
Multiband effect and electron-hole asymmetry in the transport properties of iron pnictides

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The different families of pnictides

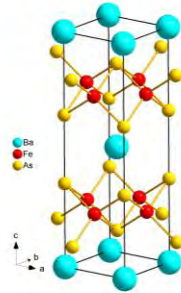
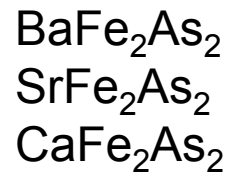
1111



Highest $T_c = 57\text{K}$

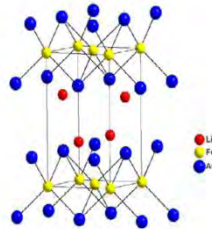
Synthesis of single crystals difficult
Few results
Chemical composition

122



A very rich family
Good single crystals

111



Good single crystals
Stoichiometric → low defect content
Very sensitive to air

Chalcogenides

The different families of pnictides

1111



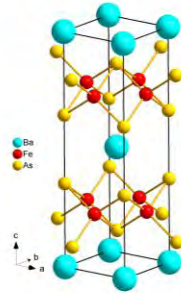
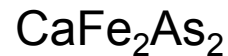
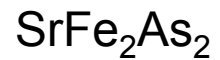
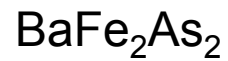
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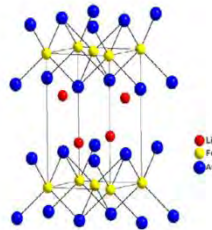
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Analysis of the resistivity and Hall effect measurements in a minimal two-band model: electron doped BaFe_2As_2 with Co and Ni substitution

Strong electron-hole asymmetry in the scattering rates

Comparison with ARPES results

Isovalent substitution: Ru/Fe

Resistivity, Hall effect and magnetoresistance

Evidence of a Fermi liquid behavior in a two band model

LiFeAs Resistivity, Hall effect and magnetoresistance

Analysis in relation with ARPES and quantum oscillations data

Beyond the two-band model

The 122 Phase AFe_2As_2 ($A=Ba, Sr, Ca, Eu$)

Possibility to get large single crystals

Pr^{3+}, La^{3+}

$Ba^{2+}, Sr^{2+}, Ca^{2+}, ..$

K^+

Electron-doped

$T_c^{max} \sim 45K$ in $CaFe_2As_2$

Hole-doped

$T_c^{max} = 37K$

Ru on Fe site
P on As Site

Isovalent substitution

● Ba
● Fe
● As

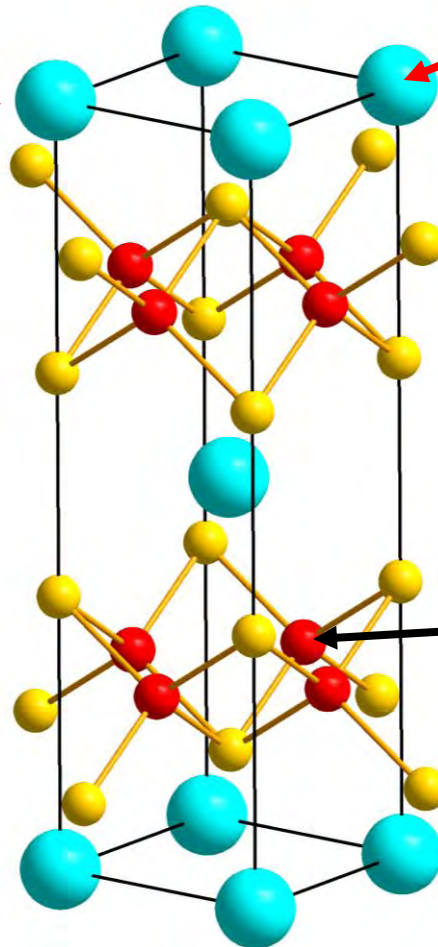
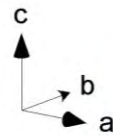
Co : $3d^7$
Ni : $3d^8$

Rh : $4d^7$
Ir : $5d^9$

Electron-doped

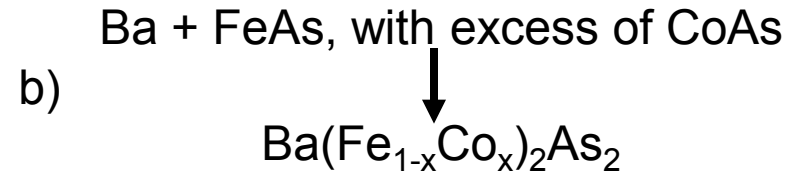
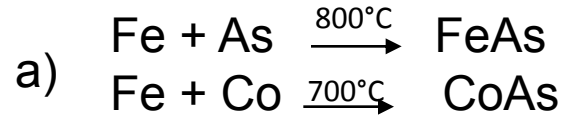
$T_c^{max} = 25K$

26 Fe	27 Co	28 Ni
44 Ru		



122 Phase BaFe_2As_2 : crystal growth

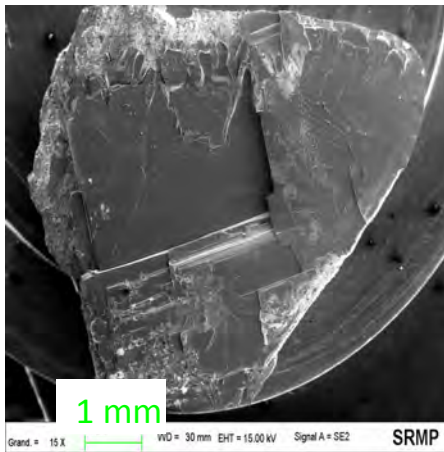
Self flux method



Synthesis in quartz tubes, sealed under vacuum
Kept at 1180°C for 4h and cooling at 5°C/h \longrightarrow

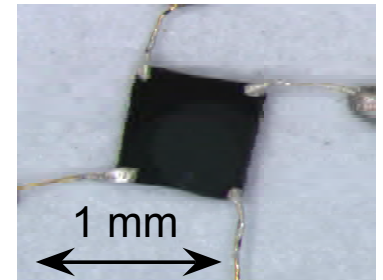
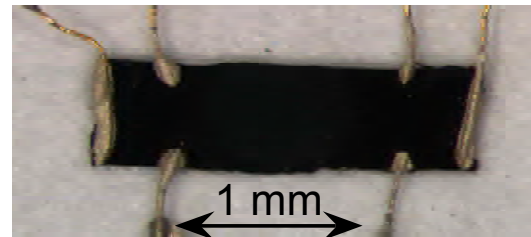
1000°C for 6h
Cooled down to room T

Mechanical extraction of
platelets single crystals,



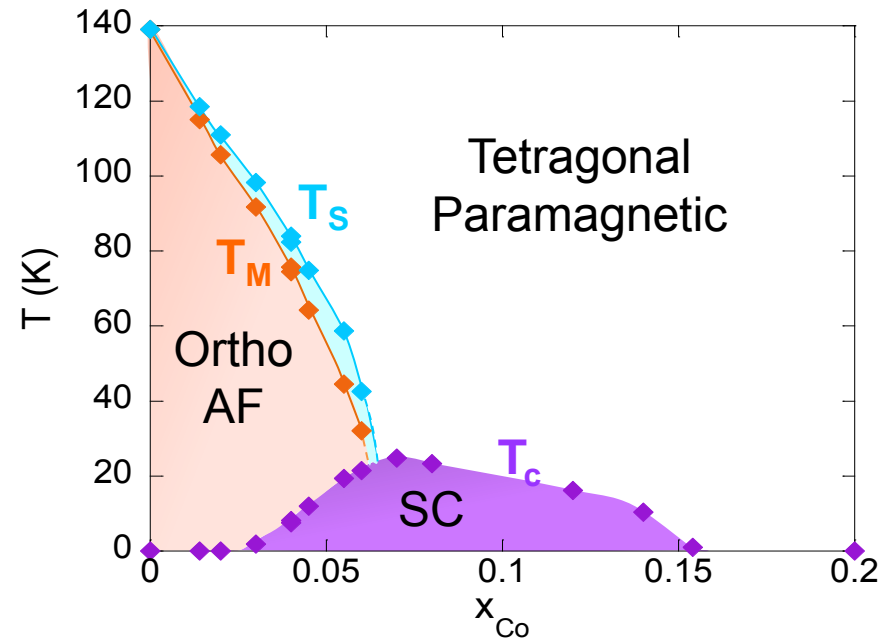
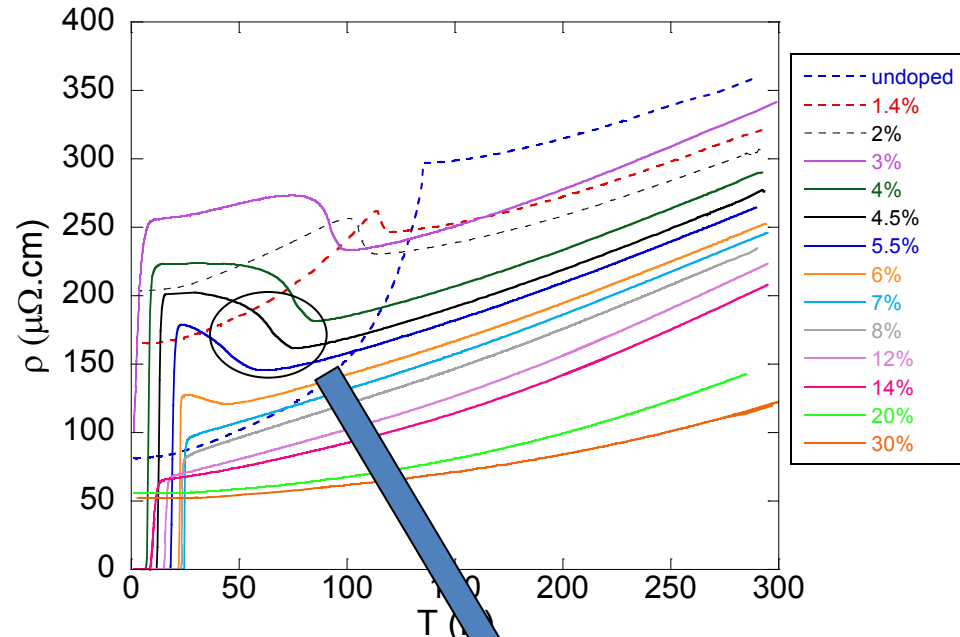
Measurements of
Transport properties

Thickness between 10 to 30 μm



Co content determined by wave length
dispersive X-ray spectroscopy

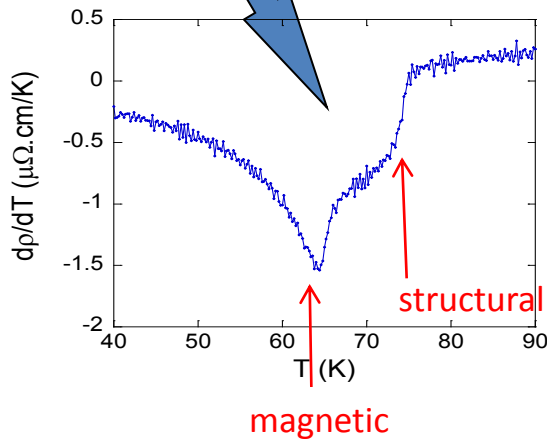
Transport properties of $\text{Ba}(\text{Fe}_{1-x}\text{Co}_x)_2\text{As}_2$ single crystals



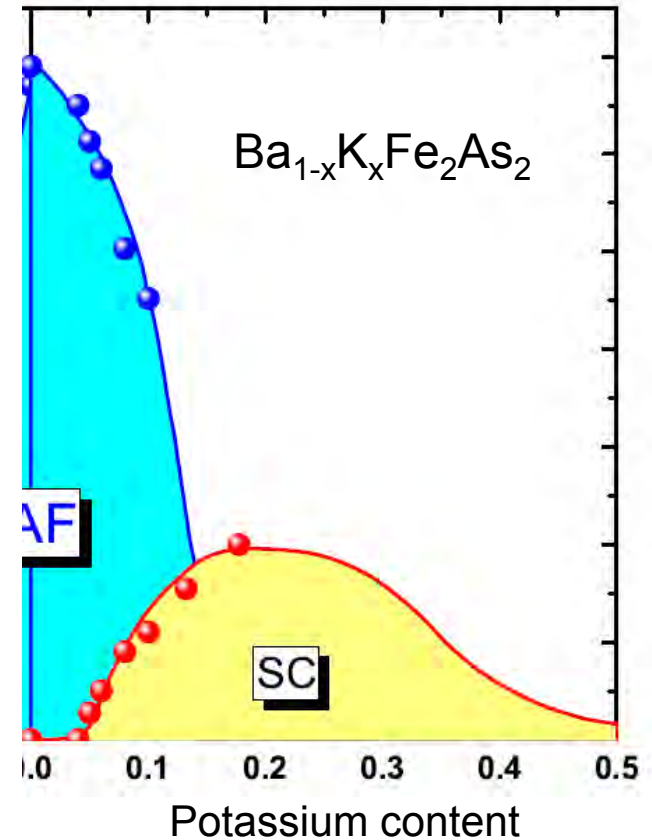
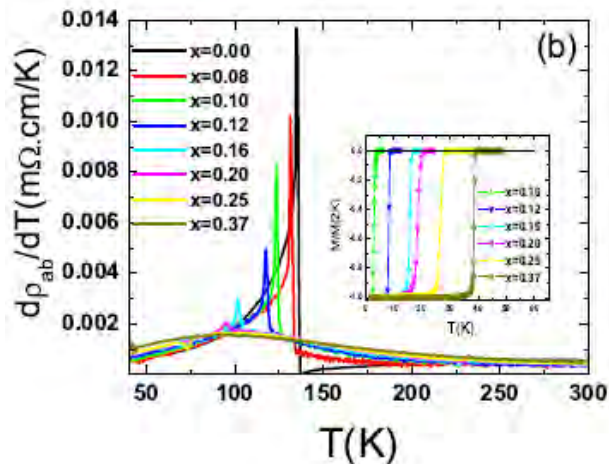
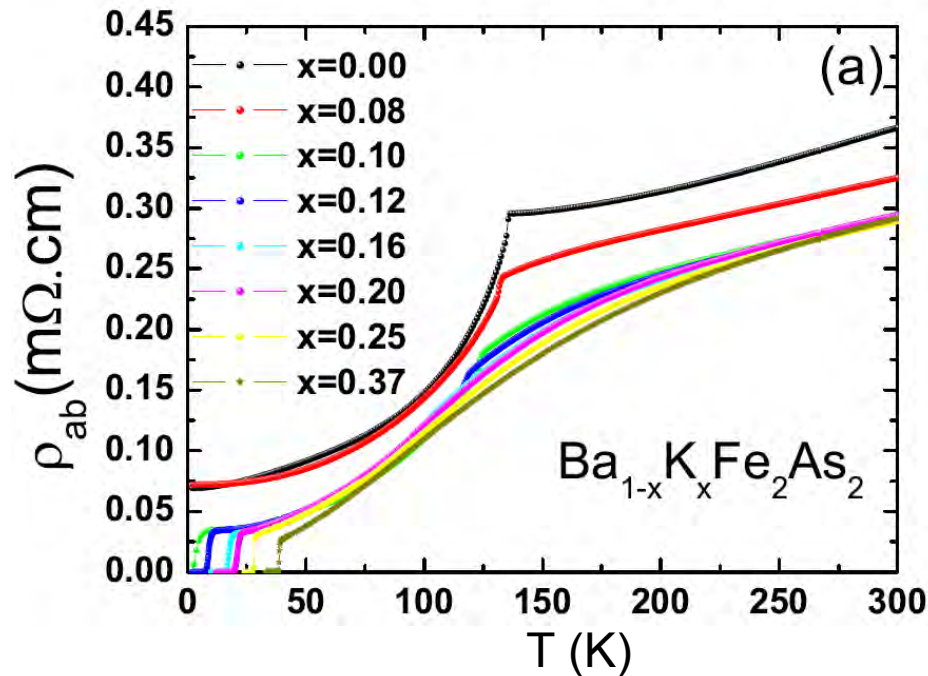
Coexistence of AF and SC states at local scale

NMR : Y. Laplace et al. PRB (2009)

Structural and magnetic transitions

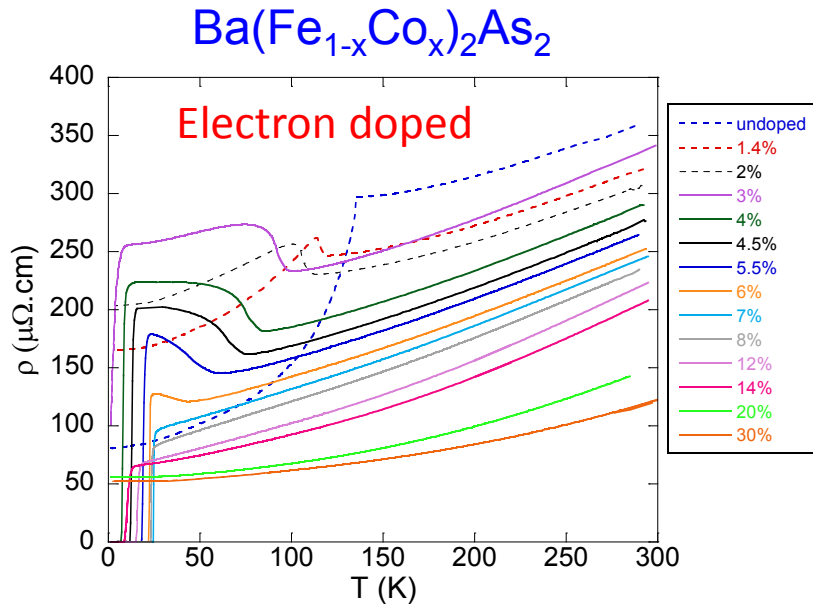


Transport properties of $\text{Ba}_{1-x}\text{K}_x\text{Fe}_2\text{As}_2$ single crystals

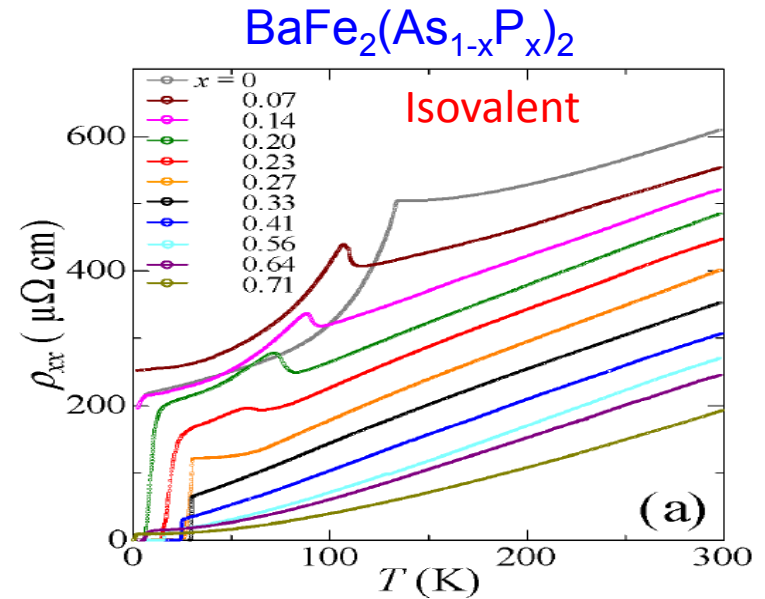


No splitting between structural and AF transitions

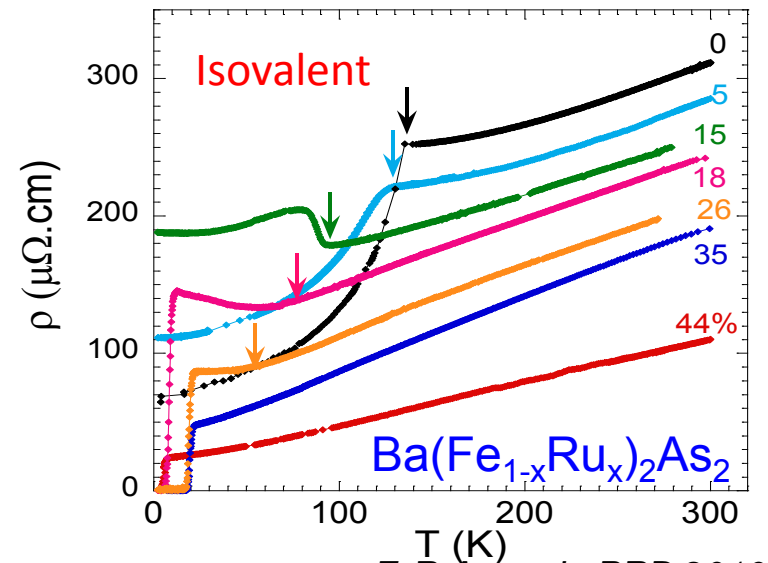
Resistivity evolution by substitution in the BaFe_2As_2 family



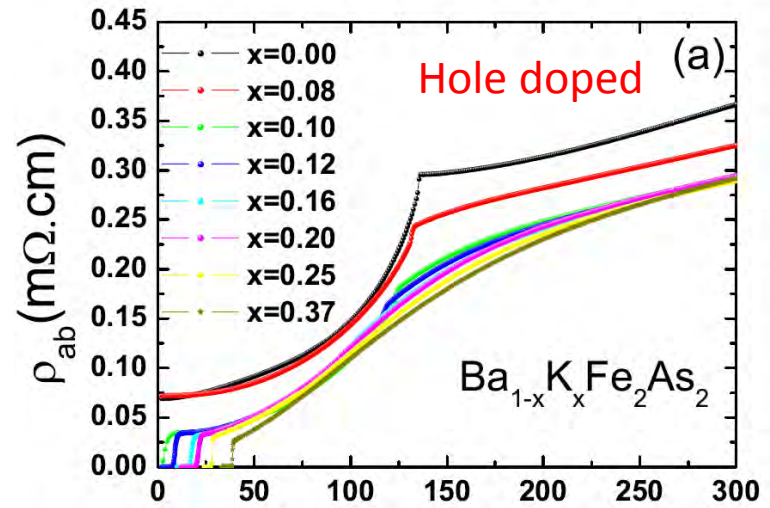
F. R.A. et al., PRL 2009



S. Kasahara et al., PRB 2010

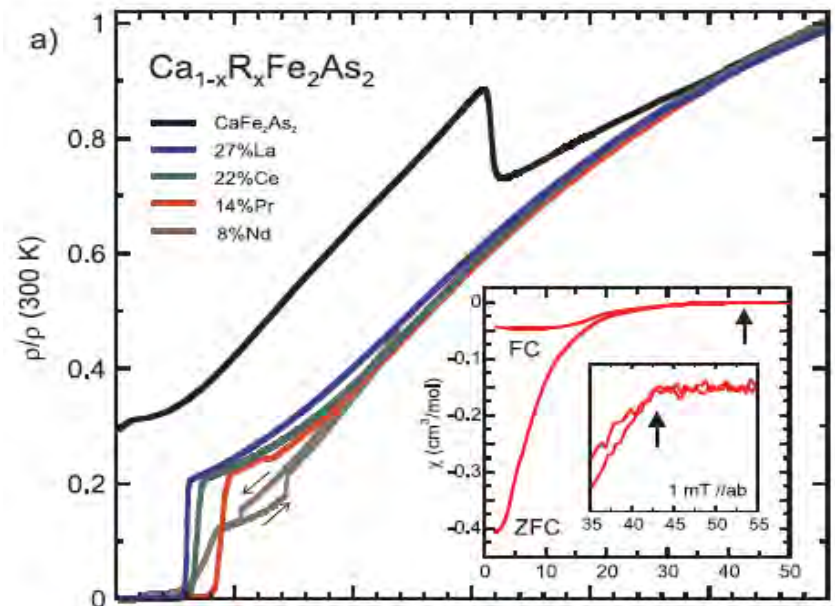
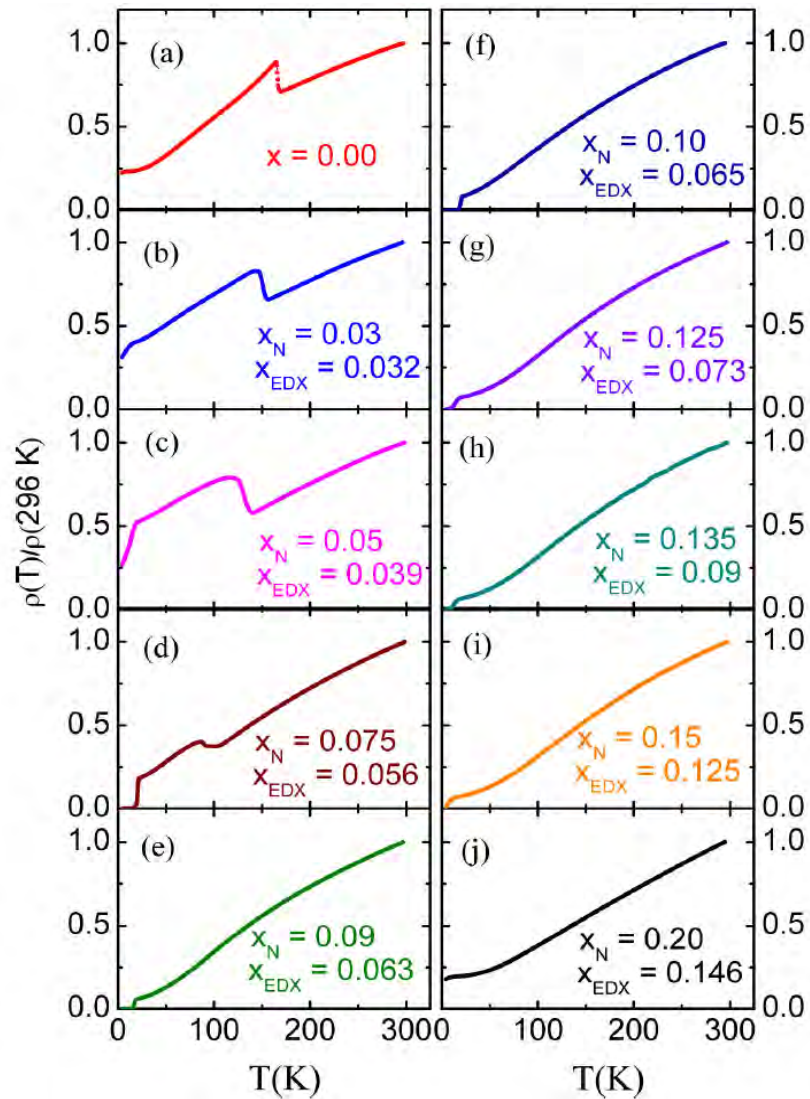
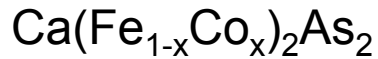


F. R.A. et al., PRB 2010



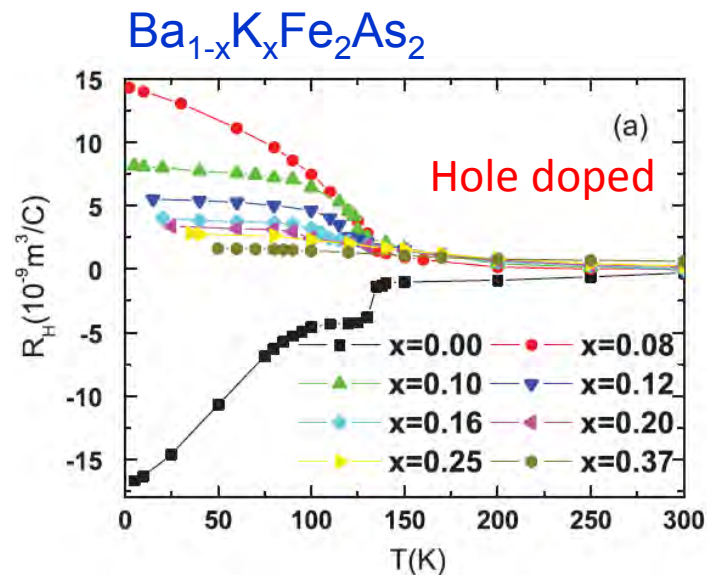
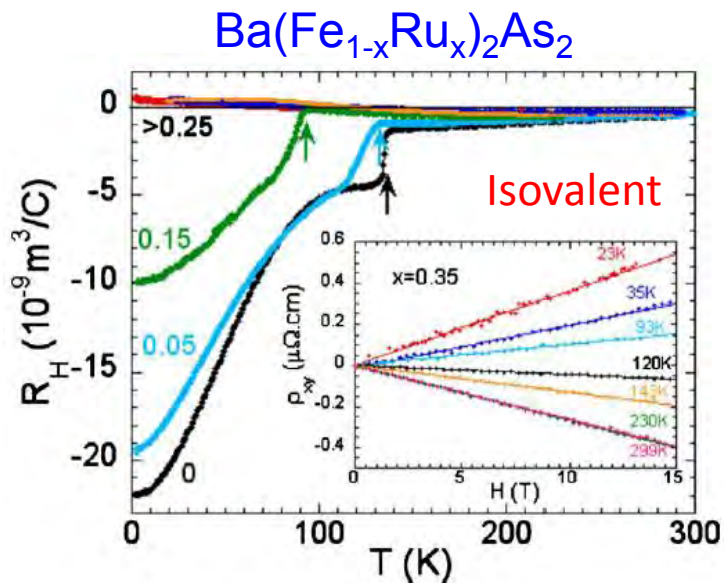
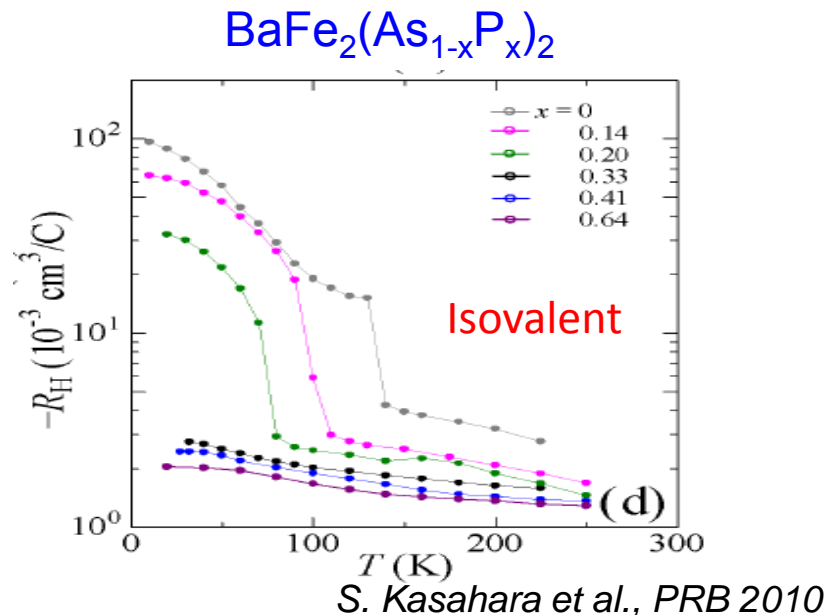
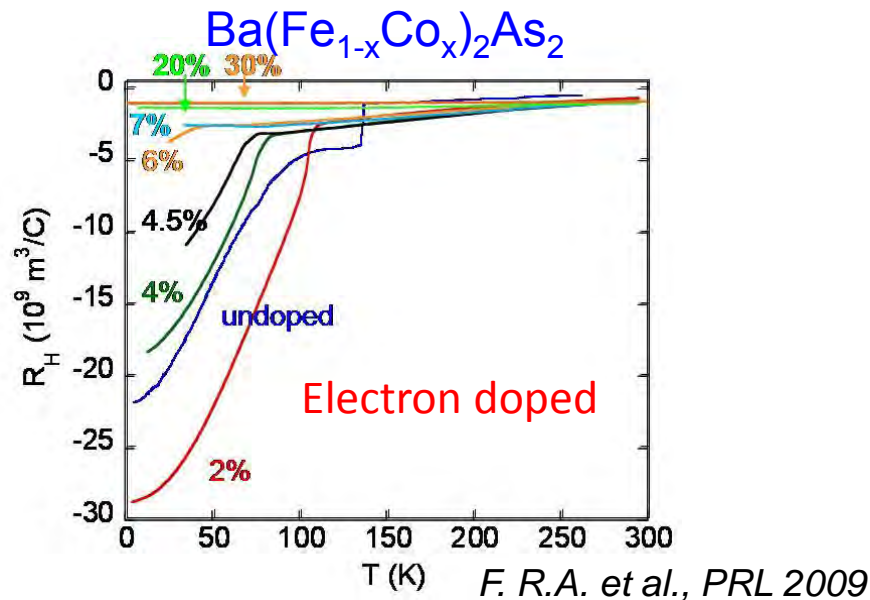
B. Shen et al., PRB (2011)

The CaFe_2As_2 family



S.R. Saha et al., arXiv 1105.4798

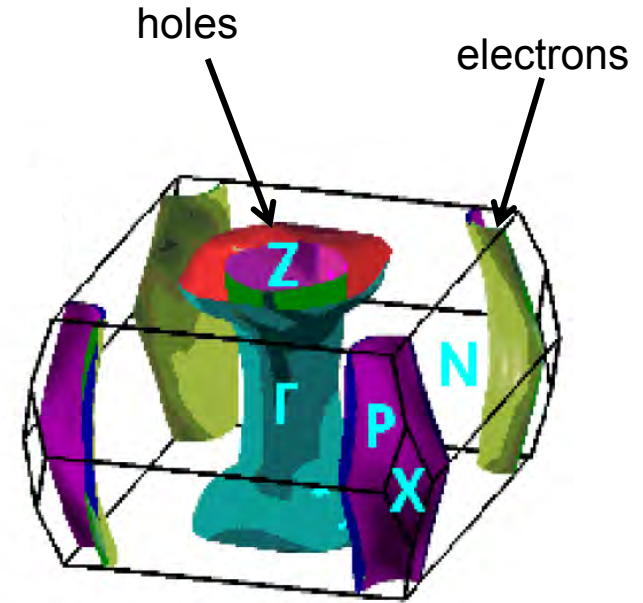
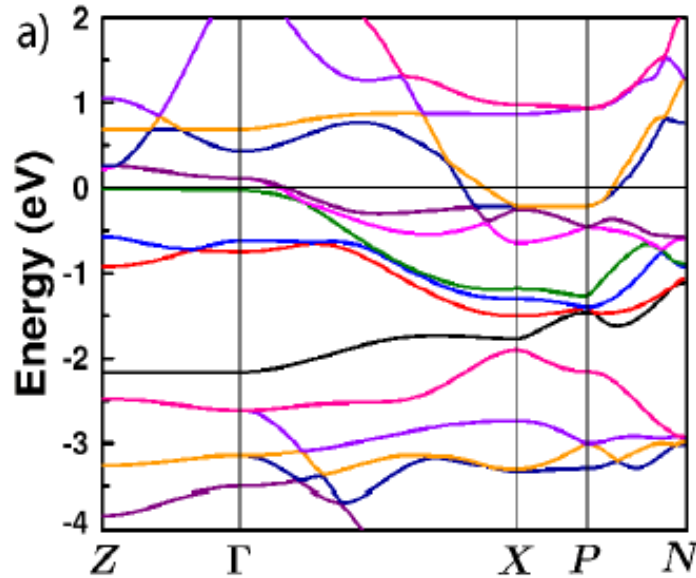
Hall effect : Large variation with temperature



The Iron based superconductors: multiband structure

Electronic structure: band structure calculation

BaFe₂As₂



Compensated semimetal: $n_e = n_h$

Multiband effects



Strong influence on transport properties

Modification by doping

Resistivity, Hall effect, magnetoresistance, ...

One-band versus multiband

Single band metal

Drude formula

$$\rho = \frac{m^*}{ne^2\tau}$$

Hall effect

$$R_H = \frac{1}{ne}$$

Iron-based superconductors

Minimal two-band model

$$\sigma = \sigma_e + \sigma_h$$
$$1/\rho = 1/\rho_e + 1/\rho_h \quad \rho_{e,h} = \frac{m_{e,h}^*}{n_{e,h}e^2\tau_{e,h}}$$

$$R_H = \frac{1}{e} \frac{(-n_e\mu_e^2 + n_h\mu_h^2)}{(n_e\mu_e + n_h\mu_h)^2}$$

$$\frac{\delta\rho}{\rho(0)} = \frac{\sigma_e\sigma_h(\mu_e + \mu_h)^2}{(\sigma_e + \sigma_h)^2} H^2$$

$$\mu_{e,h} = \frac{e\tau_{e,h}}{m_{e,h}^*}$$

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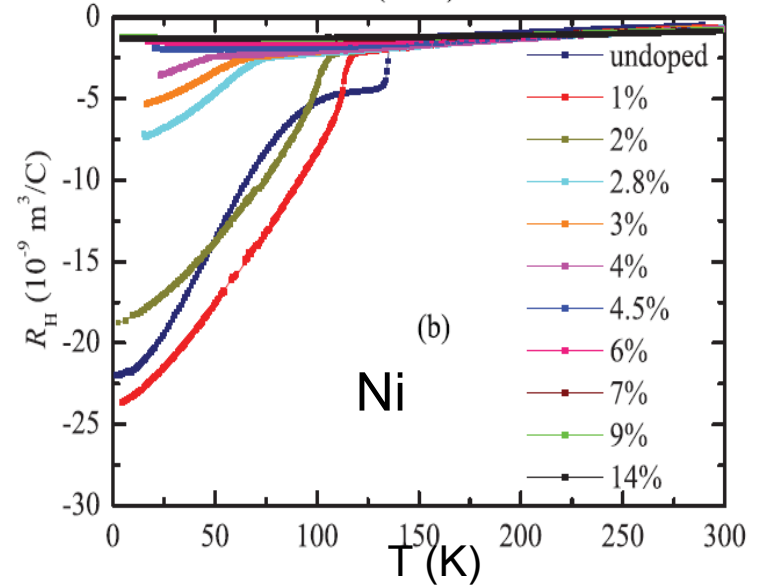
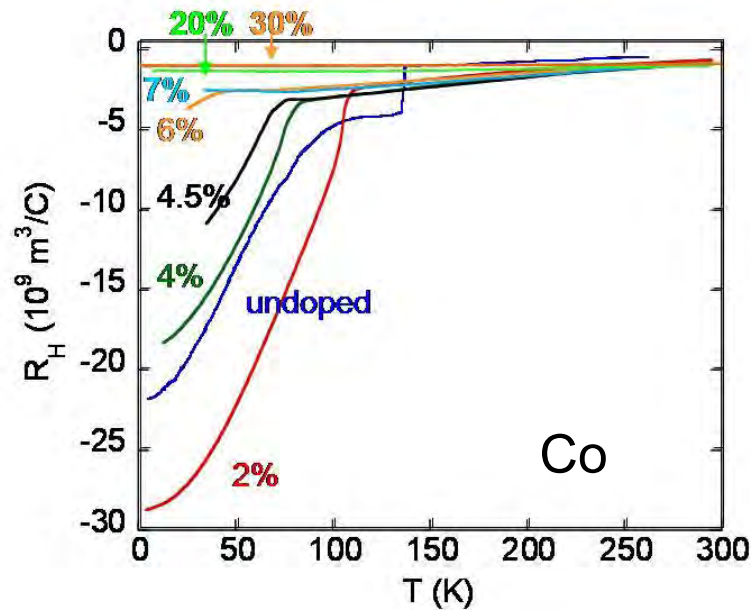
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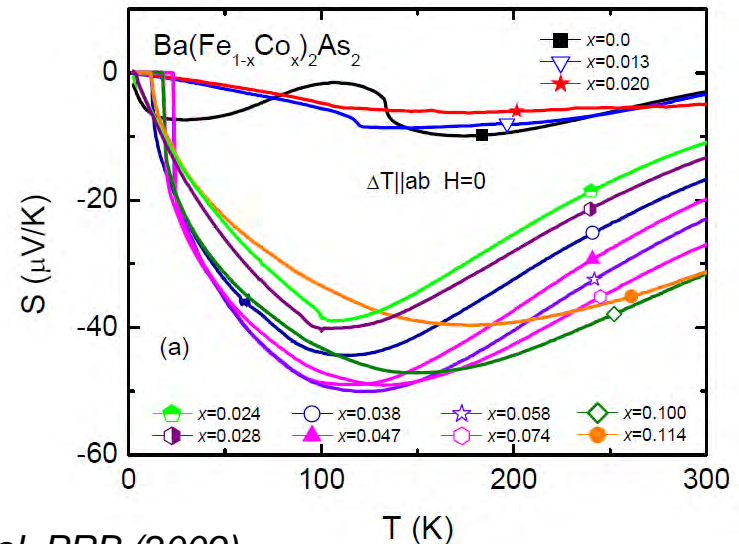
Beyond the two-band model

Hall effect: $\text{Ba}(\text{Fe}_{1-x}\text{Co}_x)_2\text{As}_2$ and $\text{Ba}(\text{Fe}_{1-x}\text{Ni}_x)_2\text{As}_2$

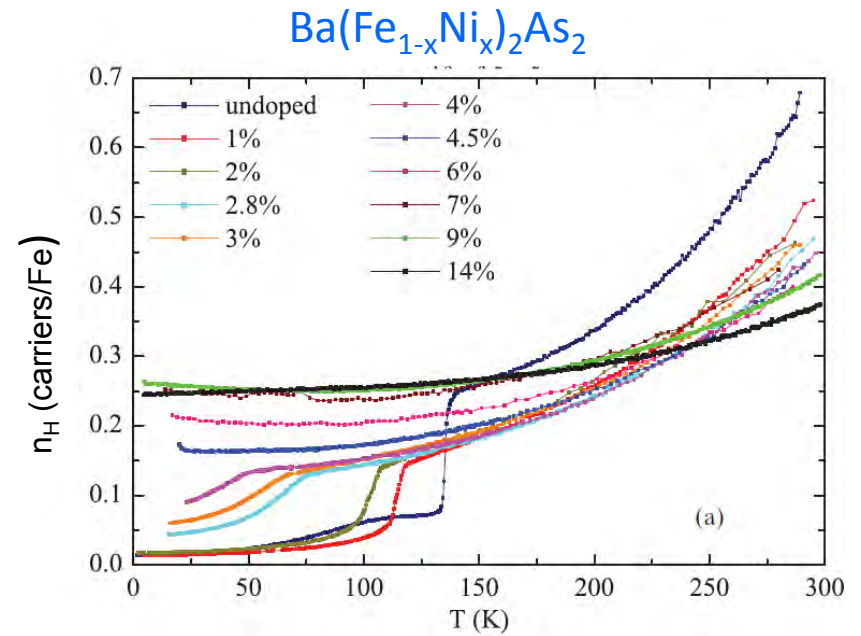
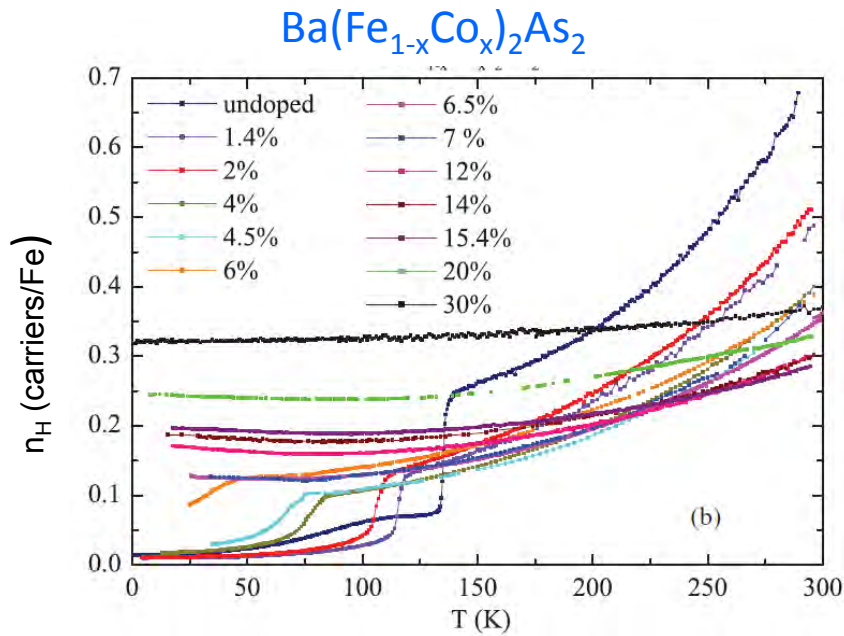


Transport dominated by the electrons all over the phase diagram

Drastic change of the Hall coefficient
in the AFM state near $x \sim 0.02$
Seen also by thermopower measurements



Hall number: $\text{Ba}(\text{Fe}_{1-x}\text{Co}_x)_2\text{As}_2$ and $\text{Ba}(\text{Fe}_{1-x}\text{Ni}_x)_2\text{As}_2$



In a two-band model:

$$\frac{1}{n_H e} = |R_H| = \frac{|-R_e \sigma_e^2 + R_h \sigma_h^2|}{(\sigma_e + \sigma_h)^2}$$

electrons

$$R_e = \frac{1}{n_e e}$$

holes

$$R_h = \frac{1}{n_h e}$$



n_H upper bound for n_e

Analysis in a two-band model

One electron band and one hole band : 4 unknown quantities : n_e , n_h , τ_e , τ_h

Experimentally : 3 equations

Conductivity: $\sigma = \sigma_e + \sigma_h$

Hall effect: $\frac{1}{n_H e} = |R_H| = \frac{|-R_e \sigma_e^2 + R_h \sigma_h^2|}{(\sigma_e + \sigma_h)^2}$

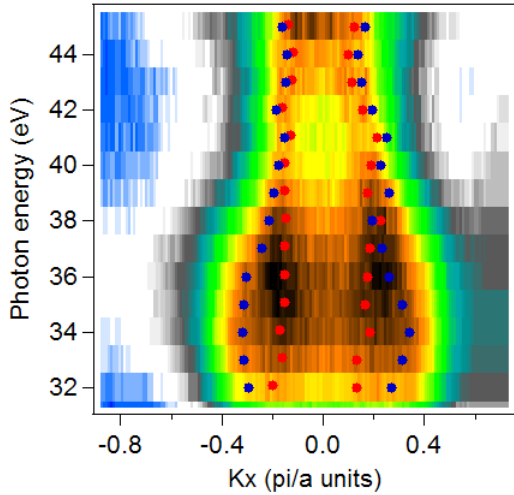
Charge conservation: $n_e = n_h + x_{Co}$

➔ Estimate of the electron density : ARPES measurements

Using ARPES measurements to determine n_e and n_h

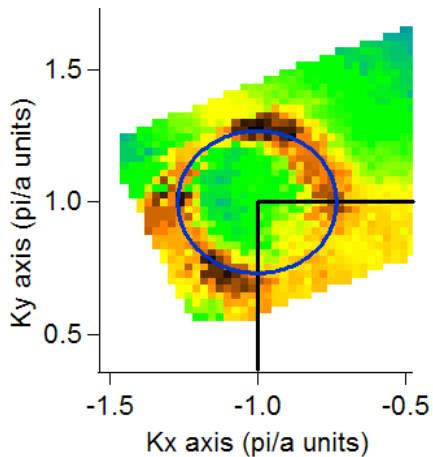
V. Brouet et al., PRB (2009)

Hole bands: two 2D, one 3D

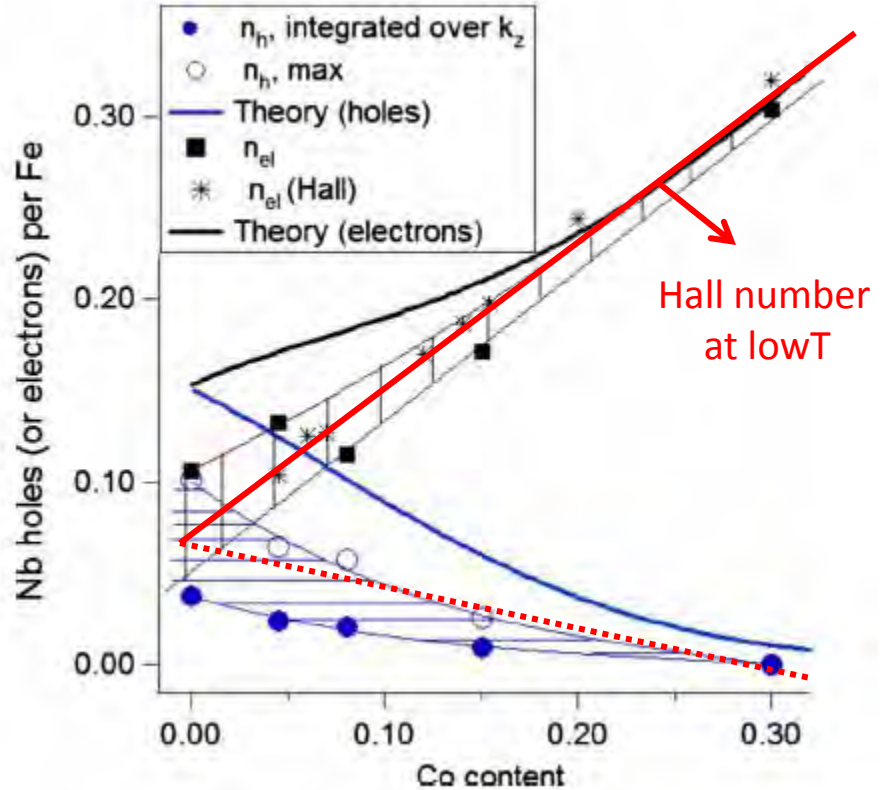


4% Co

2 degenerate electron bands: 2D



7% Co



Hall number at lowT

$$n_e - n_h = x_{Co}$$

Good agreement between ARPES and transport data



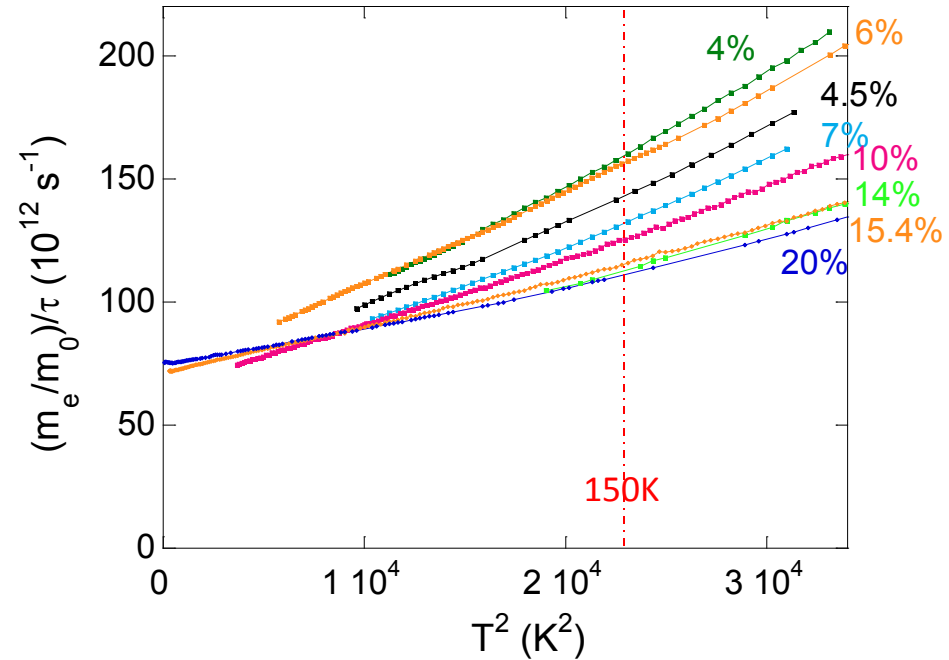
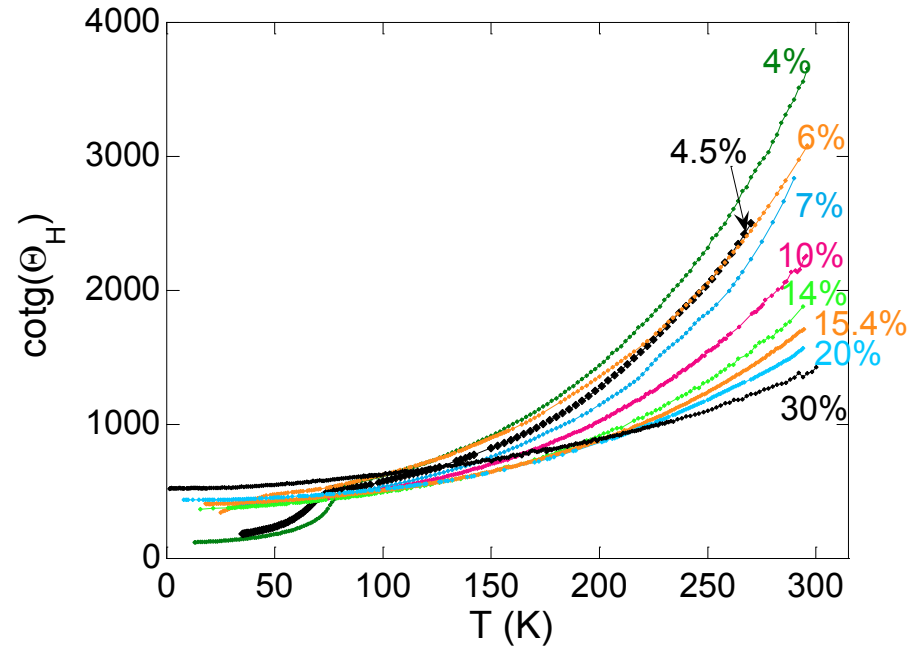
Holes not directly visible in the transport properties

Electronic scattering rates

$$n_e(T) = n_H(T)$$

$$\rho(T) = \frac{m_e}{n_e(T)e^2\tau(T)}$$

$$\cot(\Theta_H) = \rho / |R_H| = m_e / e\tau$$



Up to 150K: $1/\tau = 1/\tau_0 + BT^2$ ➔

Electron-electron interaction
Fermi liquid behavior

$1/\tau_0$ independent on Co doping for $x \leq 20\%$

No contribution of phonons nor spin fluctuations to the electronic scattering rates

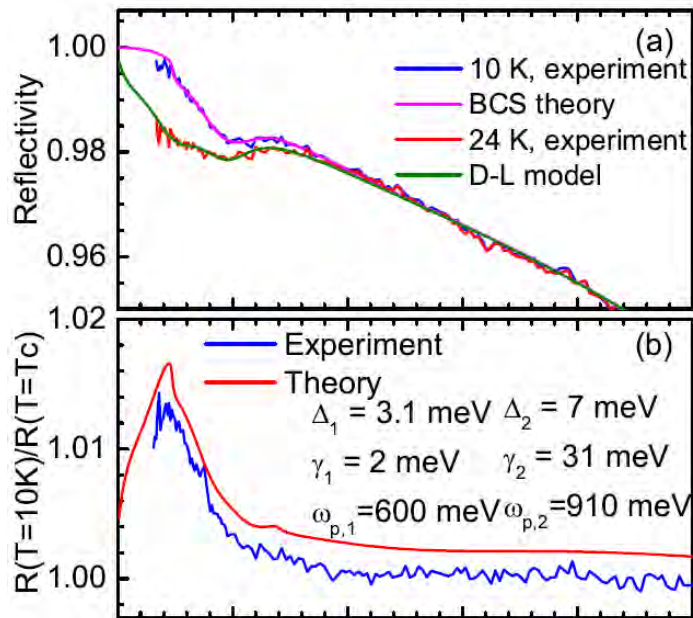
Electron and hole mobilities in $\text{Ba}(\text{Fe}_{1-x}\text{Co}_x)_2\text{As}_2$

➔ Holes more strongly scattered than electrons

Optical measurements

E. Van Heumen et al., arXiv 0912.0636

Superconducting state for $x=7\%$

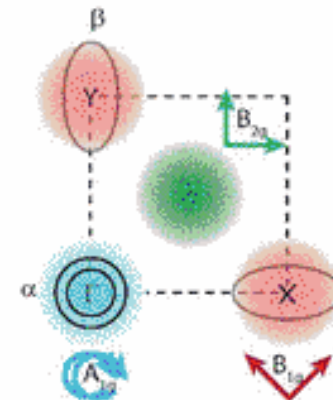
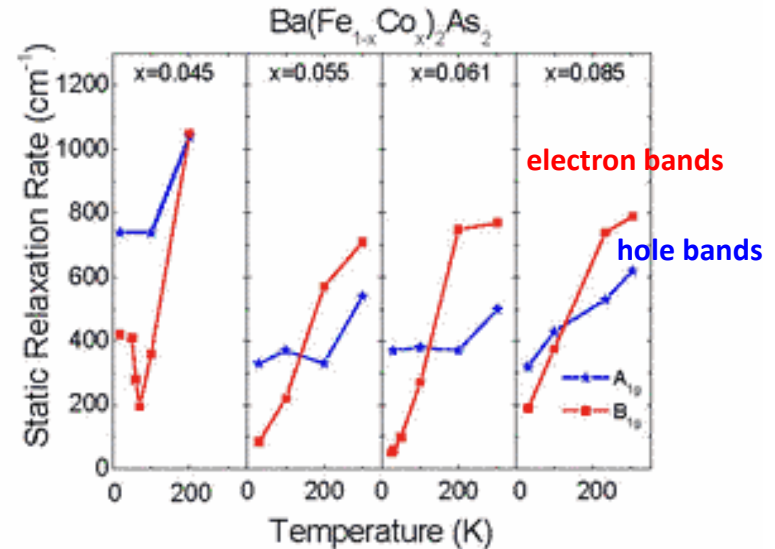


The holes are 15 times more scattered than the electrons

Interband scattering more efficient for the holes than for the electrons

Electronic Raman scattering

R. Hackl et al.

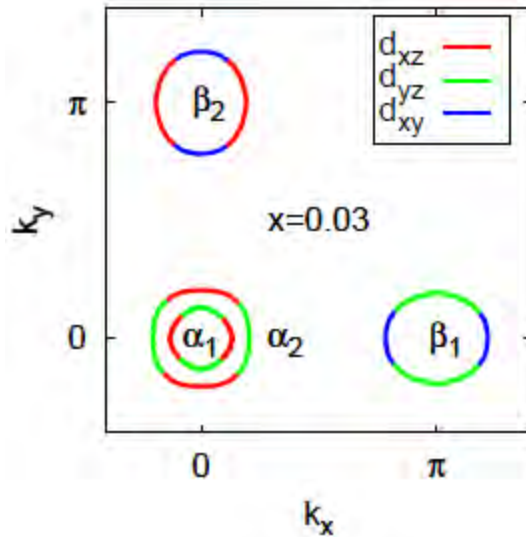


Anisotropy of the scattering rates along the FS sheets

Multiorbital composition of each Fermi pocket

A.F. Kemper et al. PRB 2011

d_{xz} , d_{yz} , d_{xy}



→ Anisotropy of the effective mass and Fermi velocity

Scattering by spin fluctuations

One particle scattering rates affected by

- the orbital character of the initial and final states
- momentum dependence of the spin susceptibility

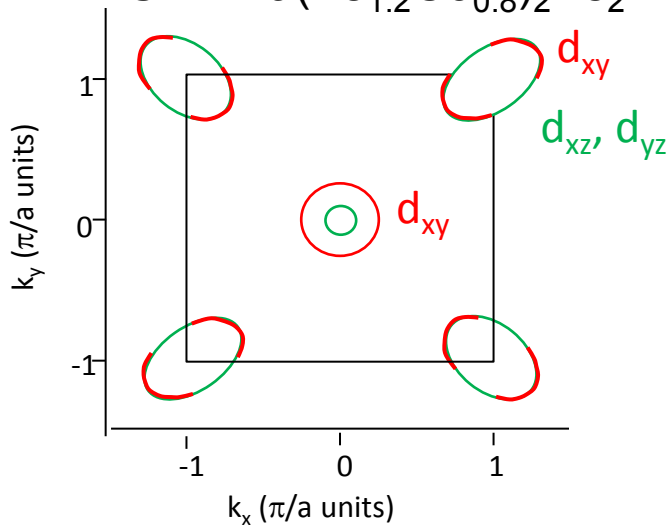
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Strong anisotropy of the scattering rates

Transport dominated by small parts of the electron FS sheets with d_{xy} character with long lifetimes and large v_F

Anisotropy of the scattering rates along the FS sheets

ARPES: in $\text{Ba}(\text{Fe}_{1.2}\text{Co}_{0.8})_2\text{As}_2$:

V. Brouet et al., arXiv 1105.5604



lifetimes about twice longer
on the d_{xy} parts of the SF sheets
both for the hole and electrons pockets

$$1/l = 1/v_F \tau$$

	k_F (π/a)	v_F (eV.Å)	m^*/m_b	$\delta k(E_F)$ (π/a)	n (carr/Fe)	\hbar/τ (meV)
hole d_{xz}/d_{yz}	0.06	0.5	2.7	0.11	0.006	44
hole d_{xy}	0.22	0.4	2.3	0.07	0.038	22
electron d_{xz}/d_{yz}	0.25	0.6	2.4	0.09	0.12	43
electron d_{xy}	0.30	0.7/1.2	5/2.9	0.04		22/38

Seems difficult to reconcile with the Hall effect dominated by the electrons

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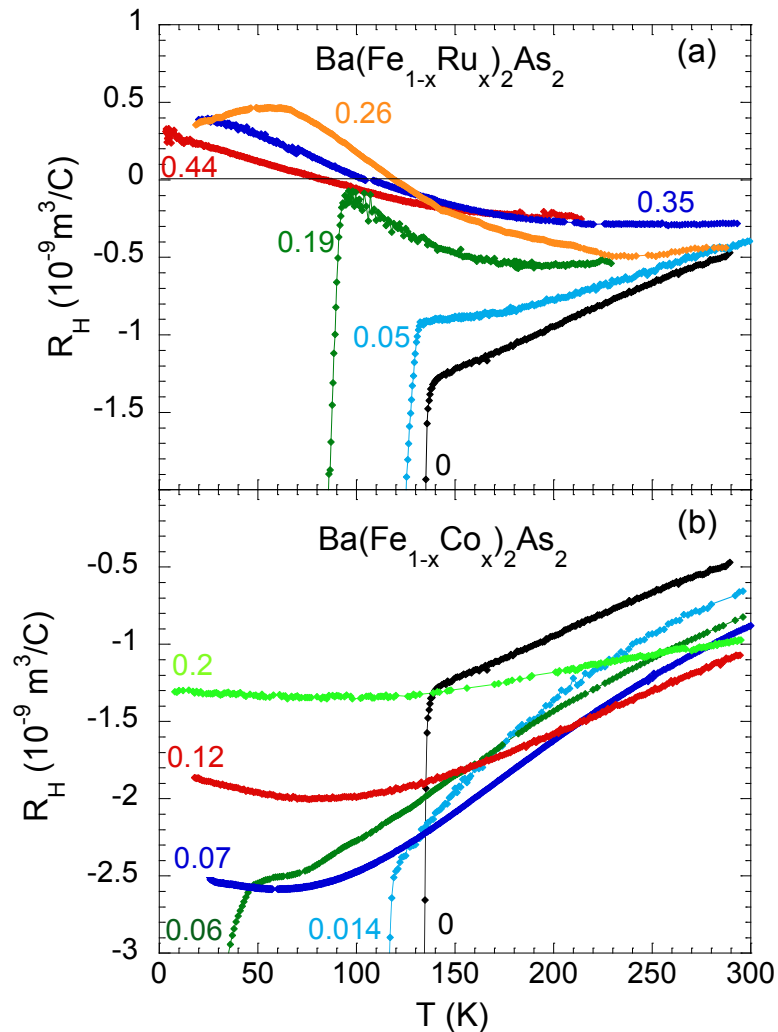
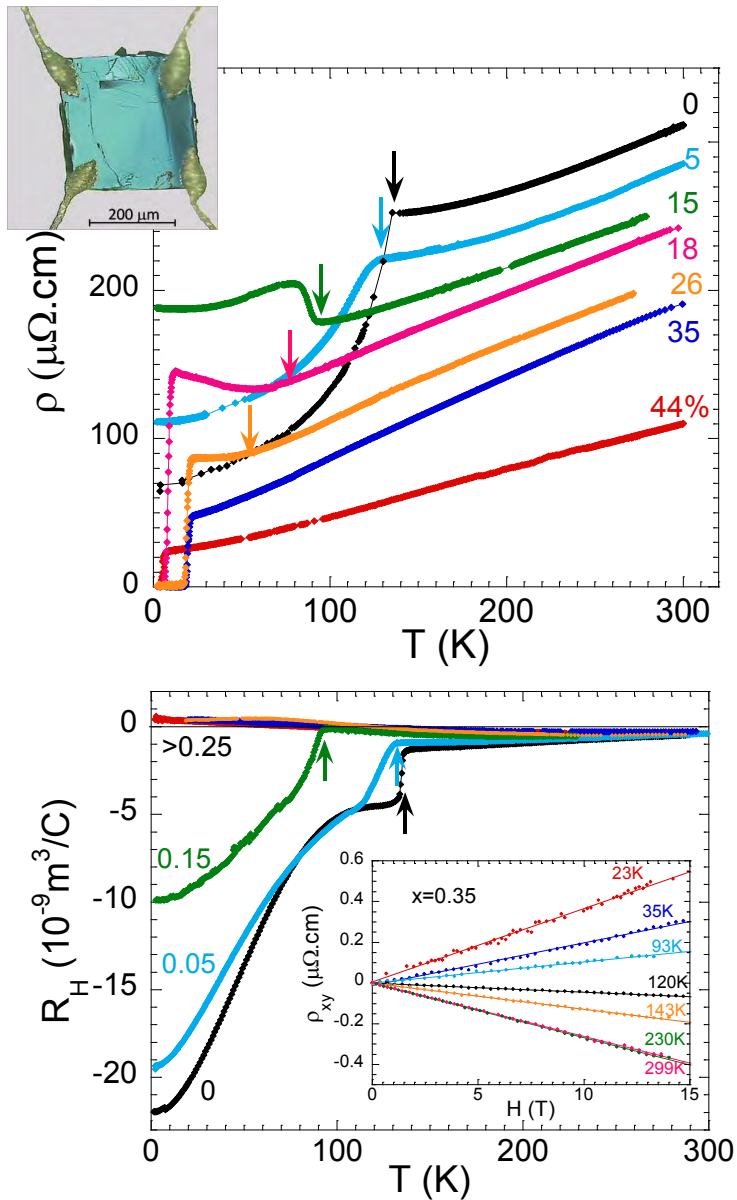
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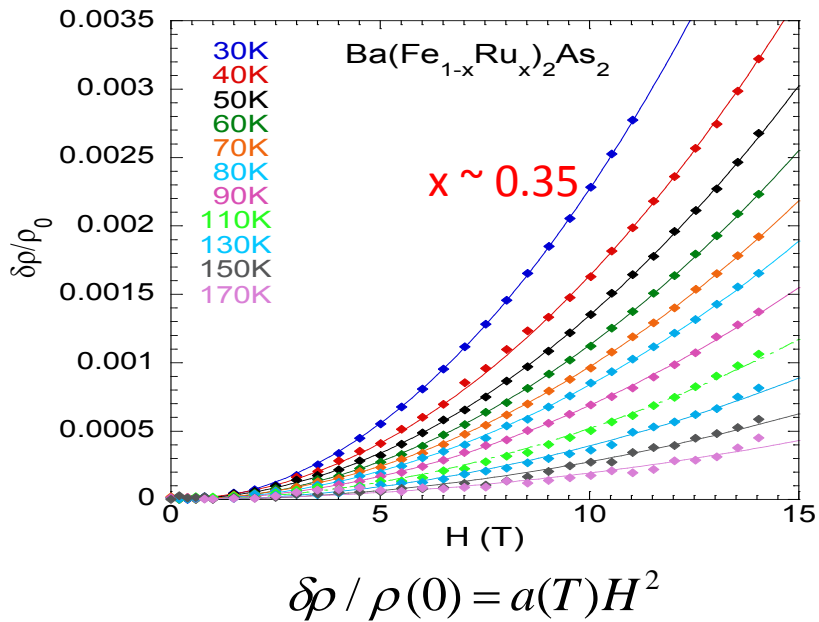
Beyond the two-band model

Transport properties of Ru-substituted BaFe_2As_2

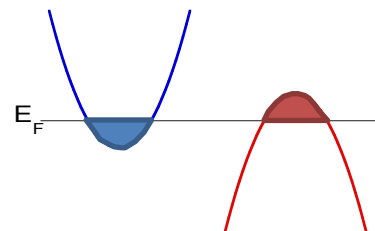


In a two band model :
mobility of holes overcomes that of electrons
at low T for $x > 0.25$

Ba(Fe_{1-x}Ru_x)₂As₂ : ρ , R_H , Magnetoresistance



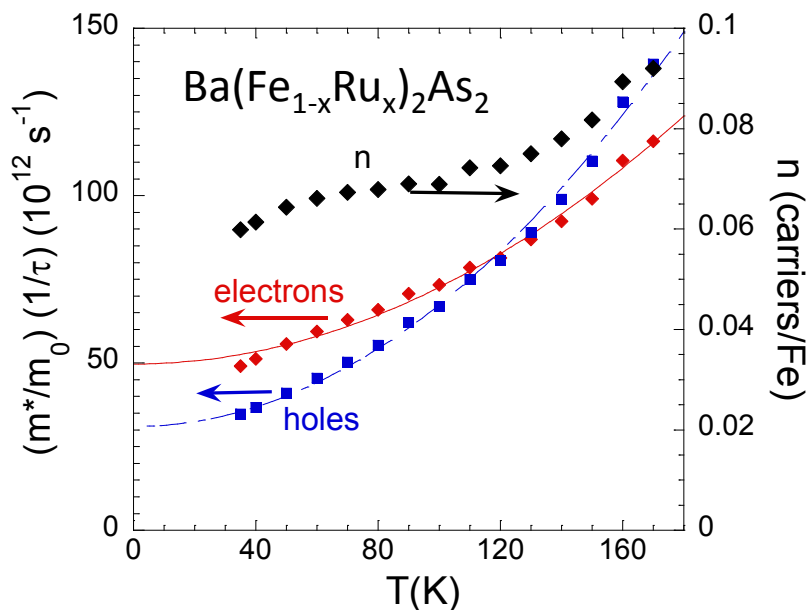
For a compensated semi-metal: $n_e = n_h = n$



$$R_H / \rho = \mu_h - \mu_e$$

$$\frac{\delta\rho}{\rho(0)} = \mu_e \mu_h H^2 \quad \longrightarrow \quad n, \mu_e, \mu_h$$

$$1 / \rho = (\mu_h + \mu_e) n e$$

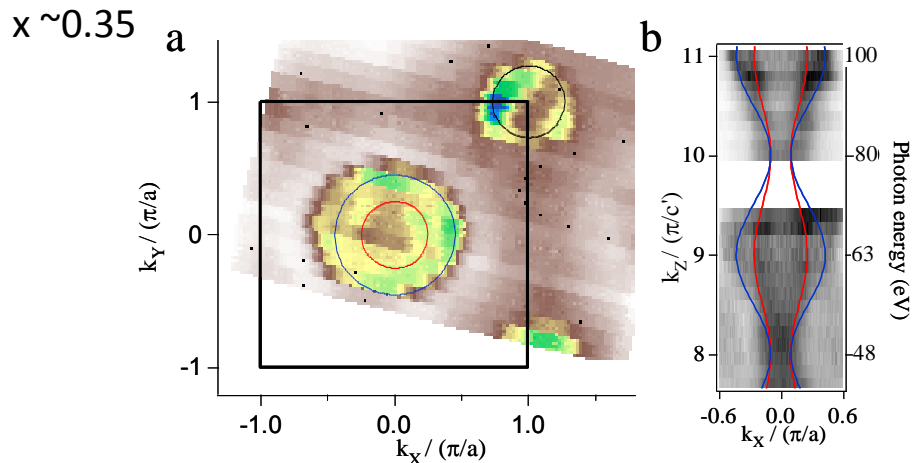


T² behavior for both holes and electrons

↓
Fermi liquid

Comparison with ARPES measurements

V. Brouet et al, PRL (2010)



Modification of the band structure

For the hole pockets:
FS significantly warped along k_z

Ru : isovalent substitution

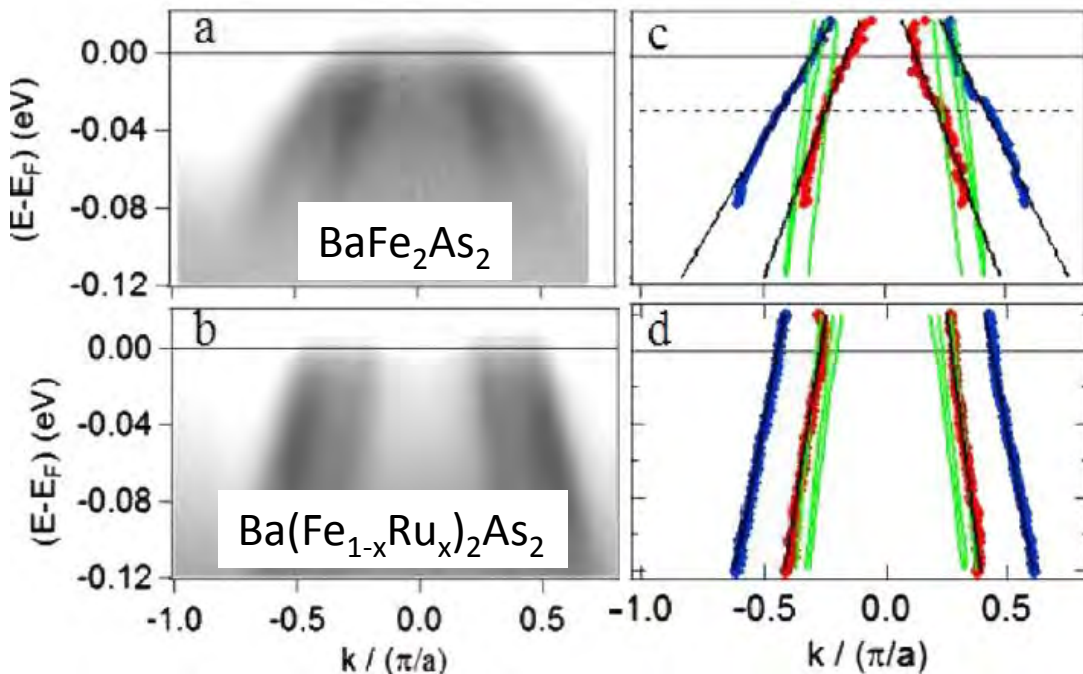
$$n_e = n_h = n \approx 0.11 / \text{Fe}$$

But other ARPES studies show
no change of the carrier concentration

R. Dhaka et al. PRL(2011), N. Xu et al. PRL (2012)

Electronic correlations
strongly reduced
 v_F nearly three times larger

Might explain the increase of
the hole mobility



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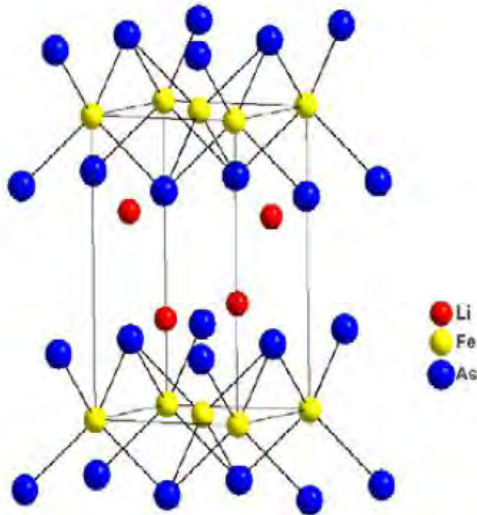
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Beyond the two-band model

LiFeAs



$$a=b=3.7715(2)\text{\AA}, c=6.3574(3)\text{\AA}$$

Mechanism of superconductivity?

- Absence of nesting: Spin fluctuations?
- Orbital fluctuations : s_{++}

$T_c \sim 17-18\text{K}$

Stoichiometric

Nearly compensated semi-metal

Defect free

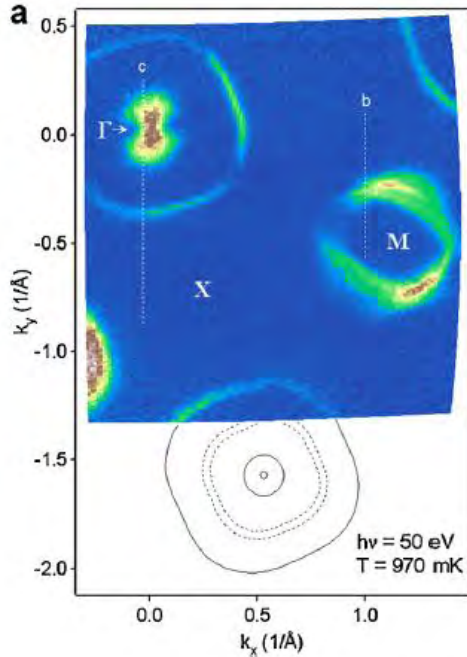
Small upper-critical fields

Quantum oscillation measurements

Non polar surface : ARPES

Drawback: Very sensitive to air

ARPES results on LiFeAs

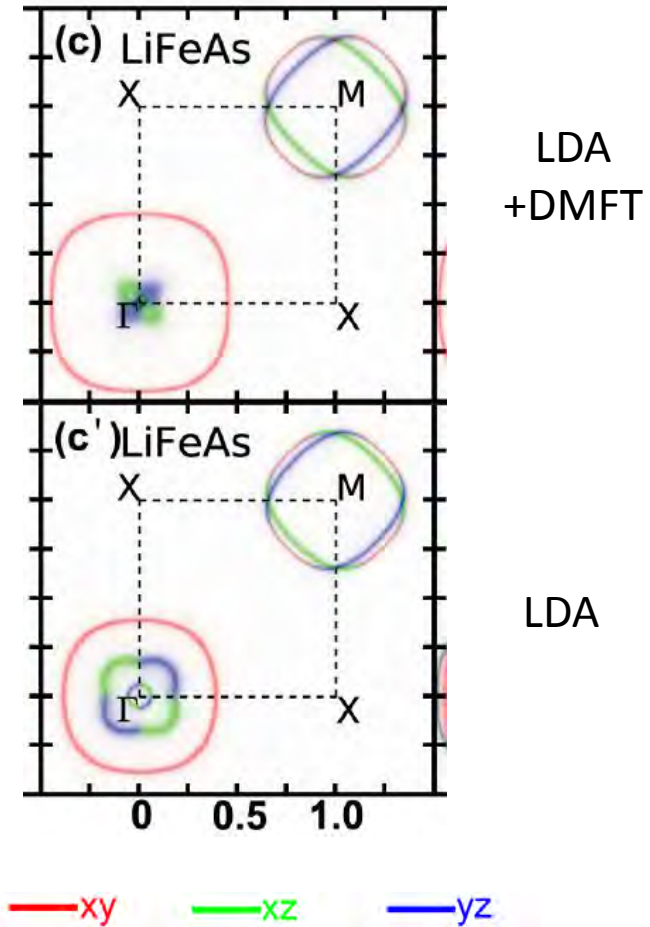


Good agreement with calculations if correlations are taken into account

Borisenko et al. PRL (2010)

Two hole bands
Two electron bands

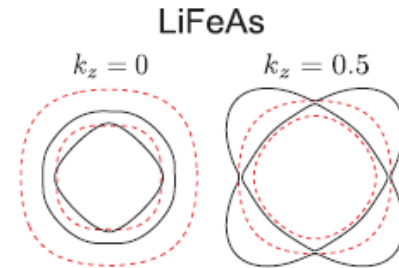
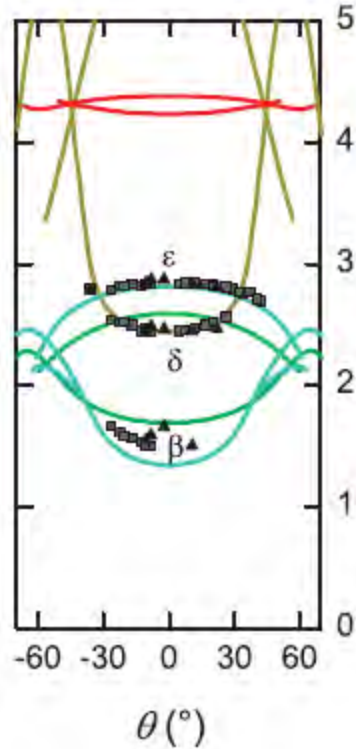
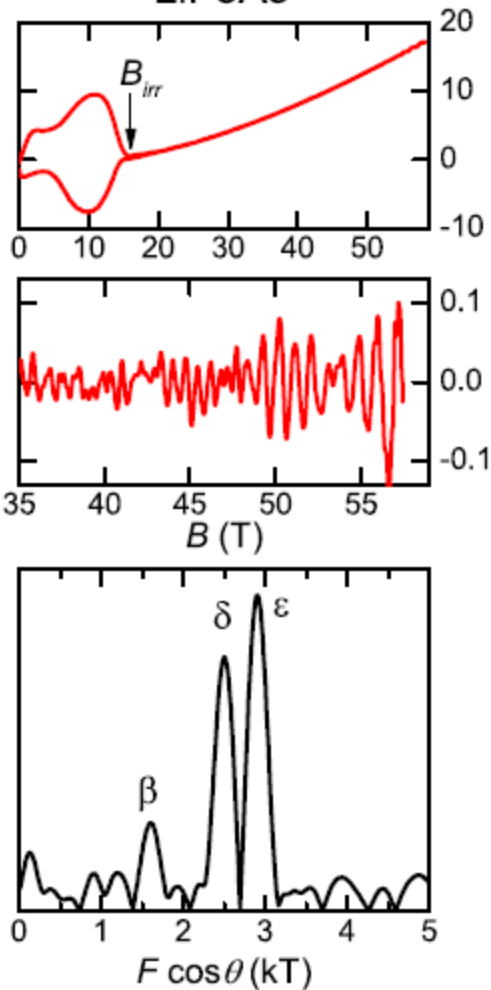
$$n_e \approx n_h \approx 0.2 \text{ carriers/Fe}$$



Yin et al., Nature materials

De Haas van Alphen experiments on LiFeAs

C. Putzke et al., PRL (2011)



3 different frequencies



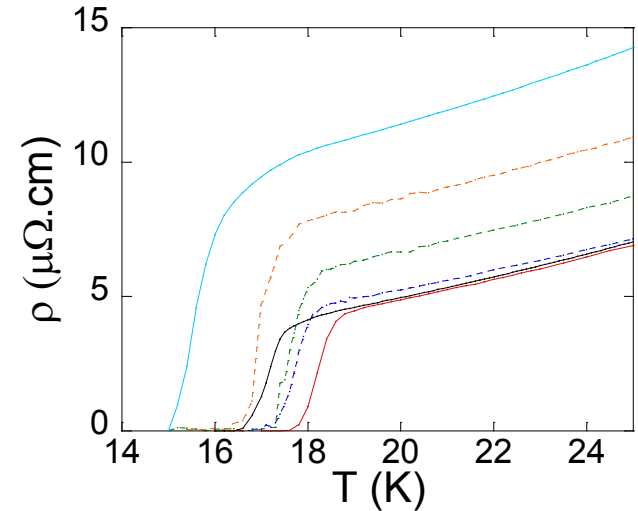
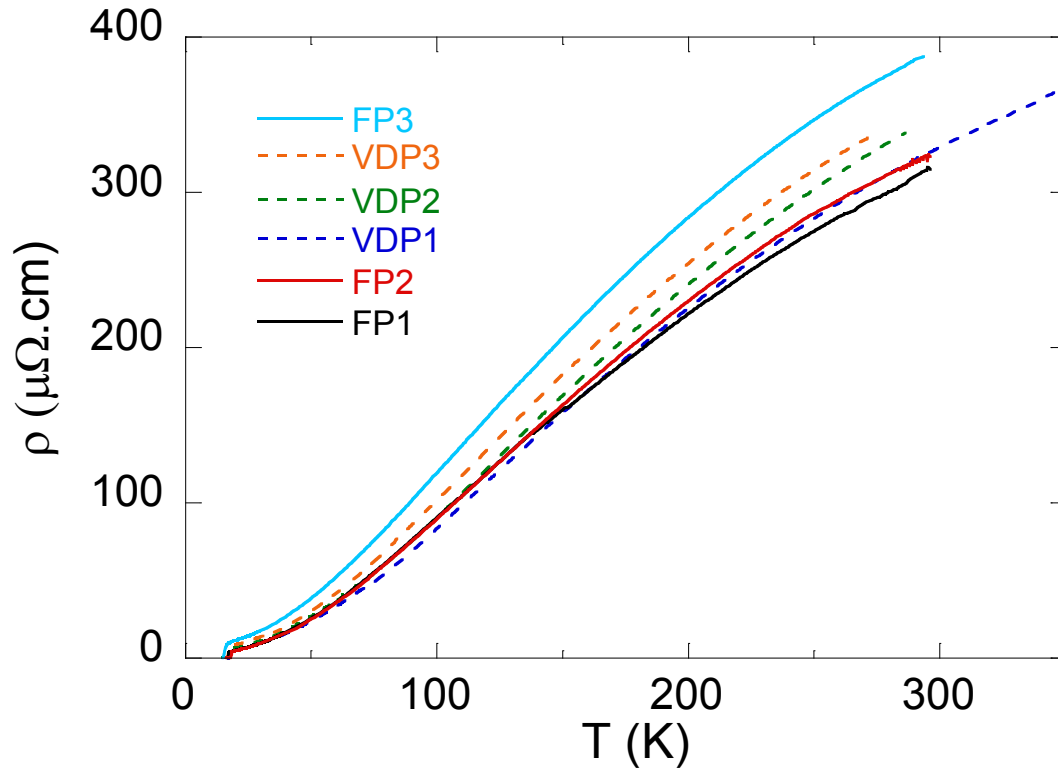
area of the FS sheets

Electron bands

Inner band : ~ 0.08 el/Fe

Outer band: ~ 0.11 el/Fe

Transport properties of LiFeAs : resistivity

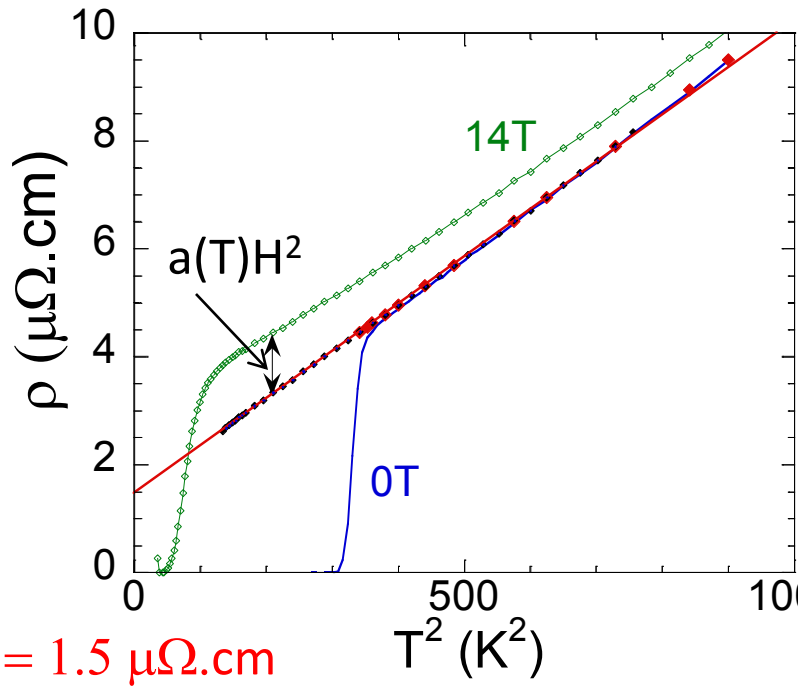


Correlation between T_c and resistivity at T_c

Compared to previous reports

Lower residual resistivity
larger RRR = $\rho(300\text{K})/\rho_0 \sim 250$

LiFeAs : a Fermi liquid compound?



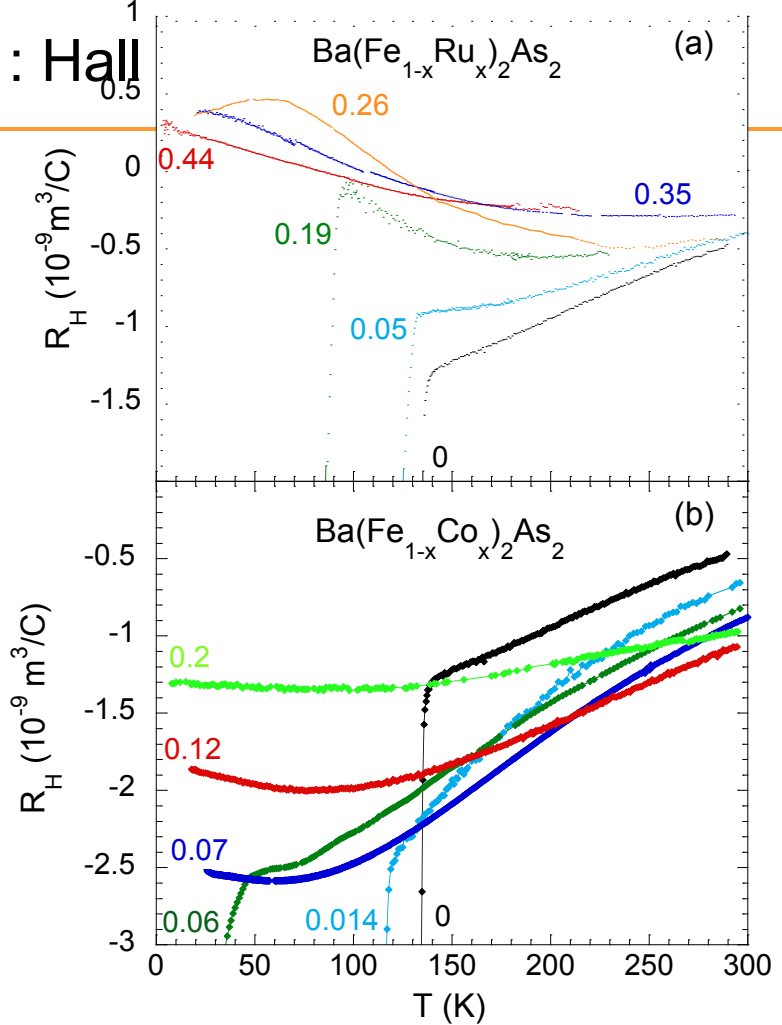
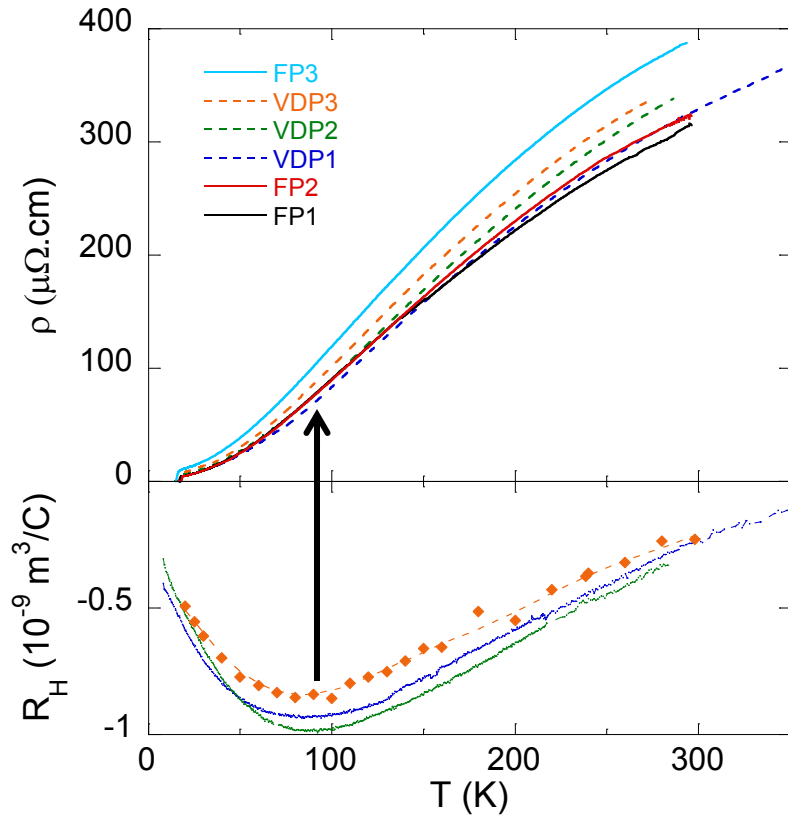
$$\rho(T) = \rho_0 + AT^2$$

$$A \approx 9 n\Omega.cm / K^2$$

Sample	T_c (K)	ρ_0 ($\mu\Omega.cm$)	RRR	A ($n\Omega.cm/K^2$)
VDP1	17.75	4.15	80	9.7
VDP2	17.6	2.93	119	9.3
VDP3	17	4.55	79	9.5
FP1	18.2	1.47	225	8.8
FP2	17.15	1.21	263	9.3
FP3	15.6	6.34	62	12.7
Ref.[2]	17.75	15.2	38	22
Ref.[3]	17.3	13	53	20

A increases with ρ_0 \longrightarrow Matthiessen's rule is not obeyed
Multiband material

Transport properties of LiFeAs : Hall

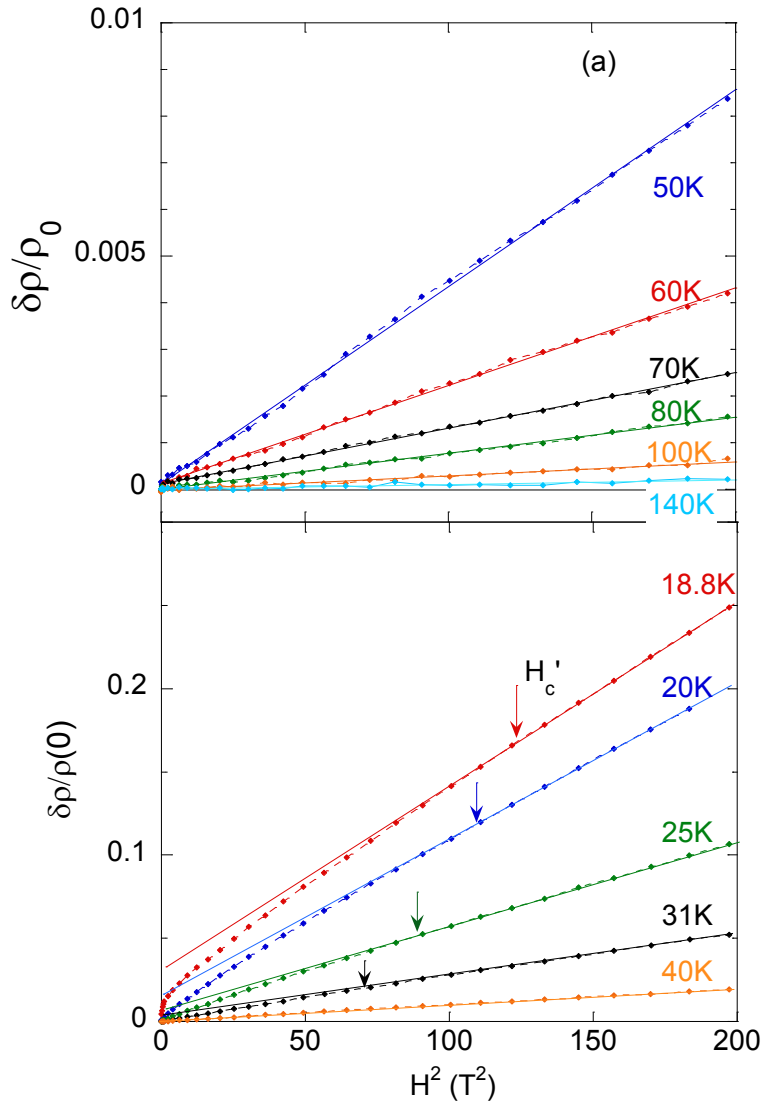


The Hall coefficient is negative

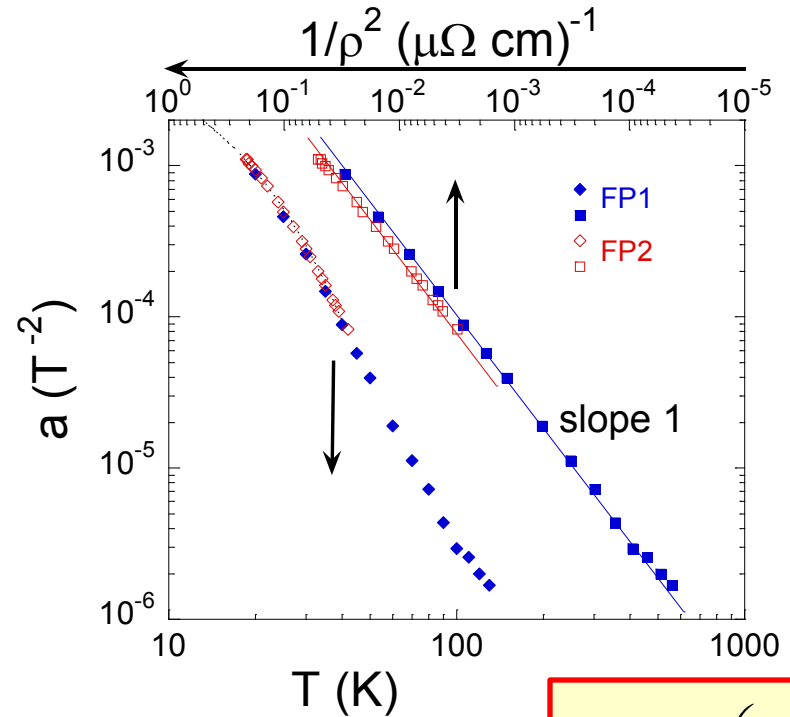
Transport dominated by the electrons as in BaFe_2As_2

Small minimum in R_H : corresponds to the change in curvature of $\rho(T)$

Transport properties of LiFeAs : Magnetoresistance



$$\frac{\delta\rho}{\rho(0)} = \frac{\rho(H) - \rho(0)}{\rho(0)} = a(T)H^2$$

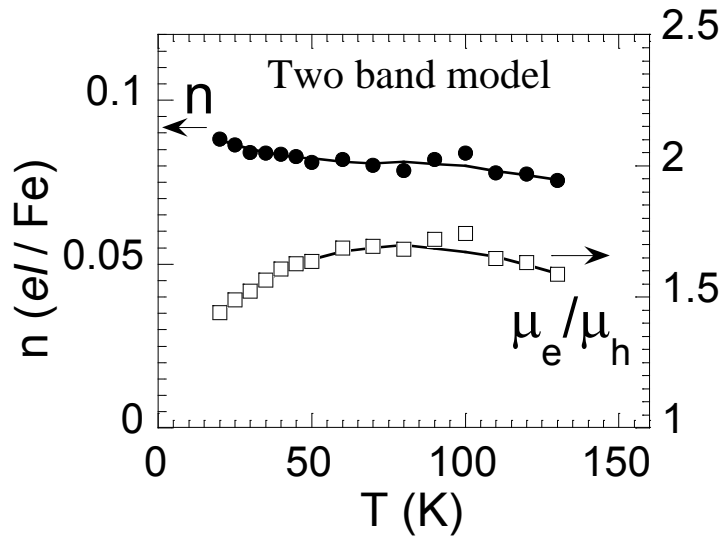


$$\frac{\delta\rho}{\rho(0)} \propto \left(\frac{H}{\rho}\right)^2$$

Kohler's rule

In a single band metal : $\delta\rho / \rho(0) \approx (\mu H)^2 \propto (H / \rho)^2$

Interpretation of the transport data in a two-band model



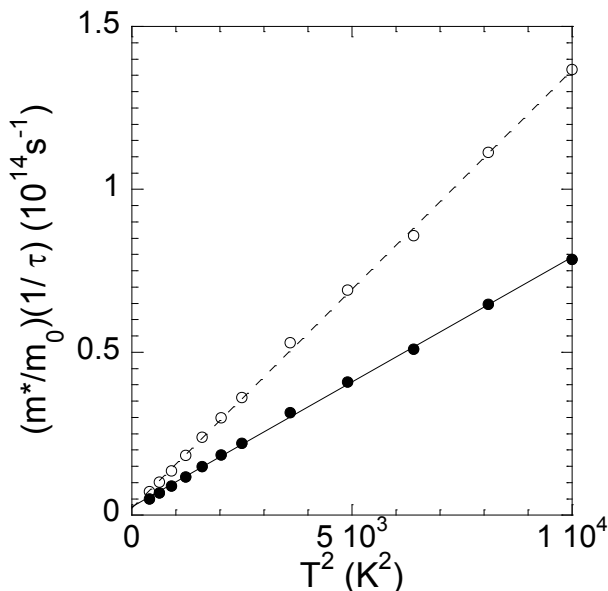
$$R_H / \rho = \mu_h - \mu_e$$

$$\frac{\delta\rho}{\rho(0)} = \mu_e \mu_h H^2$$

Ratio $\mu_e / \mu_h \approx 1.5$

Scattering rates $1/\tau \propto T^2$

➔ Fermi liquid behavior for both carriers

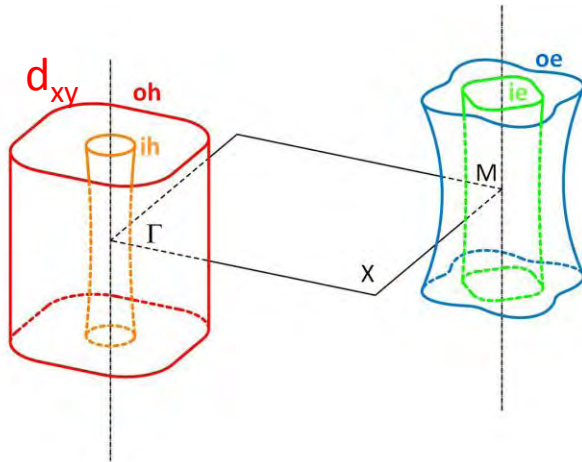


BUT : $n \approx 0.08$ carriers/Fe

More than twice less than given by
ARPES or Quantum oscillations experiments

Beyond the two band-model - 1

4 electronic bands with 4 different mobilities



Number of carriers taken from ARPES and dHvA data

$$\frac{1}{H^2} \frac{\delta\rho}{\rho(0)} = a(T) = \frac{\sigma_e \sigma_h (\mu_e + \mu_h)^2}{(\sigma_e + \sigma_h)^2} + \frac{\sigma_h}{(\sigma_e + \sigma_h)} A_h + \frac{\sigma_e}{(\sigma_e + \sigma_h)} A_e \quad A_h = \frac{\sigma_{ih} \sigma_{oh} (\mu_{ih} - \mu_{oh})^2}{\sigma_h^2}$$

Effective hole band

Effective mobility

$$\mu_h = \frac{\sigma_{ih}}{\sigma_h} \mu_{ih} + \frac{\sigma_{oh}}{\sigma_h} \mu_{oh}$$

Effective number of carriers

$$\sigma_h = n_h^{eff} e \mu_h$$

Can vary with temperature

Effective electron band

→ Two-band model with a Hall coefficient

$$R_H = \frac{(-n_e^{eff} \mu_e^2 + n_h^{eff} \mu_h^2) e}{(\sigma_e + \sigma_h)^2}$$

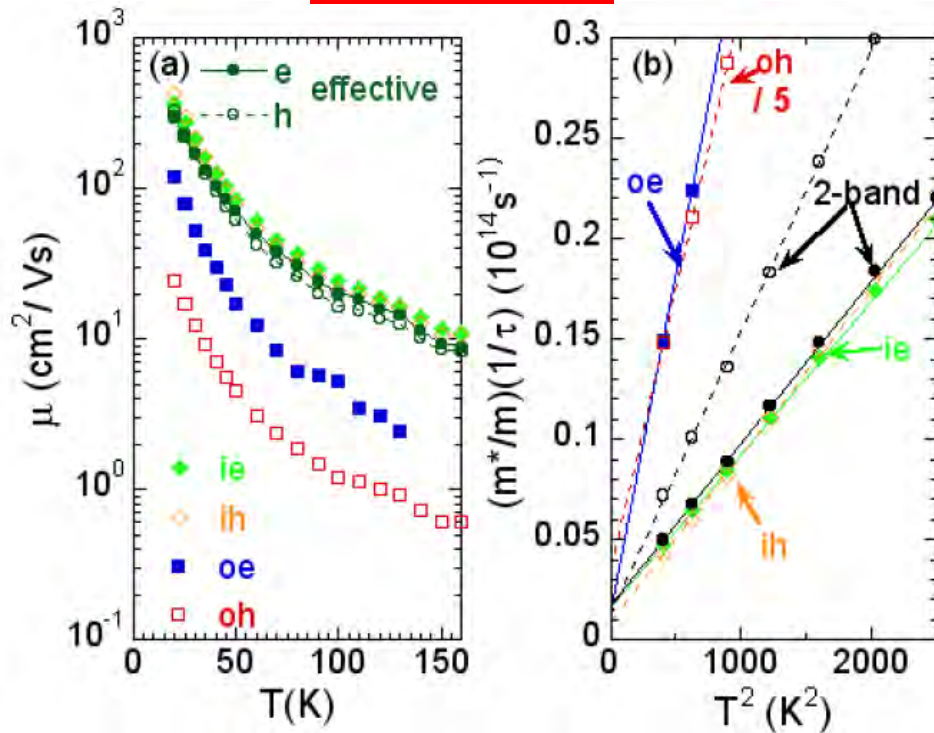
Beyond the two band-model - 2

Four different unknown parameters and only three experimental equations

➔ A unique solution cannot be acquired

Good solution must be such as $\mu_{ie} \geq \mu_{oe}$ and $0.02 h/Fe \leq n_h^{eff} \leq 0.06 h/Fe$

For instance : $n_h^{eff} = 0.05 h/Fe$ independent of T



$$\mu_{ie} \cong \mu_{ih}$$

electrons $\mu_{ie} / \mu_{oe} \cong 3-6$
 holes $\mu_{ih} / \mu_{oh} \cong 17$

$$\underline{n_e^{eff} = 0.13 e/Fe}$$

$$\mu_e^{eff} \cong \mu_h^{eff} \rightarrow \text{Kohler's rule}$$

For all the carriers :

$$1/\tau \propto T^2$$

For the outer hole band

ARPES measurements give a MDC half width: $\delta k = 18$ meV (A.A. Kordyuk et al. PRB 2011)

$(m^*/m_0)(1/\tau) \sim 8 \cdot 10^{13} \text{ s}^{-1}$ comparable with $\sim 2 \cdot 10^{13} \text{ s}^{-1}$ found here by transport

Measurements of the lifetimes for the other pockets should be very interesting

Conclusion

- Multiband description of the transport properties
- For the compounds with isovalent substitution Ru/Fe -
Description in a two band model seems reasonable
 - Number of carriers compatible with ARPES (Ru/Fe)
 - Fermi liquid behavior for both the holes and the electrons
 - Mobilities of holes and electrons comparable
- LiFeAs: low defect content
 - Comparison with ARPES and quantum oscillation measurements
 - Description more realistic of the transport properties by taking into account four different bands.
- Co substitution in BaFe_2As_2
 - Strong electron-hole asymmetry in the scattering rates: holes not visible in the transport properties
 - Difficult to reconcile with the ARPES data on the different lifetimes
 - Many theoretical questions....

Collaborators

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P. Bonville (Mössbauer)

LPS Véronique Brouet (ARPES)
H. Alloul
J. Bobroff , Y. Laplace, Y. Texier (RMN)

Co : Transport: FRA et al., Phys. Rev. Lett. **103** (2009).
ARPES: V. Brouet et al., Phys. Rev. B **80**, (2009).
RMN: Y. Laplace et al., PRB **80** (2009).
Mössbauer: P. Bonville, FRA et al. Europhys. Lett. **89** (2010)

Ni : Transport: A. Olariu, FRA et al., Phys. Rev. B **83** (2011).
Mössbauer: A. Olariu et al. arXiv 1106.1332.

Ru : Transport: FRA et al., Phys. Rev. B **81** (2010).
ARPES: V. Brouet, FRA et al., Phys. Rev. Lett. **105** (2010)

LiFeAs: Transport: FRA et al, submitted (2012).

