

# What are the indications for the pseudogap in the frequency and temperature dependence of the optical conductivity ?

D. van der Marel

Collaboration with

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E. van Heumen<sup>1,3</sup>, M. K. Chan<sup>4</sup>, X. Zhao<sup>4</sup>, Y. Li<sup>4</sup>, M. Greven<sup>4</sup>,  
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# Optical Conductivity of *c* Axis Oriented $\text{YBa}_2\text{Cu}_3\text{O}_{6.70}$ : Evidence for a Pseudogap

C. C. Homes, T. Timusk, R. Liang, D. A. Bonn, and W. N. Hardy, PRL 71 (1993)

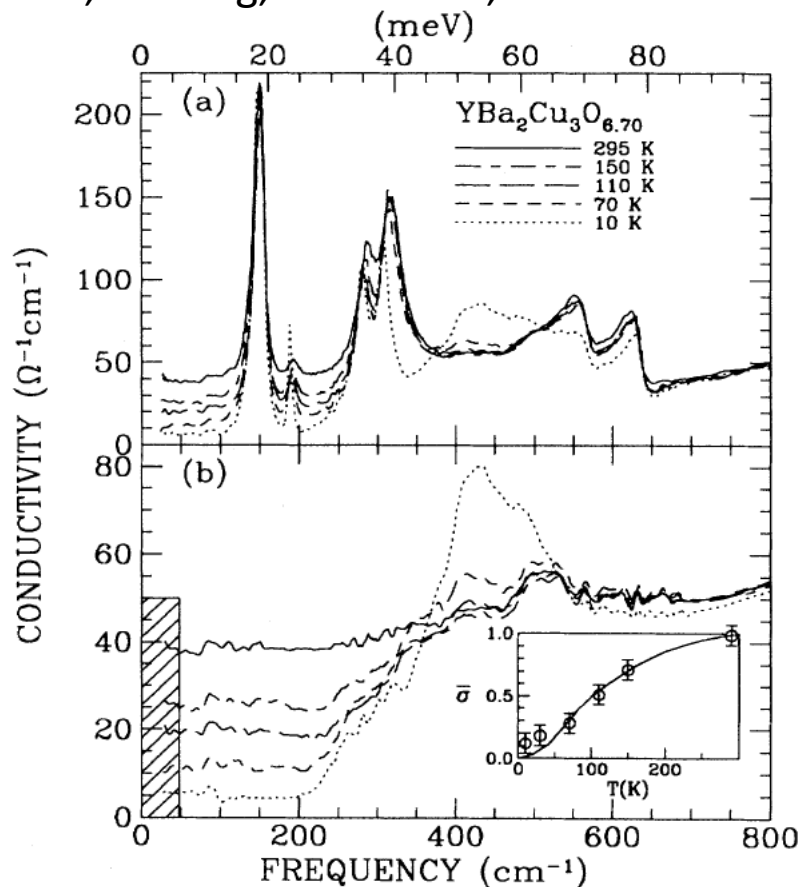
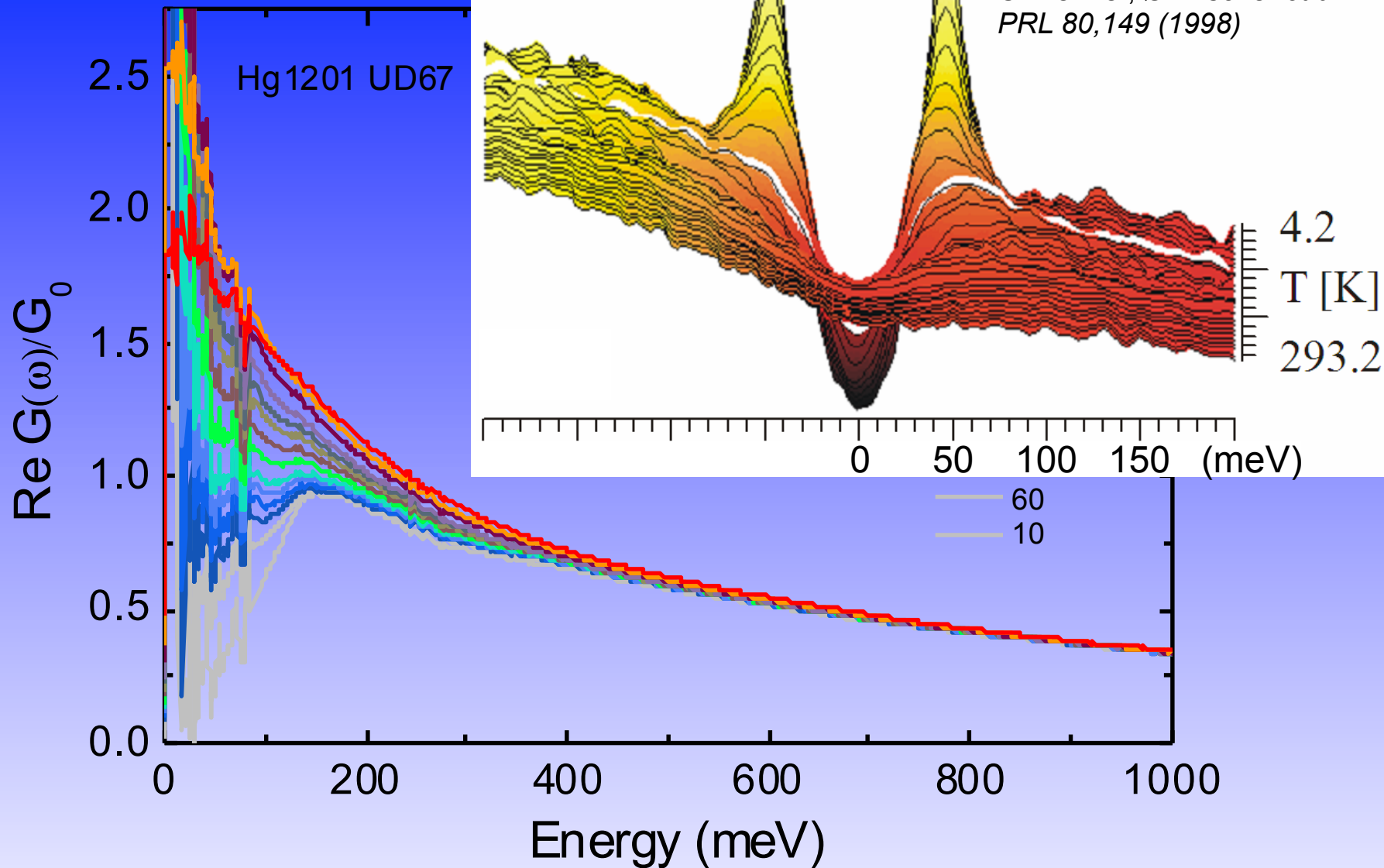
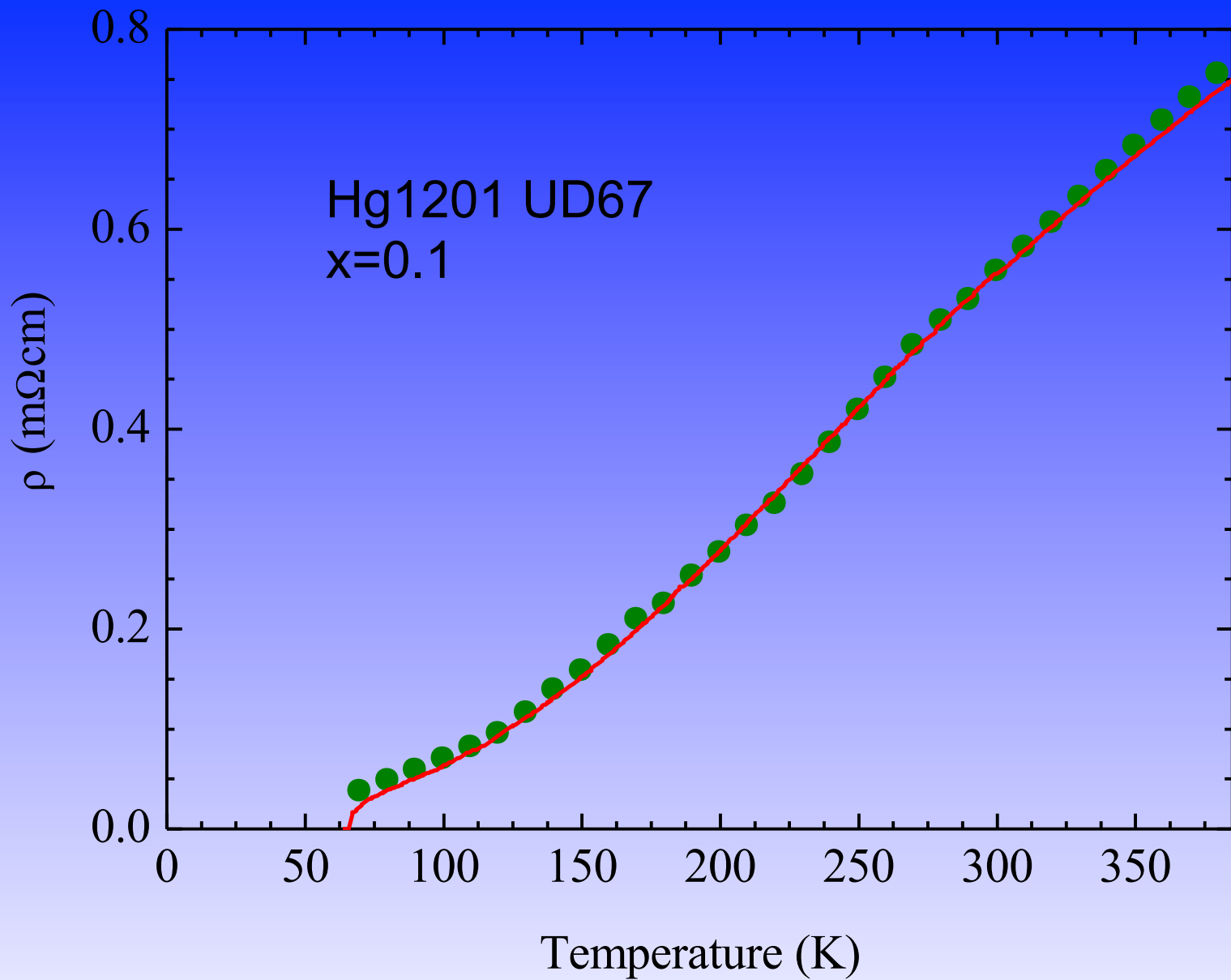


FIG. 2. The optical conductivity of  $\text{YBa}_2\text{Cu}_3\text{O}_{6.70}$  along the *c* axis from  $\approx 25$  to  $800 \text{ cm}^{-1}$  obtained by a Kramers-Kronig analysis of the reflectance with (a) the phonons at 150, 192, 286, 317, 557, and  $630 \text{ cm}^{-1}$  present and (b) subtracted to yield the electronic background. Note that the formation of a pseudogap is clearly visible well above  $T_c$  (63 K). The shaded area represents the spectral weight of the





# Optical response of a classical charged fluid

$$\sigma(\omega) = \frac{ne^2 / m}{\tau^{-1} - i\omega}$$

Generalization to interacting electrons in cuprates

$$\frac{G(\omega, T)}{G_0} = \frac{\pi K}{\tau^{-1}(\omega) - i\omega m^*(\omega)}$$

K = Integrated spectral weight

W Götze & P Wölfle, PRB 6, 1226 (1972)

JW Allen & JC Mikkelsen, PRB 15, 2952 (1977)

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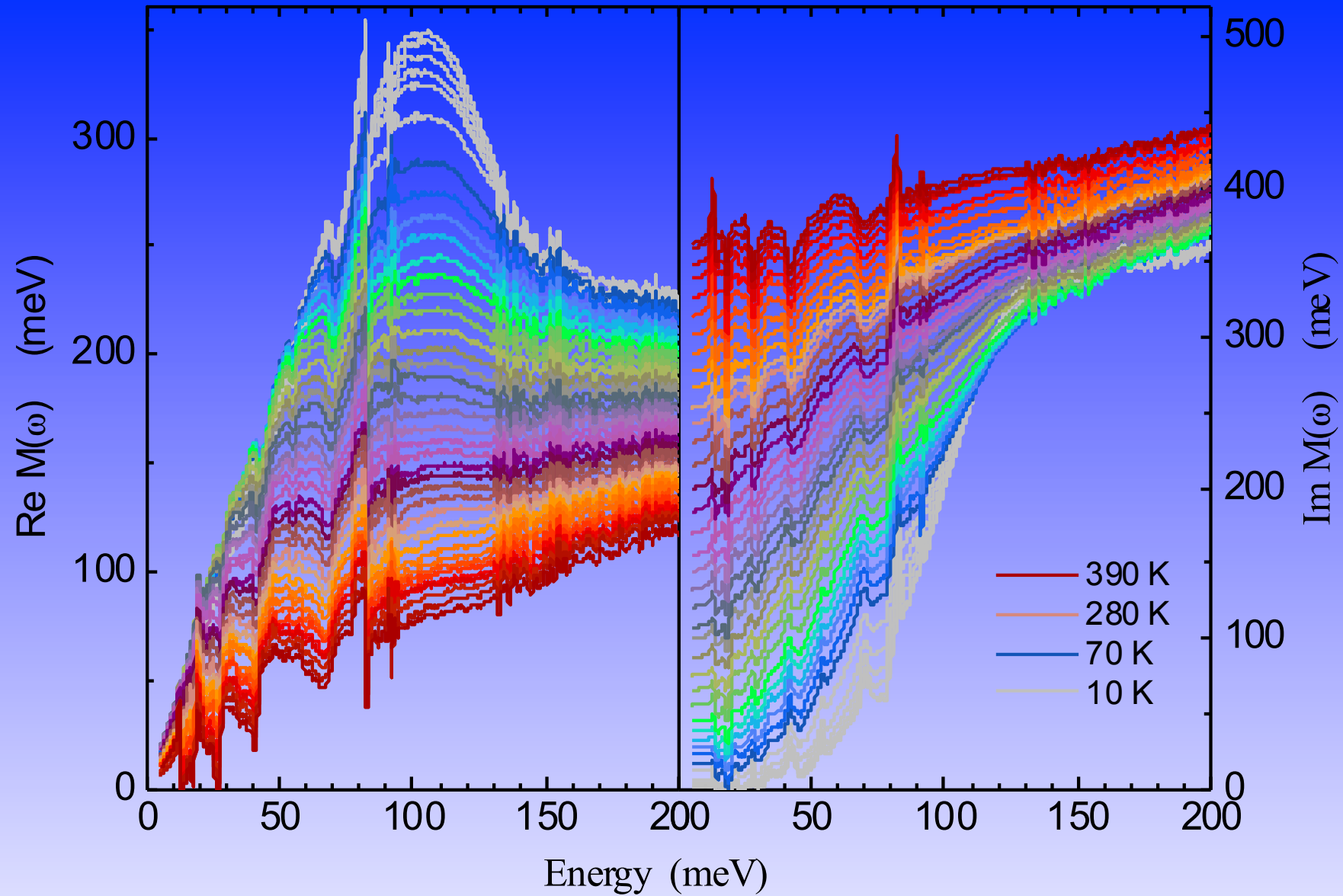
$$\frac{G(\omega, T)}{G_0} = \frac{\pi K}{M_2(\omega) - iM_1(\omega) - i\omega}$$

Straightforward inversion of the experimental data:

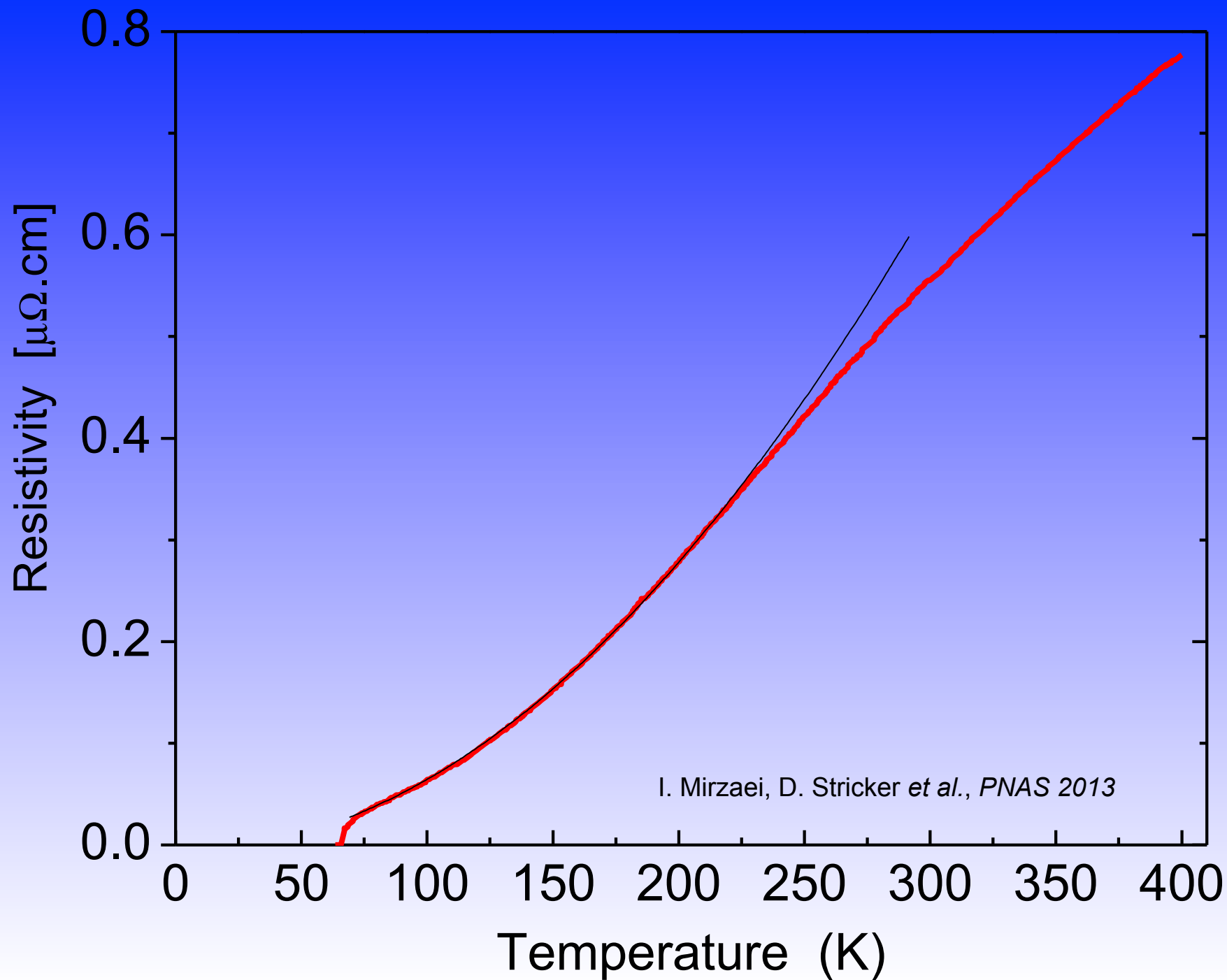
$$M_2(\omega) = \pi K \cdot \text{Re} \frac{G_0}{G(\omega, T)}$$

W Götze & P Wölfle, PRB 6, 1226 (1972)

JW Allen & JC Mikkelsen, PRB 15, 2952 (1977)







I. Mirzaei, D. Stricker *et al.*, *PNAS* 2013

# Fermi liquid

Single particle life time:  $\tau_{sp}(\varepsilon, T) \propto \left[ \varepsilon^2 + \pi^2 (k_B T)^2 \right]^{-1}$

Optical relaxation rate:  $\left\{ \begin{array}{l} 1 / \tau_{opt}(\omega, T) \propto (\hbar\omega)^2 + (p\pi k_B T)^2 \\ p = 2 \end{array} \right\}$

**R. N. Gurzhi, Sov. Phys. JETP 35, 673 (1959)**

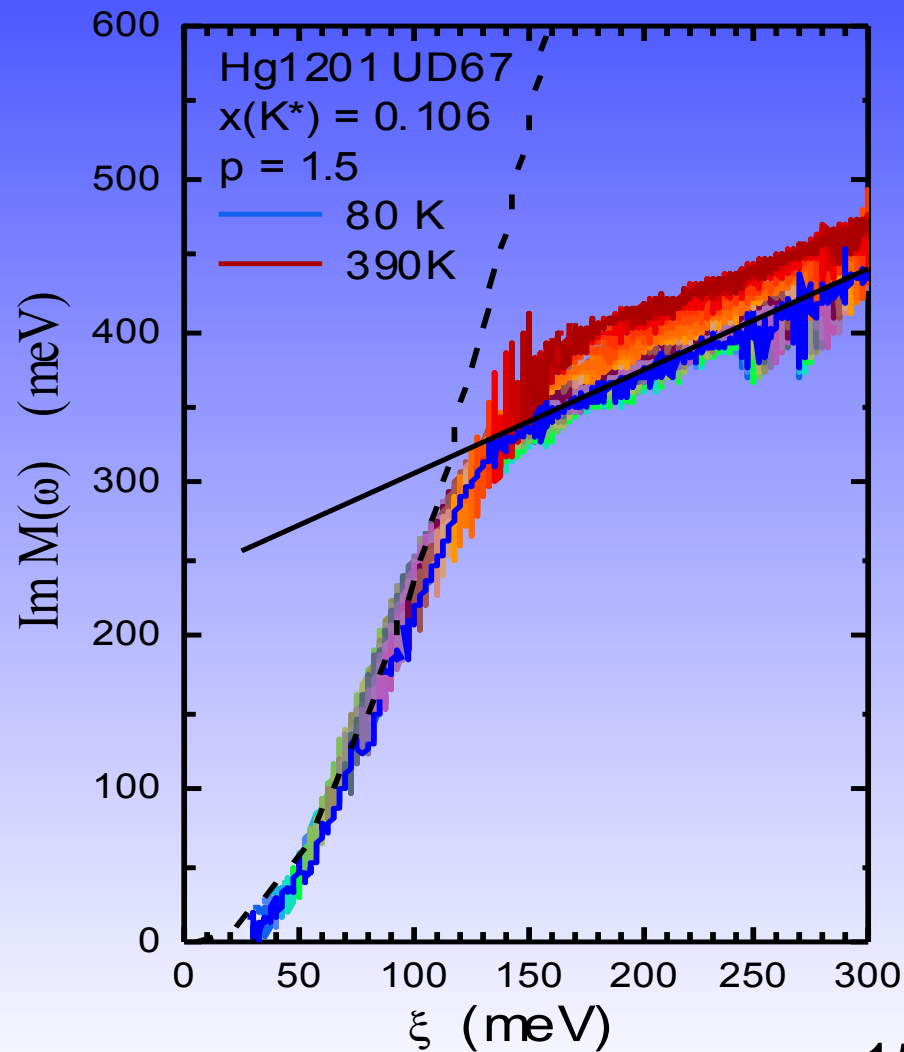
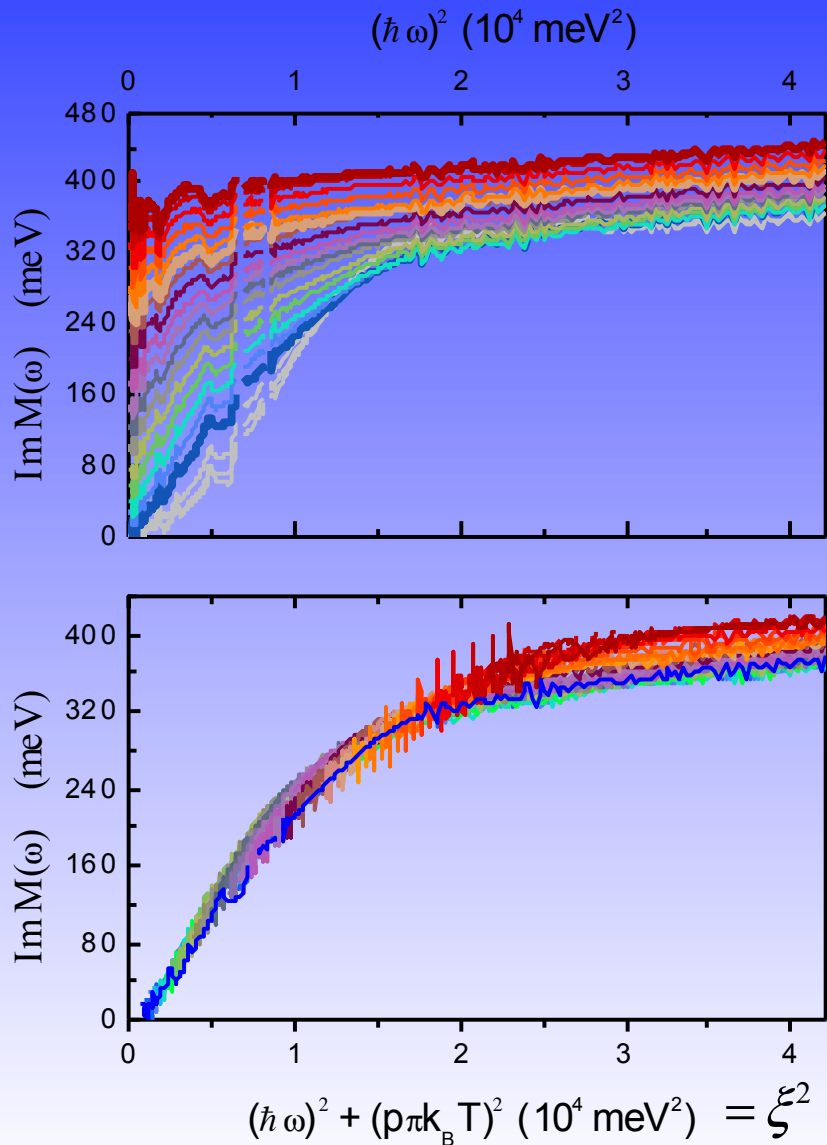
**D. L. Maslov & A. V. Chubukov, PRB 86, 155137 (2012)**

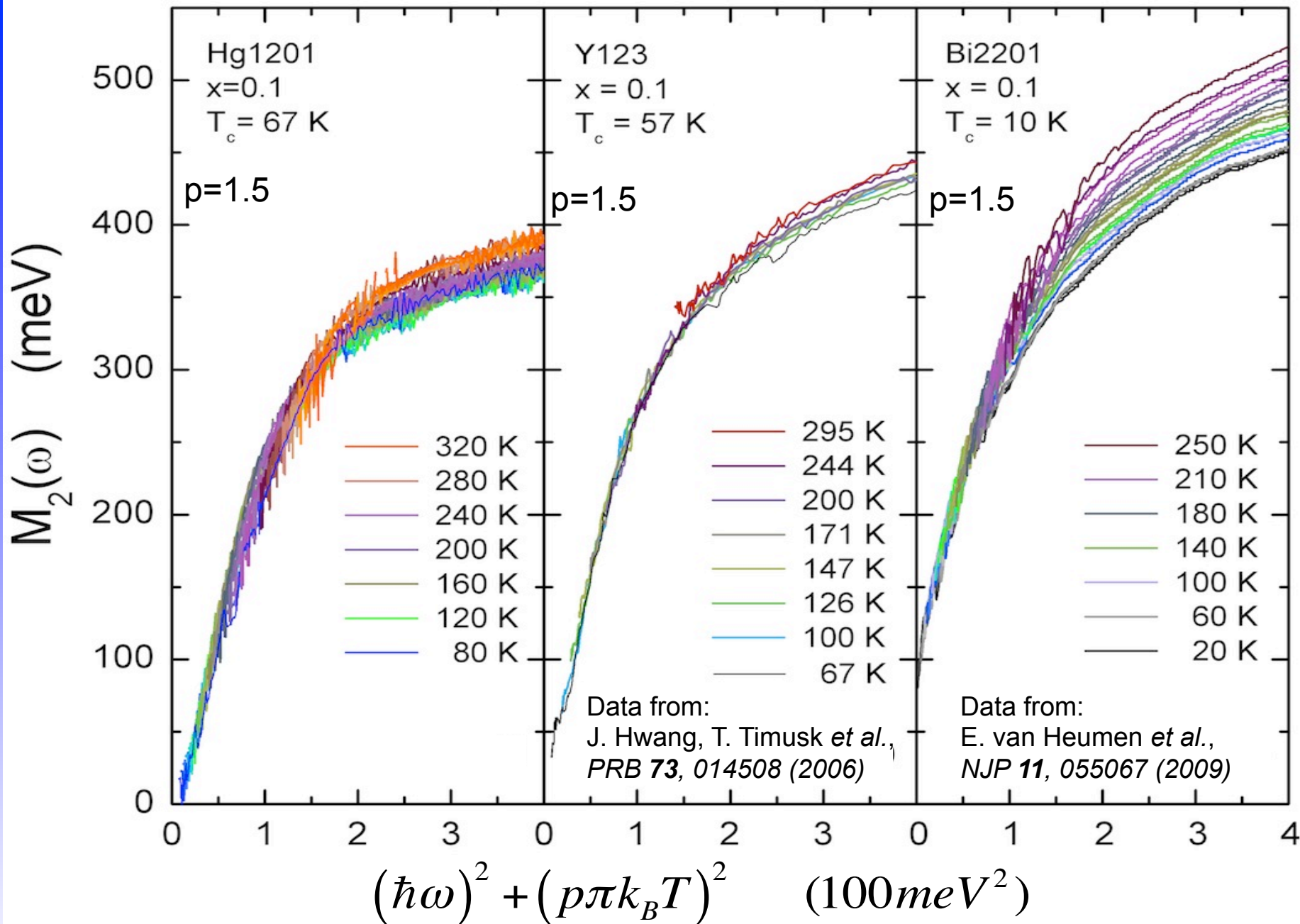
**C. Berthod *et al*, PRB 87, 115109 (2013)**



# Fermi-liquid

## Optical signature: scaling collapse





# Strong coupling theory

P.B. Allen  
PRB 1971

$$\frac{G(\omega, T)}{G_0} = \frac{i\pi K}{\omega + M(\omega, T)}$$

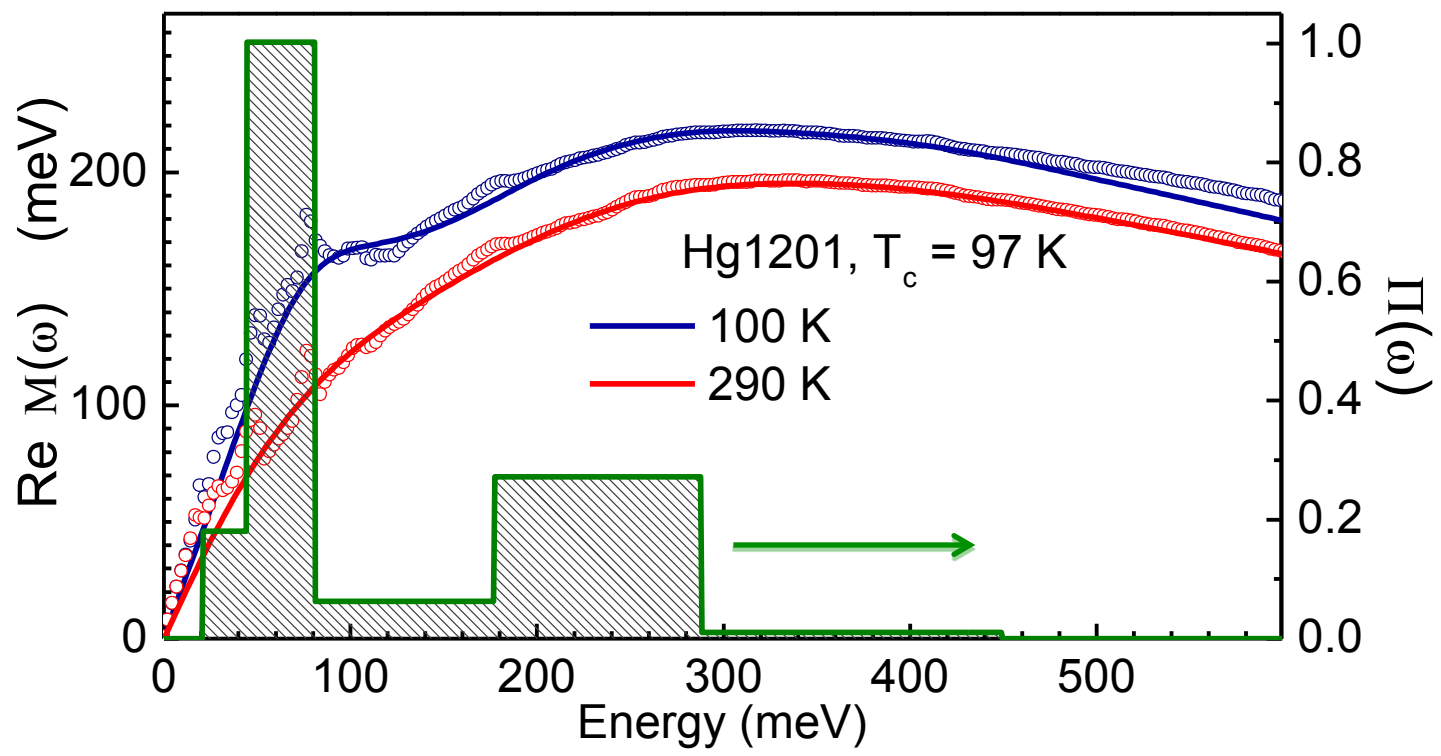
$$M(\omega, T) = \omega \left[ \int_{-\infty}^{\infty} \frac{n_F(\omega + \varepsilon, T) - n_F(\varepsilon, T)}{\omega - \Sigma(\omega + \varepsilon, T) + \Sigma^*(\varepsilon, T)} d\varepsilon \right]^{-1} - \omega$$

$$\Sigma(\omega, T) = \int d\varepsilon \int d\omega' \tilde{\Pi}(\omega') \left[ \frac{n_B(\omega') + n_F(\varepsilon)}{\omega - \varepsilon + \omega' + i\delta} + \frac{n_B(\omega') + 1 - n_F(\varepsilon)}{\omega - \varepsilon - \omega' - i\delta} \right]$$

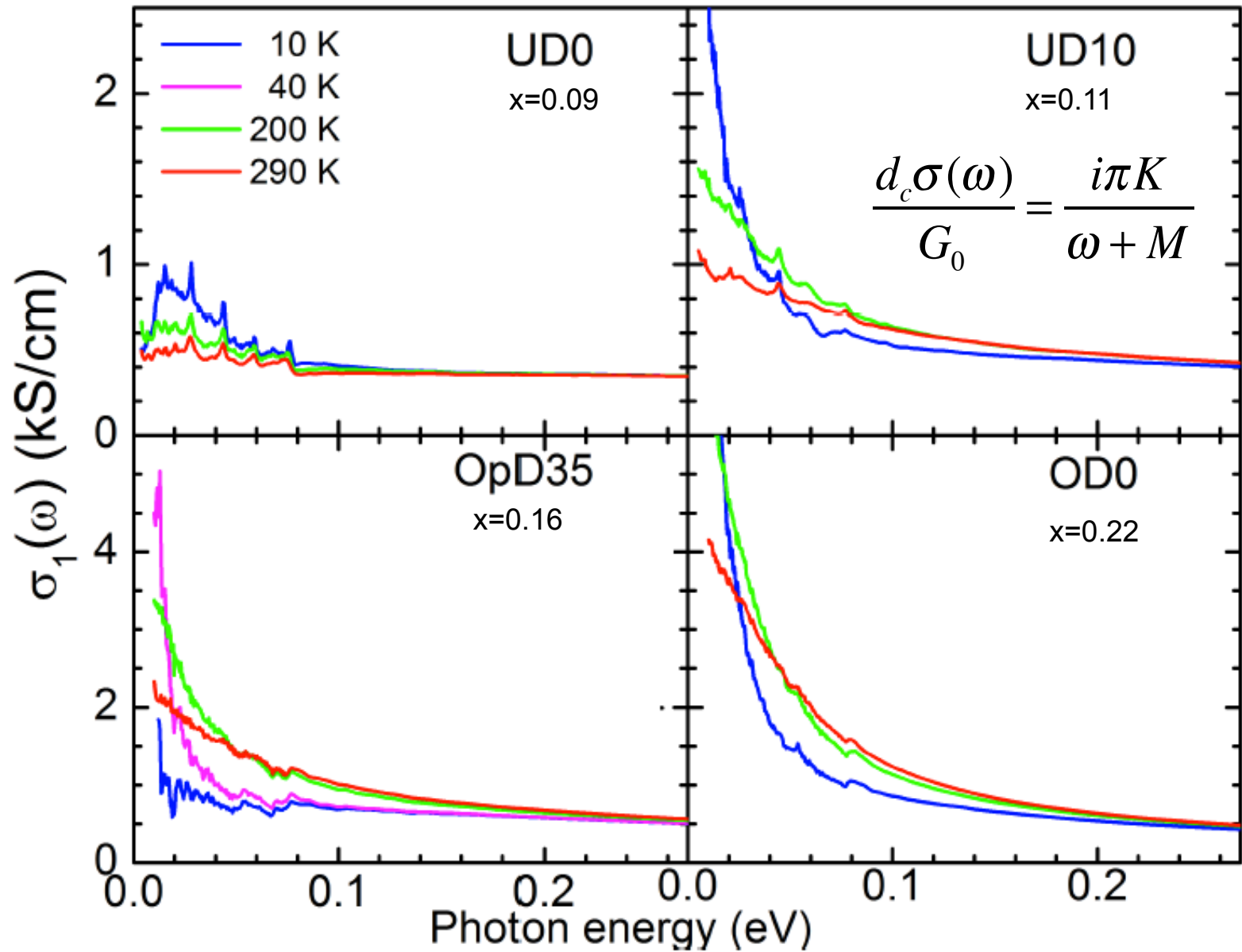
**The glue function:  $\tilde{\Pi}(x)$  can be fitted to experimental data**

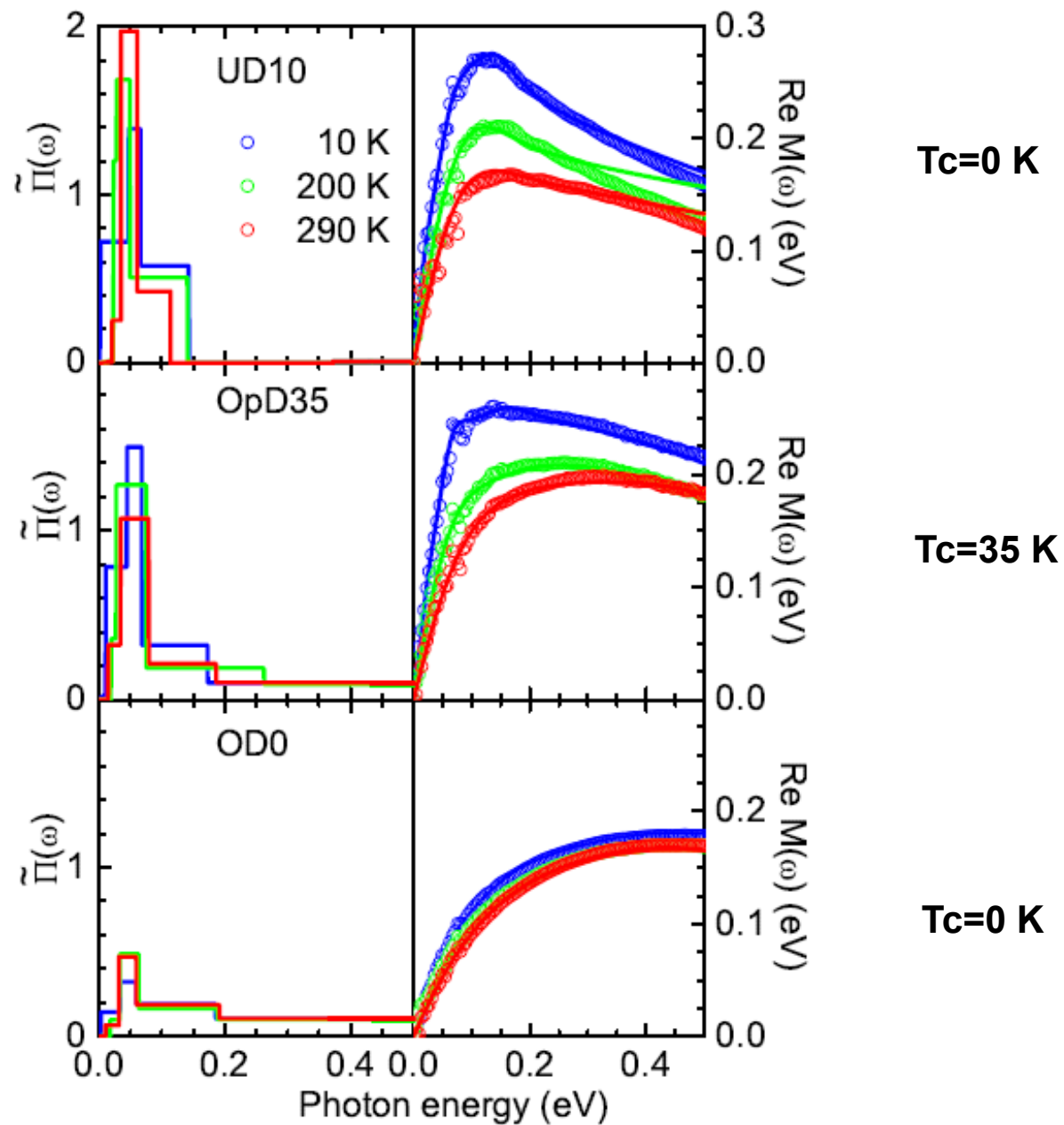
# Optical Memory function

Fit of  $\Pi(\omega)$  to the experimental data at 100 and 290 K simultaneously

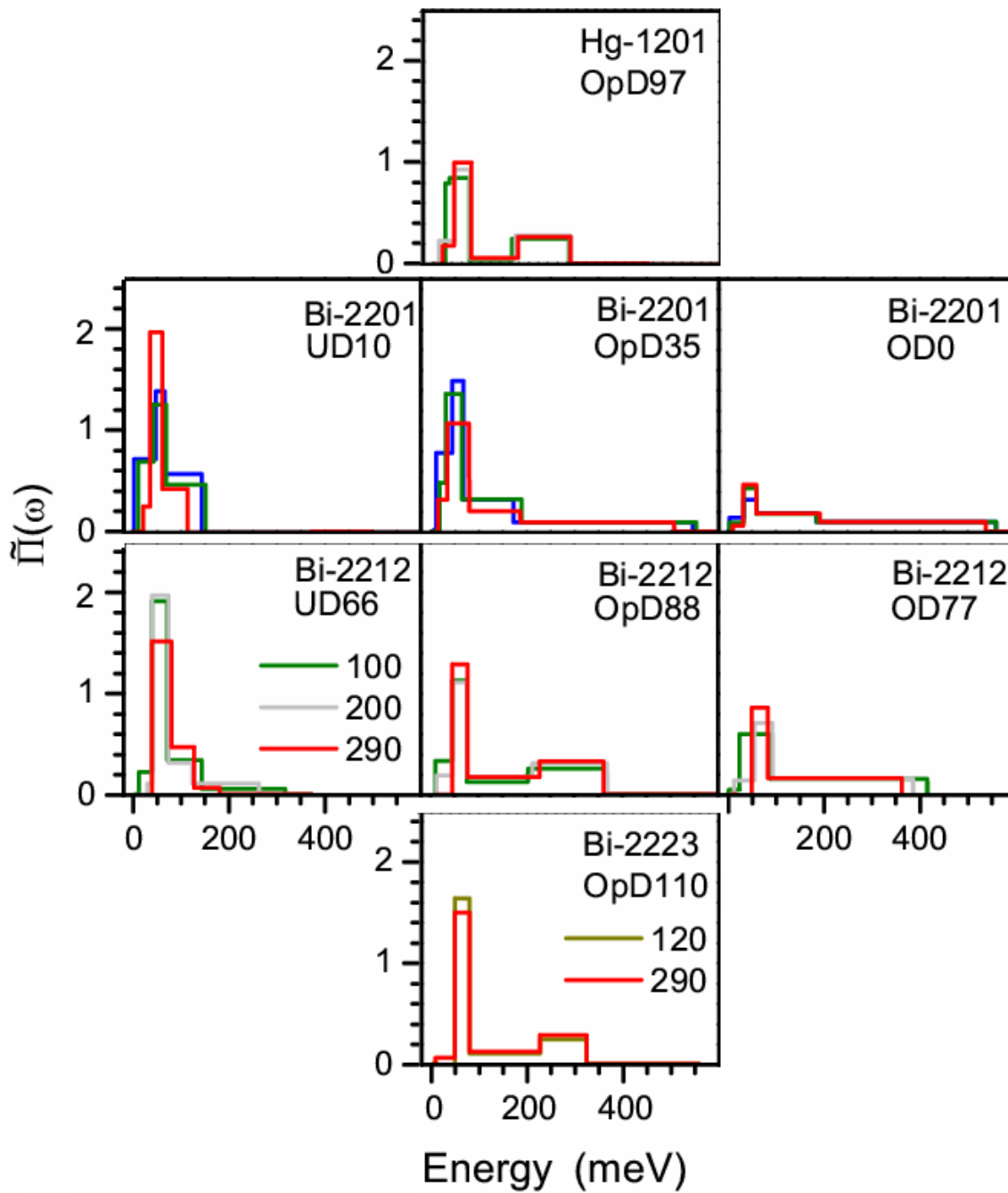


# Bi2201

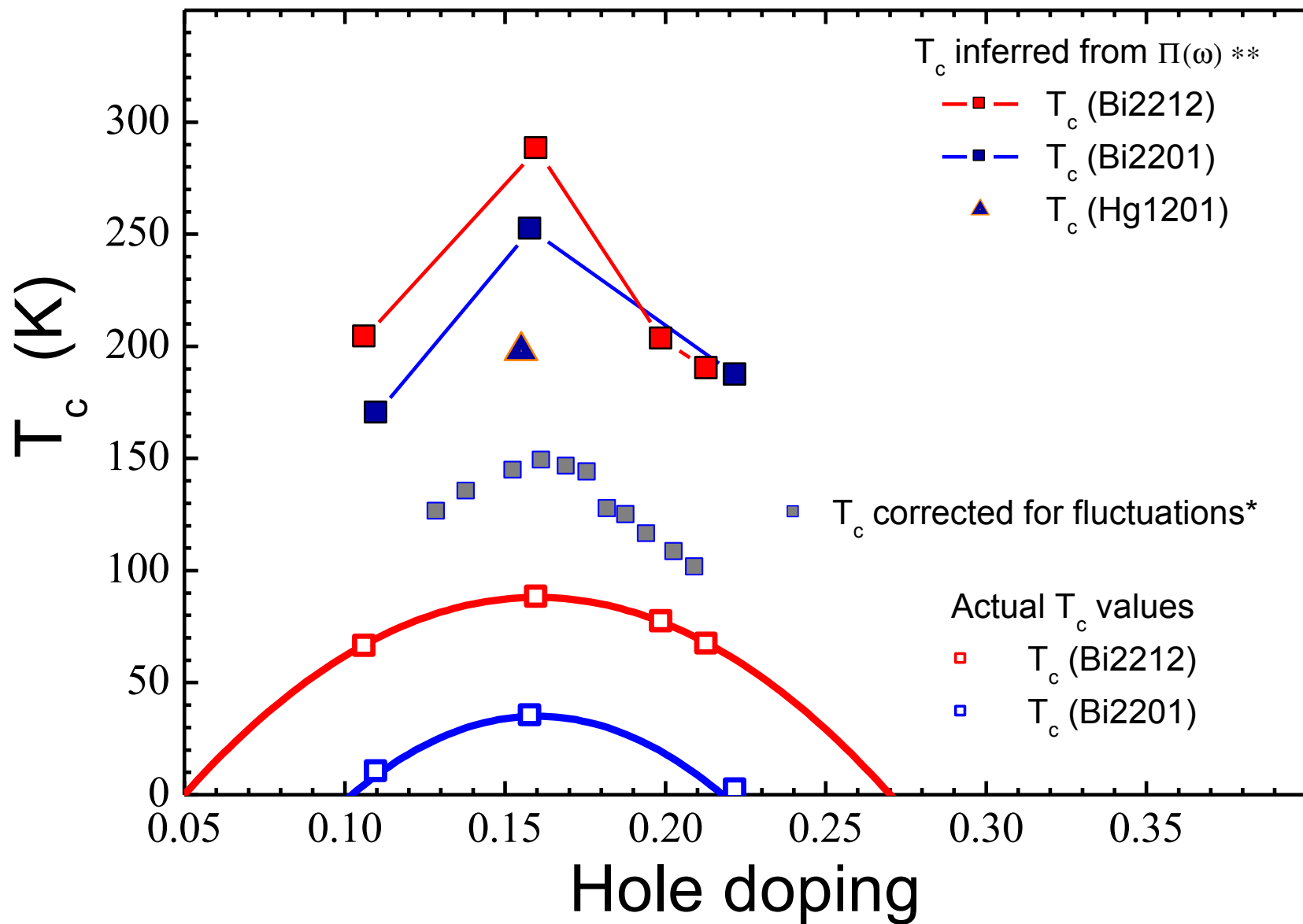






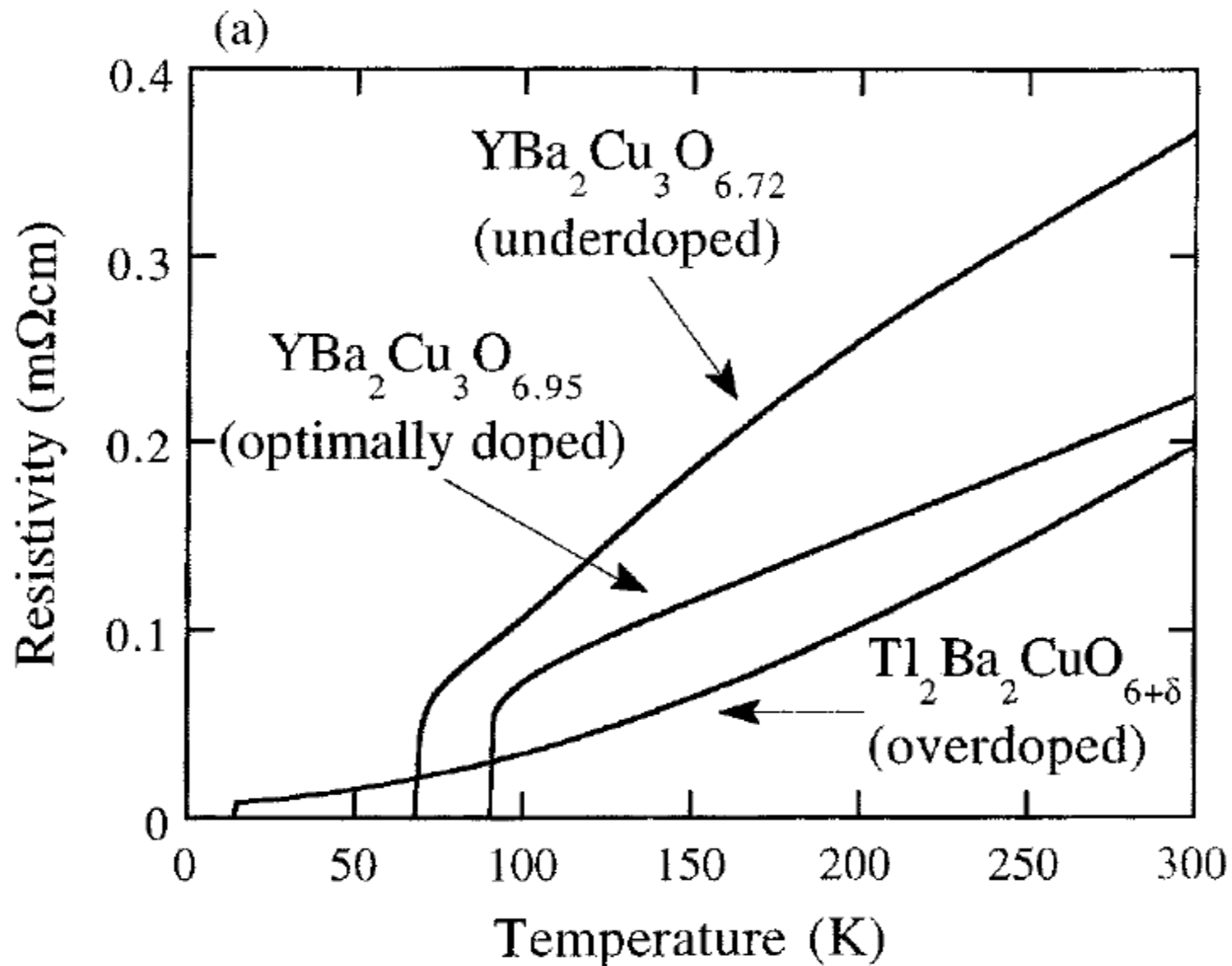


# Possible relation of fluctuation spectrum to $T_c$

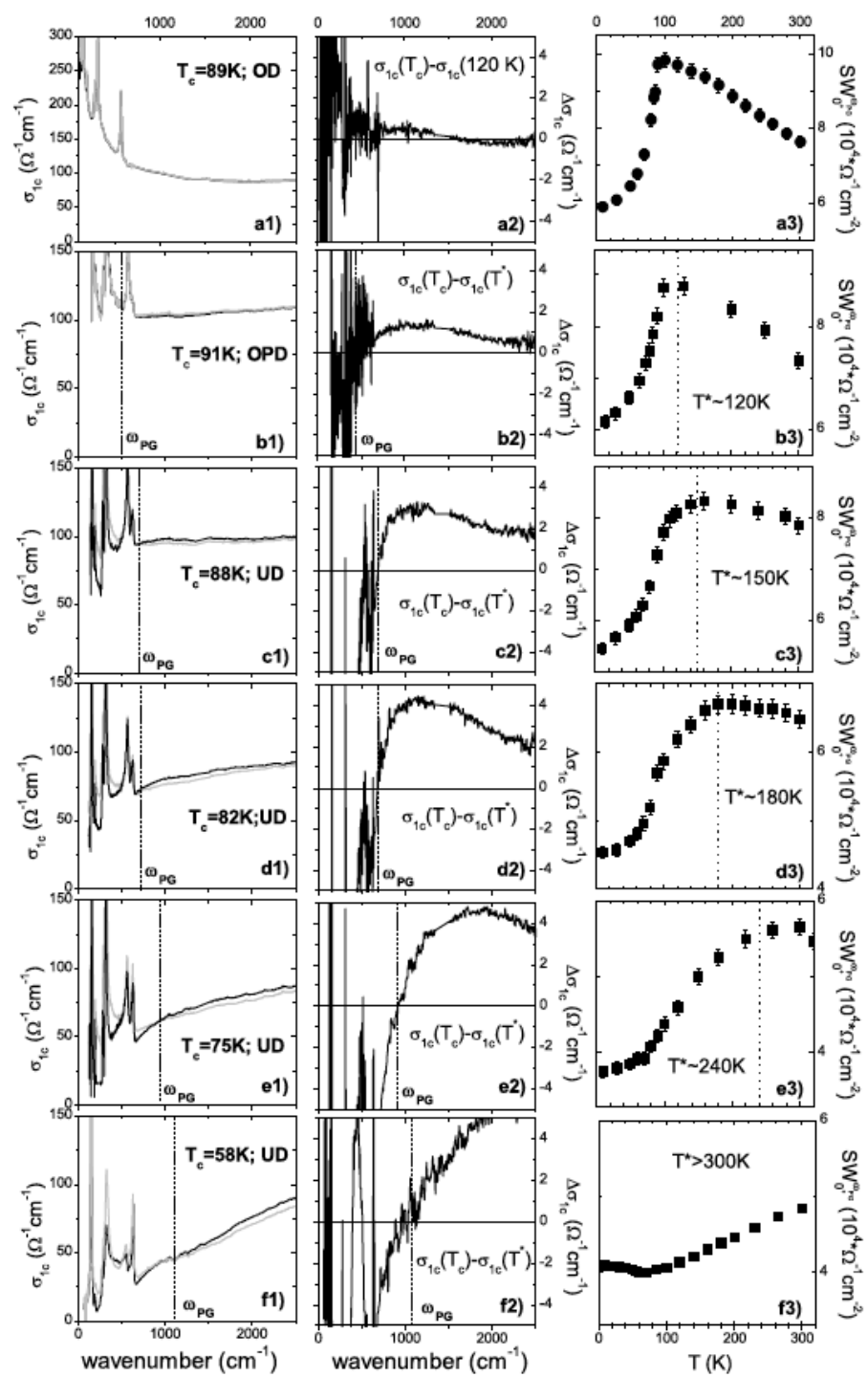
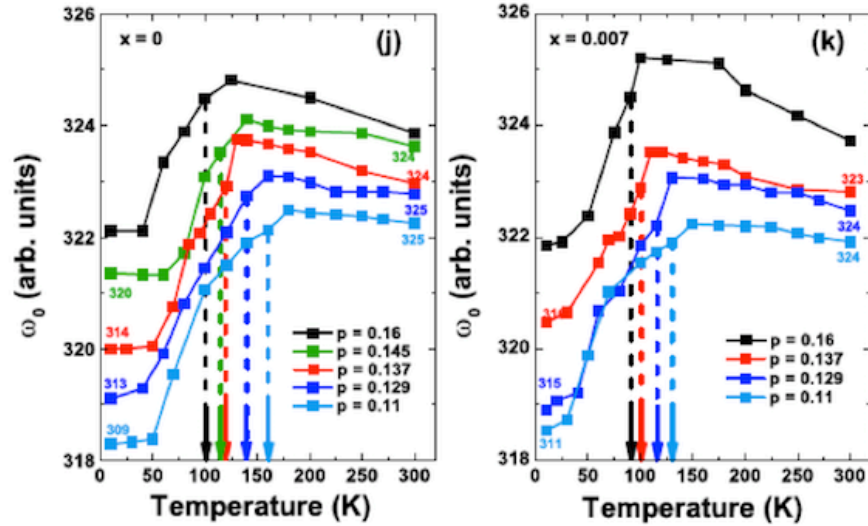


\*\* E. van Heumen et al., PRB 79, 184512 (2009)

\* J. Tallon et al., PRB 83, 092502 (2011)

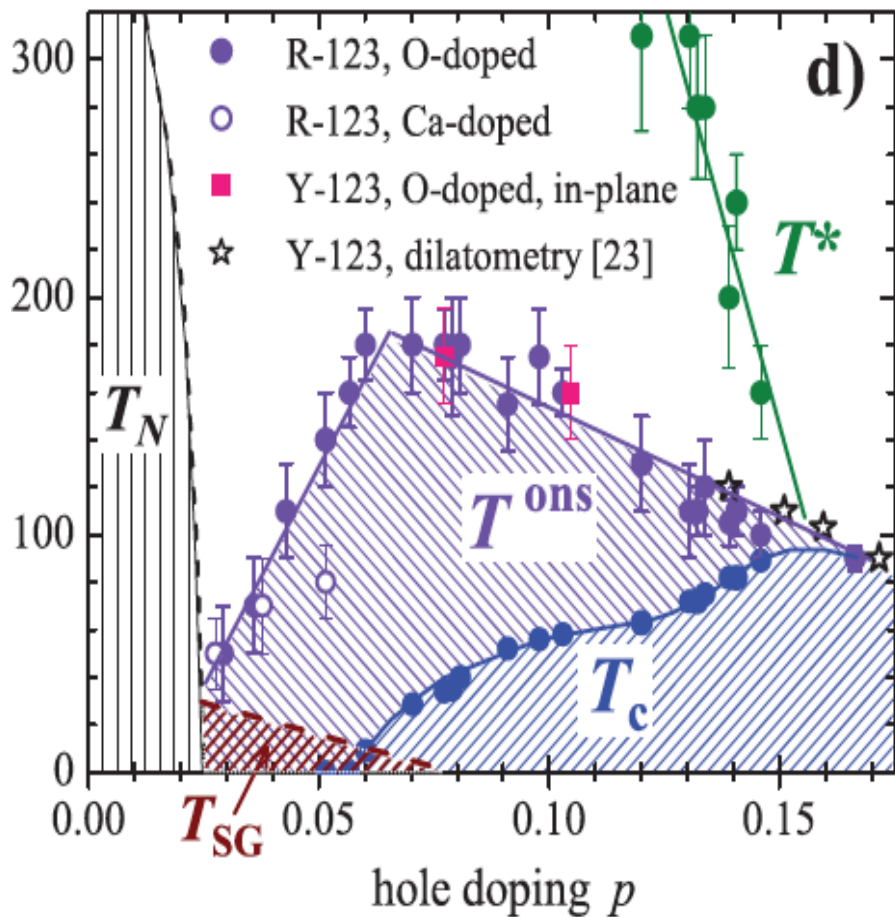


A. Mackenzie et al., *PRB* 53, 5848 (1996)

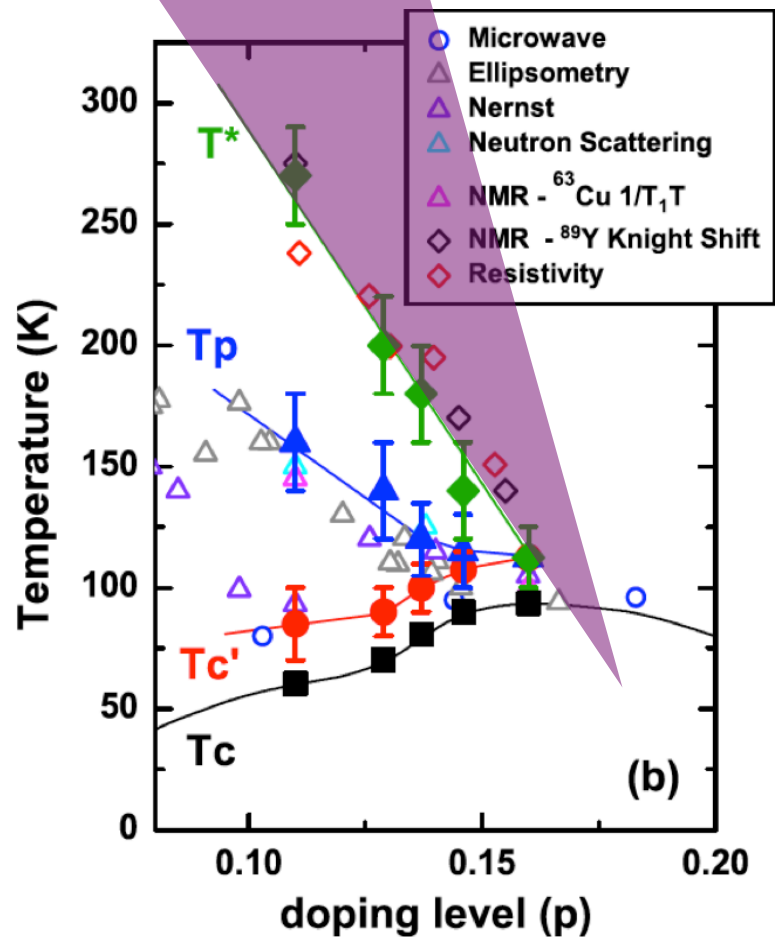


## Persistence of the Superconducting Condensate Far above the Critical Temperature

Ece Uykur, Kiyohisa Tanaka,  
 Takahiko Masui, Shigeki  
 Miyasaka, and Setsuko Tajima,  
 PRL 112, 127003 (2014)



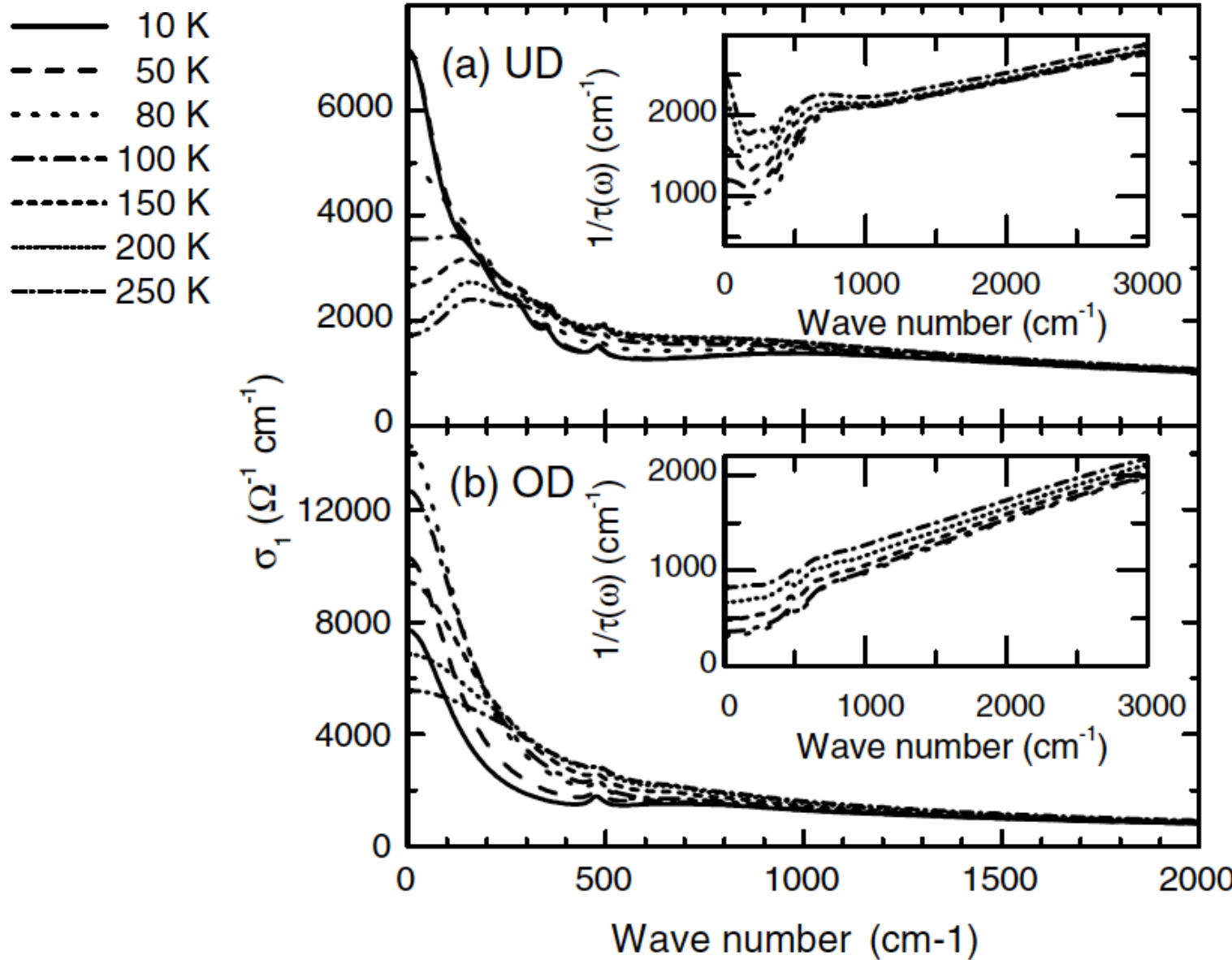
A Dubroka, M Rossle, KW Kim, VK Malik, D Munzar, DN Basov, AA Schafgans, SJ Moon, CT Lin, D Haug, V Hinkov, B Keimer, Th Wolf, JG Storey, JL Tallon, C Bernhard  
*PRL* **106**, 047006 (2011)



Ece Uykur, Kiyohisa Tanaka, Takahiko Masui, Shigeki Miyasaka, Setsuko Tajima,  
*PRL* **112**, 127003 (2014)

# Absence of a Loss of In-Plane Infrared Spectral Weight in the Pseudogap Regime of Bi-2212

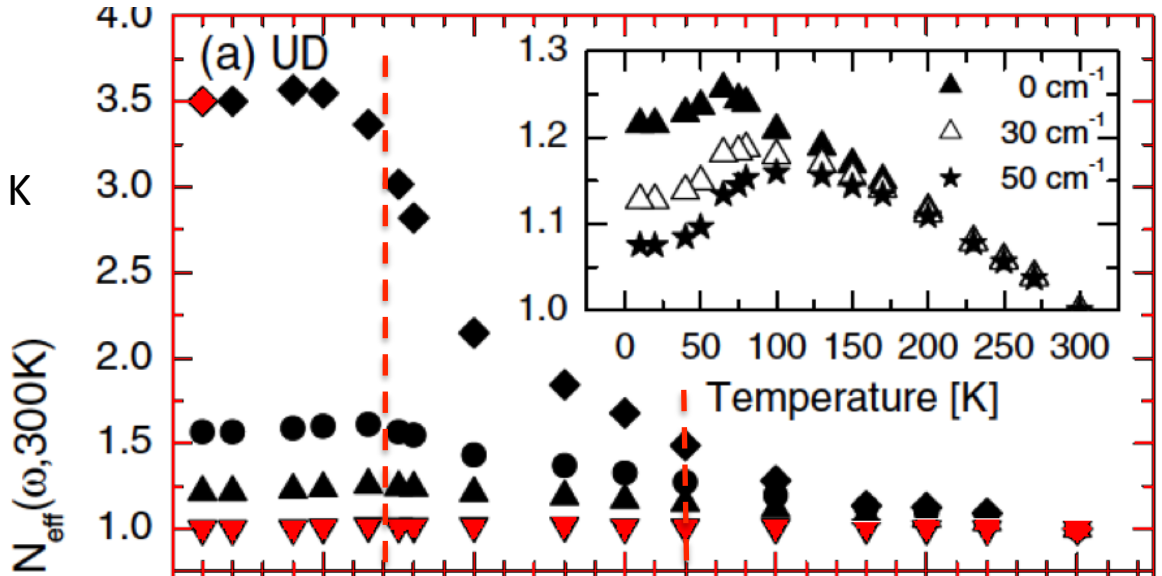
AF Santander-Syro, RPSM Lobo, N Bontemps, Z Konstantinovic, Z Li, and H Raffy, PRL 88 (2002)



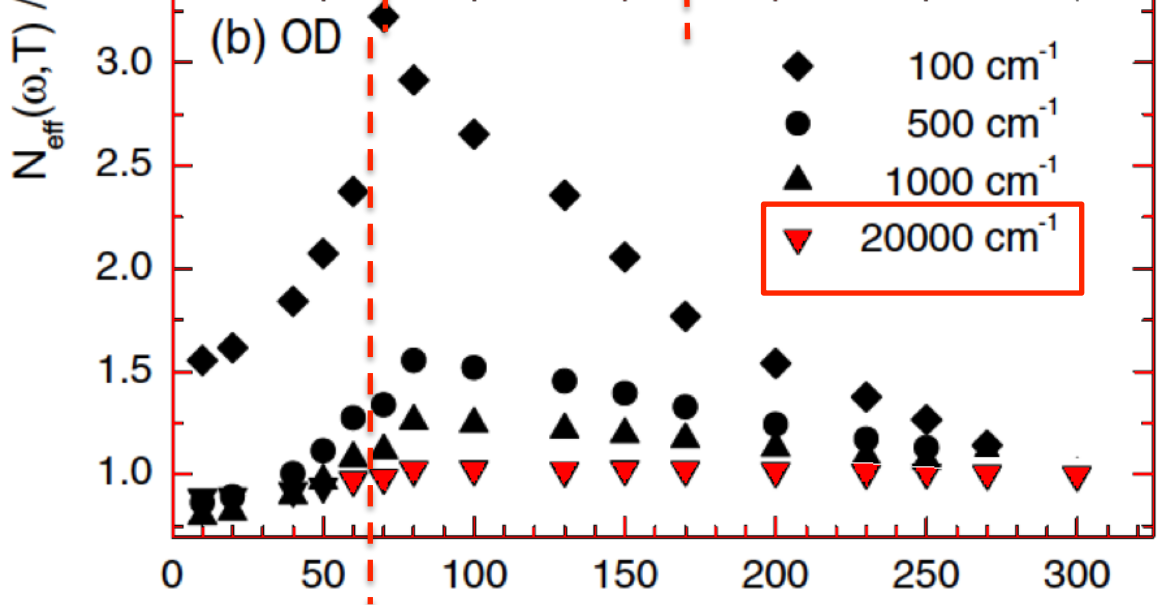
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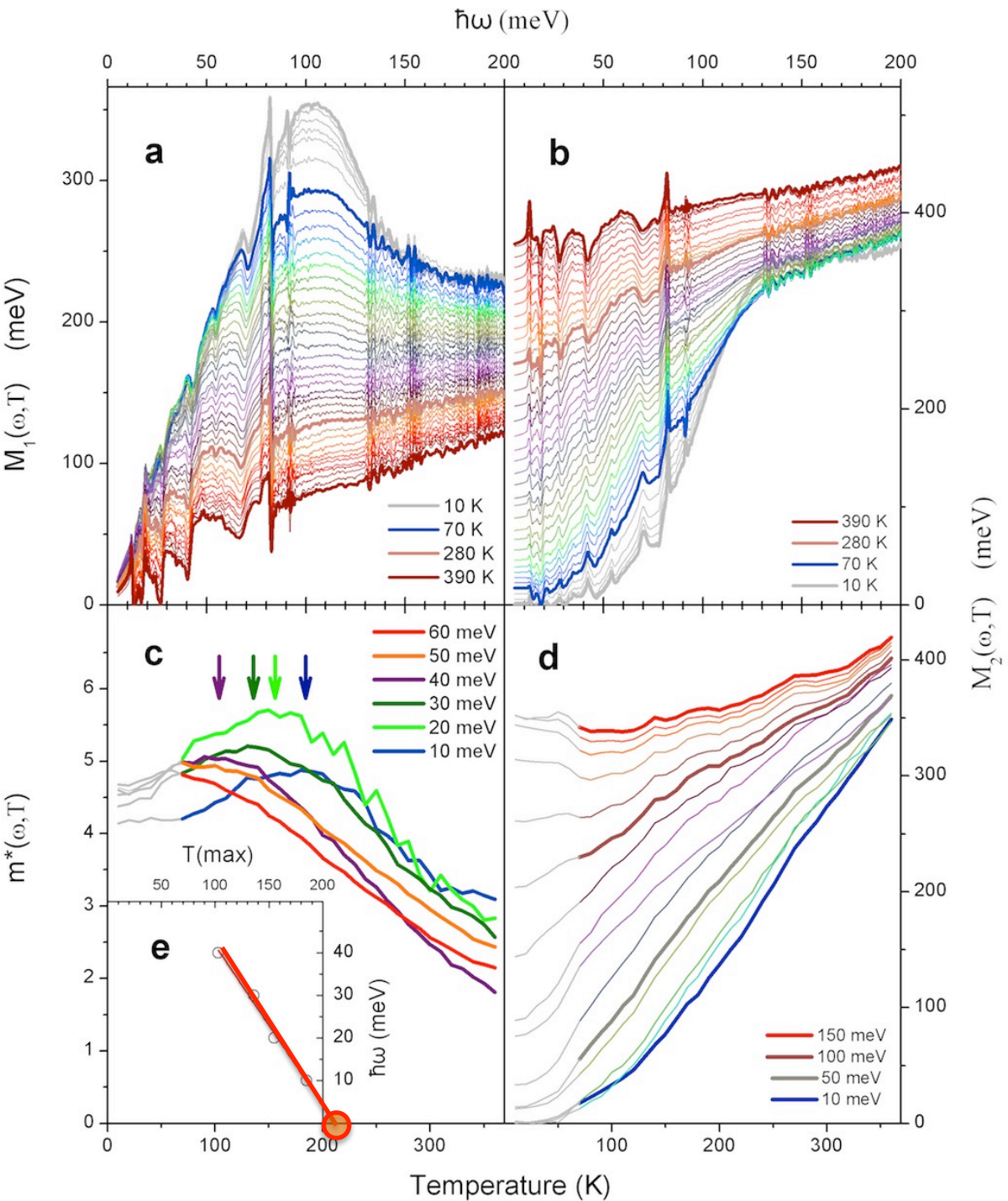
AF Santander-Syro, RPSM Lobo, N Bontemps, Z Konstantinovic, Z Li, and H Raffy, PRL 88 (2002)

$T_c = 70 \text{ K}$ ,  $T^* = 170 \text{ K}$



$T_c = 63 \text{ K}$





$$\frac{G(\omega, T)}{G_0} = \frac{i\pi K}{\omega + M(\omega, T)}$$

$$\tau^{-1}(\omega, T) = M_2(\omega, T)$$

$$m^*(\omega, T) = 1 + M_1(\omega, T) / \omega$$

Maximum of  $m^*(\omega, T)$   
for  $\omega \rightarrow 0$  gives:

$T(\text{crossover}) = 220 \text{ K}$

(  $T^* = 350 \text{ K}$  )



# Conclusions

- Pseudogap shows up clearly in frequency dependence along a and c-axis
- An important part of the pseudogap spectral features have to do with the incoherent part of the optical response function and.... consistency with  $T_c$  !
- Temperature dependence:  $T(\text{cross-over})$  shows up in the relaxation rate and in the c-axis conductivity.
- No discernable effects at higher frequencies.
- Yet, the effects of the superconducting  $T_c$  are quite visible for energies as high as 1 eV.

# sumrules

$$\int_0^{\infty} \omega^{-1} L(q, \omega) d\omega = \frac{\pi}{2}$$

$$\int_0^{\infty} L(q, \omega) d\omega = 2\pi E_C(q)$$

$$\int_0^{\infty} \omega L(q, \omega) d\omega = \frac{2\pi^2 n e^2}{m}$$

$$\int_0^{\infty} \omega^3 L(q, \omega) d\omega = \frac{2\pi^2}{m^2} \sum_G G^2 U_G \langle \rho_{-G} \rangle$$

$$\int_0^{\infty} \text{Im } \omega \epsilon(\omega) d\omega = -\pi^2 e^2 a^2 \sum_{\langle i,j \rangle} t \langle c_i^t c_j \rangle$$

