

# Transmission Eigenvalues in Inverse Scattering Theory

Fioralba Cakoni

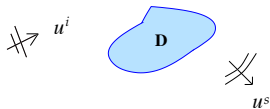
Department of Mathematical Sciences  
University of Delaware  
Newark, DE 19716, USA  
email: [cakoni@math.udel.edu](mailto:cakoni@math.udel.edu)

Jointly with D. Colton, D. Gintides, H. Haddar and A. Kirsch

Research supported by a grant from AFOSR and NSF



# Scattering by an Inhomogeneous Media



$$\Delta u + k^2 n(\mathbf{x})u = 0 \quad \text{in } \mathbb{R}^d, \quad d = 2, 3$$

$$u = u^s + u^i$$

$$\lim_{r \rightarrow \infty} r^{\frac{d-1}{2}} \left( \frac{\partial u^s}{\partial r} - iku^s \right) = 0$$

We assume that  $n - 1$  has compact support  $\bar{D}$  and  $n \in L^\infty(D)$  is such that  $\Re(n) \geq \gamma > 0$  and  $\Im(n) \geq 0$  in  $\bar{D}$ . Here  $k > 0$  is the wave number proportional to the frequency  $\omega$ .

**Question:** Is there an incident wave  $u^i$  that does not scatter?

The answer to this question leads to the [transmission eigenvalue problem](#).

# Transmission Eigenvalues

If there exists a nontrivial solution to the **homogeneous interior transmission problem**

$$\begin{aligned} \Delta w + k^2 n(x)w &= 0 && \text{in } D \\ \Delta v + k^2 v &= 0 && \text{in } D \\ w &= v && \text{on } \partial D \\ \frac{\partial w}{\partial \nu} &= \frac{\partial v}{\partial \nu} && \text{on } \partial D \end{aligned}$$

such that  $v$  can be extended outside  $D$  as a solution to the Helmholtz equation  $\tilde{v}$ , then the scattered field due to  $\tilde{v}$  as incident wave is identically zero.

Values of  $k$  for which this problem has non trivial solution are referred to as **transmission eigenvalues** and the corresponding nontrivial solution  $w, v$  as **eigen-pairs**.

# Transmission Eigenvalues

In general such an extension of  $v$  does not exist!

Since **Herglotz wave functions**

$$v_g(x) := \int_{\Omega} e^{ikx \cdot d} g(d) ds(d), \quad \Omega := \{x : |x| = 1\},$$

are **dense** in the space

$$\{v \in L^2(D) : \Delta v + k^2 v = 0 \text{ in } D\}$$

at a transmission eigenvalue there is an **incident field that produces arbitrarily small scattered field**.

# Motivation

Two important issues:

- Real transmission eigenvalues can be **determined** from the scattered data.
- Transmission eigenvalues carry **information** about material properties.

Therefore, transmission eigenvalues can be used

- to **quantify the presence** of **abnormalities inside homogeneous media** and use this information to test the integrity of materials.

How are **real transmission eigenvalues** seen in the scattering data?

# Measurements

We assume that  $u^i(x) = e^{ikx \cdot d}$  and the far field pattern  $u_\infty(\hat{x}, d, k)$  of the scattered field  $u^s(x, d, k)$  is available for  $\hat{x}, d \in \Omega$ , and  $k \in [k_0, k_1]$

$$\text{where} \quad u^s(x, d, k) = \frac{e^{ikr}}{r^{\frac{d-1}{2}}} u_\infty(\hat{x}, d, k) + O\left(\frac{1}{r^{3/2}}\right)$$

as  $r \rightarrow \infty$ ,  $\hat{x} = x/|x|$ ,  $r = |x|$ .

Define the **far field operator**  $F : L^2(\Omega) \rightarrow L^2(\Omega)$  by

$$(Fg)(\hat{x}) := \int_{\Omega} u_\infty(\hat{x}, d, k) g(d) ds(d), \quad \left( S = I + \frac{ik}{\sqrt{2\pi k}} e^{-i\pi/4} F \right)$$

# The Far Field Operator

## Theorem

*The far field operator  $F : L^2(\Omega) \rightarrow L^2(\Omega)$  is injective and has dense range if and only if  $k$  is not a transmission eigenvalue such that for a corresponding eigensolution  $(w, v)$ ,  $v$  takes the form of a [Herglotz wave function](#).*

For  $z \in D$  the **far field equation** is

$$(Fg)(\hat{x}) = \Phi_{\infty}(\hat{x}, z, k), \quad g \in L^2(\Omega)$$

where  $\Phi_{\infty}(\hat{x}, z, k)$  is the far field pattern of the fundamental solution  $\Phi(x, z, k)$  of the Helmholtz equation  $\Delta v + k^2 v = 0$ .

# Computation of Real TE

Theorem (Cakoni-Colton-Haddar, *Comp. Rend. Math.* 2010)

Assume that either  $n > 1$  or  $n < 1$  and  $z \in D$ .

- If  $k^2$  is not a transmission eigenvalue then for every  $\epsilon > 0$  there exists  $g_{z,\epsilon,k} \in L^2(\Omega)$  satisfying  $\|Fg_{z,\epsilon,k} - \Phi_\infty\|_{L^2(\Omega)} < \epsilon$  and

$$\lim_{\epsilon \rightarrow 0} \|v_{g_{z,\epsilon,k}}\|_{L^2(D)} \quad \text{exists.}$$

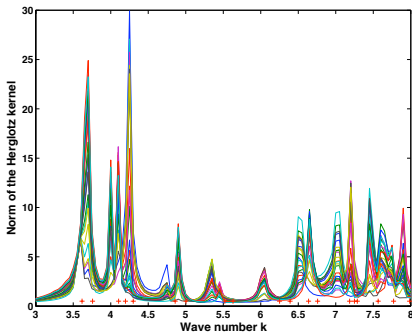
- If  $k^2$  is a transmission eigenvalue for any  $g_{z,\epsilon,k} \in L^2(\Omega)$  satisfying  $\|Fg_{z,\epsilon,k} - \Phi_\infty\|_{L^2(\Omega)} < \epsilon$  and for almost every  $z \in D$

$$\lim_{\epsilon \rightarrow 0} \|v_{g_{z,\epsilon,k}}\|_{L^2(D)} = \infty.$$

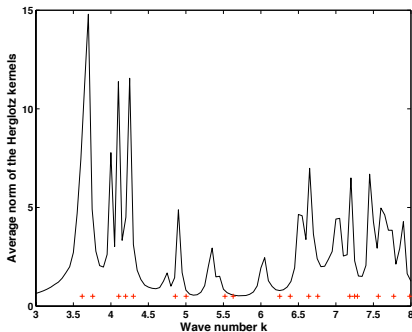
If  $g$  is the computed **Tikhonov regularized solution**, the second part still holds, whereas the first part is proven only for the scalar case *Arens, Inverse Problems (2004)*.



# Computation of Real TE



A composite plot of  $\|g_{z_j}\|_{L^2(\Omega)}$   
against  $k$  for 25 random points  $z_j \in D$



The average of  $\|g_{z_j}\|_{L^2(\Omega)}$   
over all choices of  $z_j \in D$ .

Computation of the transmission eigenvalues from the far field equation for the unit square  $D$ .

# Transmission Eigenvalue Problem

Recall the transmission eigenvalue problem

$$\begin{aligned}
 \Delta w + k^2 n(x) w &= 0 && \text{in } D \\
 \Delta v + k^2 v &= 0 && \text{in } D \\
 w &= v && \text{on } \partial D \\
 \frac{\partial w}{\partial \nu} &= \frac{\partial v}{\partial \nu} && \text{on } \partial D
 \end{aligned}$$

It is a nonstandard eigenvalue problem

$$\int_D (\nabla w \cdot \nabla \bar{\psi} - k^2 n(x) w \bar{\psi}) \, dx = \int_D (\nabla v \cdot \nabla \bar{\phi} - k^2 v \bar{\phi}) \, dx$$

- If  $n = 1$  the interior transmission problem is degenerate
- If  $\Im(n) > 0$  in  $\bar{D}$ , there are no **real** transmission eigenvalues.

# Historical Overview

- The transmission eigenvalue problem in scattering theory was introduced by *Kirsch (1986)* and *Colton-Monk (1988)*
- Research was focused on the discreteness of transmission eigenvalues for variety of scattering problems:  
*Colton-Kirsch-Päivärinta (1989)*, *Rynne-Sleeman (1991)*,  
*Cakoni-Haddar (2007)*, *Colton-Päivärinta-Sylvester (2007)*,  
*Kirsch (2009)*, *Cakoni-Haddar (2009)*, *Hickmann (to appear)*.

In the above work, it is always assumed that either  $n - 1 > 0$  or  $1 - n > 0$ .

## Historical Overview, cont.

- The first proof of existence of at least one transmission eigenvalues for large enough contrast is due to *Päivärinta-Sylvester (2009)*.
- The existence of an infinite set of transmission eigenvalues is proven by *Cakoni-Gintides-Haddar (2010)* under only assumption that either  $n - 1 > 0$  or  $1 - n > 0$ . The existence has been extended to other scattering problems by *Kirsch (2009)*, *Cakoni-Haddar (2010)*, *Cakoni-Kirsch (2010)*, *Bellis-Cakoni-Guzina (2011)*, *Cossonniere (Ph.D. thesis)* etc.
- *Hitrik-Krupchyk-Ola-Päivärinta (2010)*, in a series of papers have extended the transmission eigenvalue problem to a more general class of differential operators with constant coefficients.

## Historical Overview, cont.

- *Finch* has connected the discreteness of the transmission spectrum to a uniqueness question in thermo-acoustic imaging for which  $n - 1$  can change sign.
- *Cakoni-Colton-Haddar (2010)* and then *Cossonniere-Haddar (2011)* have investigated the case when  $n = 1$  in  $D_0 \subset D$  and  $n - 1 > \alpha > 0$  in  $D \setminus \overline{D_0}$ .
- Recently *Sylvester (to appear)* has shown that the set of transmission eigenvalues is at most discrete if  $n - 1$  is positive (or negative) only in a neighborhood of  $\partial D$  but otherwise could change sign inside  $D$ . A similar result is obtained by *Bonnet Ben Dhia - Chesnel - Haddar (2011)* using T-coercivity and *Lakshitanov-Vainberg (to appear)*, for the case when there is contrast in both the main differential operator and lower term.

# Scattering by a Spherically Stratified Medium

We consider the **interior eigenvalue problem** for a ball of radius  $a$  with index of refraction  $n(r)$  being a function of  $r := |x|$

$$\begin{aligned} \Delta w + k^2 n(r) w &= 0 && \text{in } B \\ \Delta v + k^2 v &= 0 && \text{in } B \\ w &= v && \text{on } \partial B \\ \frac{\partial w}{\partial r} &= \frac{\partial v}{\partial r} && \text{on } \partial B \end{aligned}$$

where  $B := \{x \in \mathbb{R}^3 : |x| < a\}$ .

# Scattering by a Spherically Stratified Medium

Look for solutions in polar coordinates  $(r, \theta, \varphi)$

$$v(r, \theta) = a_\ell j_\ell(kr) P_\ell(\cos \theta) \quad \text{and} \quad w(r, \theta) = a_\ell Y_\ell(kr) P_\ell(\cos \theta)$$

where  $j_\ell$  is a spherical Bessel function and  $Y_\ell$  is the solution of

$$Y_\ell'' + \frac{2}{r} Y_\ell' + \left( k^2 n(r) - \frac{\ell(\ell+1)}{r^2} \right) Y_\ell = 0$$

such that  $\lim_{r \rightarrow 0} (Y_\ell(r) - j_\ell(kr)) = 0$ . There exists a **nontrivial solution of the interior transmission problem** provided that

$$d_\ell(k) := \det \begin{pmatrix} Y_\ell(a) & -j_\ell(ka) \\ Y_\ell'(a) & -kj_\ell'(ka) \end{pmatrix} = 0.$$

Values of  $k$  such that  $d_\ell(k) = 0$  are the **transmission eigenvalues**.  $d_\ell(k)$  are entire function of  $k$  of finite type and bounded for  $k > 0$ .

# Transmission Eigenvalues

Assume that  $\Im(n) = 0$  and  $n \in C^2[0, a]$ .

- If either  $n(a) \neq 1$  or  $n(a) = 1$  and  $\int_0^a \sqrt{n(\rho)} d\rho \neq a$ .
  - The set of all transmission eigenvalue is discrete.
  - There exists an infinite number of real transmission eigenvalues accumulating only at  $+\infty$ .
- For a subclass of  $n(r)$  there exist infinitely many complex transmission eigenvalues, *Leung-Colton, (to appear)*.

## Inverse spectral problem

- All transmission eigenvalues uniquely determine  $n(r)$  provide  $n(0)$  is given and either  $n(r) > 1$  or  $n(r) < 1$ .  
*Cakoni-Colton-Gintides, SIAM Journal Math Analysis, (2010).*
- If  $n(r) < 1$  then transmission eigenvalues corresponding to spherically symmetric eigenfunctions uniquely determine  $n(r)$ .  
*Aktosun-Gintides-Papanicolaou, Inverse Problems, (2011).*



# Transmission Eigenvalue Problem

Recall the transmission eigenvalue problem

$$\begin{aligned}\Delta w + k^2 n(x)w &= 0 && \text{in } D \\ \Delta v + k^2 v &= 0 && \text{in } D \\ w &= v && \text{on } \partial D \\ \frac{\partial w}{\partial \nu} &= \frac{\partial v}{\partial \nu} && \text{on } \partial D\end{aligned}$$

Let  $u = w - v$ , we have that

$$\Delta u + k^2 n u = k^2 (n - 1) v.$$

Then eliminate  $v$  to get an equation only in terms of  $u$  by applying  $(\Delta + k^2)$

# Transmission Eigenvalues

Let  $n \in L^\infty(D)$ , and denote  $n^* = \sup_{x \in D} n(x)$  and  $0 < n_* = \inf_{x \in D} n(x)$ .

To fix our ideas assume  $n_* > 1$  (similar analysis if  $n^* < 1$ ).

Let  $u := w - v \in H_0^2(D)$ . The transmission eigenvalue problem can be written for  $u$  as an eigenvalue problem for the fourth order equation:

$$(\Delta + k^2) \frac{1}{n-1} (\Delta + k^2 n) u = 0$$

i.e. in the variational form

$$\int_D \frac{1}{n-1} (\Delta u + k^2 n u) (\Delta \varphi + k^2 \varphi) dx = 0 \quad \text{for all } \varphi \in H_0^2(D)$$

**Definition:**  $k \in \mathbb{C}$  is a **transmission eigenvalue** if there exists a nontrivial solution  $v \in L^2(D)$ ,  $w \in L^2(D)$ ,  $w - v \in H_0^2(D)$  of the homogeneous interior transmission problem.

# Transmission Eigenvalues

Obviously we have

$$0 = \int_D \frac{1}{n-1} |(\Delta u + k^2 n u)|^2 dx + k^2 \int_D (|\nabla u|^2 - k^2 n |u|^2) dx.$$

Poincare inequality yields the Faber-Krahn type inequality for the first transmission eigenvalue (not isoperimetric)

$$k_{1,D,n}^2 > \frac{\lambda_1(D)}{n^*}.$$

where  $\lambda_1(D)$  is the first Dirichlet eigenvalue of  $-\Delta$  in  $D$ .

In particular there are no real transmission eigenvalues in the interval  $(0, \lambda_1(D)/n^*)$ .

# Transmission Eigenvalues

Letting  $k^2 := \tau$ , the transmission eigenvalue problem can be written as a **quadratic pencil operator**

$$u - \tau K_1 u + \tau^2 K_2 u = 0, \quad u \in H_0^2(D)$$

with **selfadjoint compact operators**  $K_1 = T^{-1/2} T_1 T^{-1/2}$  and  $K_2 = T^{-1/2} T_2 T^{-1/2}$  where

$$(Tu, \varphi)_{H^2(D)} = \int_D \frac{1}{n-1} \Delta u \Delta \varphi \, dx \quad \text{coercive}$$

$$(T_1 u, \varphi)_{H^2(D)} = - \int_D \frac{1}{n-1} (\Delta u \varphi + n u \Delta \varphi) \, dx$$

$$(T_2 u, \varphi)_{H^2(D)} = \int_D \frac{n}{n-1} u \varphi \, dx \quad \text{non-negative.}$$

# Transmission Eigenvalues

The transmission eigenvalue problem can be transformed to the eigenvalue problem

$$(\mathbb{K} - \xi \mathbb{I})U = 0, \quad U = \begin{pmatrix} u \\ \tau K_2^{1/2} u \end{pmatrix}, \quad \xi := \frac{1}{\tau}$$

for the **non-selfadjoint compact operator**

$\mathbb{K}: H_0^2(D) \times H_0^2(D) \rightarrow H_0^2(D) \times H_0^2(D)$  given by

$$\mathbb{K} := \begin{pmatrix} K_1 & -K_2^{1/2} \\ K_2^{1/2} & 0 \end{pmatrix}.$$

However from here one can see that the **transmission eigenvalues form a discrete** set with  $+\infty$  as the only possible accumulation point.

# Transmission Eigenvalues

To obtain existence of transmission eigenvalues and isoperimetric Faber-Krahn type inequalities we rewrite the transmission eigenvalue problem in the form

$$(\mathbb{A}_\tau - \tau \mathbb{B})u = 0 \quad \text{in } H_0^2(D)$$

$$(\mathbb{A}_\tau u, \varphi)_{H^2(D)} = \int_D \frac{1}{n-1} (\Delta u + \tau u)(\Delta \varphi + \tau \varphi) dx + \tau^2 \int_D u \cdot \varphi dx$$

$$(\mathbb{B}u, \varphi)_{H^2(D)} = \int_D \nabla u \cdot \nabla \varphi dx$$

Observe that

- The mapping  $\tau \rightarrow \mathbb{A}_\tau$  is continuous from  $(0, +\infty)$  to the set of **self-adjoint coercive operators** from  $H_0^2(D) \rightarrow H_0^2(D)$ .
- $\mathbb{B} : H_0^2(D) \rightarrow H_0^2(D)$  is self-adjoint, compact and non-negative.

# Transmission Eigenvalues

Now we consider the **generalized eigenvalue problem**

$$(\mathbb{A}_\tau - \lambda(\tau)\mathbb{B})u = 0 \quad \text{in } H_0^2(D)$$

**Note** that  $k^2 = \tau$  is a transmission eigenvalue if and only if satisfies  $\lambda(\tau) = \tau$

For a fixed  $\tau > 0$  there exists an increasing sequence of eigenvalues  $\lambda_j(\tau)_{j \geq 1}$  such that  $\lambda_j(\tau) \rightarrow +\infty$  as  $j \rightarrow \infty$ .

These eigenvalues satisfy

$$\lambda_j(\tau) = \min_{W \subset \mathcal{U}_j} \left( \max_{u \in W \setminus \{0\}} \frac{(\mathbb{A}_\tau u, u)}{(\mathbb{B}u, u)} \right).$$

# Transmission Eigenvalues

Hence, if there exists two positive constants  $\tau_0 > 0$  and  $\tau_1 > 0$  such that

- $\mathbb{A}_{\tau_0} - \tau_0 \mathbb{B}$  is positive on  $H_0^2(D)$ ,
- $\mathbb{A}_{\tau_1} - \tau_1 \mathbb{B}$  is non positive on a  $m$  dimensional subspace of  $H_0^2(D)$

then each of the equations  $\lambda_j(\tau) = \tau$  for  $j = 1, \dots, m$ , has at least one solution in  $[\tau_0, \tau_1]$  meaning that there exists  $m$  transmission eigenvalues (counting multiplicity) within the interval  $[\tau_0, \tau_1]$ .

It is now obvious that determining such constants  $\tau_0$  and  $\tau_1$  provides the existence of transmission eigenvalues as well as the desired isoperimetric inequalities.



# Transmission Eigenvalues

- $(\mathbb{A}_\tau \mathbf{u} - \tau \mathbb{B} \mathbf{u}, \mathbf{u})_{\mathcal{U}_0(D)} \geq \alpha \|\mathbf{u}\|_{\mathcal{U}_0(D)}$  for all  $0 < \tau < \frac{\lambda_1(D)}{n^*}$ .
- Take  $\tau_1 := k^2(B_r)$  the first eigenvalue a ball  $B_r \subset D$  and  $n(x) = n_*$  constant,  $u_r$  the corresponding eigenfunction and denote  $\tilde{u}_r \in H_0^2(D)$  its extension by zero to the whole of  $D$ . Then

$$(\mathbb{A}_{\tau_1} \tilde{u}_r - \tau_1 \mathbb{B} \tilde{u}_r, \tilde{u}_r)_{\mathcal{U}_0(D)} \leq 0.$$

If the radius of the ball is such that  $m(r)$  disjoint balls can be included in  $D$ , the above condition is satisfied in a  $m(r)$ -dimensional subspace of  $H_0^2(D)$

Thus there exists  $m(r)$  transmission eigenvalues (counting multiplicity). As  $r \rightarrow 0$ ,  $m(r) \rightarrow \infty$  and since the multiplicity of an eigenvalue is finite we prove the **existence of an infinite set of real transmission eigenvalues**.

# Faber-Krahn Inequalities

Theorem (Cakoni-Gintides-Haddar, SIMA (2010))

Assume that  $1 < n_*$ . Then, there exists an infinite discrete set of *real transmission eigenvalues* accumulating at infinity  $+\infty$ . Furthermore

$$k_{1,n^*,B_1} \leq k_{1,n^*,D} \leq k_{1,n(x),D} \leq k_{1,n_*,D} \leq k_{1,n_*,B_2}.$$

where  $B_2 \subset D \subset B_1$ .

One can prove that, for  $n$  constant, the first transmission eigenvalue  $k_{1,n}$  is continuous and strictly monotonically decreasing with respect to  $n$ . In particular, this shows that the *first transmission eigenvalue determine uniquely the constant index of refraction*, provided that it is known a priori that either  $n > 1$ .

Similar results can be obtained for the case when  $0 < n^* < 1$ .

# Detection of Anomalies in an Isotropic Medium

What does the first transmission eigenvalue say about the inhomogeneous media  $n(x)$ ?

We find the constant  $n_0$  such that the first transmission eigenvalue of

$$\begin{aligned} \Delta w + k^2 n_0 w &= 0 && \text{in } D \\ \Delta v + k^2 v &= 0 && \text{in } D \\ w &= v && \text{on } \partial D \\ \frac{\partial w}{\partial \nu} &= \frac{\partial v}{\partial \nu} && \text{on } \partial D \end{aligned}$$

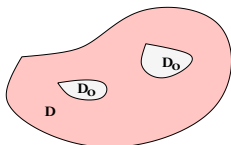
is  $k_{1,n(x)}$  (which can be determined from the measured data).

Then from the previous discussion we have that  $n_* \leq n_0 \leq n^*$ .

**Open Question:** Find an exact formula that connect  $n_0$  to  $n(x)$  and  $D$ .

# The Case with Cavities

Can the assumption  $n > 1$  or  $0 < n < 1$  in  $D$  be relaxed?



$$n = 1 \text{ in } D_0$$

$$n - 1 \geq \delta > 0 \text{ in } D \setminus D_0$$

The case when there are regions  $D_0$  in  $D$  where  $n = 1$  (i.e. cavities) is more delicate. The same type of analysis can be carried through by looking for solutions of the transmission eigenvalue problem

$v \in L^2(D)$  and  $w \in L^2(D)$  such that  $w - v$  is in

$$V_0(D, D_0, k) := \{u \in H_0^2(D) \text{ such that } \Delta u + k^2 u = 0 \text{ in } D_0\}.$$

*Cakoni-Colton-Haddar, SIMA (2010)*

# The Case with Cavities

In particular if  $n > 1$  and  $k(D_0, n(x))$  is the first eigenvalue for a fixed  $D$ , one has the following properties:

- The **Faber Krahn inequality**

$$0 < \frac{\lambda_1(D)}{n^*} \leq k(D_0, n(x)).$$

- Monotonicity with respect to the index of refraction

$$k(D_0, n(x)) \leq k(D_0, \tilde{n}(x)), \quad \tilde{n}(x) \leq n(x).$$

- Monotonicity with respect to voids

$$k(D_0, n(x)) \leq k(\tilde{D}_0, n(x)), \quad D_0 \subset \tilde{D}_0.$$

where  $\lambda_1(D)$  is the first Dirichlet eigenvalue of  $-\Delta$  in  $D$ .

## The Case of $n - 1$ Changing Sign

Recently, progress has been made in the case of the contrast  $n - 1$  changing sign inside  $D$  with state of the art result by *Sylvester (to appear)*. Roughly speaking he shows that transmission eigenvalues form a discrete (possibly empty) set provided  $n - 1$  has fixed sign **only in a neighborhood** of  $\partial D$ . There are two aspects in the proof:

- **Fredholm property.** Sylvester considers the problem in the form

$$\Delta u + k^2 n u = k^2 (n - 1) v, \quad \Delta v + k^2 v = 0, \quad u \in H_0^2(D), \quad v \in H^1(D)$$

and uses the concept of upper-triangular compact operators. This property can also be obtained via variational formulation (*Kirsch*) or integral equation formulation (*Cossonniere-Haddar*).

- **Find a  $k$  that is not a transmission eigenvalues.** This requires careful estimates for the solution inside  $D$  in terms of its values in a neighborhood of  $\partial D$ .

The existence of transmission eigenvalues under such weaker assumptions is still open.

# Complex Eigenvalues

Current results on complex transmission eigenvalues for media of general shape are limited to **identifying eigenvalue free zones in the complex plane**.

- The first result for homogeneous media is given in *Cakoni-Colton-Gintides SIMA (2010)*.
- The **best result to date** is due *Hitrik-Krupchyk-Ola-Päivärinta, Math. Research Letters (2011)*, where they show that almost all transmission eigenvalues are confined to a parabolic neighborhood of the positive real axis. More specifically they show

## Theorem (Hitrik-Krupchyk-Ola-Päivärinta)

For  $n \in C^\infty(\bar{D}, \mathbb{R})$  and  $1 < \alpha \leq n \leq \beta$ , there exists a  $0 < \delta < 1$  and  $C > 1$  both independent of  $\alpha, \beta$  such that all transmission eigenvalues  $\tau := k^2 \in \mathbb{C}$  with  $|\tau| > C$  satisfies  $\Re(\tau) > 0$  and  $\Im(\tau) \leq C|\tau|^{1-\delta}$ .

# Absorbing-Dispersive Media

$$\Delta w + k^2 \left( \epsilon_1 + i \frac{\gamma_1}{k} \right) w = 0 \quad \text{in } D$$

$$\Delta v + k^2 \left( \epsilon_0 + i \frac{\gamma_0}{k} \right) v = 0 \quad \text{in } D$$

where  $\epsilon_0 \geq \alpha_0 > 0$ ,  $\epsilon_1 \geq \alpha_1 > 0$ ,  $\gamma_0 \geq 0$ ,  $\gamma_1 \geq 0$  are bounded functions.

For the corresponding spherically stratified case we have:

## Theorem

If

$$\frac{\gamma_0 a}{\sqrt{\epsilon_0}} = \int_0^a \frac{\gamma_1(r)}{\sqrt{\epsilon_1(r)}} dr \quad \text{and} \quad \sqrt{\epsilon_0} a \neq \int_0^a \sqrt{\epsilon_1(r)} dr$$

there exist an infinite number of real transmission eigenvalues. If the first condition is not met then there exist an infinite number of complex eigenvalues.



# Absorbing-Dispersive Media

In the general case we have proven *Cakoni-Colton-Haddar* (to appear):

- The set of transmission eigenvalues  $k \in \mathbb{C}$  in the right half plane is discrete, provided  $\epsilon_1(x) - \epsilon_0(x) > 0$ .
- Using the stability of a finite set of eigenvalues for closed operators we have shown that if  $\sup_D(\gamma_0 + \gamma_1)$  is small enough there exists at least  $\ell > 0$  transmission eigenvalues each in a small neighborhood of the first  $\ell$  real transmission eigenvalues corresponding to  $\gamma_0 = \gamma_1 = 0$ .
- For the case of  $\epsilon_0, \epsilon_1, \gamma_0, \gamma_1$  constant, we have identified eigenvalue free zones in the complex plane

The existence of transmission eigenvalues for general media if absorption is present is still open.

# Anisotropic Media

The corresponding **transmission eigenvalue problem** is to find  $v, w \in H^1(D)$  such that

$$\begin{aligned} \nabla \cdot A \nabla w + k^2 n w &= 0 && \text{in } D \\ \Delta v + k^2 v &= 0 && \text{in } D \\ w &= v && \text{on } \partial D \\ \nu \cdot A \nabla w &= \nu \cdot \nabla v && \text{on } \partial D. \end{aligned}$$

This transmission eigenvalue problem has a more complicated nonlinear structure than quadratic.

The existence has been shown in *Cakoni-Gintides-Haddar, SIAM J. Math. Anal. (2010)* and *Cakoni-Kirsch, IJCSM (2010)*.

# Existence of Transmission Eigenvalues

Set  $u = w - v \in H_0^1(D)$ . Find  $v = v_u$  by solving a Neuman type problem: For every  $\psi \in H^1(D)$

$$\int_D (A - I) \nabla v \cdot \nabla \bar{\psi} - k^2(n-1)v\bar{\psi} \, dx = \int_D A \nabla u \cdot \nabla \bar{\psi} - k^2 n u \bar{\psi} \, dx.$$

Having  $u \rightarrow v_u$ , we require that  $v := v_u$  satisfies  $\Delta v + k^2 v = 0$ .

Thus we define  $\mathbb{L}_k : H_0^1(D) \rightarrow H_0^1(D)$

$$(\mathbb{L}_k u, \phi)_{H_0^1(D)} = \int_D \nabla v_u \cdot \nabla \bar{\phi} - k^2 v_u \cdot \bar{\phi} \, dx, \quad \phi \in H_0^1(D).$$

Then the **transmission eigenvalue problem is equivalent** to

$$\mathbb{L}_k u = 0 \quad \text{in} \quad H_0^1(D) \quad \text{which can be written}$$

$$(\mathbb{I} + \mathbb{L}_0^{-1/2} \mathbb{C}_k \mathbb{L}_0^{-1/2}) u = 0 \quad \text{in} \quad H_0^1(D)$$

$\mathbb{L}_0$  self-adjoint positive definite and  $\mathbb{C}_k$  self-adjoint compact.

# Existence of Transmission Eigenvalues

- If  $n(x) \equiv 1$  and the contrast  $A - I$  is either positive or negative in  $D$  then there exists an infinite discrete set of real transmission eigenvalues accumulating at  $+\infty$ .
- If the contrasts  $A - I$  and  $n - 1$  have the same fixed sign, then there exists an infinite discrete set of real transmission eigenvalues accumulating at  $+\infty$ .
- If the contrasts  $A - I$  and  $n - 1$  have the opposite fixed sign, then there exists at least one real transmission eigenvalue providing that  $n$  is small enough.

# Discreteness of Transmission Eigenvalues

The strongest result on the discreteness of transmission eigenvalues for this problem is due to *Bonnet Ben Dhia - Chesnel - Haddar*, *Comptes Rendus Math.* (2011) (using the concept of **T-coercivity**).

In particular, the discreteness of transmission eigenvalues is proven under either one of the following assumptions (weaker than for the existence):

- Either  $A - I > 0$  or  $A - I < 0$  in  $D$ , and  $\int_D (n - 1) dx \neq 0$  or  $n \equiv 1$ .
- The contrasts  $A - I$  and  $n - 1$  have the same fixed sign only in a neighborhood of the boundary  $\partial D$ .

# Numerical Example: Homogeneous Anisotropic Media

We consider  $D$  to be the unit square  $[-1/2, 1/2] \times [-1/2, 1/2]$ ,  
 $n \equiv 1$  and

$$A_1 = \begin{pmatrix} 2 & 0 \\ 0 & 8 \end{pmatrix} \quad A_2 = \begin{pmatrix} 6 & 0 \\ 0 & 8 \end{pmatrix} \quad A_{2r} = \begin{pmatrix} 7.4136 & -0.9069 \\ -0.9069 & 6.5834 \end{pmatrix}$$

Matrix	Eigenvalues $a_*$ , $a^*$	Predicted $a_0$
$A_{iso}$	4, 4	4.032
$A_1$	2, 8	5.319
$A_2$	6, 8	7.407
$A_{2r}$	6, 8	6.896

*Cakoni-Colton-Monk-Sun, Inverse Problems, (2010)*

# Open Problem

- Can the existence of real transmission eigenvalues for non-absorbing media be established if the assumptions on the sign of the contrast are weakened?
- Do complex transmission eigenvalues exist for general non-absorbing media?
- Do real transmission eigenvalues exist for absorbing media?
- What would the necessary conditions be on the contrasts that guaranty the discreteness of transmission eigenvalues?
- Can Faber-Krahn type inequalities be established for the higher eigenvalues?
- Can an inverse spectral problem be developed for the general transmission eigenvalue problem? (Completeness of eigen-solutions?)

*Cakoni - Haddar, Transmission Eigenvalues in Inverse Scattering Theory, in Inside Out 2, Uhlmann ed. MSRI Publication (to appear).*