

Cavity optomechanics using microresonators

**Stefan Weis², Samuel Deleglise², Remi Riviere², Dr. Olivier Arcizet²
Georg Anetsberger², Johannes Hofer and Emanuel Gavartin¹**

Tobias J. Kippenberg^{1,2}

**1) Tenure Track Assistant Professor, EPFL
&**

**2) Independent Max Planck Junior Research Group Leader
MPQ, (Division of T. W. Hänsch)**



College de France, Paris, 18th January 2010

**MAX-PLANCK-INSTITUT
FÜR QUANTENOPTIK
GARCHING**



Collaborators

EPFL-CMI K. Lister
J. P. Kotthaus
W. Zwerger
I. Wilson-Rae
A. Marx
J. Raedler
R. Holtzwarth (Menlo System)



European Research Council

MP-IJRG



DFG



nano_{science} Europe NanoSci-ERA

MC-EXT

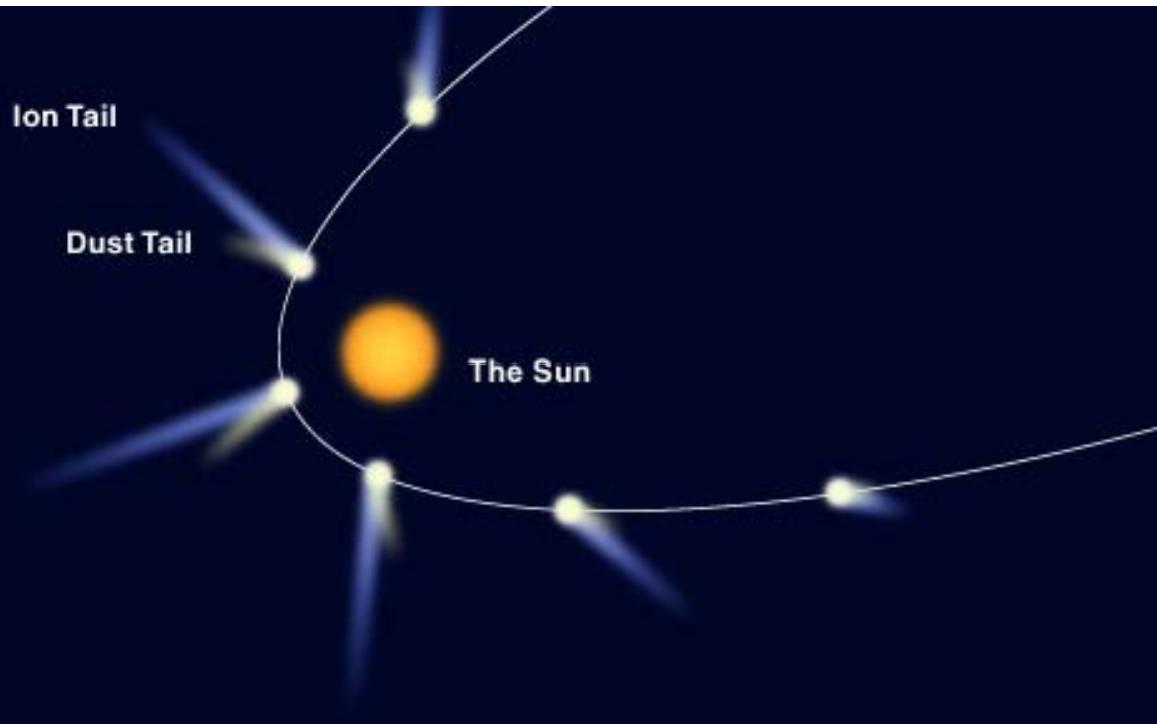
DFG-NIM

GSC

MC-IRG/IEF

NanoSci-ERA

Mechanical Effects of Light



Photons carry momentum which leads to *radiation pressure*

$$p = \hbar \cdot k$$

In the 16th century Kepler proposed that radiation pressure is responsible for comet tails pointing away from the sun

Mechanical Effects of Light on Atoms

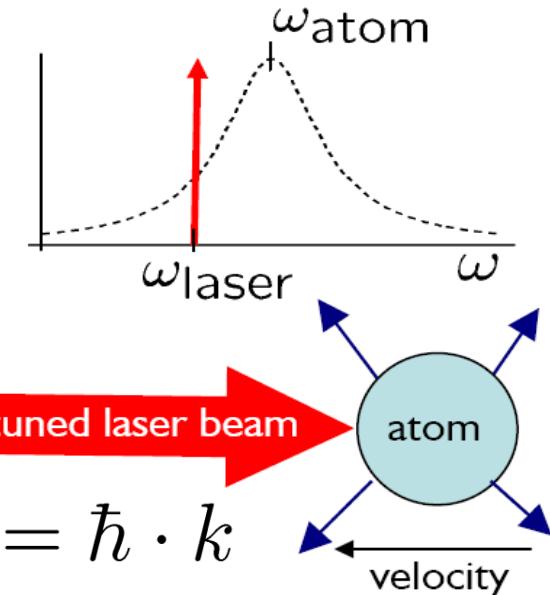
1950: Kastler “Effet luminorefrigorigique et luminocalirofique”

1970: Ashkin Trapping of Particles in Laser Light

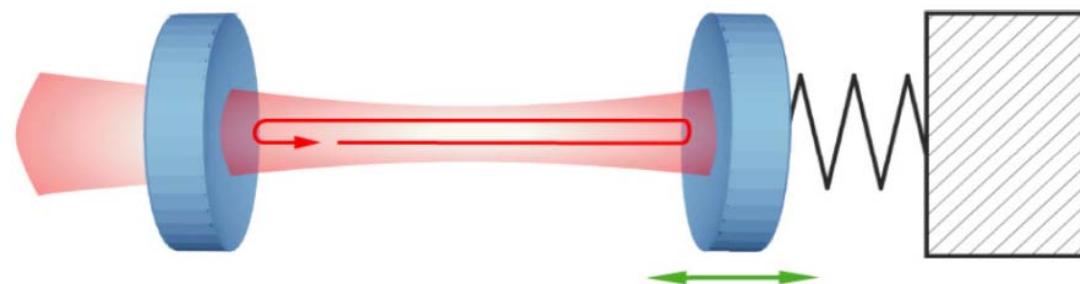
1975 Hänsch and A. Schawlow, Dehmelt and Wineland “Cooling gases by laser radiation”

**1989: “Laser Cooling to the Zero Point Energy of Motion”
(Wineland)**

.....



Is it also possible observe radiation pressure effects on mesoscopic system?

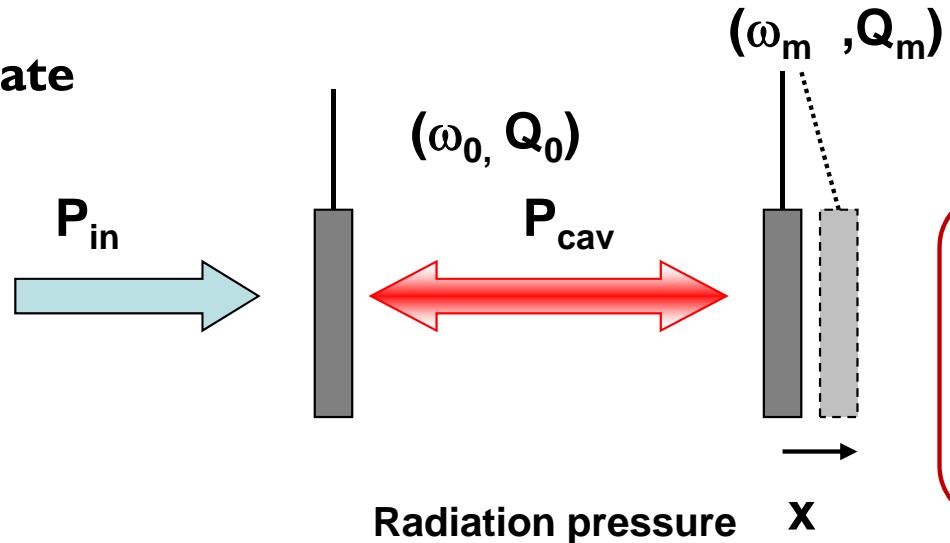


Studied first by Braginsky, Manukin (1967)

Braginsky, Manukin: *Measurement of Weak Forces in Physics Experiments* (1977)

Principle of parametric motion transduction

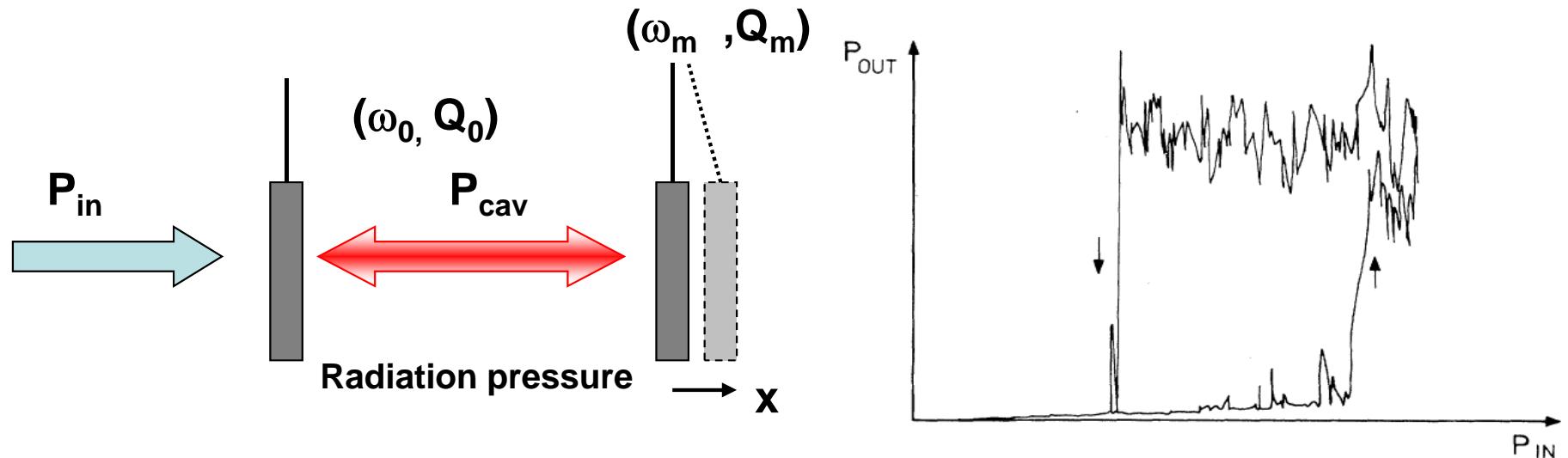
**Radiation pressure
determines the ultimate
sensitivity of an
interferometer**



$$H_{int} = -\frac{\hbar\omega}{L} a^\dagger a x$$

$$F_{RP} = \frac{\hbar\omega}{L} a^\dagger a$$

Observation of Radiation Pressure



MPQ 1983:

There is a discernible effect of radiation pressure in a high-finesse cavity with a suspended mirror.

Dorsel, Meystre, Walther et al., PRL 51, 1550 (1983)

$$H_{int} = \frac{\hbar\omega}{L} aa^\dagger x$$

Radiation Pressure Limits Position Sensitivity: Standard Quantum Limit

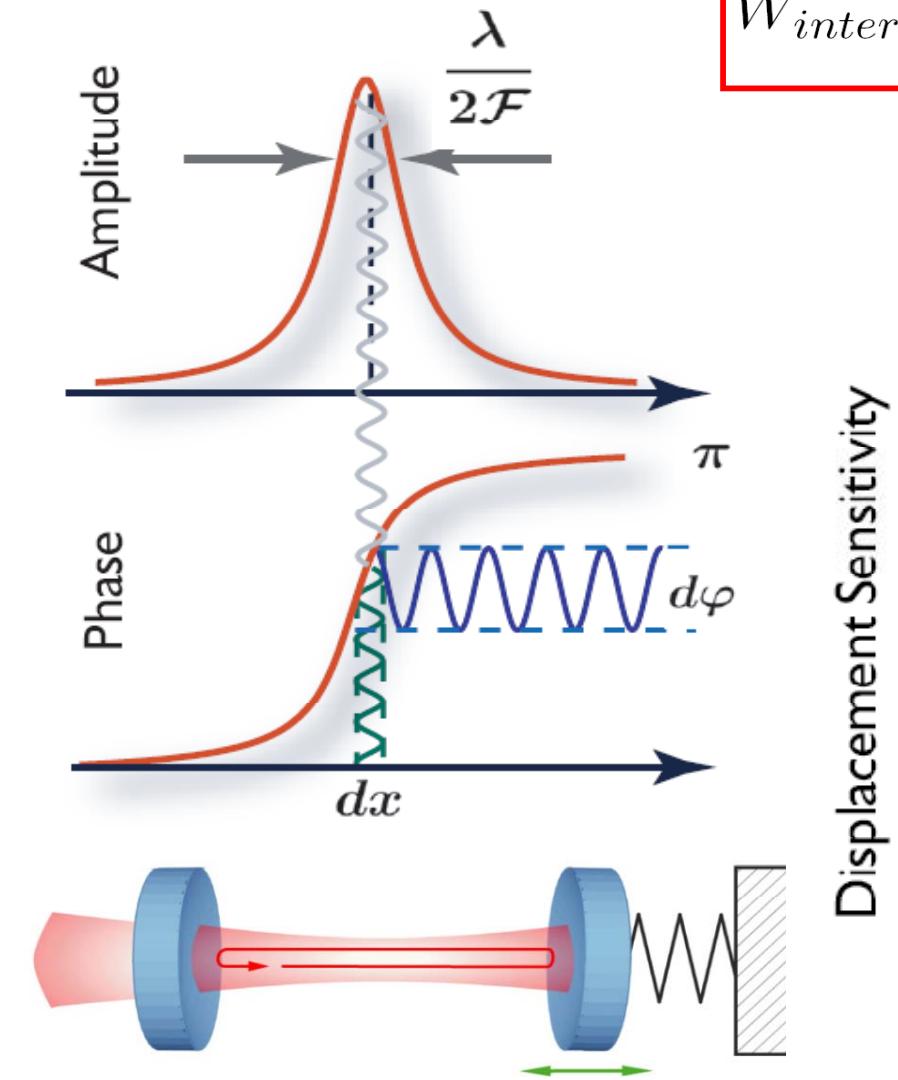
$$S_{SQL} = \sqrt{\frac{\hbar}{2m\Omega_m\gamma}}$$

Braginsky, Manukin: Measurement of Weak Forces in Physics Experiments (1977)

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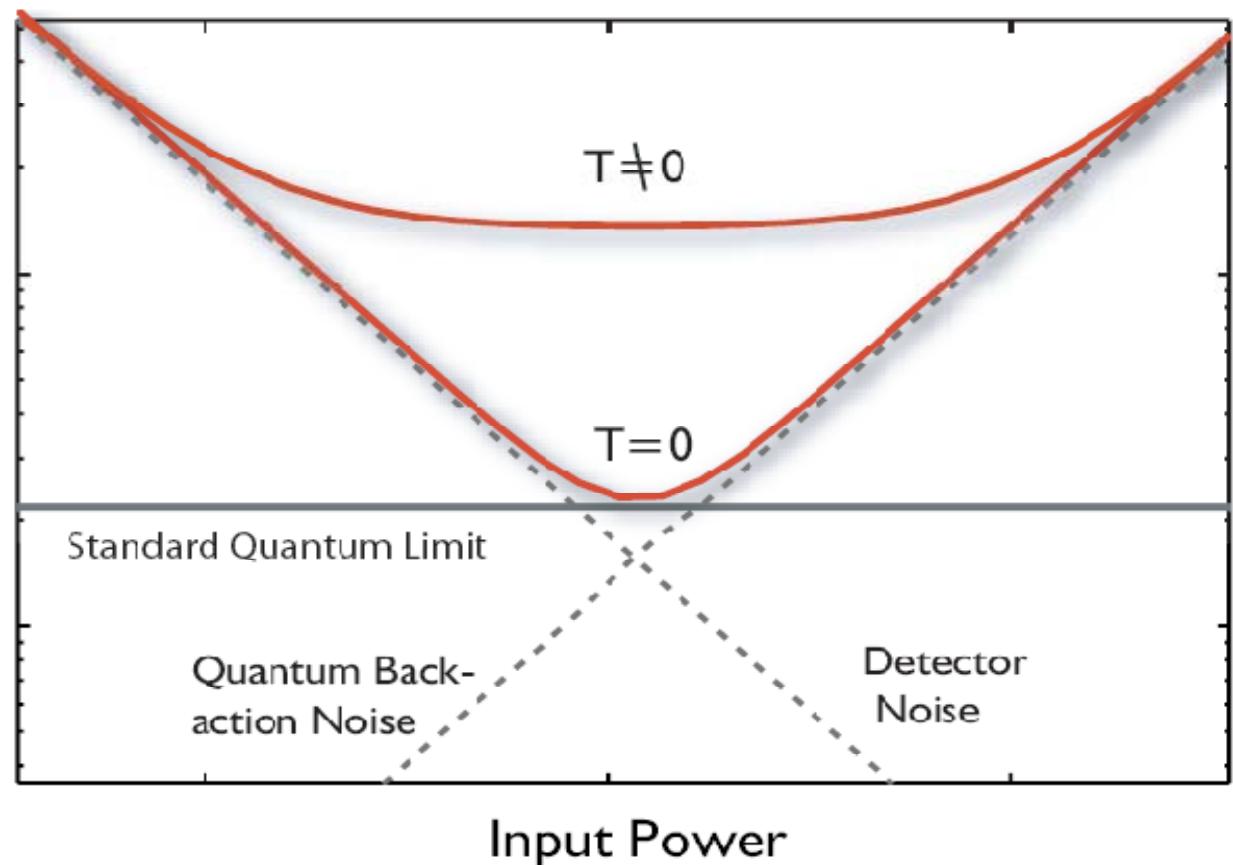
Standard Quantum Limit of Motion Detection



$$W_{\text{interaction}} = - \frac{x}{L} |a_p|^2$$

$$\Rightarrow \ddot{x} = \frac{|a_p|^2}{mL}$$

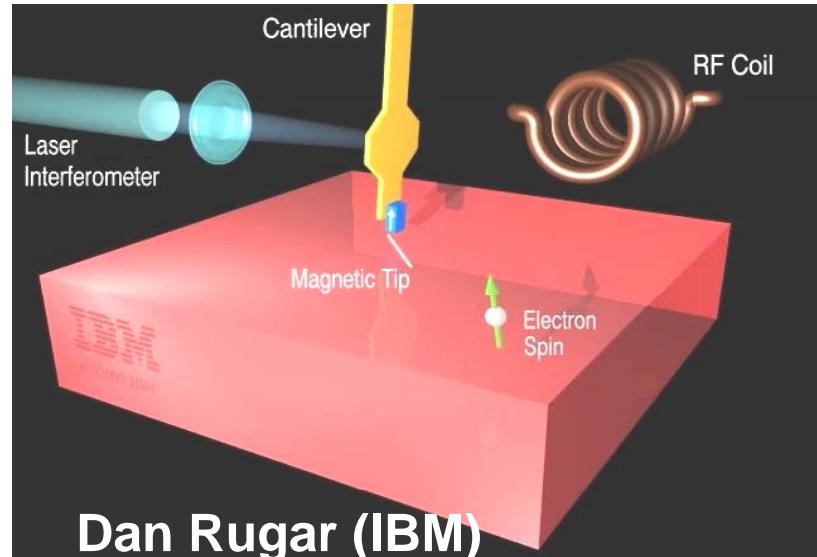
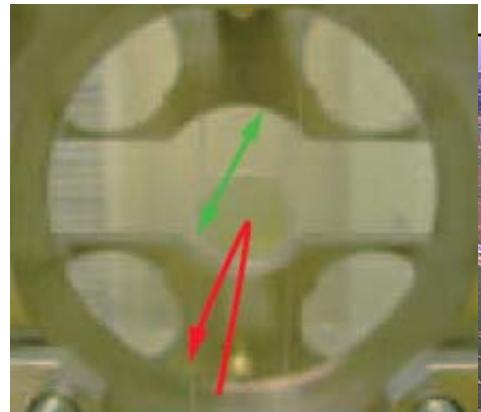
$$\Rightarrow \dot{a}_p = -i\omega_o \frac{x}{L} a_p$$



Caves, Phys. Rev. D (1980)

Braginsky, Manukin: Measurement of Weak Forces in Physics Experiments (1977)

Coupling light with mechanical oscillators



Gravitational wave interferometric
Detection (VIRGO)

Weak force detection (IBM, San Jose)

Radiation Pressure Quantum Fluctuations limit Position Sensitivity:
Standard Quantum Limit

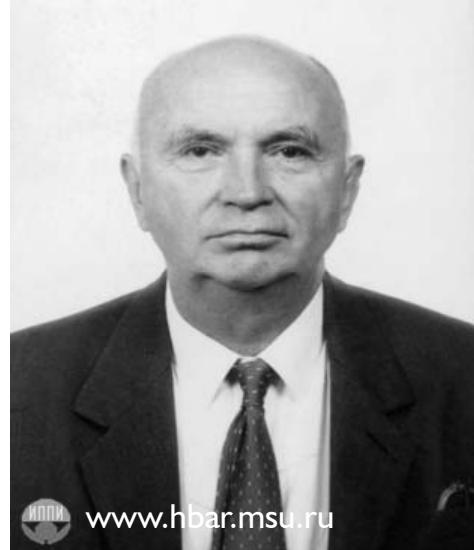
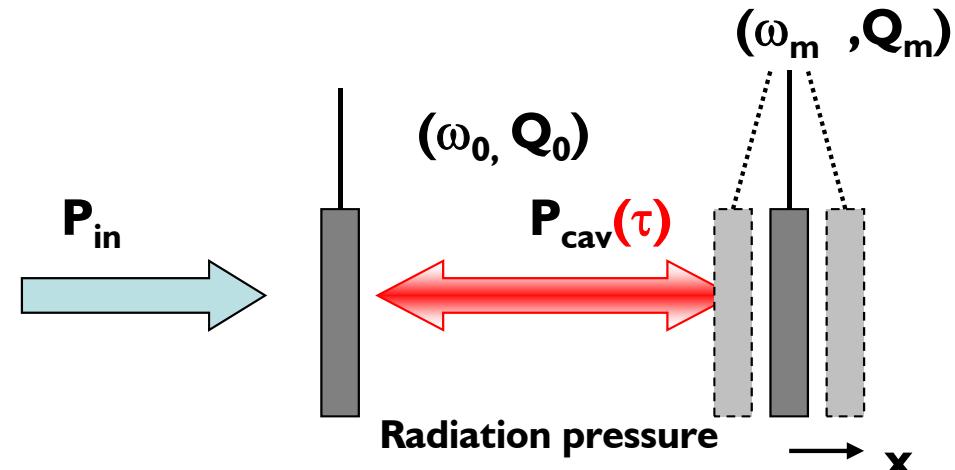
$$S_{SQL} = \sqrt{\frac{\hbar}{m\Omega_m \gamma_m}}$$

→ Quantum effects with macroscopic mechanical oscillators ?

$$n = \frac{k_B T}{\hbar \Omega} \quad \Omega = 1 MHz, T_Q = 50 \mu K$$

Dynamical Backaction

The mutual coupling of optical and mechanical modes was first theoretically studied by V.B. Braginsky, “Measurement of Weak Forces”



$$\frac{d^2x}{dt^2} + \frac{1}{2\tau_m} \frac{dx}{dt} + \omega_m^2 x = \frac{F_{rp}(x(t - \tau))}{m_{eff}}$$

V.B. Braginsky

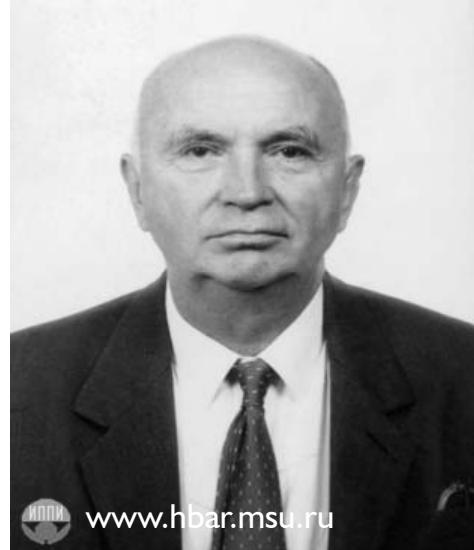
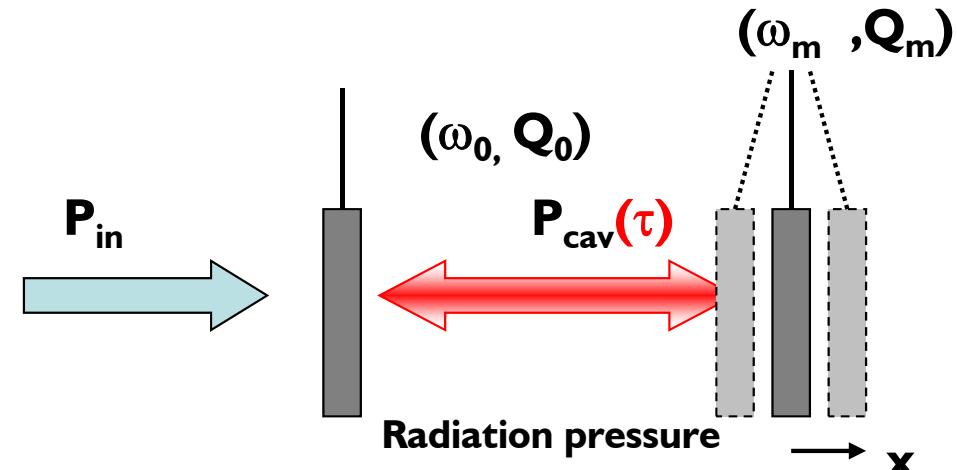
Taylor Expansion yields:

$$\Delta\omega = \frac{\frac{d}{dx} F_{rp}(x)}{m_{eff}} \quad \text{Position dependent term}$$

$$\Delta\beta = -\tau \frac{\frac{d}{dx} F_{rp}(x)}{m_{eff}} \quad \text{Velocity dependent term}$$

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V.B. Braginsky

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SOVIET PHYSICS JETP

VOLUME 31, NUMBER 5

NOVEMBER 1970

INVESTIGATION OF DISSIPATIVE PONDEROMOTIVE EFFECTS OF ELECTROMAGNETIC RADIATION

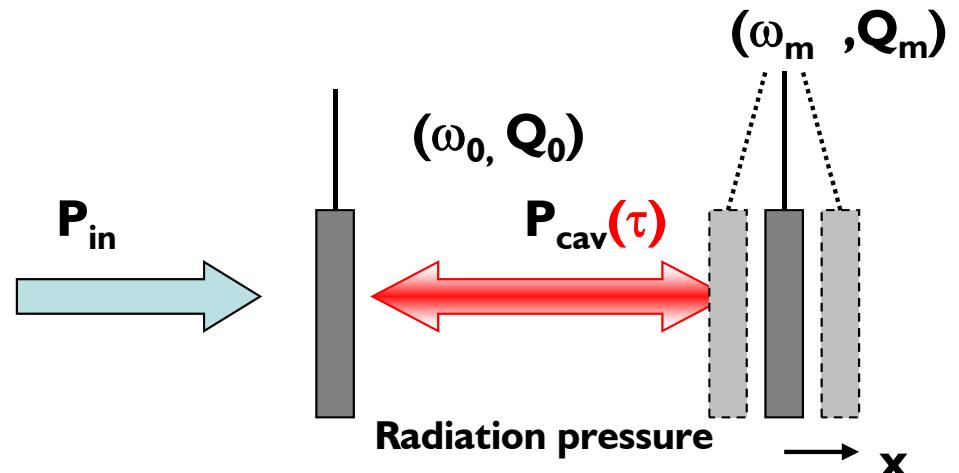
V. B. BRAGINSKII, A. B. MANUKIN, and M. Yu. TIKHONOV

Moscow State University

Submitted October 17, 1969

Dynamical Backaction

The mutual coupling of optical and mechanical modes was first theoretically studied by V.B. Braginsky, "Measurement of Weak Forces"



$$\frac{d^2x}{dt^2} + \frac{1}{2\tau_m} \frac{dx}{dt} + \omega_m^2 x = \frac{F_{rp}(x(t - \tau))}{m_{eff}}$$

Predicted for more than 30 years,
but only recently observed.

incident term

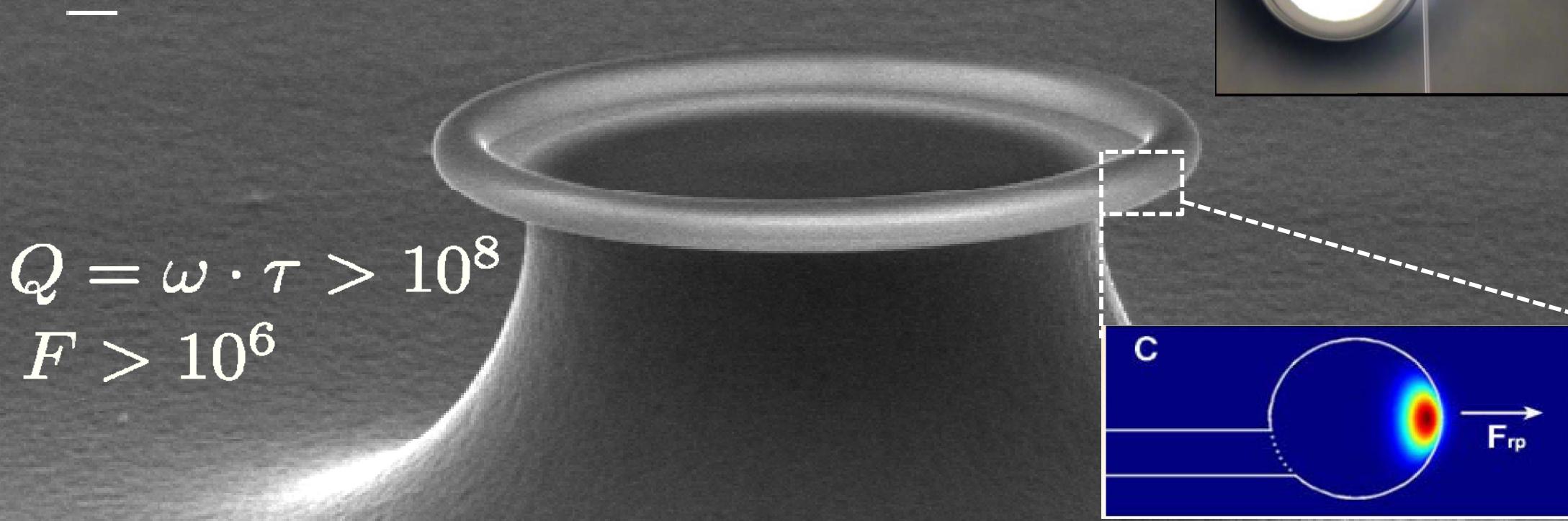
*incident term;
Blue detuning
red detuning*

$\Delta\gamma < 0 \Rightarrow$ Amplification

$\Delta\gamma > 0 \Rightarrow$ Cooling

Experiments (1977)

Structures Fabricated at CMI-EPFL

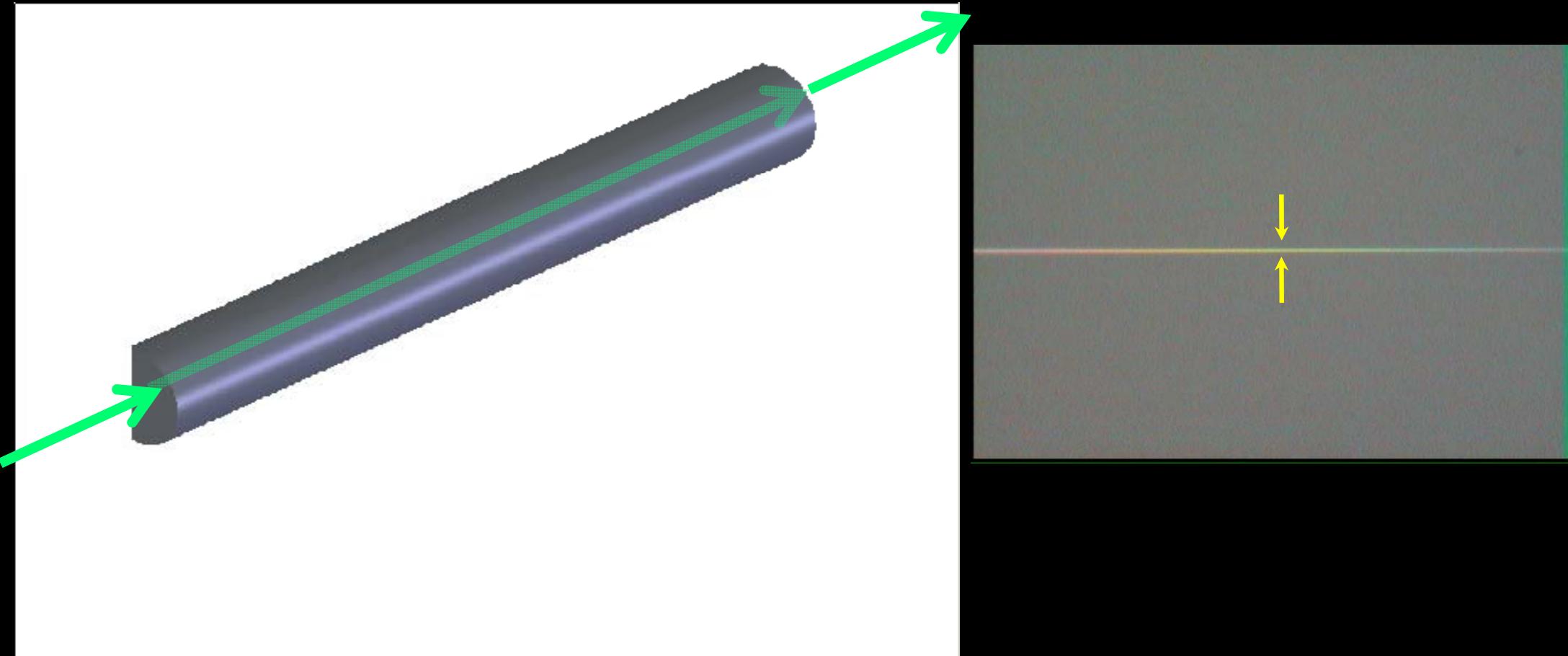


$$Q = \omega \cdot \tau > 10^8$$
$$F > 10^6$$

* D. K. Armani, T. J. Kippenberg, S. M. Spillane, K. J. Vahala.
Nature 421, 925-928 (2003).

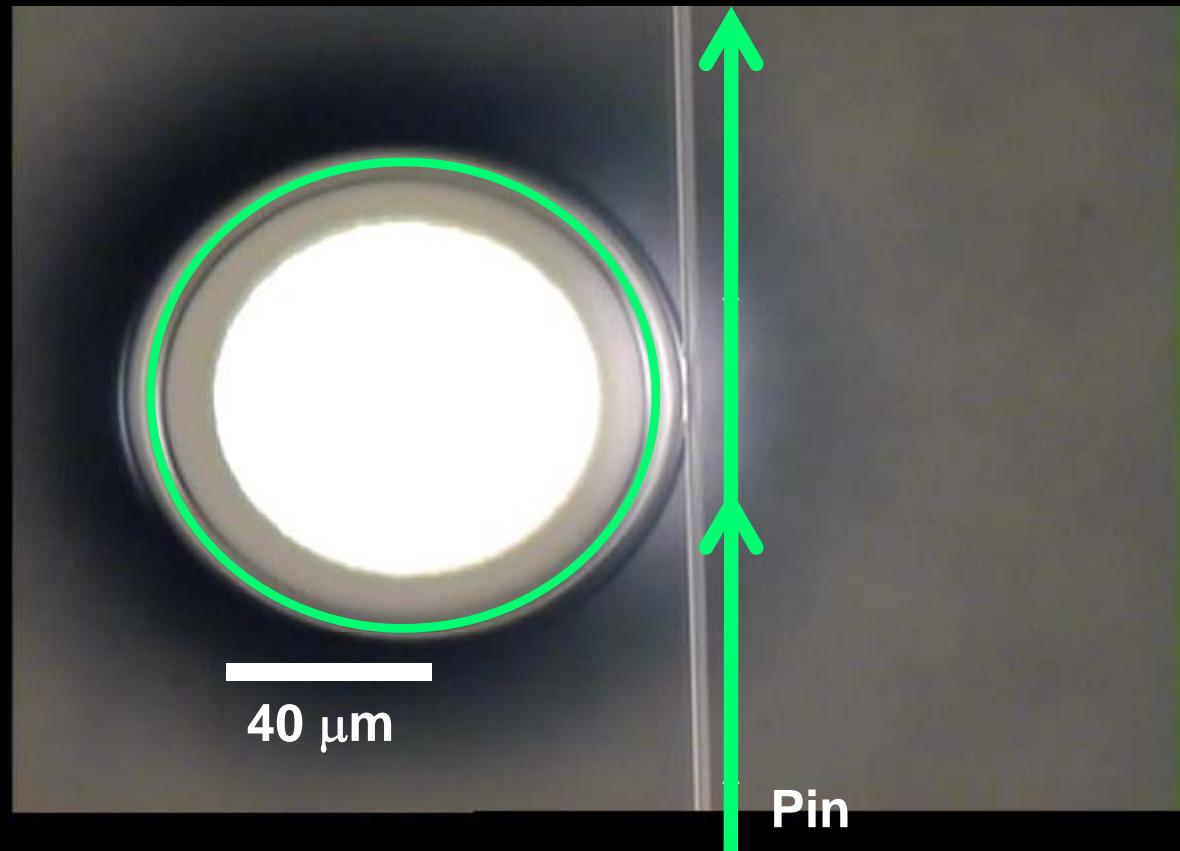
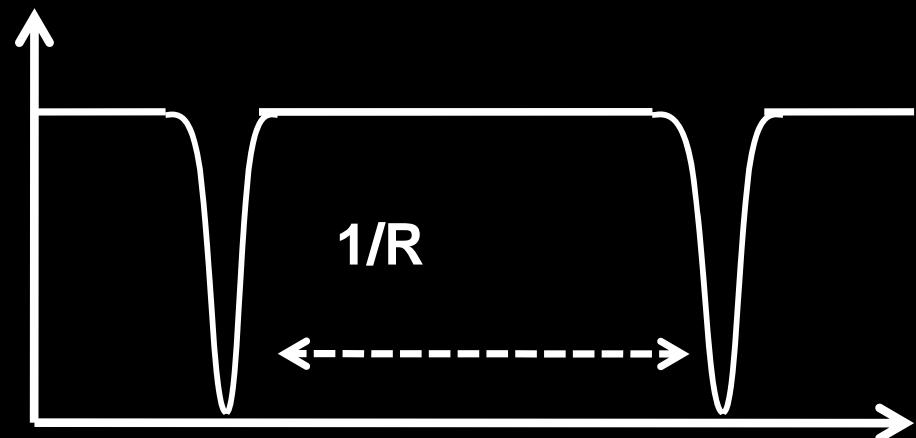


Optical fiber coupling



M. Cai, O.J. Painter, K. J. Vahala. *Phys. Rev. Lett.* (2002).

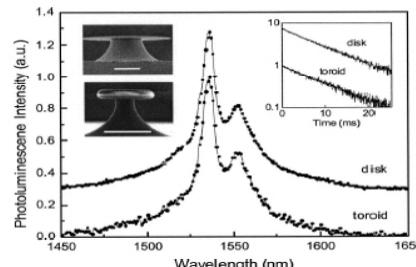
Optical Microresonators with Giant Photon Lifetimes



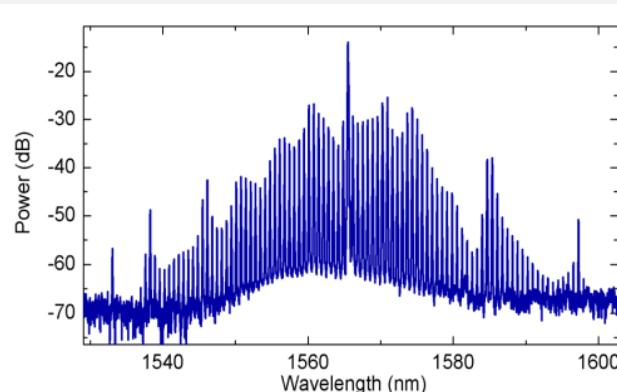
S. M. Spillane, T. J. Kippenberg, O.J. Painter, K. J. Vahala. Phys. Rev. Lett. (2003).
T.J. Kippenberg, S.M. Spillane, K.J. Vahala, Optics Letters, (2002).

Ultra-high Q Physics

Narrow Linewidth Laser sources



Frequency Combs on a Chip



Nonlinear Optics at ultra low power

Kippenberg, Spillane, Vahala.
Phys. Rev. Lett. (2004).

Spillane, Kippenberg, Vahala.
Nature, 621-623 (2002).

Cavity QED strong coupling

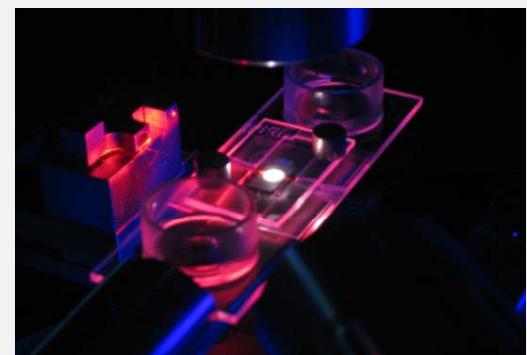


Kimble Group
Nature (2006), Science (2008)

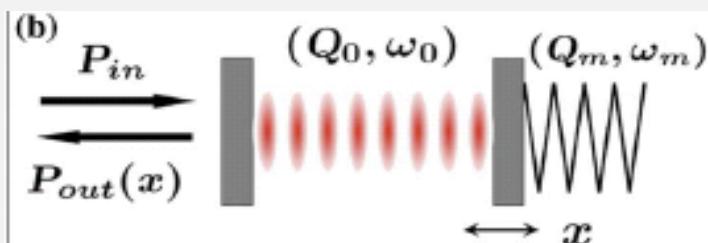
Applications in Science, Technology, Metrology



Molecule Sensing

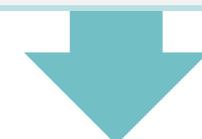


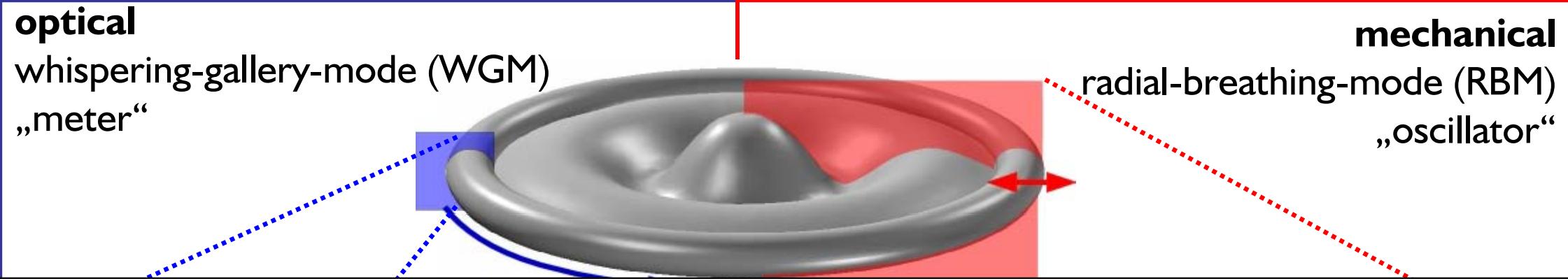
Cavity Optomechanics



Kippenberg, Vahala Science (2008)

Fundamental Research





$$\hat{H} = \hbar\omega_L \hat{a}_p^\dagger \hat{a}_p + \hbar\omega_m \hat{a}_m^\dagger \hat{a}_m + \hbar g_m \hat{a}_p^\dagger \hat{a}_p (\hat{a}_m + \hat{a}_m^\dagger)$$

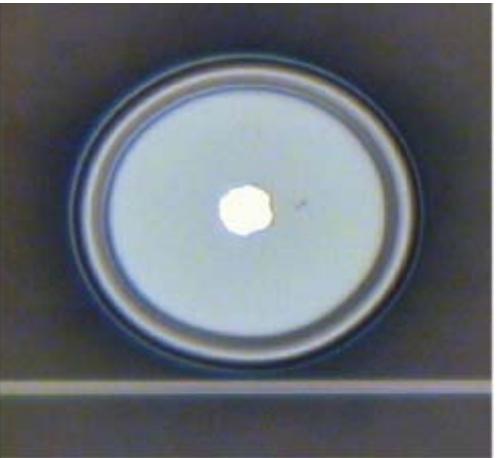
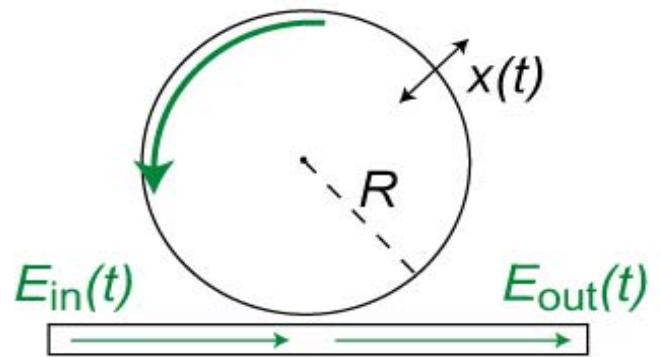
$$g_m = \frac{x_0}{L} \omega_0$$

Coupling strength

$$x_0 = \sqrt{\frac{\hbar}{2m\omega}}$$

Zero point motion

Optomechanical Coupling

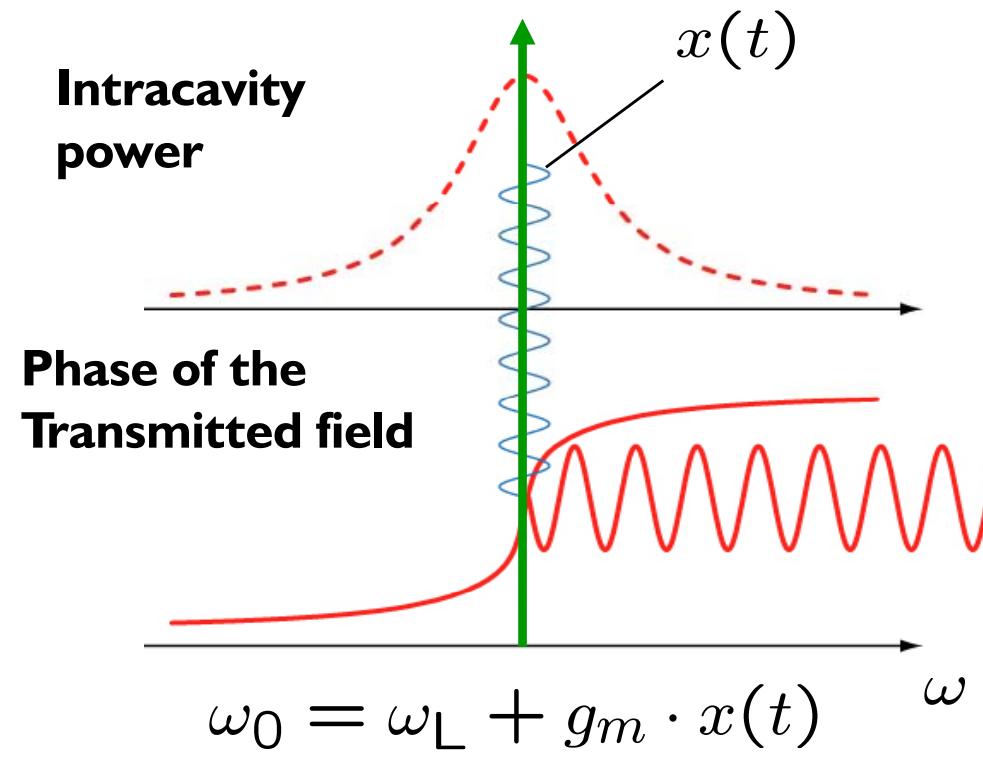


$$H_{int} = \hbar \omega a^\dagger a (1 - g_m x)$$

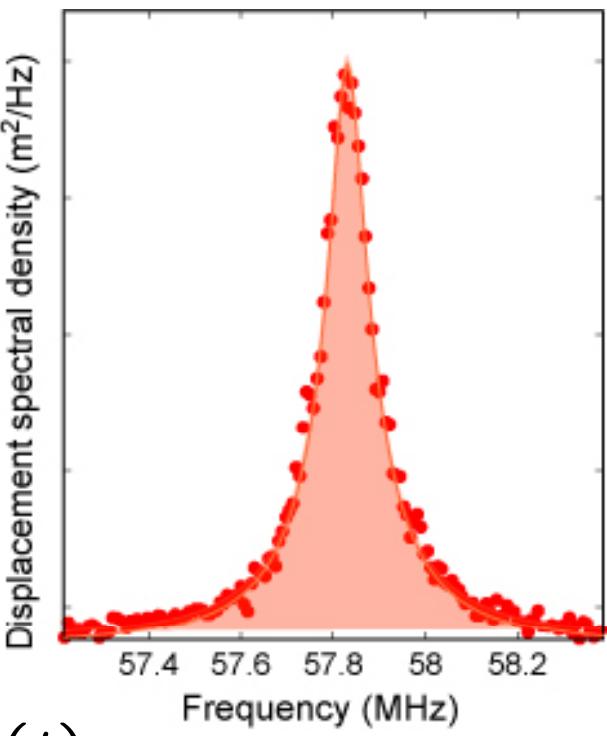
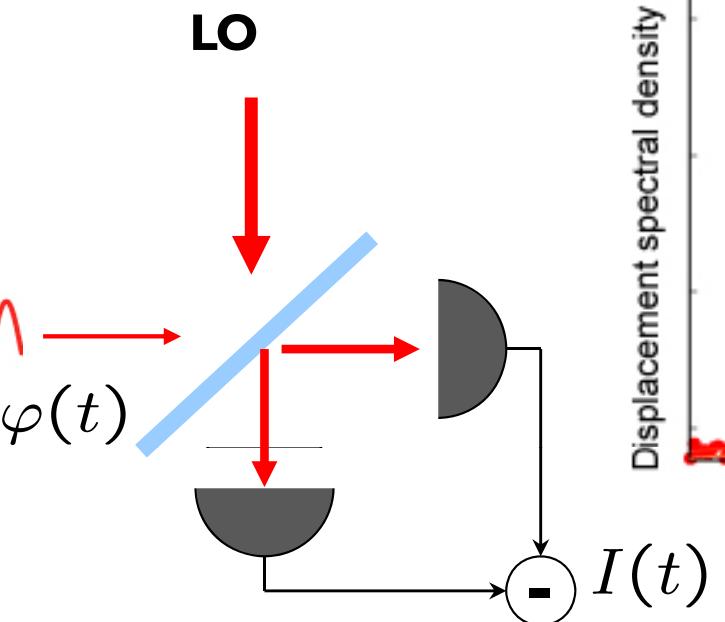
$$g_m = \frac{\omega}{R} \approx 10\text{GHz/nm}$$

$$F_{rad} = \hbar \cdot g_m$$

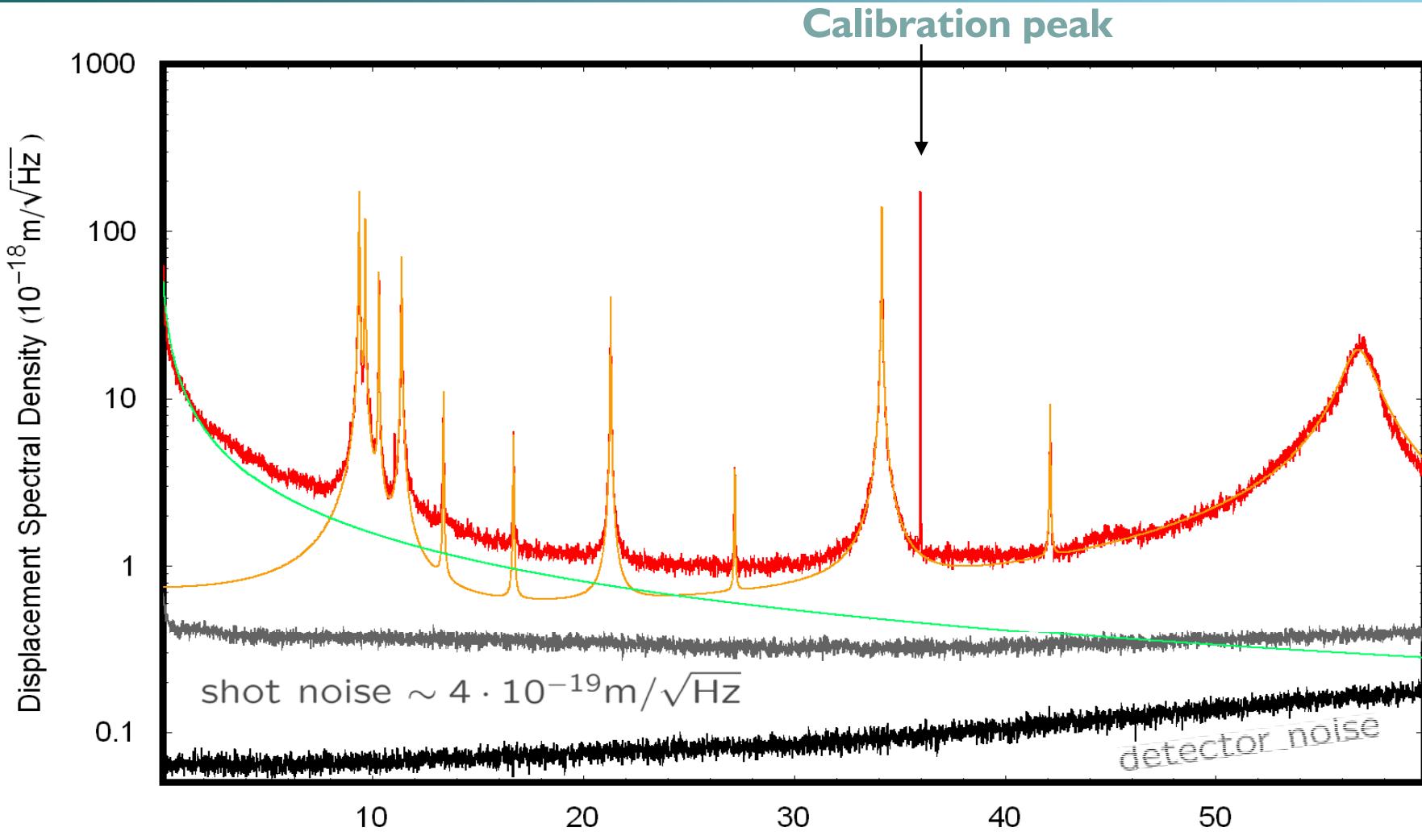
$$T_{eff} = \frac{1}{k_B} \int m_{eff} \delta \Omega \cdot x [\Omega]^2 \cdot \Omega^2$$



Homodyne Detection



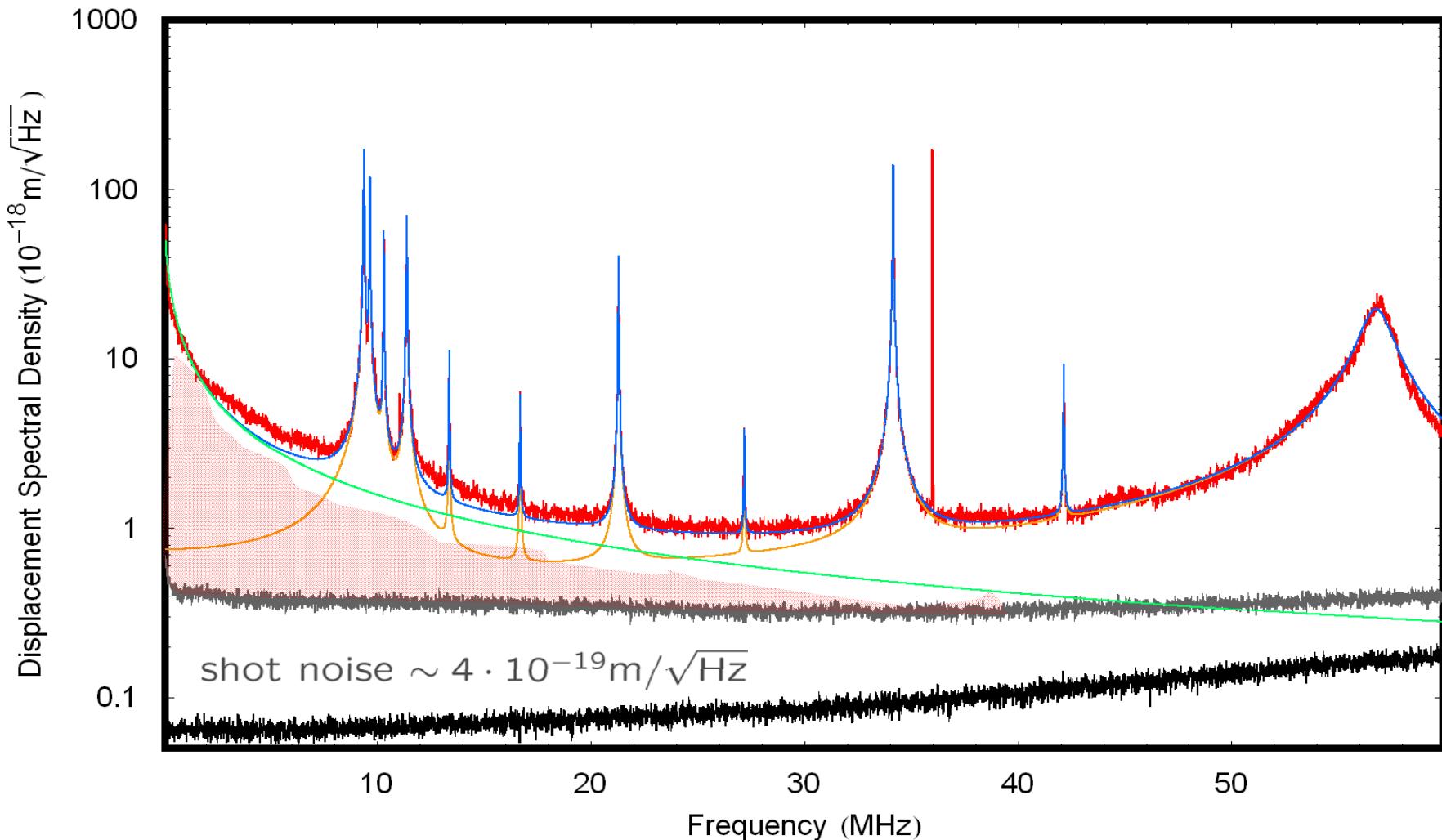
Quantitative Analysis of Noise Spectra



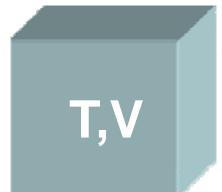
$$\sqrt{S_x^{\text{shot}}(\Omega)} \approx \frac{\lambda}{16\pi\eta_C\mathcal{F}} \cdot \sqrt{\frac{1 + \left(\frac{\Omega}{\kappa/2}\right)^2}{P_{\text{in}}/\hbar\omega}}$$

Schliesser, Anetsberg, Riviere, Arcizet, Kippenberg NJP, 10, 095007

Quantitative Analysis of Noise Spectra



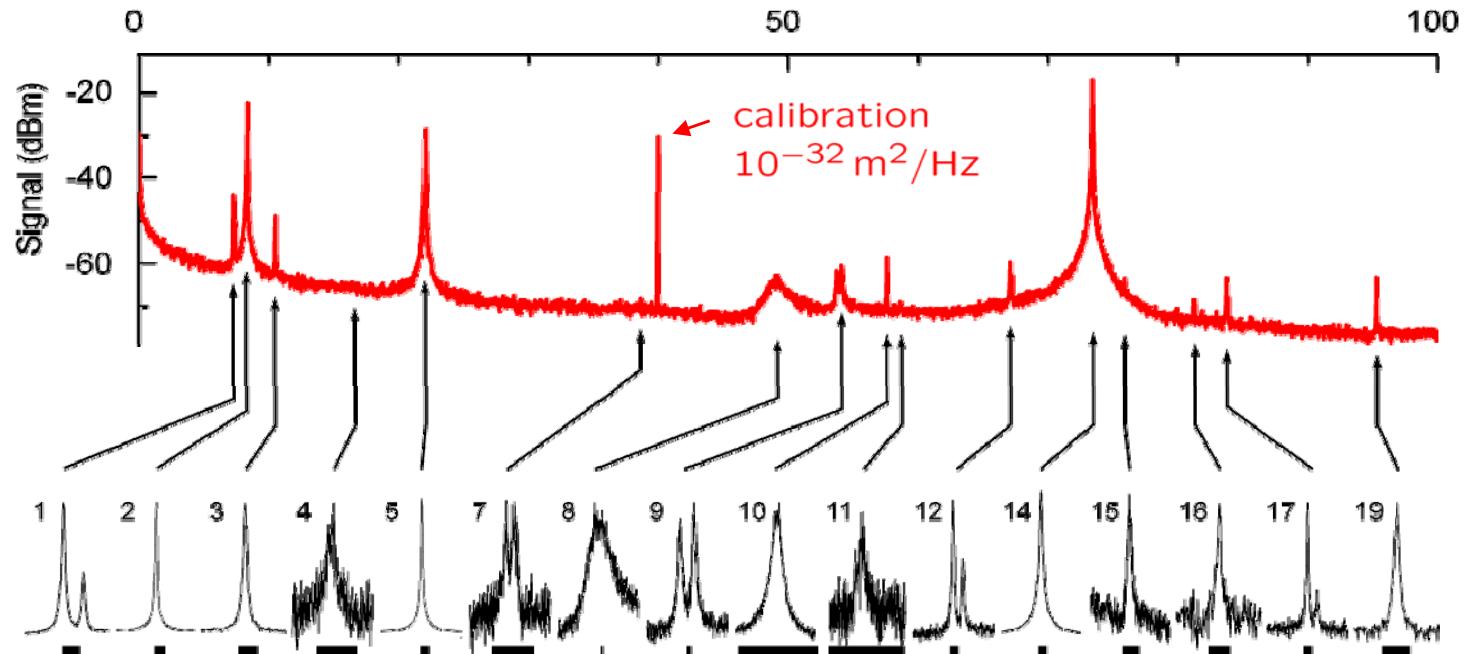
Thermorefractive noise caused by temperature fluctuations



$$\langle \Delta T^2 \rangle = \frac{k_B T^2}{\rho C_p V}$$

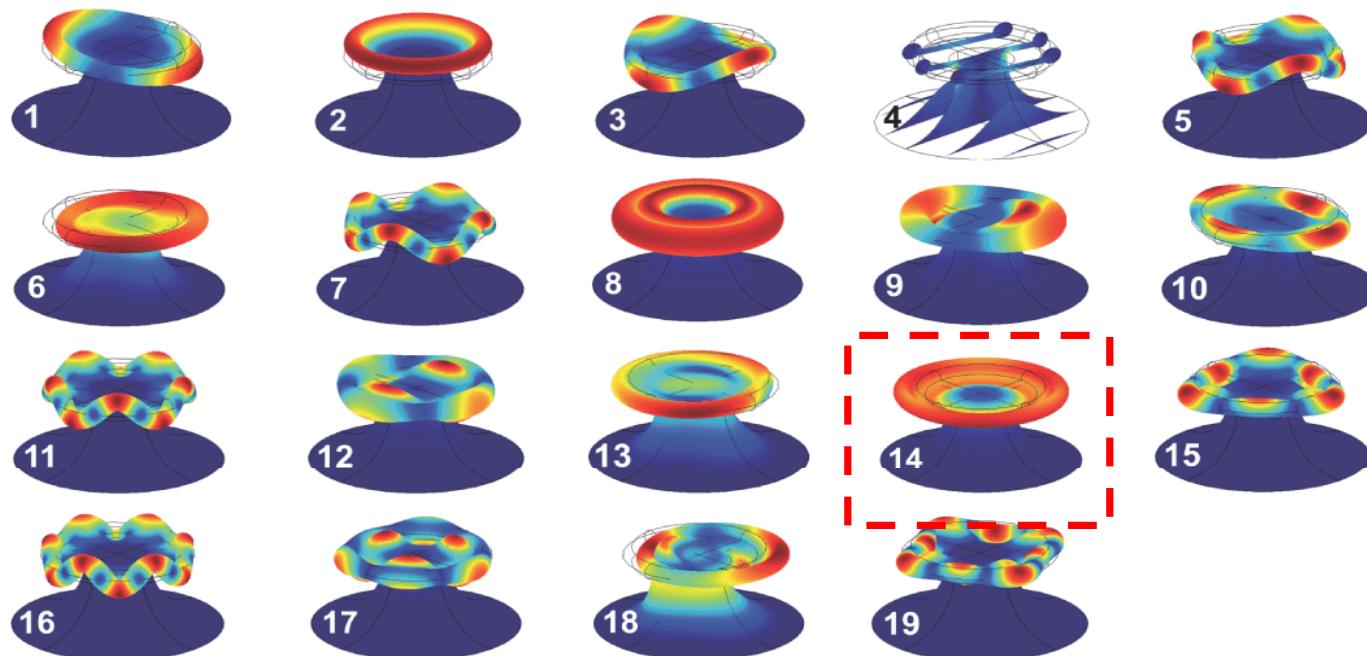
V. B. Braginsky, M. L. Gorodetsky, and S. P. Vyatchanin, "Thermodynamical fluctuations and photo-thermal shot noise in gravitational wave antennae," Phys. Lett. A 264, 1–10 (1999).

Mechanical Modes



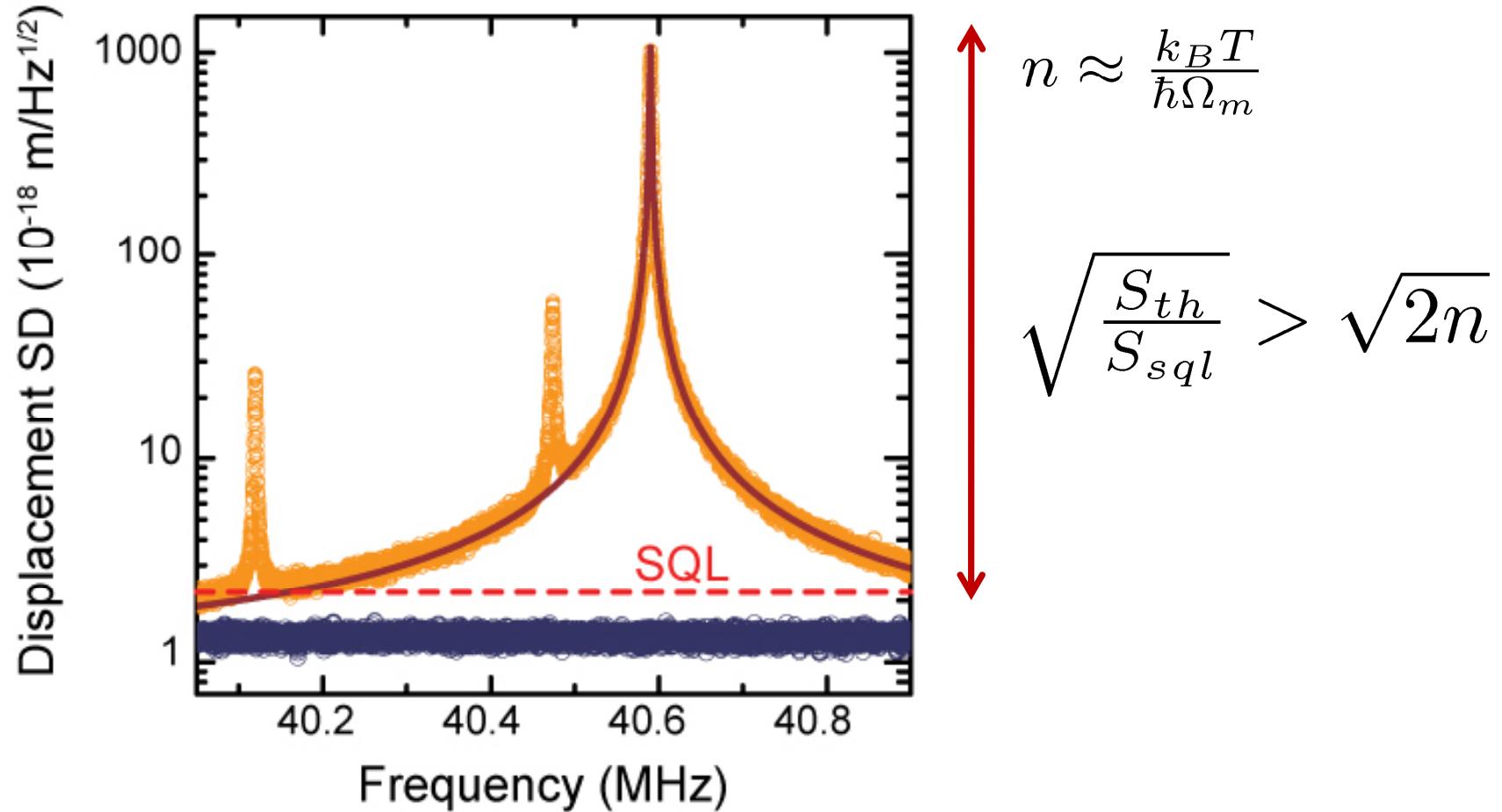
measured
mechanical
spectrum

zoom on
individual peaks



mode patterns
obtained from
finite element
modeling

Sensitivity of the Displacement Measurement

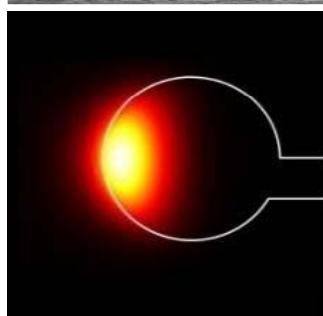
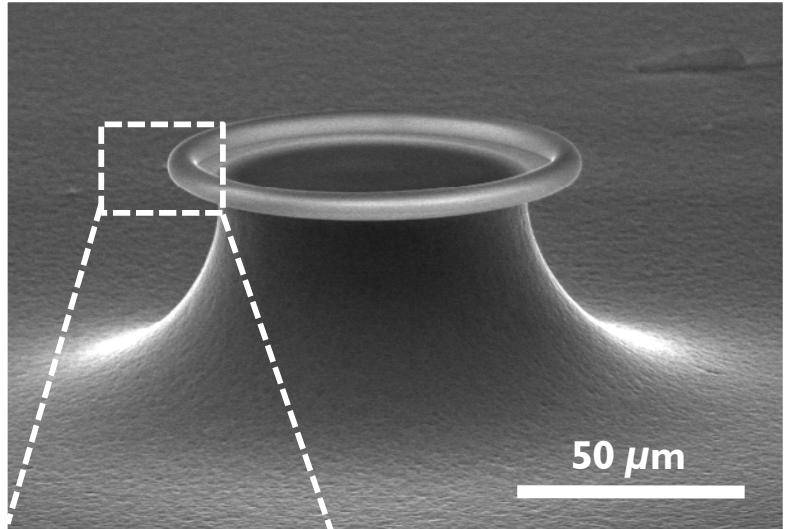


$$\frac{\sqrt{S_x^{sig}[\Omega_m]}}{\sqrt{S_x^{zpm}[\Omega_m]}} = \sqrt{2n}$$

From signal to background one can deduce that the sensitivity is *below* the SQL

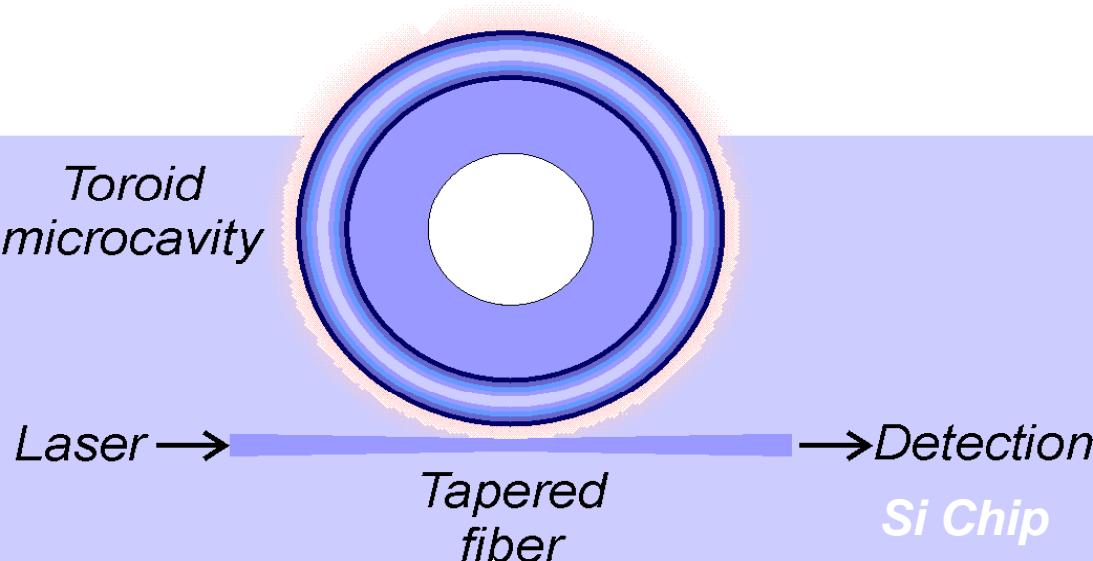
Nano-Optomechanics

Toroid microcavity:



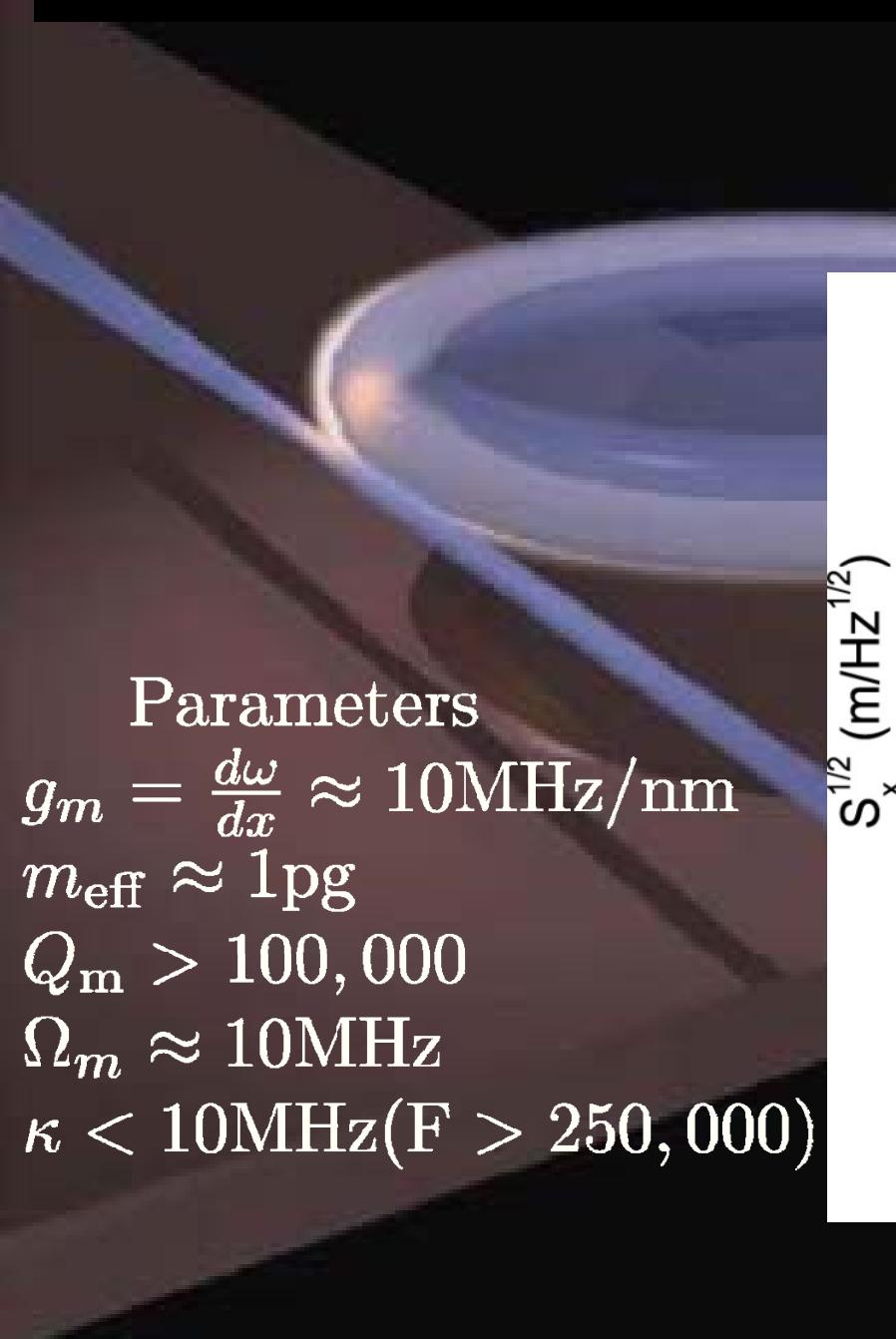
Finesse $> 10^6$
Linewidth $< 1 \text{ MHz}$

Armani et al. Nature 421, 925 (2003).
Kippenberg et al., APL 85, 6113 (2004)



Optomechanical near-field interaction

High Q SiN nanomechanical beams



Parameters

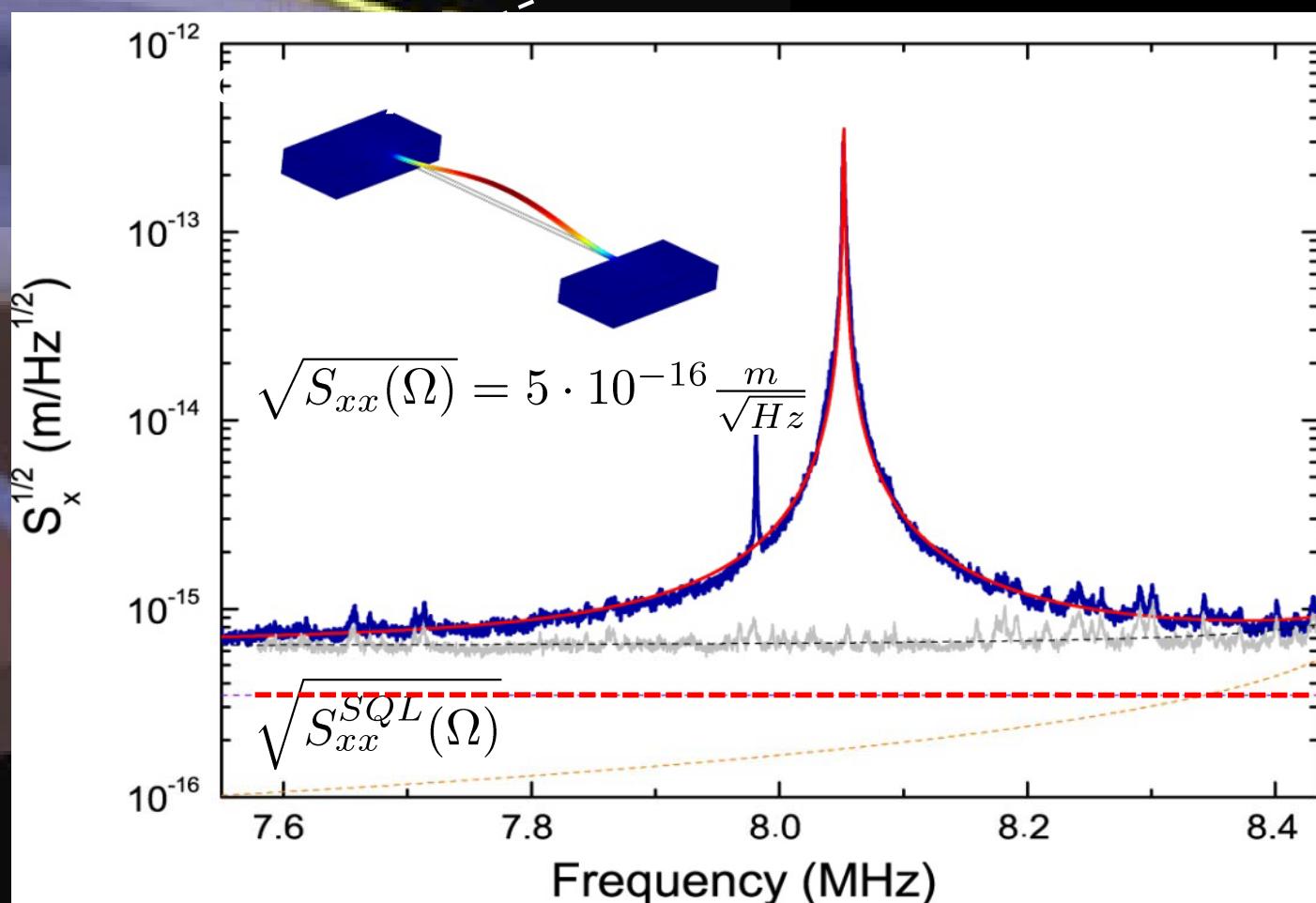
$$g_m = \frac{d\omega}{dx} \approx 10 \text{MHz/nm}$$

$$m_{\text{eff}} \approx 1 \text{pg}$$

$$Q_m > 100,000$$

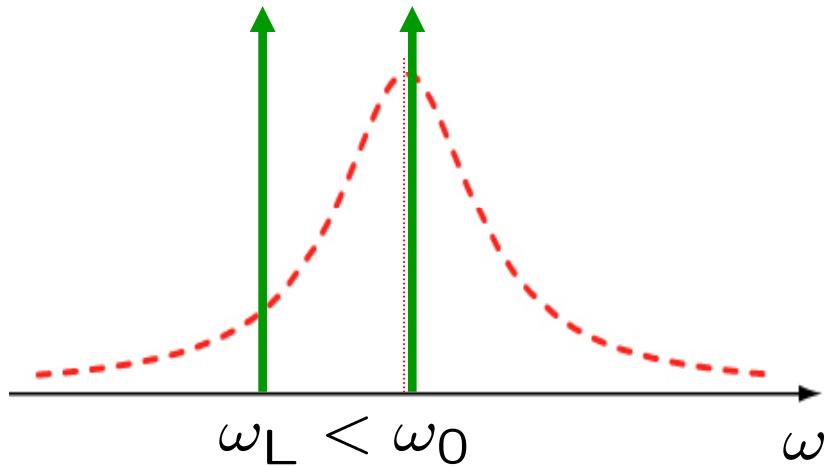
$$\Omega_m \approx 10 \text{MHz}$$

$$\kappa < 10 \text{MHz} (\text{F} > 250,000)$$

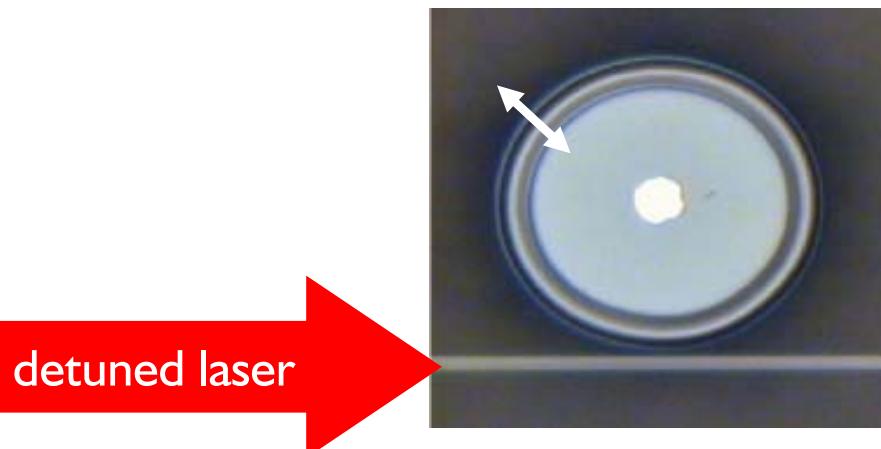


Cooling by Dynamical Backaction at MPQ

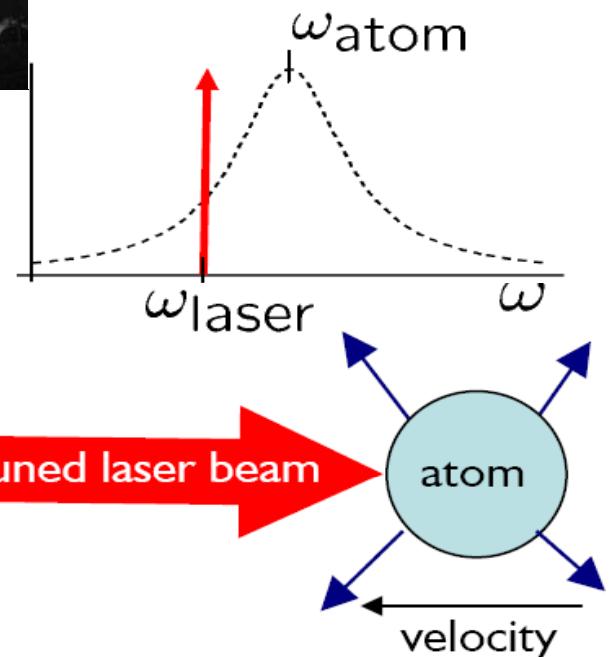
detuned laser induces viscous force on cavity boundary



V.B. Braginsky



.... similar to
laser cooling of
atoms

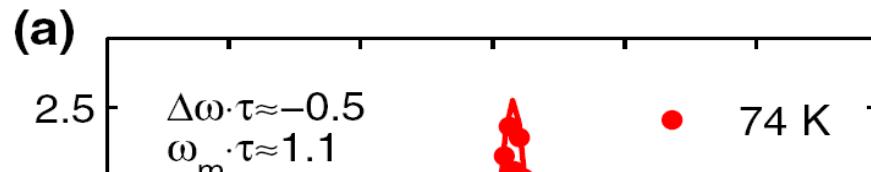


Braginsky, Manukin, "Ponderomotive effects of electromagnetic radiation", JETP Letters 52, 986 (1967)

Hänsch, Schawlow, "Cooling of gases by laser radiation", Opt. Commun. 13, 68 (1975)

Wineland, Dehmelt, "Proposed $10^{14} \Delta\nu < \nu$ laser fluorescence spectroscopy on Ti^+ ion mono-oscillator", Bull APS 20, 637 (1975)

Demonstration of Backaction Cooling @ MPQ



Key Parameters:

- Mechanical frequency of the cooled ring: **57.8 MHz**
- Initial temperature: **300 K**
- Final effective temperature: **74 K**

INVESTIGATION OF DISSIPATIVE PONDEROMOTIVE EFFECTS OF ELECTROMAGNETIC RADIATION

V. B. BRAGINSKII, A. B. MANUKIN, and M. Yu. TIKHONOV

Moscow State University

Submitted October 17, 1969

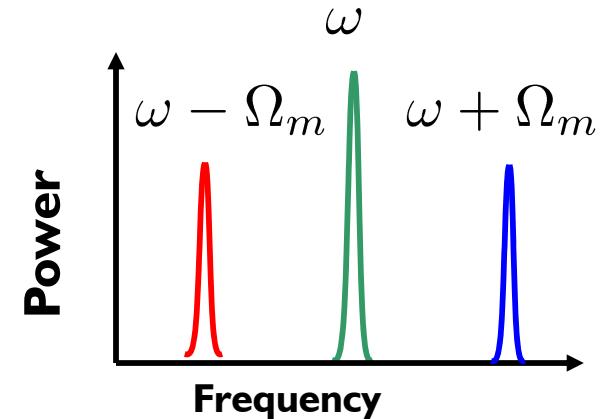
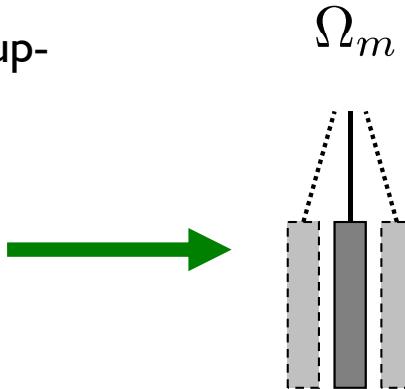
I'Haye., Nooshi,
Rev. Lett. **97**
don, Briant,
444, 71
Nature **444**,

→ **First Demonstration of Radiation Pressure Cooling as predicted by Braginsky and Dykman**

57.4 57.6 57.8 58 58.2
Frequency (MHz)

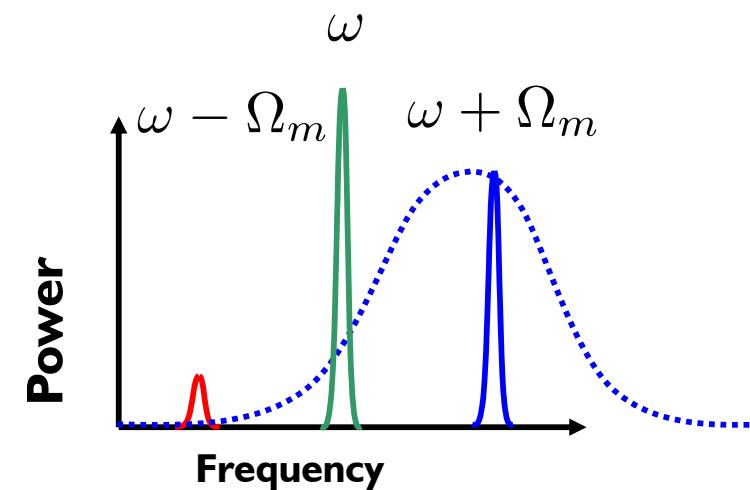
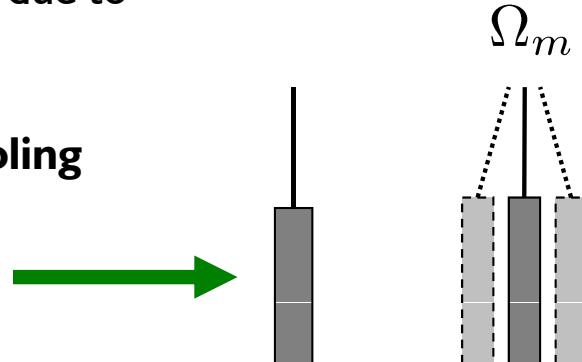
Sideband interpretation

An oscillating mirror will cause Doppler up- and down-shifted fields.



A cavity can create an imbalance due to resonant buildup

Excess anti-Stokes photons: **Cooling**



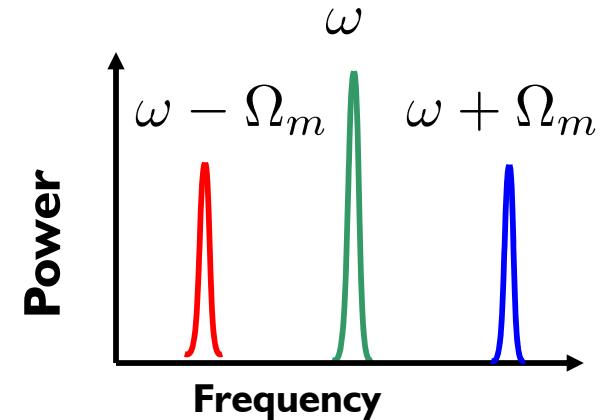
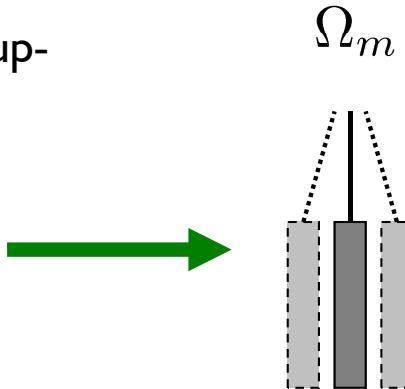
Similar mechanism to cavity cooling of atoms and molecules (coherent scattering)

V. Vuletic, S. Chu, *Phys. Rev. Lett.*, Vol. 84, No. 17 (2000)

P. Maunz, Puppe, Schuster, Syassen, Pinkse, Rempe, *Nature* (2004)

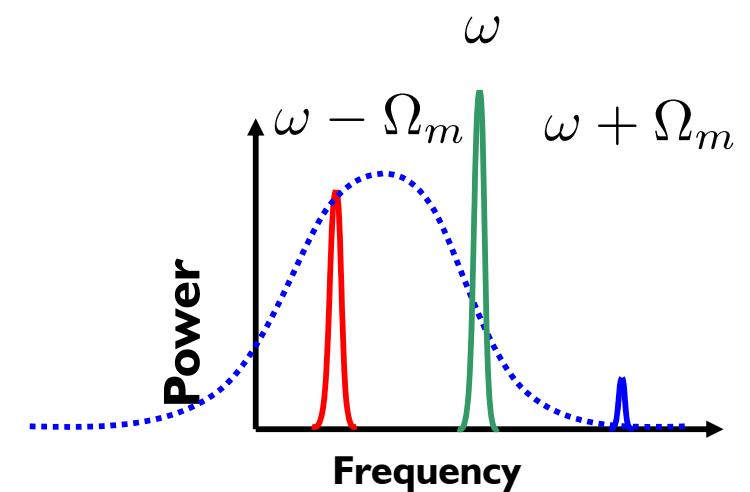
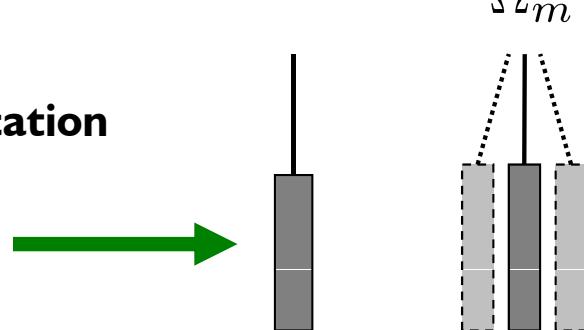
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A cavity can create an imbalance due to resonant buildup

Excess Stokes photons: **amplification**



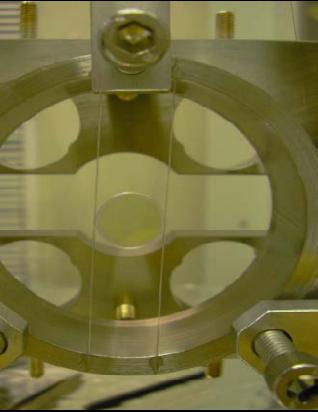
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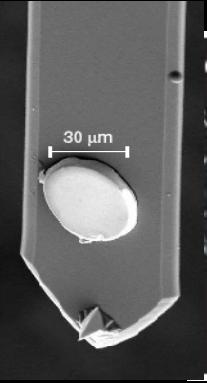
P. Maunz, Puppe, Schuster, Syassen, Pinkse, Rempe, *Nature* 428, 50 (2004).

Recent Experiments in cavity Optomechanics (2006-2009)

$$H_{int} = g_m \hbar a^\dagger a (a_m^\dagger + a_m)$$



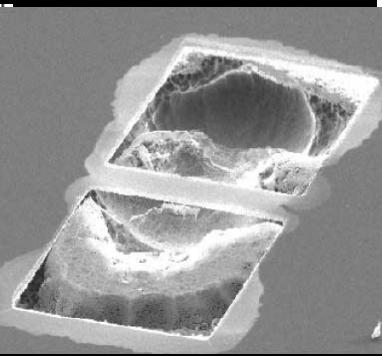
MIT/LIGO
Paris



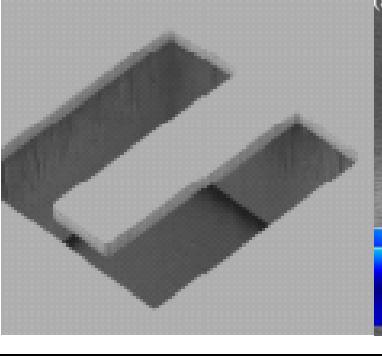
UCSB



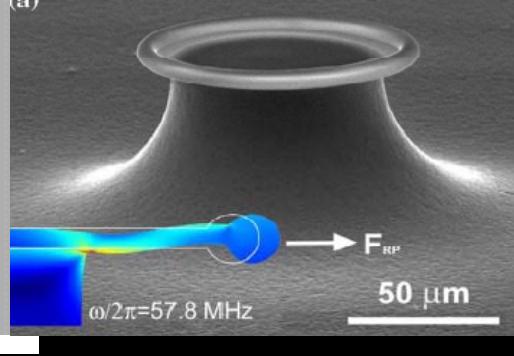
Yale



Vienna



Paris

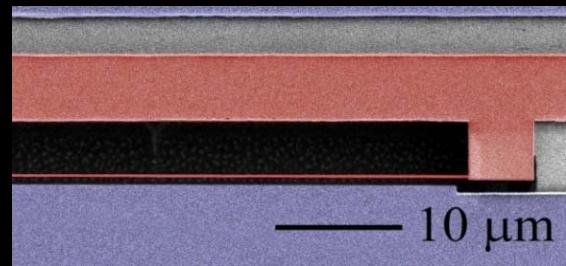


MPQ/Caltech/EPFL

$1Hz$

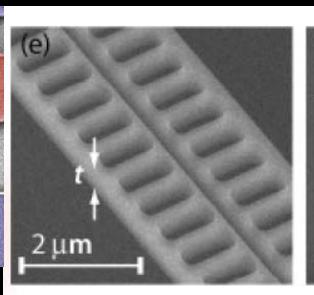


$1kHz$



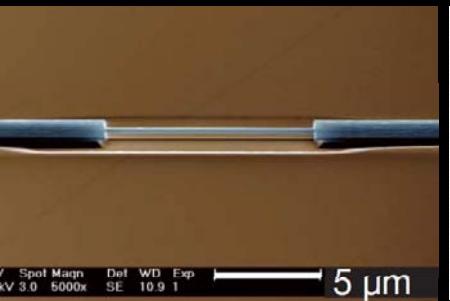
JILA, USA

$1MHz$

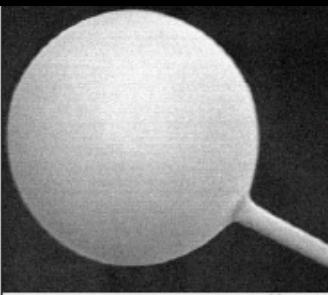


Caltech

$1GHz$



Ghent IMEC, Yale



Oregon

Quantum regime of mechanical oscillators

Prerequisites to be able observe e.g. zero point motion or quantum back-action of the measurement

- **Imprecision at the zero point motion level** $S_{SQL} = \sqrt{\frac{\hbar}{m\Omega_m\gamma_m}}$
- **Mechanical oscillators in which thermal noise is sufficiently reduced (close to quantum ground state)**

$$n \approx \frac{k_B T}{\hbar \Omega} \quad \Omega_m = 1\text{MHz}, T_Q = 50\mu\text{K}$$

Difficult! e.g. Science Magazine, A. Cho:

Physics

Researchers Race Quantum Into Me

3 JANUARY 2003

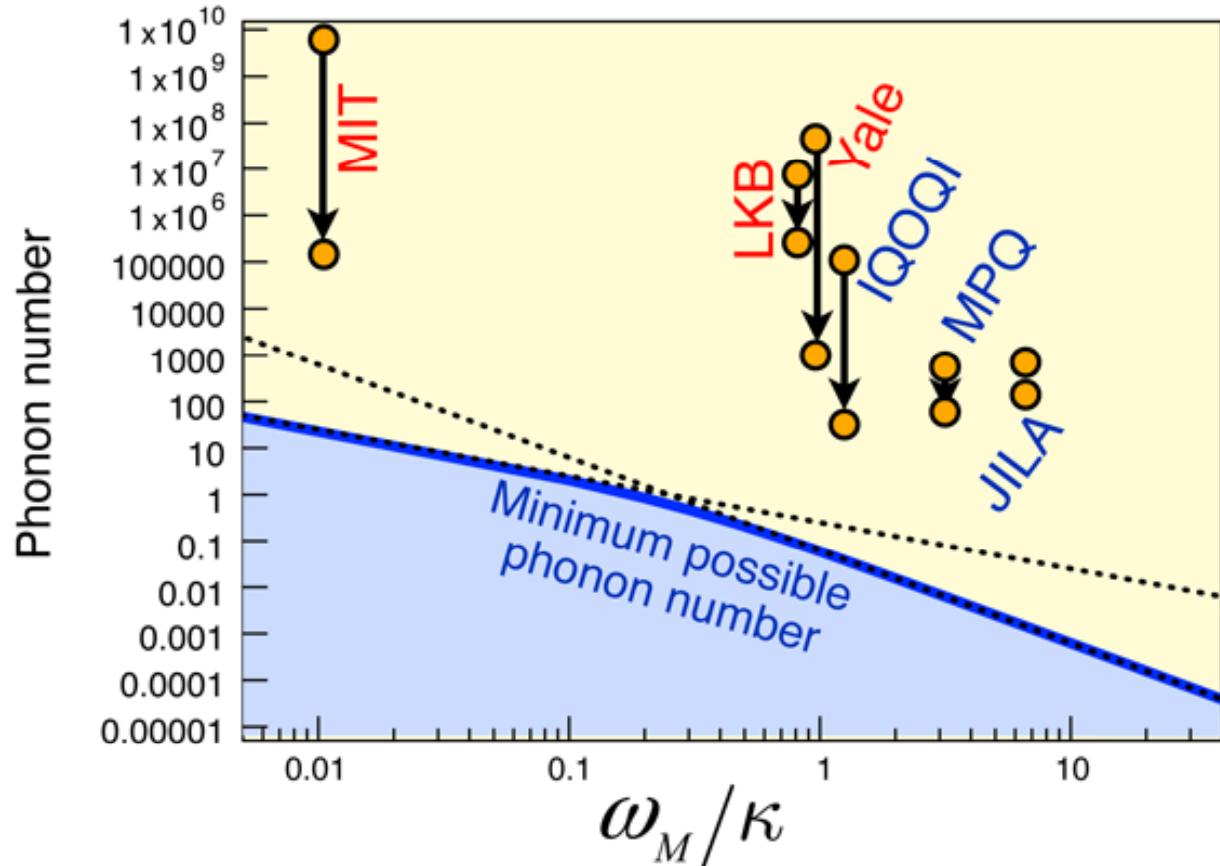
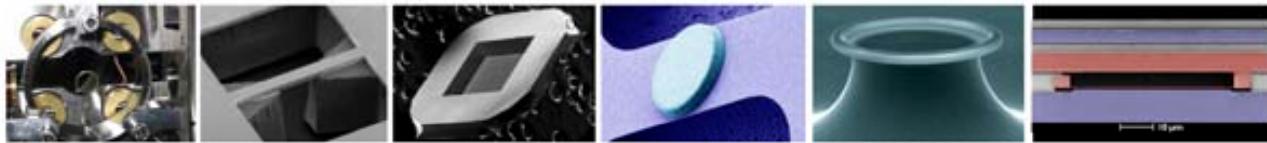
ble motion. At least four groups hope to reach the quantum limit of motion within months. The feat could open

"We don't see quantum behavior in our sense we're nics," says physicist at New to find out, putting pro devices into to observe

SCIENCE ARCHIVING MPQ

ScienceMag.org on October 17, 2003

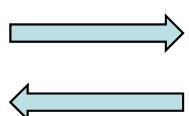
Demonstration of Backaction Cooling



Physics 2, 40 (2009)
DOI: 10.1103/Physics.2.40
F. Marquardt, Girvin, Optomechanics

Quantum Limits of Radiation Pressure Cooling

Quantum Limits of radiation pressure cooling – photon shot noise perspective



$$\Delta p = 2\hbar \cdot k$$

$$F_{RP} = \langle F_{RP} \rangle + \delta F(t)$$

$$S[\Omega]_{NN} = n \cdot \frac{\kappa}{\frac{\kappa^2}{4} + (\Omega - \Delta)^2}$$

Spectrum of Photon Number Fluctuations

$$S_{FF}[\Omega] = \left(\frac{2\hbar k}{T_{rt}}\right)^2 S_{NN}[\Omega]$$

Spectrum of Radiation Pressure Fluctuations

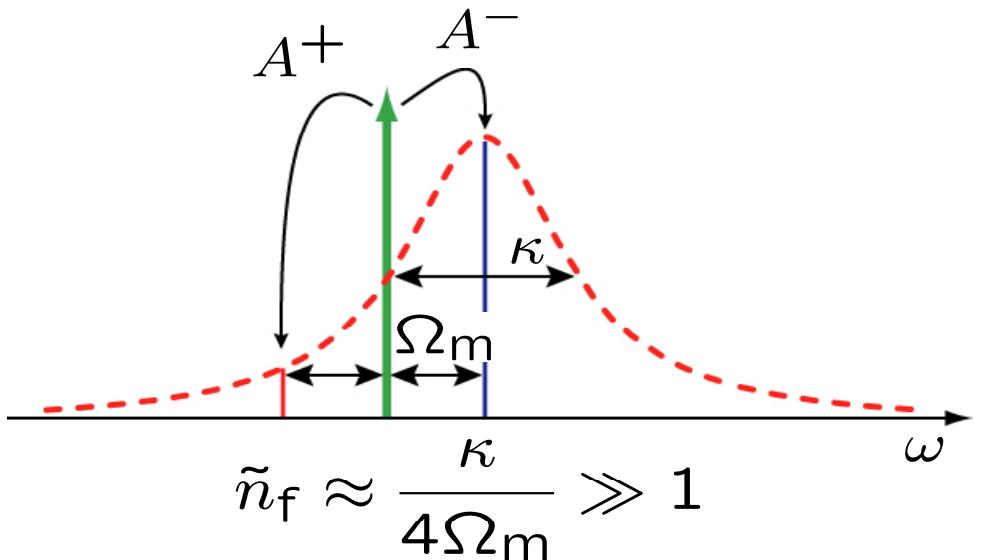
$$n_f = \frac{1}{\hbar \Omega_m} \int_{-\infty}^{\infty} S_{FF}[\Omega] \cdot \chi(\Omega)^2 \Omega^2 m_{eff} d\Omega$$

Final temperature due to
Radiation Pressure
Fluctuations

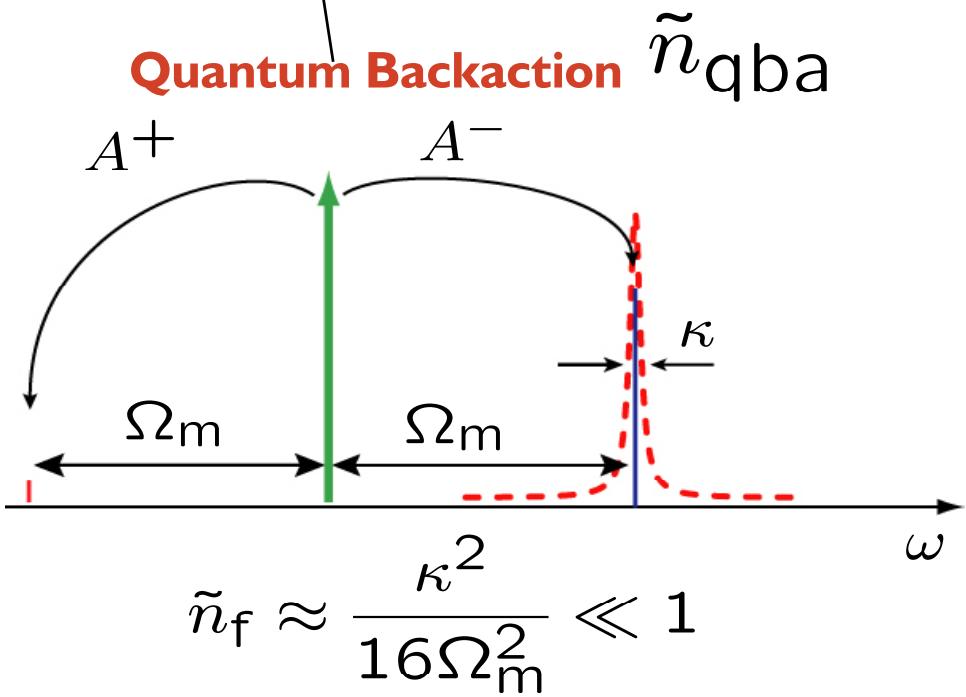
Quantum Limits of Radiation Pressure Cooling

$$n_f \approx \frac{\Gamma_m}{\Gamma_{cool}} n_i + \frac{A_+}{A_- - A_+}$$

Reservoir heating

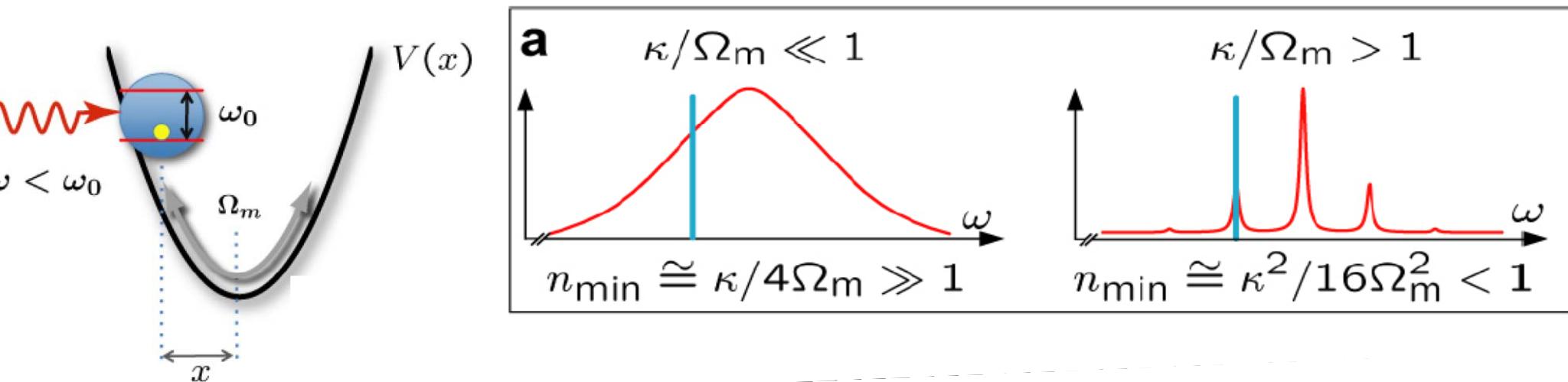


„Doppler“ limit
ground-state cooling impossible



resolved sideband cooling
ground-state cooling possible

Resolved Sideband Cooling



Lead to ground state cooling of ions:

F. Diedrich, J. C. Bergquist, W. M. Itano, D. J. Wineland, *Physical Review Letters* 62, 403 (Jan, 1989).

EL 9 Proposed $10^{14} \Delta v < v$ Laser Fluorescence Spectroscopy on Tl^+ Mono-Ion Oscillator III.
D. WINELAND and H. DEHMELT, U. of Washington.--The pulsed

EL 9 Proposed $10^{14} \Delta v < v$ Laser Fluorescence Spectroscopy on Tl^+ Mono-Ion Oscillator III.
D. WINELAND and H. DEHMELT, U. of Washington.--

oscillation parallel the λ_2 beam. Then the MIO pre-dominantly absorbs $(v_2 - v_v)$ photons. As it emits photons at all side band frequencies $v_2 + nv_v$, $|n|=0,1,2\dots$, but of average energy hv_2 the balance hv_v has to come from the oscillatory motion. The maximum cooling rate is $-hv_v/2\tau_2 \approx .2\text{eV/s}$! This only drops off for $v_m \ll v_m^*$, promising oscillatory temperatures $hv_v/2k \approx 10^{-4}\text{eV} \leq T_i \ll Mv_m^*{}^2/2k \approx .004\text{eV}$, $-\delta_D \ll 7\text{mHz}$, $-\delta_S \ll 3\text{mHz}$ and no Doppler side bands on the λ_0 resonance! Directing the λ_2 beam along $(-i + j + k)$ makes the cooling 3-dimensional.--H.D. thanks H. Walter and coworkers for stimulating discussions. We thank the National Science Foundation for its generous support.

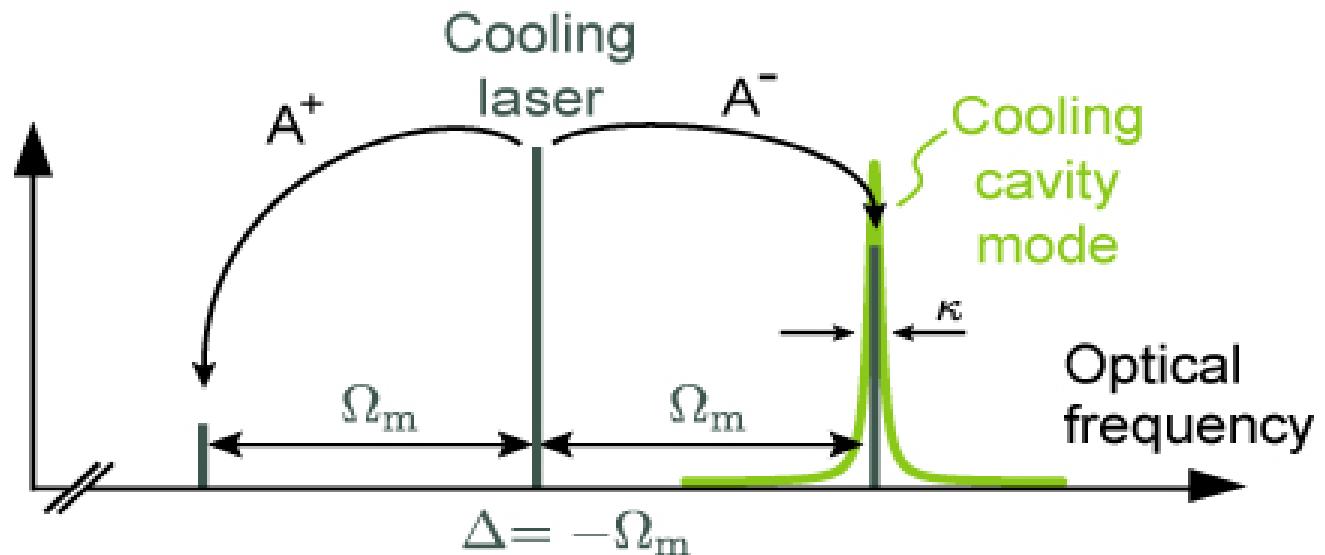
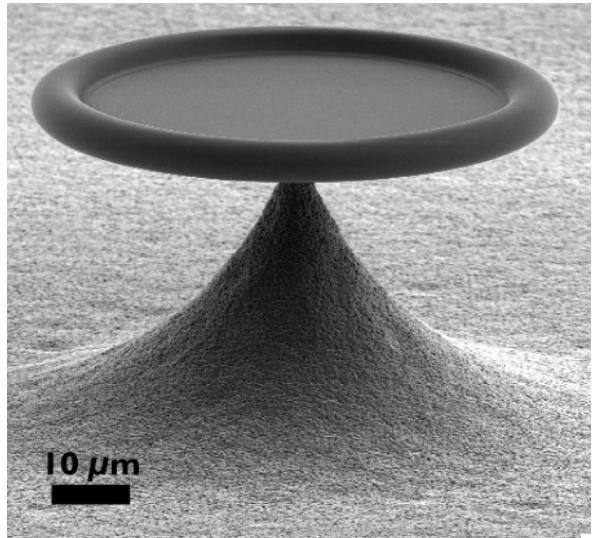
¹Dehmelt, Bull. APS 18, 1521 (1973) & 20, 60 (1975).

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FÜR QUANTENOPTIK
GARCHING



Demonstration of Resolved Sideband Cooling

$$\kappa \ll \Omega_m$$

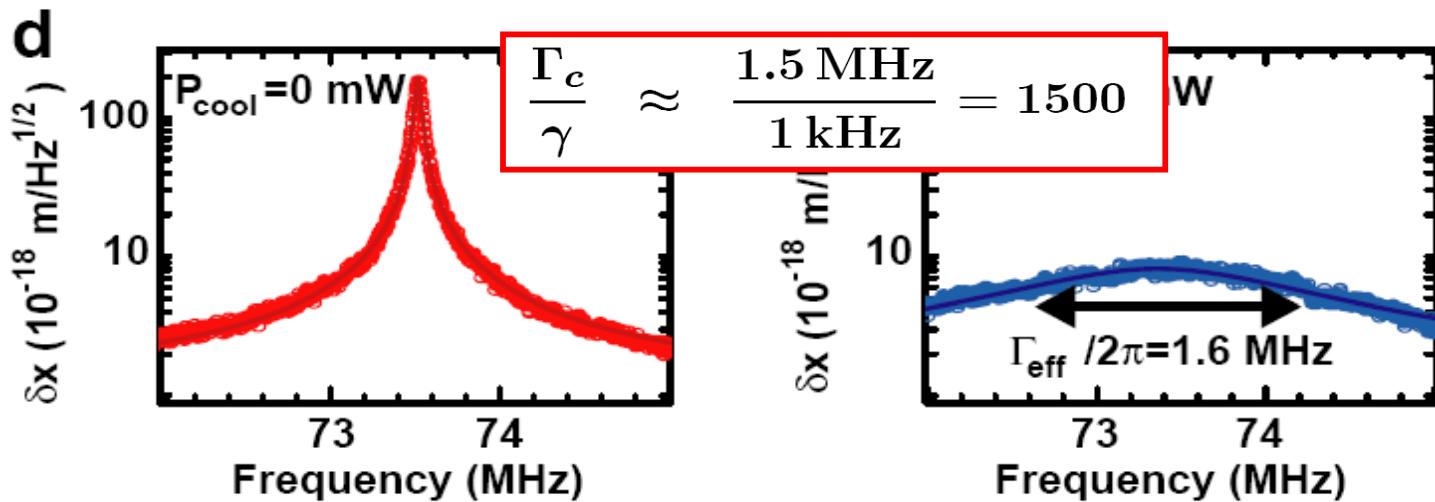


$$\Omega_m/2\pi = 73.5 \text{ MHz}$$

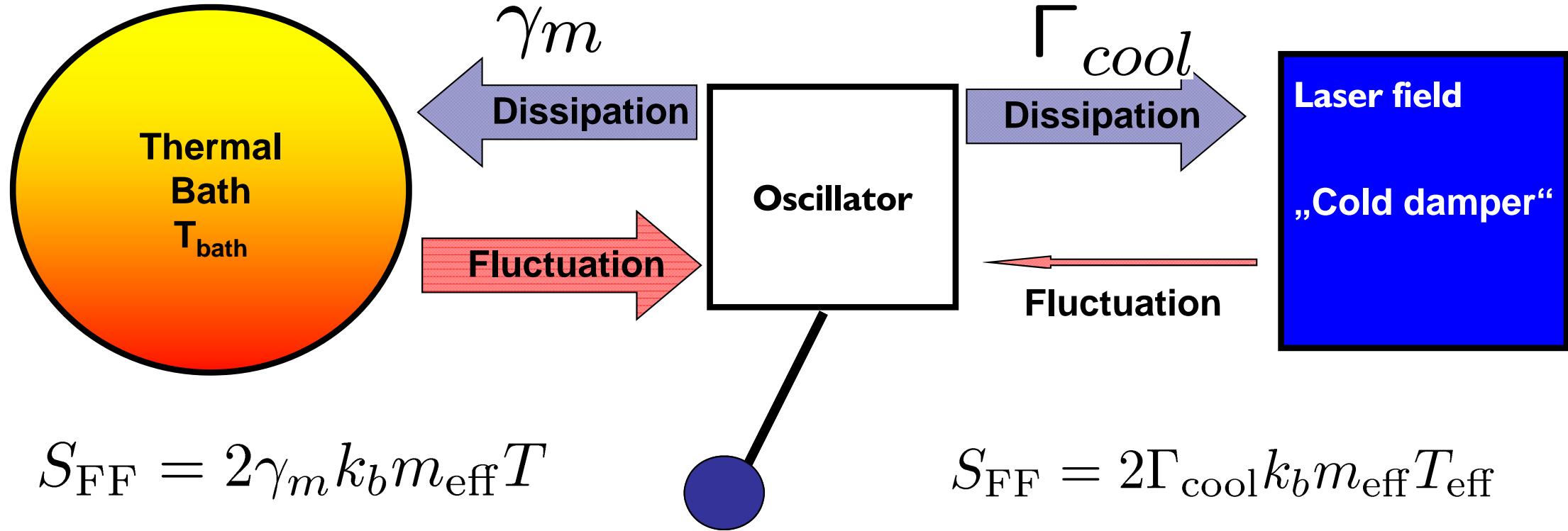
$$\kappa/2\pi = 3.2 \text{ MHz}$$

$$\mathcal{F} = 440000$$

$$\Omega_m/\kappa \approx 22$$



The challenge of ground state cooling



$$S_{FF} = 2\gamma_m k_b m_{\text{eff}} T$$

$$S_{FF} = 2\Gamma_{\text{cool}} k_b m_{\text{eff}} T_{\text{eff}}$$

$$T_{\text{eff}} = \frac{\hbar\Omega_m}{2k_b} \left(1 + \frac{A^+}{A^+ - A^-} \right)$$

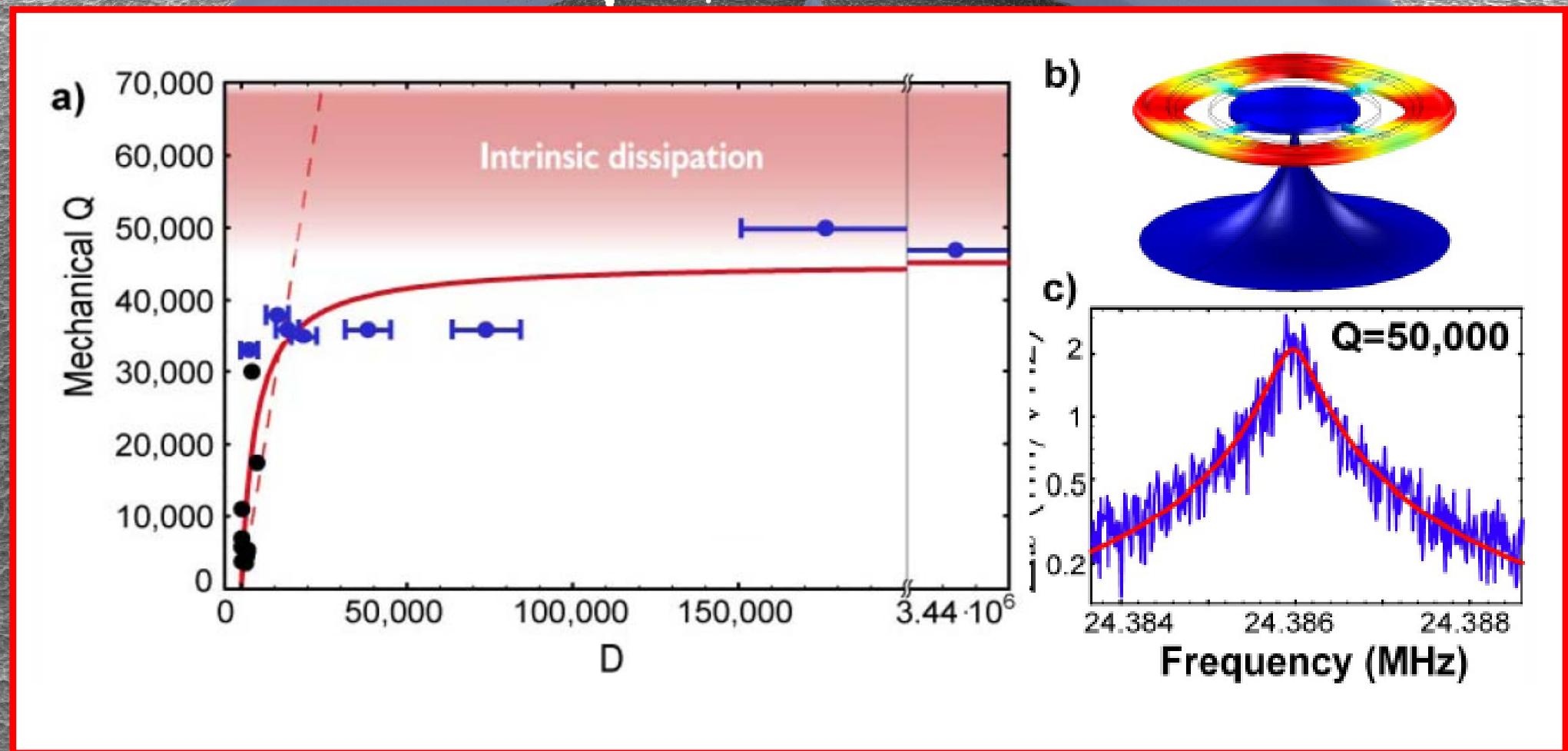
Wilson-Rae, Nooshi, Zwerger, Kippenberg, PRL **99**, 093901 (2007)
 Marquardt, Chen, Clerk, Girvin, PRL **99**, 093902 (2007)

$$T_f \approx \frac{\gamma_m}{\gamma_m + \Gamma_{\text{cool}}} \cdot T_i$$

Improving mechanical Q

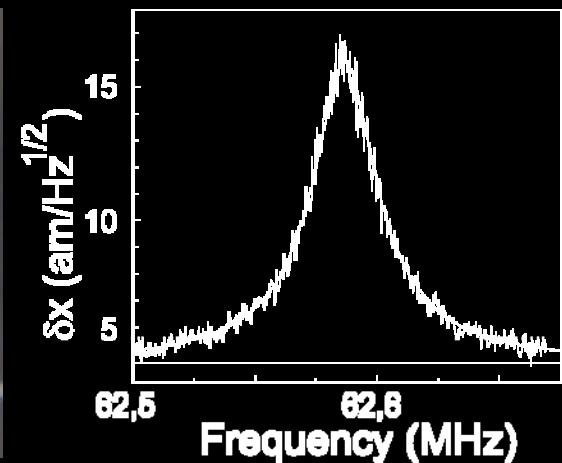
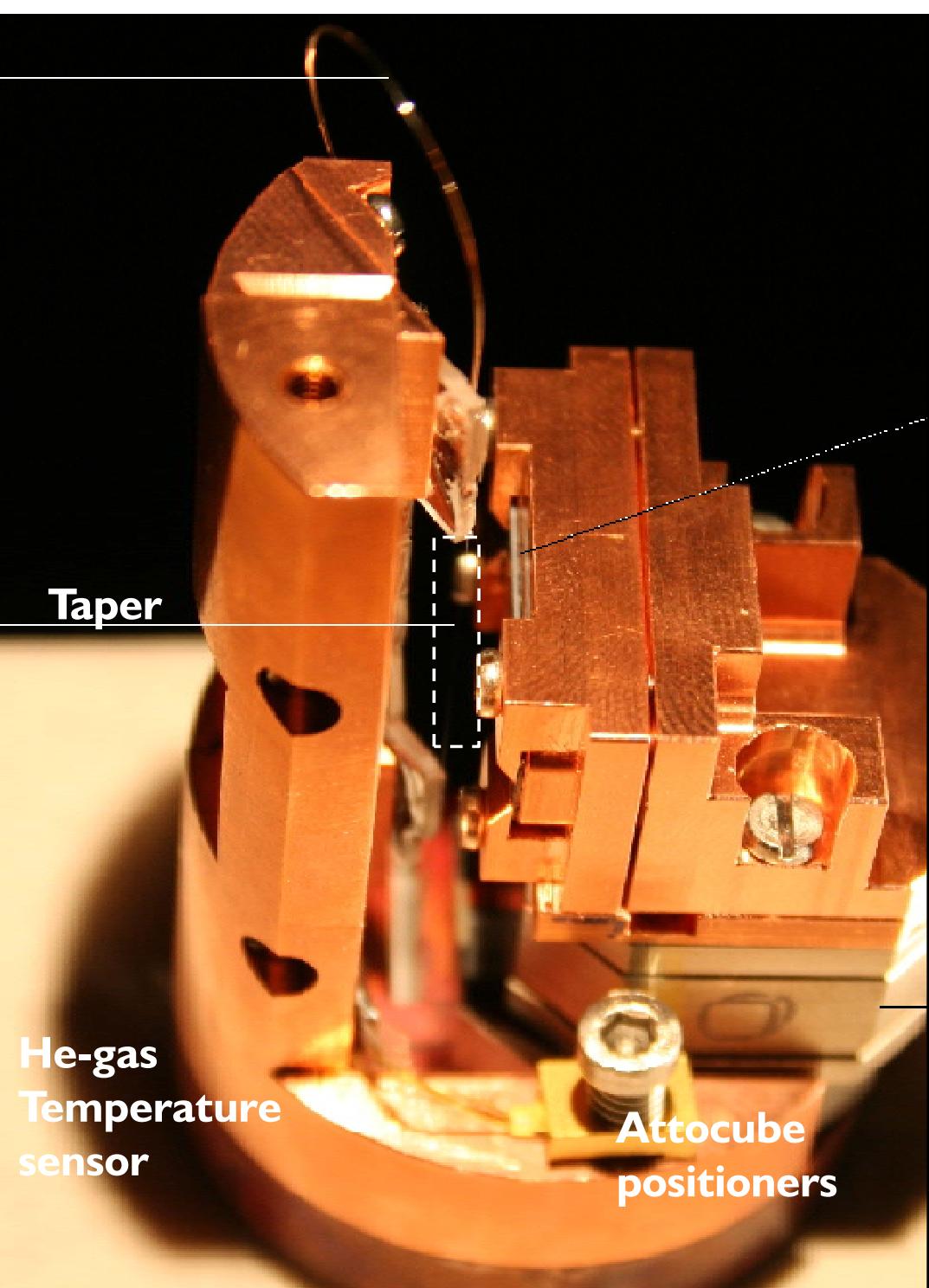
Cryogenics....

$$\nu_i = \frac{i^2 \pi}{2L^2} \sqrt{\frac{Ehw^3}{12\rho A}}$$



Acc.V
2.00 kV

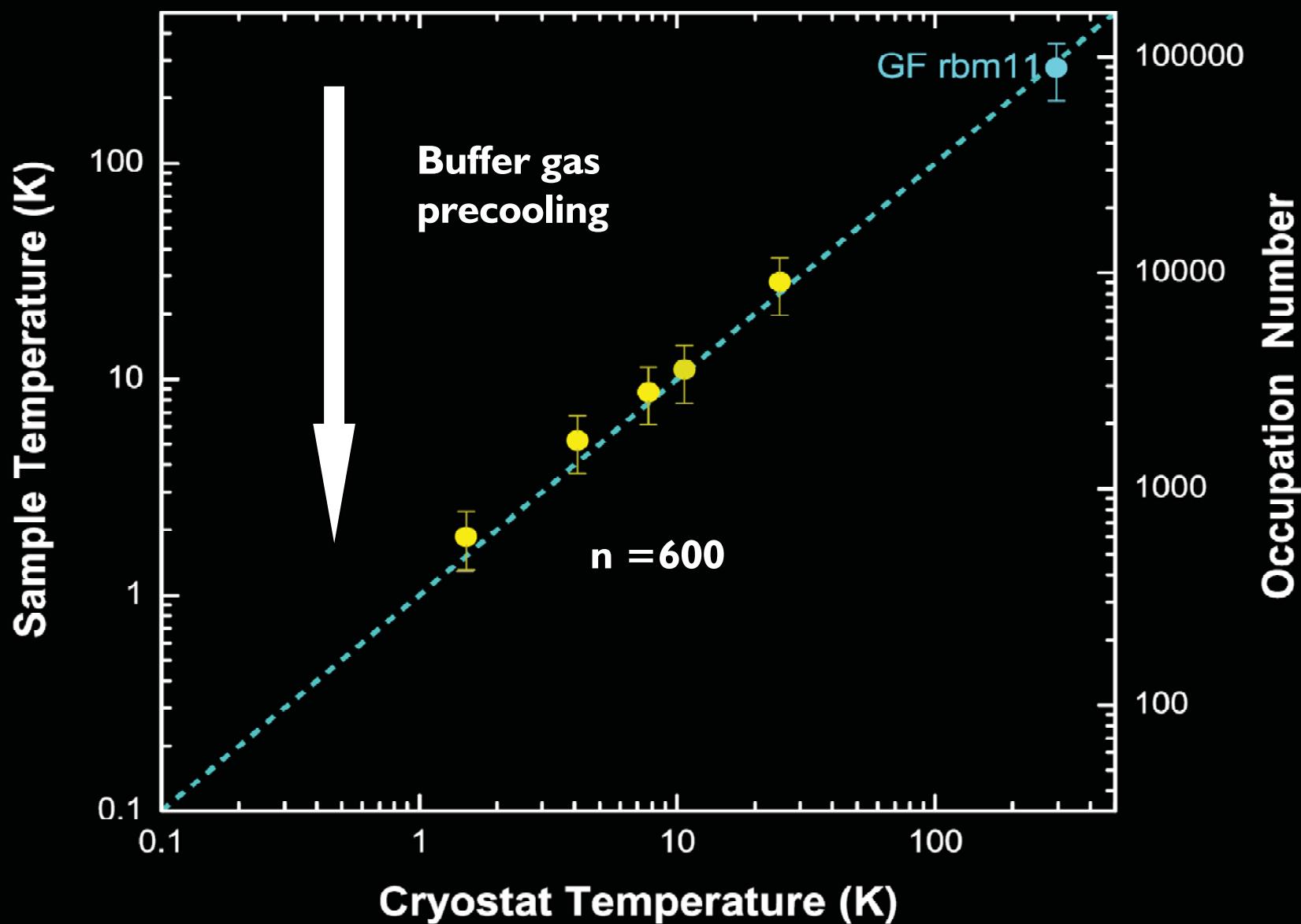
High Finesse and Mechanical Q factor in one and the same device



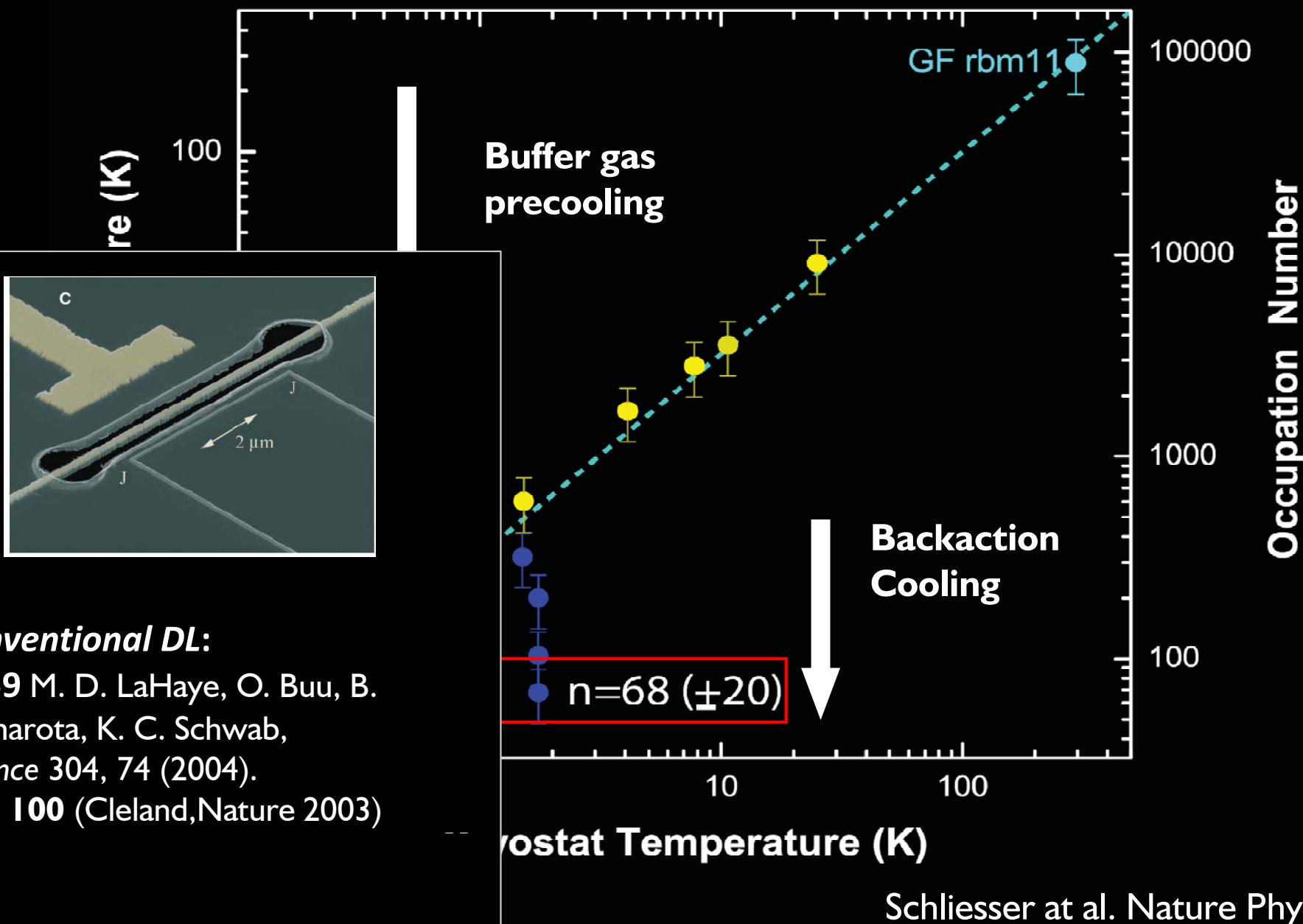
Characteristics He Buffer gas cooling
 $\Omega_m = 62\text{MHz}$ $T = 1.6K$

$$n = \frac{k_B T}{\hbar \Omega} \approx 500$$

Cryogenic Cavity Optomechanics @ MPQ



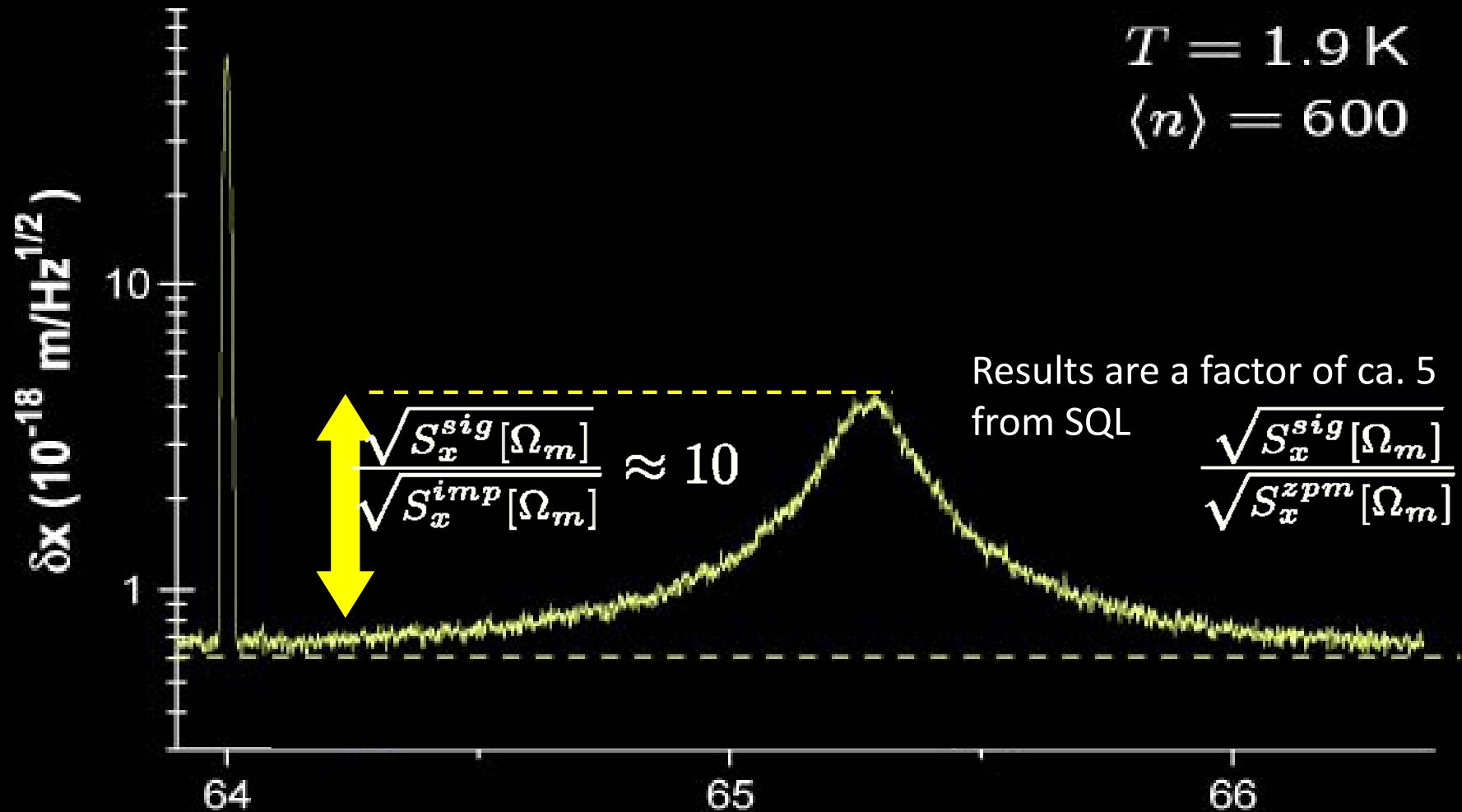
Cryogenic Cavity Optomechanics @ MPQ



Schliesser et al. *Nature Physics* 2009

Approaching the SQL using a cryogenic microresonator

Schliesser, et al. Nature Physics (2009)



-Backaction-Imprecision product only
x100 from Heisenberg uncertainty

limit

$$\sqrt{S_x S_F} \approx \frac{\hbar}{2} \cdot 100$$

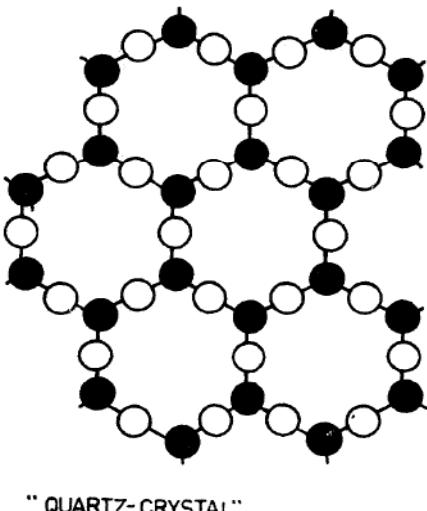
Low Temperatures: Two level fluctuators (TLS)

Anomalous Low-temperature Thermal Properties of Glasses and Spin Glasses

By P. W. ANDERSON†, B. I. HALPERIN and C. M. VARMA

Bell Laboratories, Murray Hill, New Jersey 07974

Crystalline structure

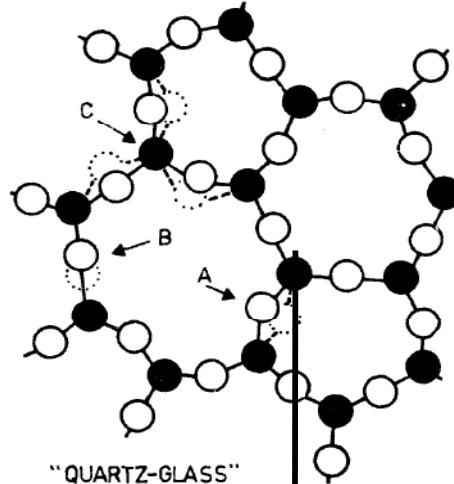


2-level system approximation [I]

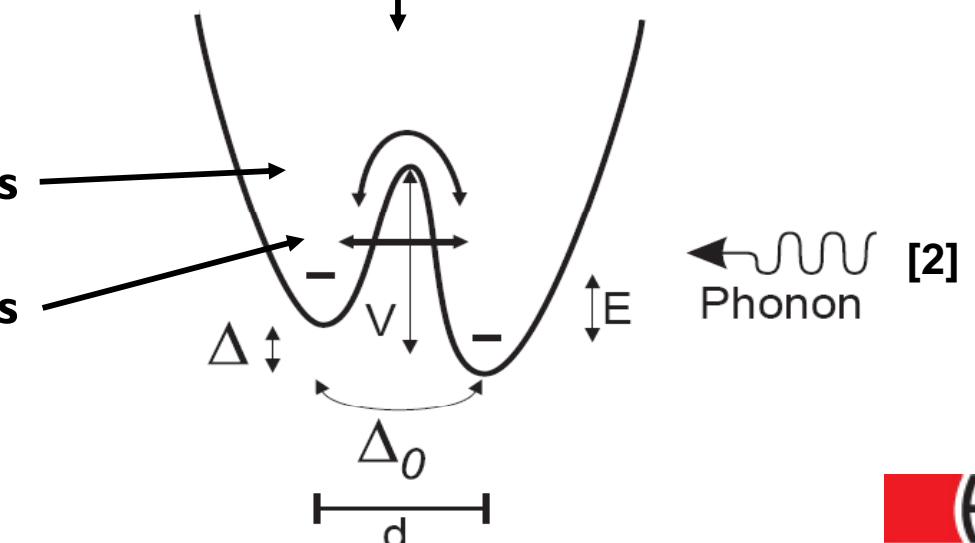
Thermally activated process

Tunelling process

[2]



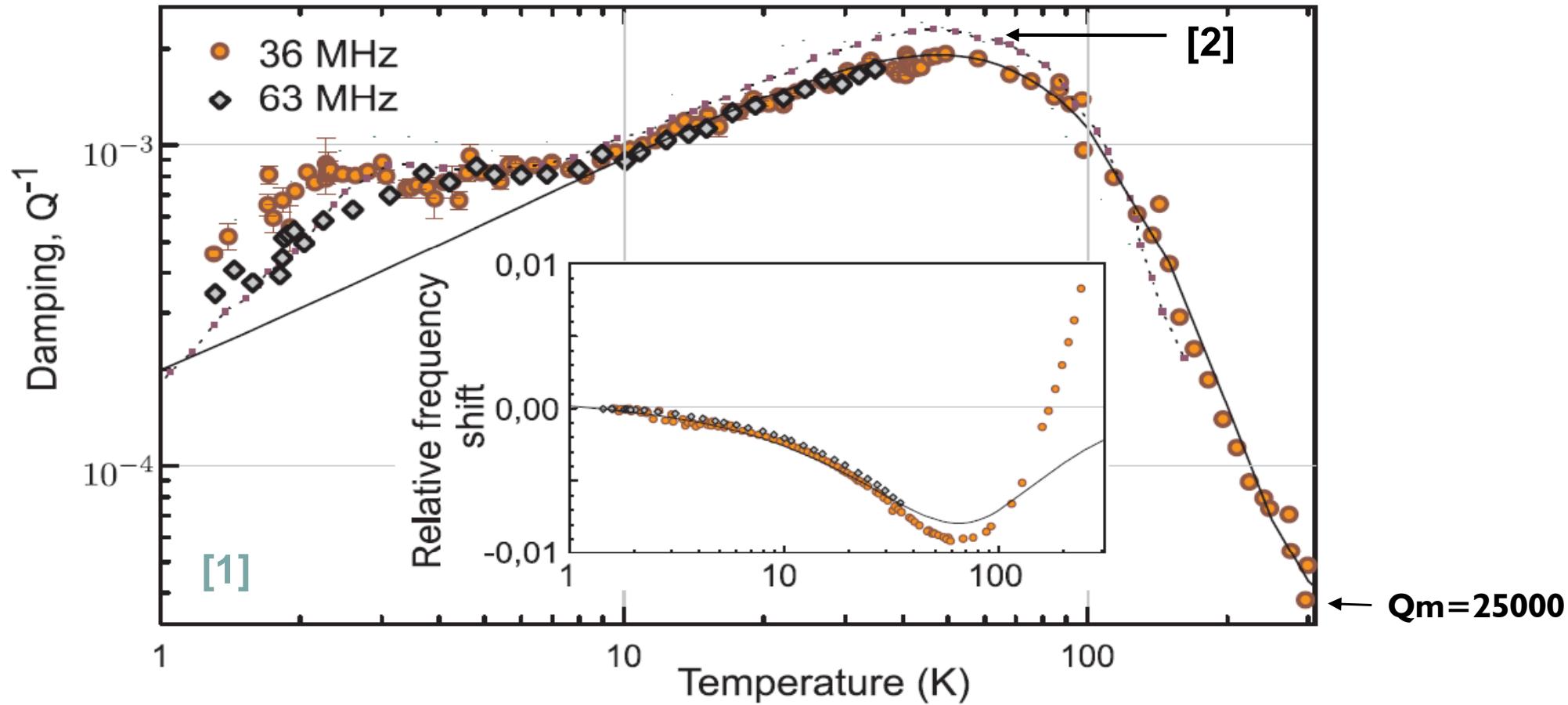
**Glassy structure:
Different possible
conformations**



[1] Vacher, Courtens, Forêt, PRB 72 214205 (2005)

[2] Jäckle, Piché, Huncklinger, J. Non-Crys. Sol. 20 365 (1976)

Low Temperatures: Two level fluctuators (TLS)



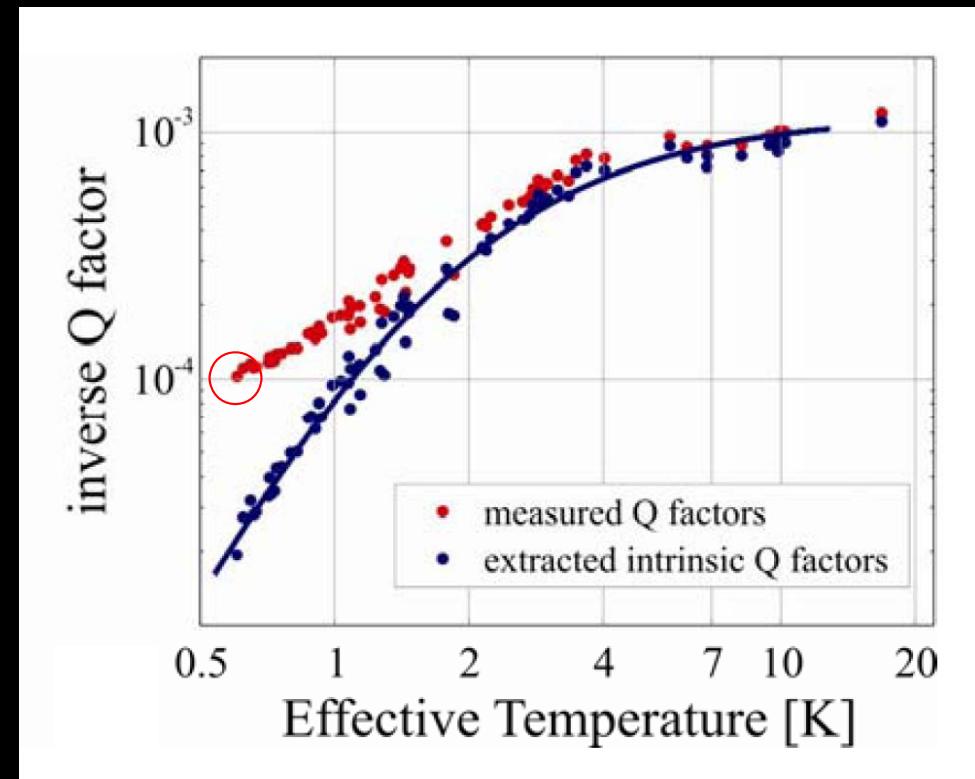
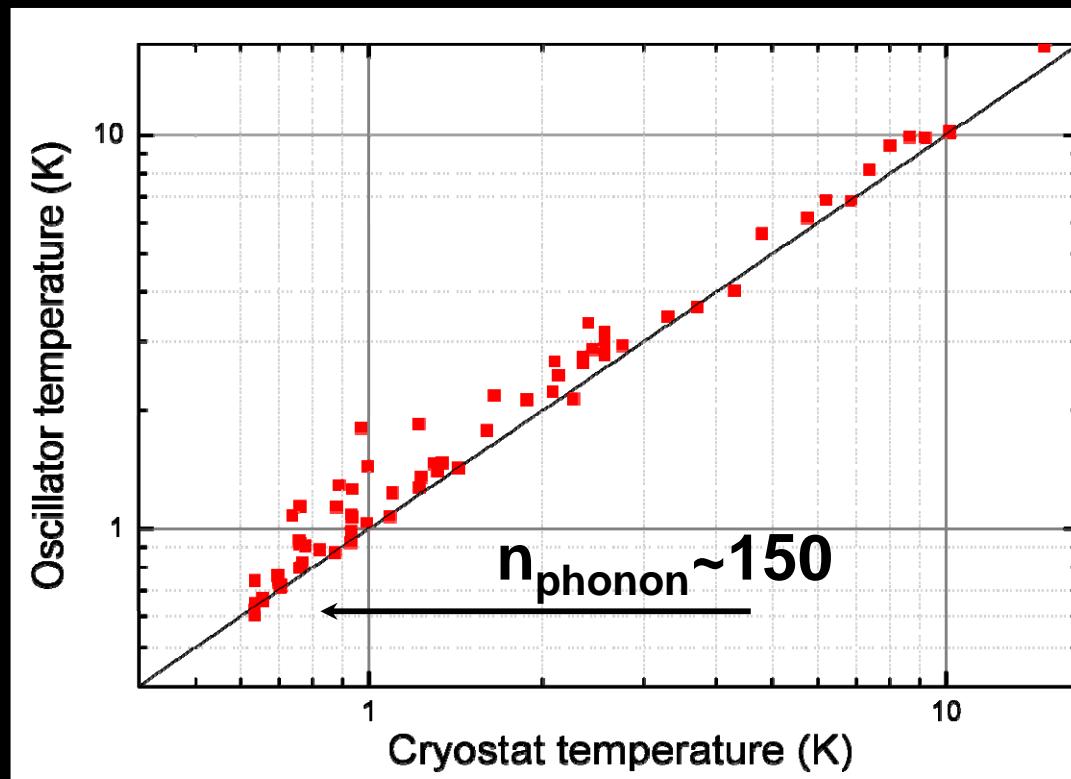
$$Q^{-1} = \mathcal{C} \operatorname{Erf} \left(\frac{\sqrt{2}T}{\Delta_C} \right) \frac{1}{T} \int_0^\infty \left(\frac{V}{V_0} \right)^{-\xi} e^{-\frac{1}{2} \frac{V^2}{V_0^2}} \frac{\Omega \tau_0 e^{V/T}}{1 + \Omega^2 \tau_0^2 e^{2V/T}} dV \quad [3]$$

[1] Arcizet, Rivière, Schliesser, Anetsberger, Kippenberg, PRA (2009)

[2] U. Bartell et al., J. Phys. (Paris) Colloq. 43, C9 (1982) II

[3] Vacher, Courtens, Forêt, PRB 72, 214205 (2005)

Optomechanics at Helium-3 Temperatures (600 mK)



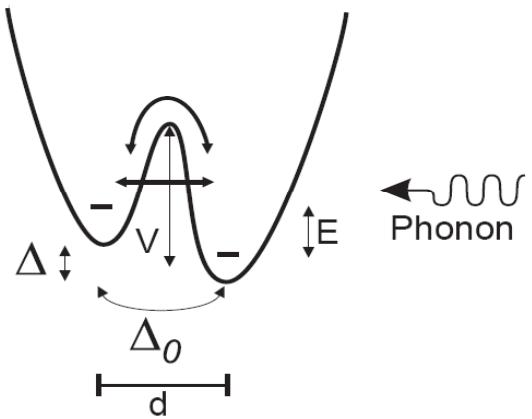
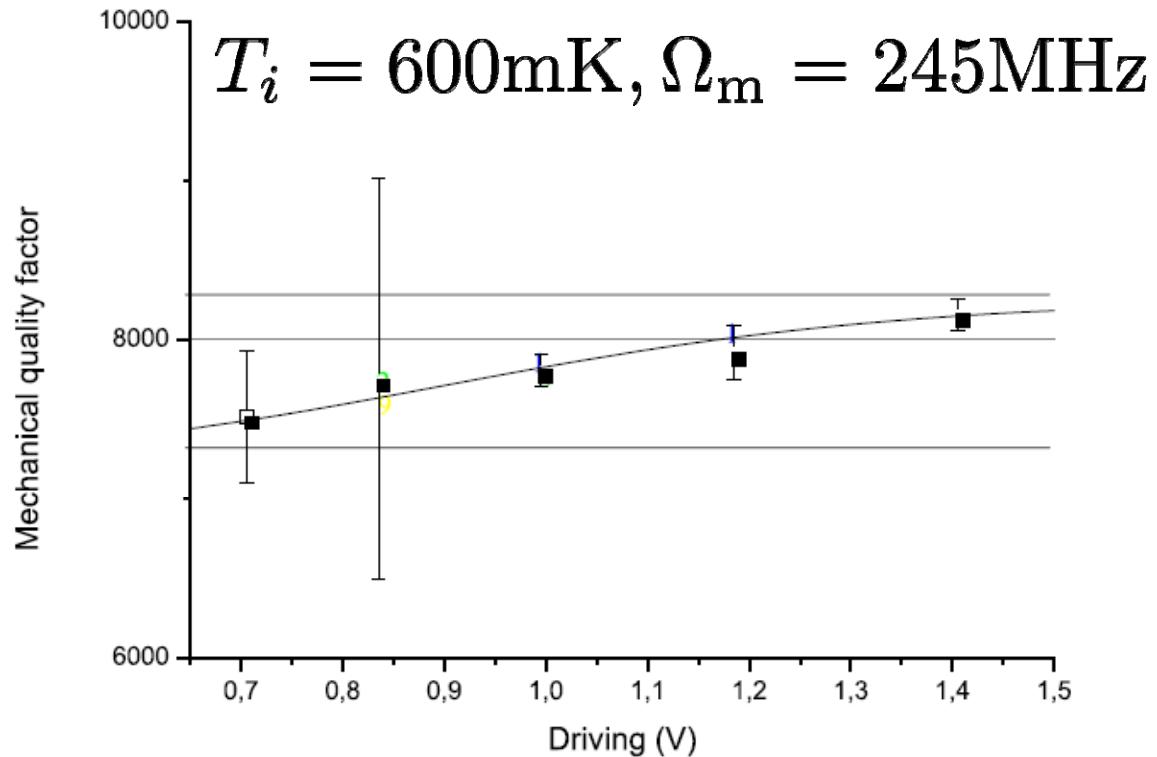
$$T_i = 600 \text{mK} : Q_m > 10^4, \Omega_m = 75 \text{MHz}$$

$$n_i = \frac{k_B T_i}{\hbar \Omega_m} \approx 150$$

Experiments fulfills are prerequisites to achieve cooling to below $n=10$ and possess imprecision close to SQL.
(O. Arcizet, S. Weis, R. Riviere unpublished)

Evidence for direct phonon absorption (Quantum friction)

2nd Order radial breathing mode



Observation: Driving of mechanical oscillator (pm amplitudes)
increases the mechanical Q factor

Explanation: Saturation of Two Level Systems due to mechanical
oscillator also termed “Quantum friction”

L. G. Remus, M. P. Blencowe, Y. Tanaka Phys. Rev. B 80, 174103 (2009)

Scientific Goals in cavity Optomechanics

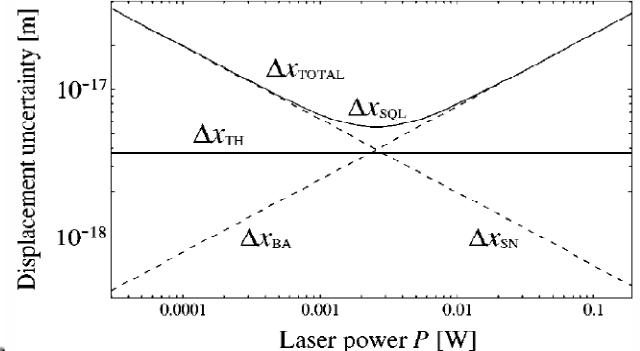
Quantum Backaction – how Nature enforces the Heisenberg uncertainty

Quantum-mechanical noise in an interferometer

Carlton M. Caves

W. K. Kellogg Radiation Laboratory, California Institute of Technology, Pasadena, California

(Received 15 August 1980)



Ground-State Cooling – true quantum *mechanics* of a macroscopic object

Analog to: Diedrich et al. PRL, 1989....

Laser Cooling to the Zero-Point Energy of Motion

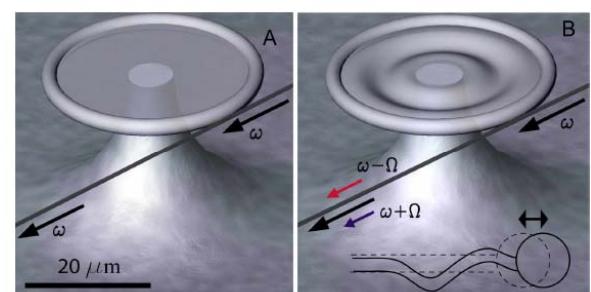
F. Diedrich,^(a) J. C. Bergquist, Wayne M. Itano, and D. J. Wineland

Time and Frequency Division, National Institute of Standards and Technology, Boulder, Colorado 80303

(Received 28 July 1988)

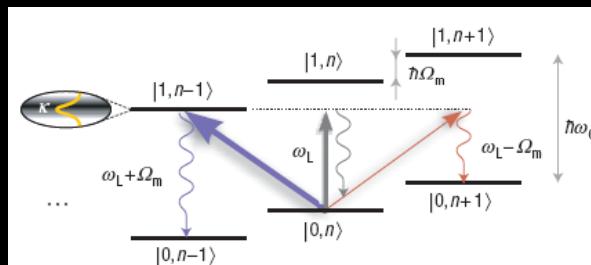
Applied Side: Radiation Pressure Driven Quartz Oscillators–
Novel form of photonic oscillators for metrology and
timekeeping

..



Conclusions

Sideband Cooling

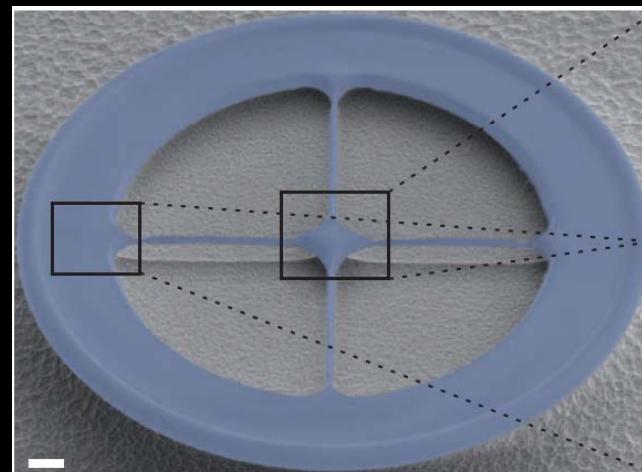


PRL, Dec. 2008

Nature Physics 4, 415 - 419 (2008)

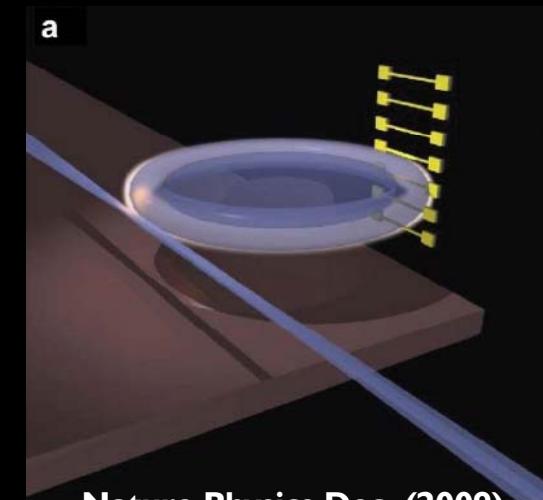
Nature Physics, 5, 509-514, (2009)

Low dissipation optomechanics



Nature Photonics 2, 627 (2008)

SQL for Nanomechanics



Nature Physics Dec. (2009)

Developed crystalline resonators



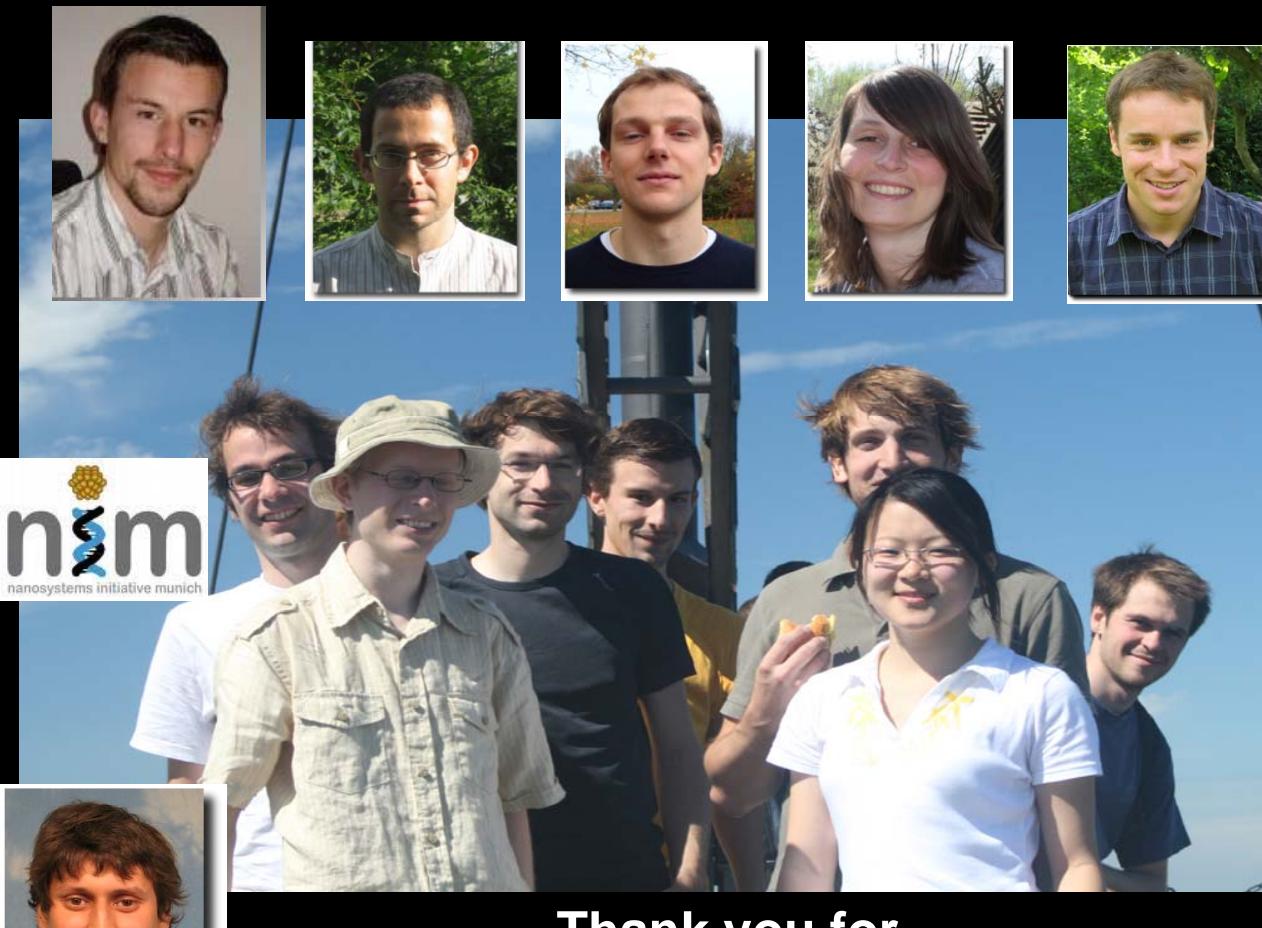
Future Directions...

Ground state cooling of a mechanical oscillator

Measurement of Zero Point Motion

Study of Quantum Friction

Acknowledgements



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your attention
www.mpq.mpg.de/k-lab

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(LMU, clean room access)
- **F. Bürsgens**
(LMU, clean room access)
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(MPQ, cryostat support)
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(EDFA and ESA support)



MP-IJRG

ERC StG

MC-EXT

GSC



MC-IRG/IEF

NanoSci-ERA