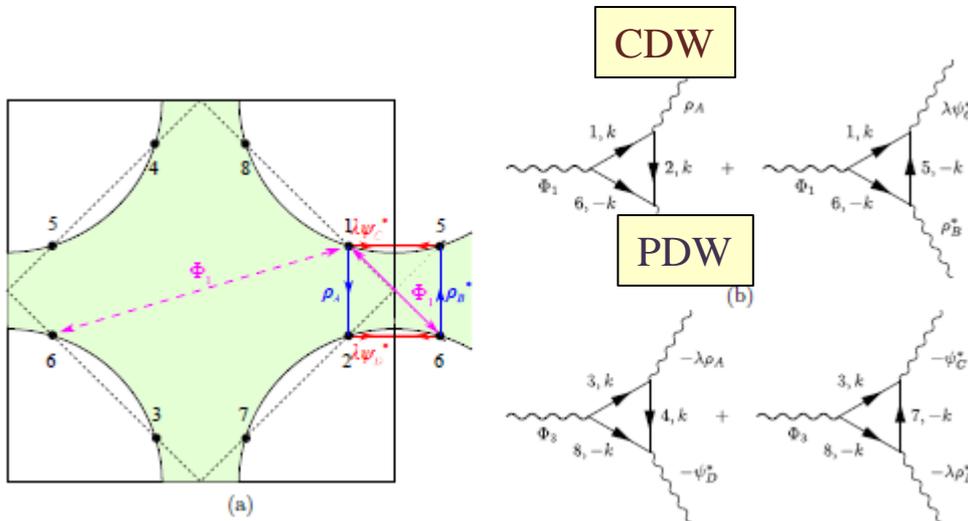
“Orthogonal” CDW and PDW orders  
carry the same momentum!

## This charge order:

- breaks  $C_4$  (CDW is a stripe along X or along Y)
- breaks time-reversal symmetry and generate incommensurate density fluctuations and incommensurate current

The reason it minimizes GL energy is rather subtle:

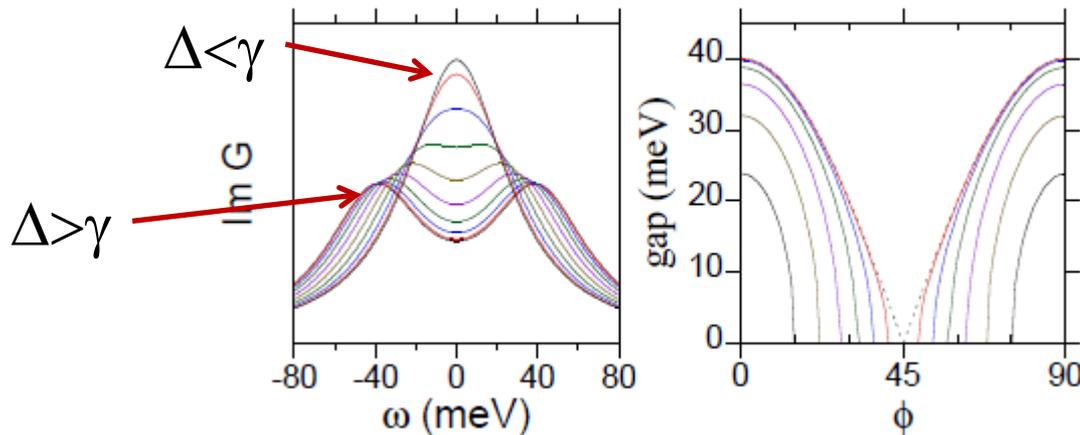
CDW and PDW, which carry the same momentum, together generate a secondary uniform s- type superconductivity



Triple coupling generates uniform SC order, and this reduces the energy of the CDW/PDW state

## This leads to the specific prediction:

- In the normal state, fermionic damping is non-zero, hence a small SC gap is difficult to detect in ARPES



Abrahams, Millis, Norman, A.C

- But in a co-existence state with a true d-wave SC with a large gap and reduced fermionic damping, additional s-wave gap component can be detectable. Two states are possible: s+d and  $s + e^{i\phi} d$ . In both cases, the gap along the nodal direction is non-zero

Search for s-wave component of SC gap in the co-existence state

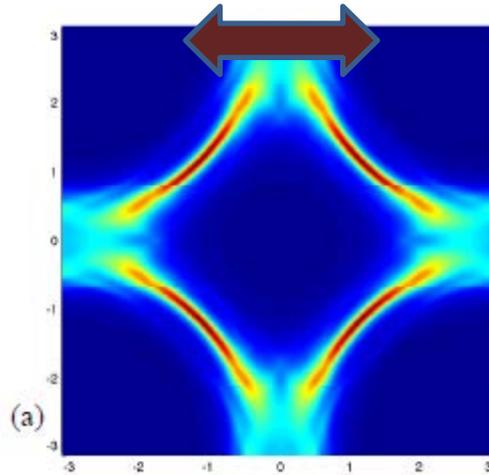
The presence of PDW also helps explain ARPES data, which without PDW component would be inconsistent with the theory

P.A. Lee

# Comparison with ARPES

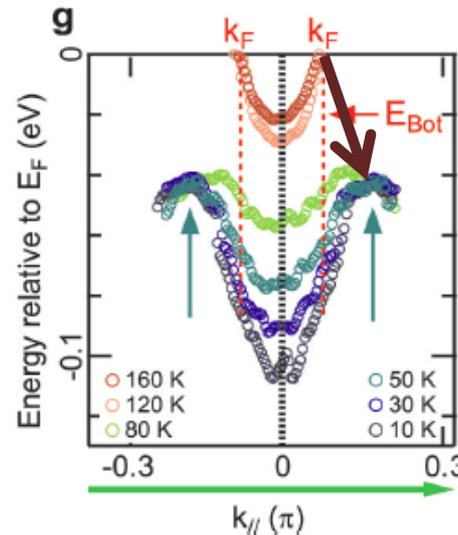
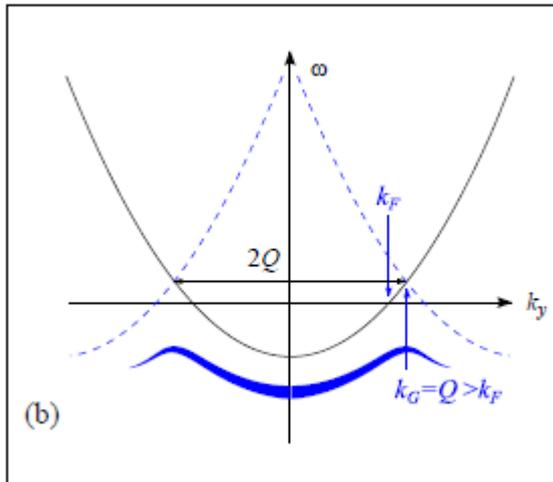
## Pure CDW state

### 1. Fermi arc



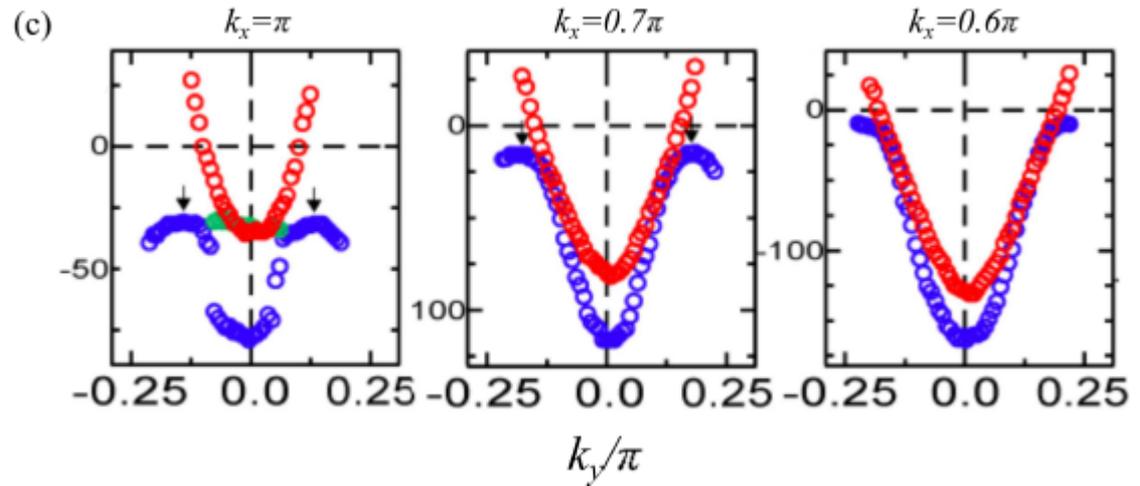
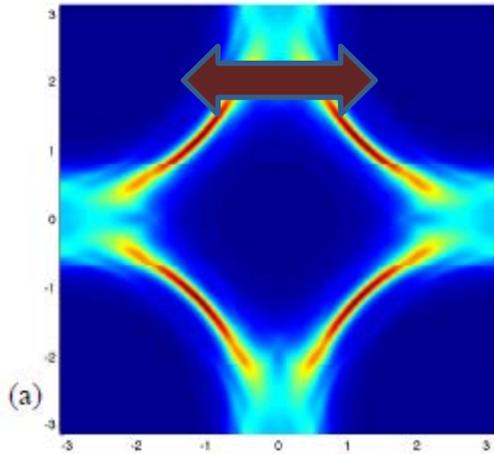
Model: domains of CDW (0,Q) and (Q,0), plus a finite damping to model a disordered CDW between  $T^*$  and  $T_{CDW}$

### 2. Shift of the position of the minimum to a larger value



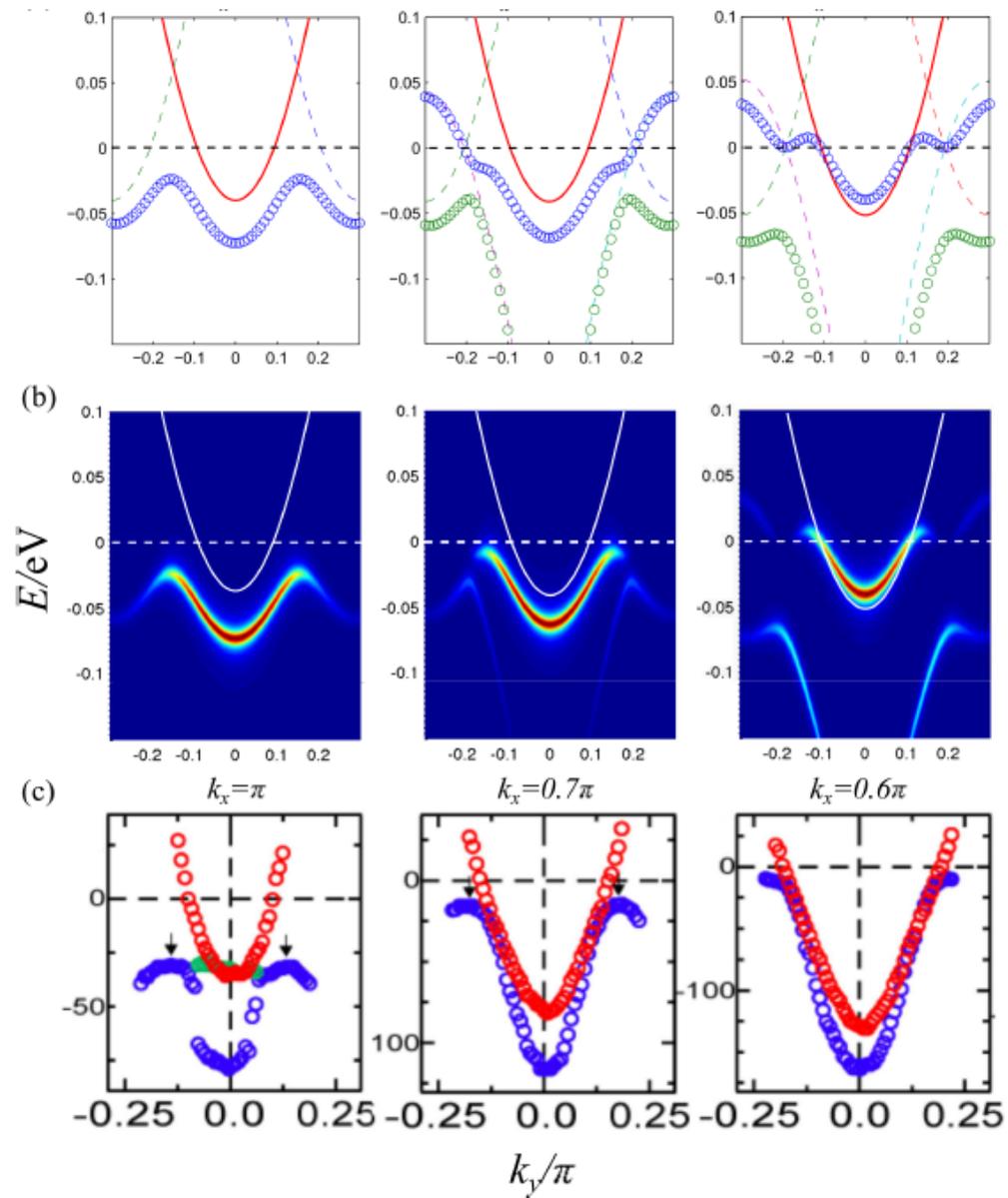
I. Vishik et al,  
Rui-Hua He et al

However, pure CDW cannot explain the behavior closer to zone diagonals

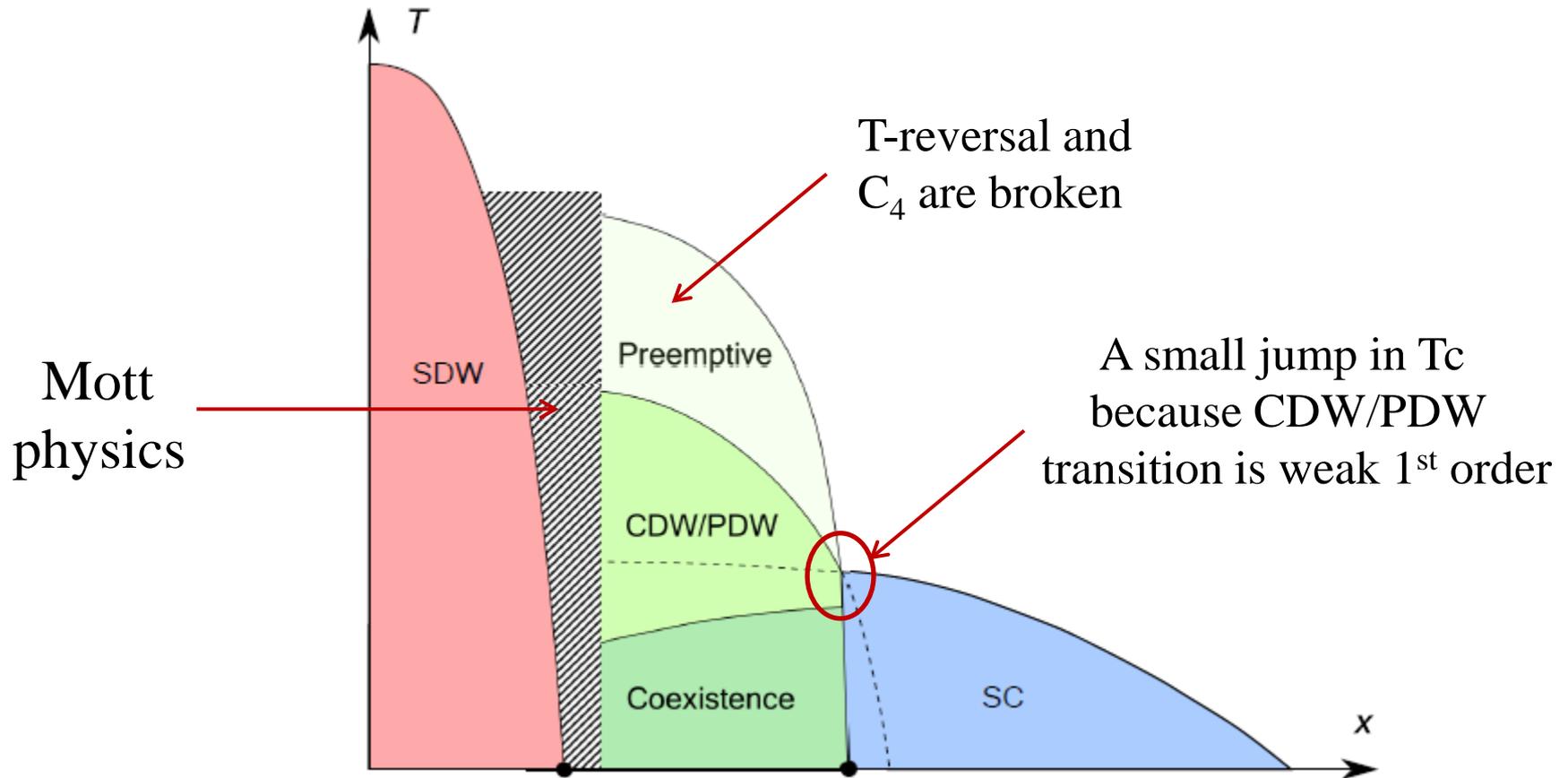


Experimental dispersion approaches Fermi level from below.  
In pure CDW description, the lower branch just keeps moving down in energy, and the upper branch lowers and hits the Fermi surface

# CDW + PDW



# Phase diagram



THANK YOU