



# A DMFT insight into the Earth's core: many-electron effects in iron under extreme conditions

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# OUTLINE

- Inner core of Earth: iron alloy at extreme conditions
- Electronic correlations in the 3 phases of iron
- Fermi-liquidness and transport in hcp-Fe at Earth core conditions
- Lattice distorsions, correlations and transport in bcc-Fe





# Solid inner core of Earth

- radius ≈ 1200 km
- density  $\approx$  13 g/cm<sup>3</sup>
  - composition:
     mainly Fe+(~5-10% Ni, Si, S, O...)
- temperature 5200÷6500 K (uncertain)
- pressure 330÷370 Gpa
- age: may be < 1 bY
- anisotropy in seismic wave propagation:
- 3-4% faster  $\vec{v}$  along spin axis (Creager Nature 1992)
  - strong anisotropy (different Fe phases?) in the inner part R < 500 km</li>
  - weak anisotropy in the outmost part; different  $\vec{v}$  in hemispheres

Hirose *et al*. Annu. Rev. Earth Planet. Sci. 2013 Tkalčić Rev. of Geophys. 2015

# State of iron in the inner core of Earth



#### **Density functional theory (DFT):**

<u>bcc</u> α: Belonoshko *et al.* Nature 2004 Luo *et al.* PNAS 2010 Belonoshko *et al.* Nat. Geo. 2017

fcc γ: Mikhaylushkin *et al.* PRL 2007 hcp ε: Vočadlo *et al.*, Nature 2003

#### **Diamond anvil cell experiments:**

**bcc** *Fe* Hrubiak et al. arXiv 2018 *Fe*<sub>0.9</sub>*Ni*<sub>0.1</sub>: Dubrovinsky *et al.* Science 2007

hcp: Tateno *et al. Fe* (370 GPa, 5800 K): Science 2010
 *Fe*<sub>0.9</sub>*Ni*<sub>0.1</sub>: Geophys. Res. Lett. 39 L12305 (2012)

# Anisotropy of Fe inside the inner core

Relative velocity vs. angle to the rotation axis



#### Deuss Ann. Rev. Ear. Pl. Sci. 2014

DFT ab initio molecular dynamics

0.98

0.97

0.96

0

20

10

50

Polar/Crystallographic angle (ξ)

70

80

90

c/a ratio (hence, T) Belonoshko et al. Nature 2004,

Science 2008 Mattesini PNAS 2010

#### hcp-Fe:

Vočaldo et al.

anisotropy

sensitive to

Earth. Pl. Sci. Lett. 2009

hcp-Fe: velocity

low anisotropy in hcp (high-T c/a ~1.633)

#### bcc-Fe:

large anisotropy [100] vs. [111]



# Iron metal at low P-T and electronic correlations

Low P-T phase diagram of Fe



 narrow (*W*~6 eV) 3d band hybridized with wide 4s

*U*(3d)≈3-5 eV < *W* but large *J<sub>H</sub>*~0.9 eV

α-Fe: moderate QP renormalization(Z~0.5) in ARPES and theory

Schäfer et al. 2005 Sànchez-Barriga et al. 2009 , Sànchez-Barriga et al., 09, Sponza et al., PRB 17, Hausoel et al., Nat. Com. 17

PM  $\alpha$ -Fe and  $\alpha \rightarrow \gamma$  transition phonon spectra of PM  $\alpha$  and  $\gamma$ Leonov et al. PRL 11, SciRep. 14, Han et al. PRL 17



- α bcc, ground state, FM
- γ fcc, T> 1185 K (P=0)
- **δ** bcc, T > 1667 K (P=0)
- hcp, P-stabilized, paramagnetic, superconducting

Are correlations still relevant at high P and T?

$$\mathsf{P} \uparrow \longrightarrow \mathsf{V} \downarrow \longrightarrow \mathsf{W} \uparrow \mathsf{and} \mathsf{U}/\mathsf{W} \downarrow$$

at high pressure correlation effects can be ignored **?** 

- U/W < 1 in Fe even at P=0 (but  $J_H$  is more important)
- high T stabilizes high-entropy (local moment) phases
- for metals el.-el. scattering  $\sim T^2 \rightarrow$  significant at high T

# self-consistent DFT+DMFT method



- hybridization-expansion CT-QMC: segment picture (mostly),
   full U: Seth *et al.* Comp. Phys. Comm. 2016
- self-consistency over the charge density  $\rho$
- conductivity: Kubo formula neglecting vertex corrections
- Calculation for the Earth core atomic volume 7.05 Å<sup>3</sup>/atom
- U in the range  $3.3 \div 5 \text{ eV}$ ,  $J_H \sim 0.9 \text{ eV}$

### Fermi-liquid vs. non-Fermi liquid behavior of bcc, fcc and hcp-Fe at extreme conditions

Inverse quasiparticle lifetime:

 $\Gamma=Z Im[\Sigma(0)] \sim T^2$ in Fermi-liquid (FL) regime

Hence  $\Gamma/T^T$  in FLs



hcp ε-Fe: Fermi-liquid behavior in the full T range up to 6000 K bcc α-Fe: a strongly non-FL behavior for T > 1000 K fcc γ-Fe: Fermi-liquid in t2g, an intermediate case for eg Pourovskii *et al.* PRB 2013

# Why is the bcc phase more correlated? DFT DOS



Large peak in bcc DOS due to a Van-Hove singularity



Suppression of hybridization at low E:  $Im[\Delta(E_F)] \approx -[\pi D(E_F)]^{-1}$ 

Enhancement of correlations

Mravlje *et al.* PRL 2011 Leonov *et al.* PRL 2015

# Uniform magnetic susceptibility in applied field



Pauli behavior
 fcc and hcp

• Curie-Weiss  $\chi$  in bcc due to a local moment

# Earth core thermal conductivity "puzzle"

Low thermal conductivity is necessary to sustain the convection

For the geodynamo (liquid iron) <u>extrapolation of shock-wave results</u> Stacey&Andereson Phys. Earth Planet. Inter. 2001 about **30-50 W/(m·K)** 

ab initio DFT calculations:

thermal conductivity in liquid Fe: > 200 W/(m·K)

in iron-rich Fe-(Ni)-(Si)-(O): ~ 150 W/(m·K)

DFT calculations of el.-ph. electrical and thermal conductivities:

Pozzo *et al.* Nature 2012, PRB 2013 Earth Planet. Sci. Lett. 2014 de Koker *et al.* PNAS 2012 Shi *et al.* J. Phys. Condens. Matter 2011

### Formation of solid core anisotropy



#### requires κ well below 200 W/(m·K) see B. Buffet Nature 2012

ab initio DFT value for el-ph κ: pure ε-Fe ~300 W/(m·K) Fe-Si alloy ~235 W/(m·K)

### **Inconclusive experiments in DAC**

Konôpková *et al.* Nature 2016 direct measurements of  $\kappa$  +extrapolation

Ohta *et al.* Nature 2016 direct measurements of  $\rho$  +WF law



### Theory:

previous ab initio DFT calculations: only electron-phonon scattering included what is the contribuiton of electron-electron scattering?

### Fermi-liquid behavior of ε-Fe: DMFT self-energy at real frequencies

Fermi-liquid behavior of  $\Gamma$  stems from  $Im[\Sigma(\omega, T)] = const * (\omega^2 + \pi^2 T^2)$ 

![](_page_13_Figure_2.jpeg)

Fermi-liquid temperature scale  $T_{FL} \approx 0.1T_0 = 14\ 000\ K$ 

### Electron-electron contribution to electrical resistivity of ε-Fe: weak

![](_page_14_Figure_1.jpeg)

The electron-phonon  $\rho_{e-ph}$  is about 5.3 · 10<sup>-5</sup> $\Omega$  · cm from Pozzo *et al.* Earth & PI. Sci. Let. 2014

# Electron-electron contribution to thermal resistivity of $\epsilon$ -Fe: significant

![](_page_15_Figure_1.jpeg)

Smaller than predicted by Wiedemann-Franz with standard Lorenz number

$$L = \frac{\pi^2}{3} \left(\frac{k_B}{e}\right)^2$$

At T~6000 K κ<sub>ee</sub> ≈ **540 W/(m·K)** 

Electron-phonon  $\kappa_{e-ph} \approx 300$  W/(m·K) from Pozzo *et al.* Earth & PI. Sci. Let. 2014

Total 
$$\kappa_{tot} = (1/\kappa_{e-ph} + 1/\kappa_{ee})^{-1} \approx 190 \text{ W/(m \cdot K)}$$

### Reduction of the Lorenz number in Fermi liquids

In the semi-classical Boltzmann formalism:

$$\sigma = e^2 \int d\epsilon \Phi(\epsilon) (-f'(\epsilon)) \tau(\epsilon) \qquad \kappa = \frac{1}{T} \int d\epsilon \epsilon^2 \Phi(\epsilon) (-f'(\epsilon)) \tau(\epsilon)$$

$$\kappa/(\sigma T) = \frac{\pi^2}{3} \left(\frac{k_B}{e}\right)^2 = L \quad \text{, in SI units} \quad 2.44 \cdot 10^{-8} W\Omega$$

assuming frequency-independent lifetime  $\boldsymbol{\tau}$ 

But in Fermi liquids frequency dependence of  $\tau_{ee} \sim 1/Im[\Sigma(\omega)]$  is strong

$$1/\tau(\epsilon) = 1/\tau(\epsilon = 0) \cdot (1 + \epsilon^2/\pi^2 T^2 k_B^2)$$

and additional  $\epsilon^2$  in  $\kappa$  enhances contributions of high frequencies.

One obtains 
$$\kappa/(\sigma \cdot T) = L/1.54 = L_{FI}$$

thus significanly reducing  $\boldsymbol{\kappa}$ 

## Electron correlations and conductivity in bcc-Fe

![](_page_17_Picture_1.jpeg)

- bcc-Fe observed(?)

   in anvil cell experiments
   Dubrovinsky et al. Science 2008
   Hrubiak et al. arXiv:1804.05109
- dynamically unstable in the harmonic approximation at high P
- can be stablized by unharmonic vibrations

Belonoshko et al. Nat. Geo 2018

![](_page_17_Figure_6.jpeg)

fixed lattice calculations not reliable for bcc!

- 1. perform *ab initio* DFT mol. dynamics (MD) simulations
- 2. run DFT+DMFT for a set of MD snapshots

#### 3x3x3 undistorted

![](_page_17_Picture_11.jpeg)

#### 3x3x3 MD snapshot

![](_page_17_Picture_13.jpeg)

## non-Fermi-liquid effects in bcc washed away by lattice distorsions!

### **bcc-Fe DFT DOS at EIC volume**

- van Hove singularity due to drives nonFL behavior of e<sub>q</sub> orbitals
- it's almost washed away by lattice vibrations

![](_page_18_Figure_4.jpeg)

### bcc-Fe: Suppression of conductivity by distortions

![](_page_19_Figure_1.jpeg)

thermal conductivity of fully distorted cells ~4 times lower than in perfect bcc

can be understood

$$r_{t} = rac{1}{\kappa_{latt}} + rac{1}{\kappa_{el.-el}}$$

where  $\kappa_{latt} \approx 280 \text{ W/(m·K)}, \kappa_{el.-el.} \approx 800 \text{ W/(m·K)}$ 

K<sub>1</sub>

### Electrical and thermal resistivity vs. distorsions

![](_page_20_Figure_1.jpeg)

 $\lambda = 0$ : transport is determined by electron-electron scattering

 $\lambda = 1$ : what determines the conductivity?

### Optical conductivity vs. distorsions in bcc Fe

optical conductivity for perfect and fully distorted case for different  $\Sigma(\omega)$ :

![](_page_21_Figure_2.jpeg)

In the distorted limit enhancing  $\langle \Sigma(\omega) \rangle$  seems to have little impact

#### the conductivity is fixed by the distorsion levels

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### References

LP, Miyake, Simak, Ruban, Dubrovinsky, Abrikosov PRB **87,** 115130 (2013) Vekilova, LP, Abrikosov, Simak, PRB **91,** 245116 (2015) LP, Mravlje, Georges, Simak, Abrikosov, New J. Phys. **19,** 073022 (2017) LP, Mravlje, Alfe ... (in preparation)

review: LP  $\Psi_k$  Highlight 11/2018 & Topical Review JPCM (in press)

# bcc Fe: hybridization suppression at low

![](_page_23_Figure_1.jpeg)

Large peak in bcc DOS due to a Van-Hove singularity Suppression of hybridization at low  $Im[ \Delta (E_F)] \approx -[\pi D (E_F)]^{-1}$