Three-terminal quantum-dot thermoelectrics

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Outline

Introduction

Quantum dots and Coulomb blockade Quantum dots and Thermoelectrics Three-terminal thermoelectrics

Coulomb-coupled conductors

Coulomb-blockade regime Chaotic cavities Resonant tunneling

Harvesting bosons

Phonons Magnons Microwave photons

Summary





Introduction



Quantum dots

Confine electrons in all three spatial directions, quasi 0-dimensional "Artificial atom"



Characteristic energy scales

- Charging energy $E_{ch} = E_C (N C_g V_g / e)^2$, $E_C = \frac{e^2}{2C}$
- Level quantization $\Delta \varepsilon$
- Tunnel couplings $\Gamma = 2\pi |t|^2 \rho$



Coulomb blockade



- Charging energy $E_{ch} = \frac{e^2}{2C} (N C_g V_g / e)^2$
- Charge fixed away from degeneracy points
- Transport only at degeneracy points
- Coulomb oscillations



Energy quantization and level scheme



- Finite level spacing in small quantum dots $E_N = E_{ch} + \sum_{i=1}^N \varepsilon_i$
- Addition energy $\Delta = 2E_C + \Delta \varepsilon$
- $\Delta \varepsilon \ll k_{\rm B}T, eV$: many levels involved \rightarrow metallic island
- $\Delta \varepsilon \gg k_{\rm B}T, eV$: only single level \rightarrow quantum dot



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Transport regimes



- Sequential tunneling
 - Tunneling of single electron
 - First order in Γ
 - Real occupation of the dot
 - Energy conservation
 - Dominant on resonance $k_{\mathsf{B}}T \gg \Gamma, |\varepsilon|$ or $eV \gg \Gamma, |\varepsilon|$



Transport regimes



- Cotunneling
 - Tunneling of two electrons
 - Second order in Γ
 - Virtual occupation of the dot
 - Energy conservation only for total process
 - Dominant off resonance $|\varepsilon| \gg k_{\rm B}T, eV, \Gamma$



Transport regimes



- Resonant tunneling
 - Many electrons tunnel
 - Nonperturbative in Γ
 - Complicated many-body effects (Kondo)
 - Dominant for strong coupling $\Gamma \gg k_{\rm B}T, eV$



Anderson model

Single-level quantum dot



- Level position ε (tunable)
- Coulomb energy U
- Tunnel coupling $\Gamma_r = 2\pi |t_r|^2 \rho$ (often tunable)
- Temperature T (tunable)
- Voltage V (tunable)



Master equation approach

- Probability P_{χ} to find quantum dot in state $\chi \in \{0, \uparrow, \downarrow, d\}$
- Occupation probabilities obey master equation

$$\dot{P}_{\chi} = \sum_{\chi'} W_{\chi\chi'} P_{\chi'}$$

• Transition rates $W_{\chi\chi'}$ from Fermi's golden rule

$$W_{\uparrow 0} = \sum_{r} \Gamma_r f_r(\varepsilon)$$

$$W_{\uparrow d} = \sum_{r} \Gamma_r [1 - f_r(\varepsilon + U)]$$

• Current $I = \sum_{\chi\chi'} W^I_{\chi\chi'} P_{\chi'}$



Experimental realizations



Kuemmeth et al., Nano Letters 2008

- 2-dimensional electron gas
- Metallic nanoparticles
- Self-assembled quantum dots



Experimental realizations



Dirks et al., Nat. Phys. 2011

Leturcq, Nat. Phys. 2009

- Carbon nanotubes
- Graphene
- Nanowires



Thermoelectrics

- 1823 Seebeck effect: Heat \rightarrow current
- 1834 Peltier effect: Current \rightarrow Cooling

Advantages of thermoelectrics:

- No moving parts
- Scalable to the nanoscale
- Heat is ubiquitous

Disadvantage

• Low efficiency, small power



Mesoscopic physics and thermoelectrics

Fundamental research

- Theoretical analysis of thermopower of quantum dots, quantum point contacts Proetto PRB 1991, Beenakker and Staring PRB 1992, Nakpathomkun et al. PRB 2010
- Experiments on quantum dots Staring et al. EPL 1993, Dzurak PRB 1997, Godijn et al. PRL 1999, Scheibner et al. PRL 2005, PRB 2007, Svensson et al. NJP 2012, 2013
- Experiments on quantum point contacts Molenkamp et al. PRL 1990

Can mesoscopic systems be useful for thermoelectric applications?

- Quantum wires and wells for thermoelectrics Hicks and Dresselhaus PRB 1993
- Sharp spectral features increase thermoelectric performance Mahan and Sofo PNAS 1996
- Electronic refrigerator Edwards et al. APL 1993, PRB 1995, Prance et al. PRL 2009



Three-terminal thermoelectrics



Sánchez, Büttiker, PRB 2011



McClure et al., PRL 2007

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- Connection to Coulomb-drag setups
- Crossed heat and charge currents
- Separation of heat source and rectifier
- Energy harvesting

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Coulomb-coupled conductors



Coulomb-blockade regime

Sánchez, Büttiker, PRB 2011



- Coulomb-coupled quantum dots
- Exchange energy but no particles
- Conductor dot: two cold reservoirs
- Gate dot: single hot reservoir



Coulomb-blockade regime



- Drive current by temperature bias
- Energy-dependent, asymmetric tunnel barriers
- Optimal heat to charge current conversion
- One energy quantum of the bath transfers one charge quantum



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Coulomb-blockade regime



- Power P = IV
- Efficiency $\eta = P/J_g$
- Device reaches Carnot efficiency η_C at stopping voltage
- Efficiency at maximum power $\eta_{maxP} = \eta_C/2$



Chaotic cavities

BS, Sánchez, Jordan and Büttiker, PRB 2012



- Capacitively coupled chaotic cavities
- Open quantum dots: Large number N of transport channels
- How do current and power vary with N?
- Asymmetric, energy-dependent transmissions $T_r = T_r^0 eT_r' \delta U$



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Chaotic cavities

• Current
$$I = \frac{\Lambda}{\tau_{RC}} k_{\mathsf{B}}(\Theta_1 - \Theta_2)$$

- Asymmetry parameter $\Lambda = \frac{G_L G'_R G_R G'_L}{(G_L + G_R)^2}$ where $G_r = \frac{e^2}{h} T_r^0$ and $G'_r = \frac{e^3}{h} T'_r$
- RC time τ_{RC} determined by effective conductance and capacitance of double cavity
- Current independent of channel number, $I \sim 0.1 \,\mathrm{nA}$
- Power scales as 1/N, similar to Coulomb-blockade for a few open channels $P \sim 1 \, {\rm fW}$
- Efficiency scales as $1/N^2$, few percent of η_C



Power versus conductance



- Coulomb blockade regime: Power grows linear with conductance
- Open contacts: Power drops as inverse conductance
- Maximal power should be achieved for single channel
- \Rightarrow Consider resonant tunneling through quantum dots



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Jordan, BS, Sánchez and Büttiker, PRB 2013



- Central cavity in thermal equilibrium with hot reservoir
- Cavity connected to two cold electronic reservoirs via quantum dots
- Quantum dots host single resonant level with width γ and energy $E_{L,R}$



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Scattering matrix approach

- Charge current $I_j = \frac{2e}{h} \int dE \ T_j(E) [f_j(E) f_{\mathsf{C}}(E)]$
- Energy current $J_j = \frac{2}{h} \int dE \ ET_j(E) [f_j(E) f_{\mathsf{C}}(E)]$

• Transmission
$$T_j(E) = \frac{\gamma^2}{(E-E_j)^2 + \gamma^2}$$

- Heat current *J* from hot reservoir
- Conservation of charge and energy

$$0 = I_{\mathsf{L}} + I_{\mathsf{R}}$$
$$0 = J + J_{\mathsf{L}} + J_{\mathsf{R}}$$



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- Numerically optimize ΔE , γ and V for maximal power
- Optimal values: $\Delta E \approx 6k_{\rm B}T$, $\gamma \approx k_{\rm B}T$
- Maximal power $P_{\text{max}} \sim 0.4 (k_{\text{B}} \Delta T)^2 / h$, about $0.1 \, \text{pW}$ at $\Delta T = 1 \, \text{K}$
- Efficiency at maximum power $\eta_{\text{maxP}} \sim 0.2 \eta_{\text{C}}$



Scaling



- Swiss cheese sandwich with self-assembled quantum dots
- Dot positions do not have to match
- Dot size of 100 nm^2 yields 10 W/cm^2 at $\Delta T = 10 \text{ K}$
- Robust with respect to fluctuations of level positions



Harvesting bosons



Phonons

Entin-Wohlman et al. PRB 2010, PRB 2012



- Quantum dot coupled to electronic reservoirs and phonon bath
- Linear-response thermoelectrics
- Left-right and particle-hole symmetry broken: $\Gamma_L(E) \neq \Gamma_R(E)$
- Flux-dependence of response coefficients in Aharonov-Bohm geometry



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Magnons

BS and Büttiker, EPL 2012



- Quantum dot coupled to ferromagnetic electrodes and ferromagnetic insulator
- Bridge between energy harvesting and spin caloritronics
- Drive pure spin current or spin-polarized charge current by magnons



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Microwave photons

Bergenfeldt, Samuelsson, BS, Flindt and Büttiker, arXiv 2013



- Double quantum dots connected via superconducting cavity
- Combines circuit QED and thermoelectrics
- Separate hot and cold part by macroscopic distance
- Reduce leakage heat currents



Summary



Summary

- Three-terminal quantum-dot thermoelectrics
- Coulomb-coupled conductors
 - Coulomb blockade: High efficiency, small power
 - Open dots: Small efficiency, small power
 - Resonant tunneling: Large power, good efficiency
- Boson-driven heat engines
 - Phonons (hard to control)
 - Magnons: spintronics
 - Microwave photons: circuit QED



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