

Synthesizing arbitrary photon states in a superconducting resonator

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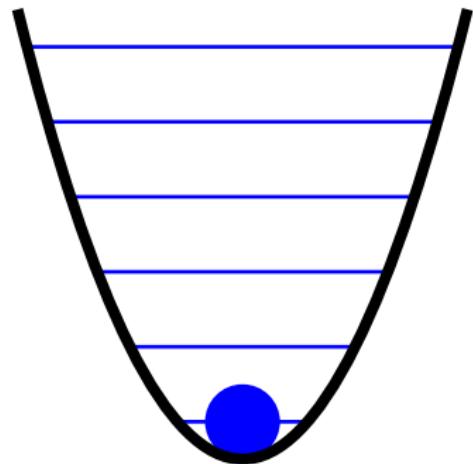
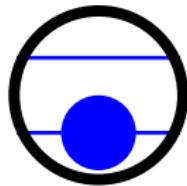


Collège de France, March 1 2010



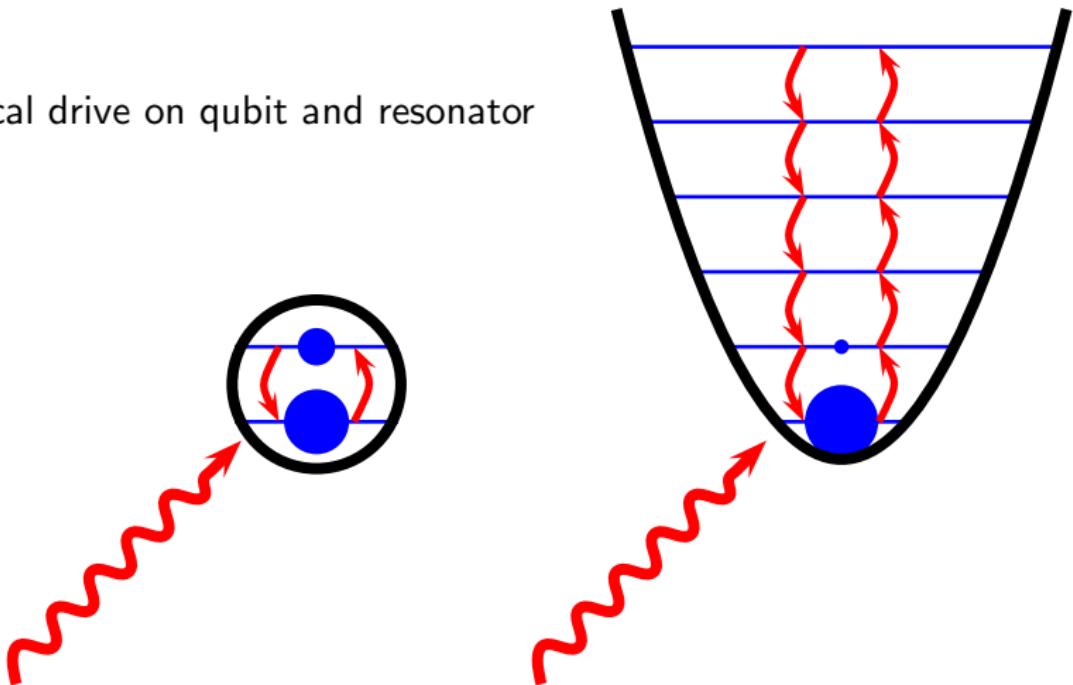
Generation of quantum states in a resonator

Classical drive on qubit and resonator



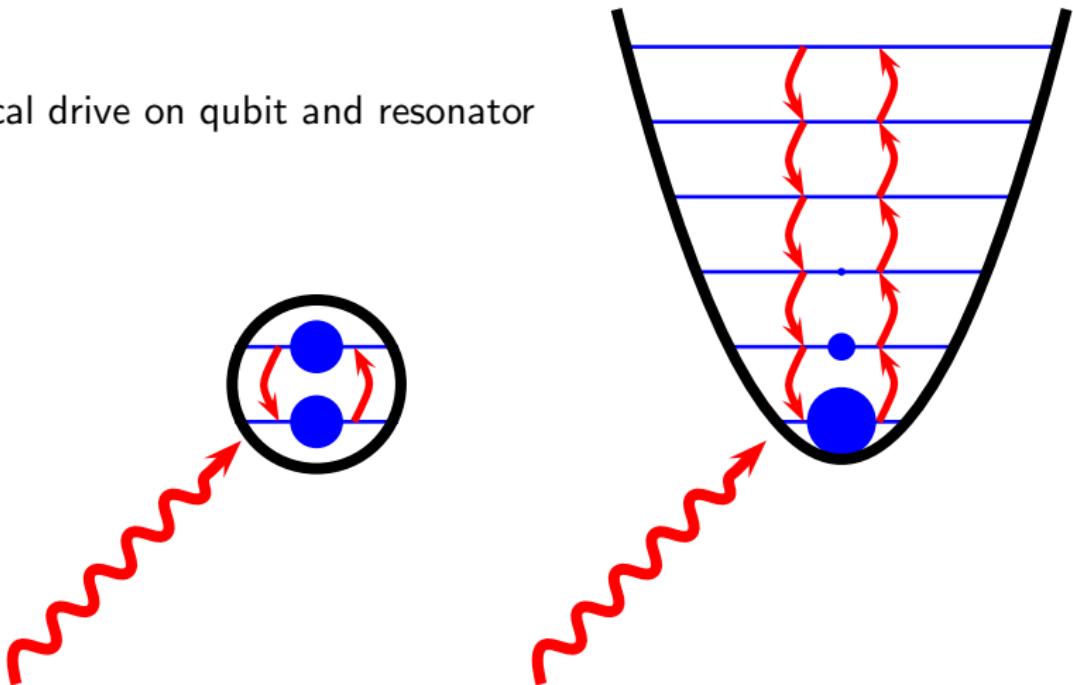
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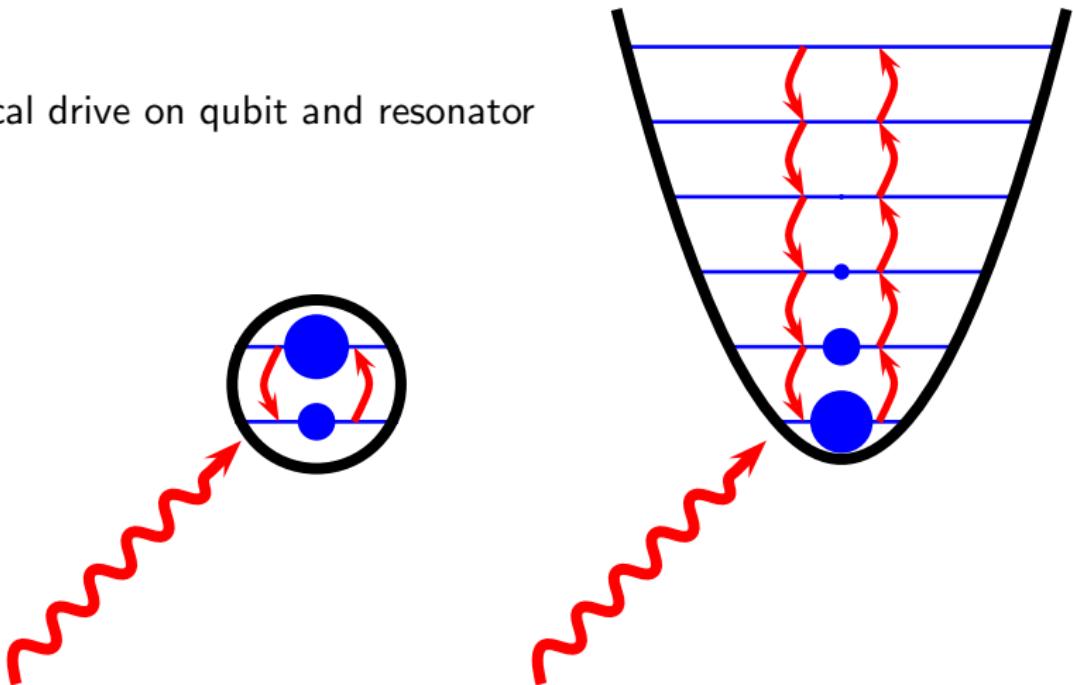
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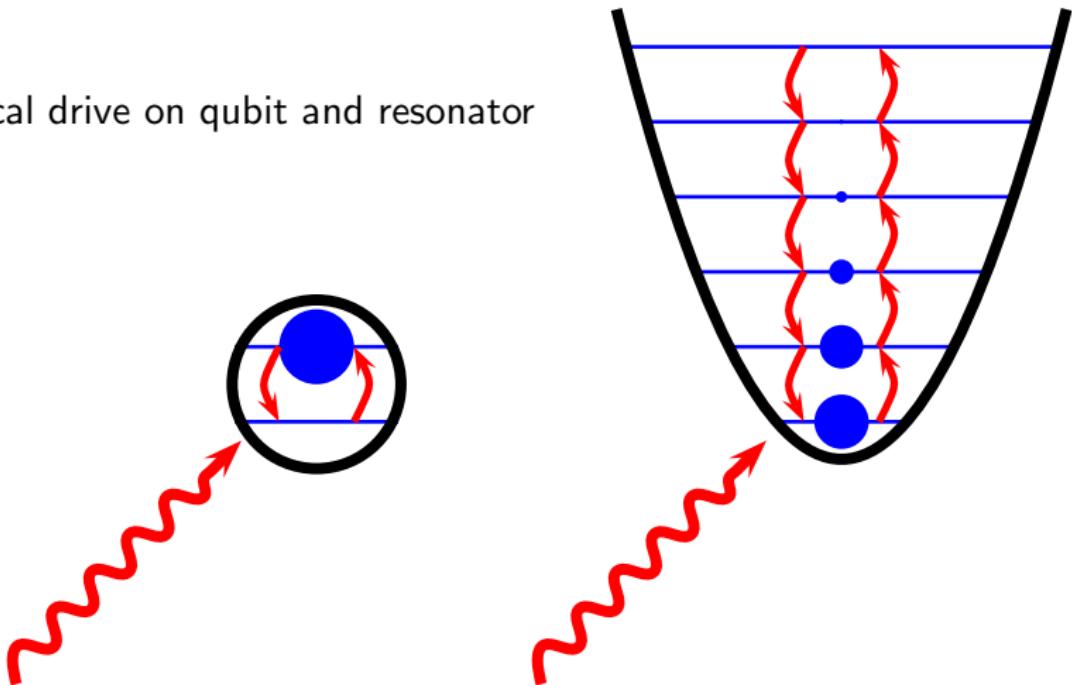
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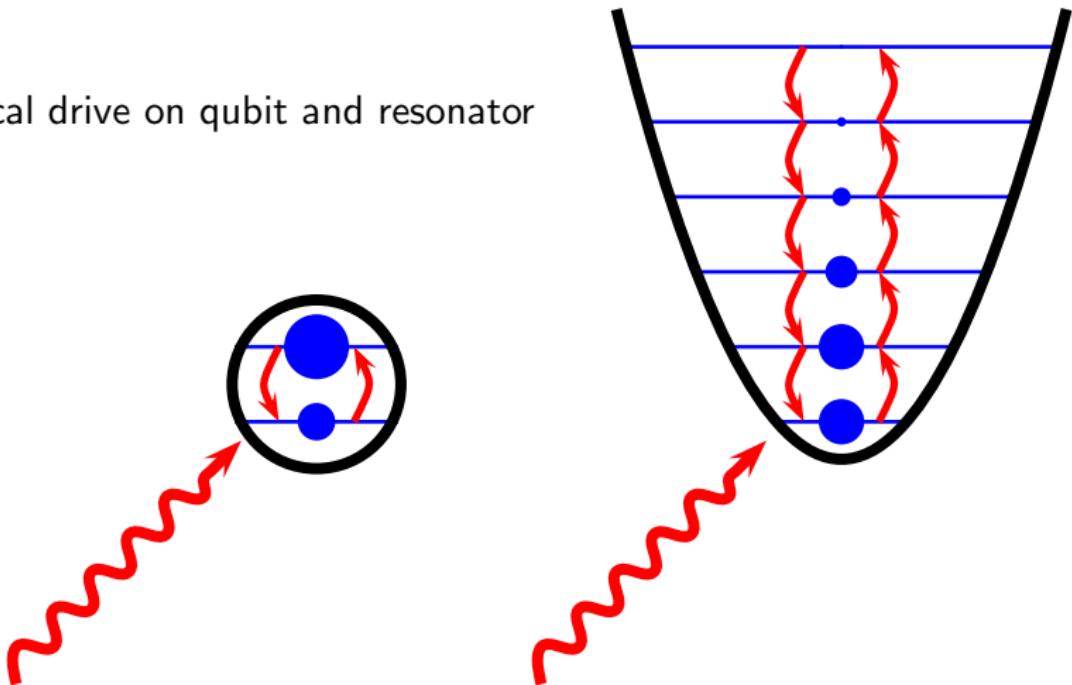
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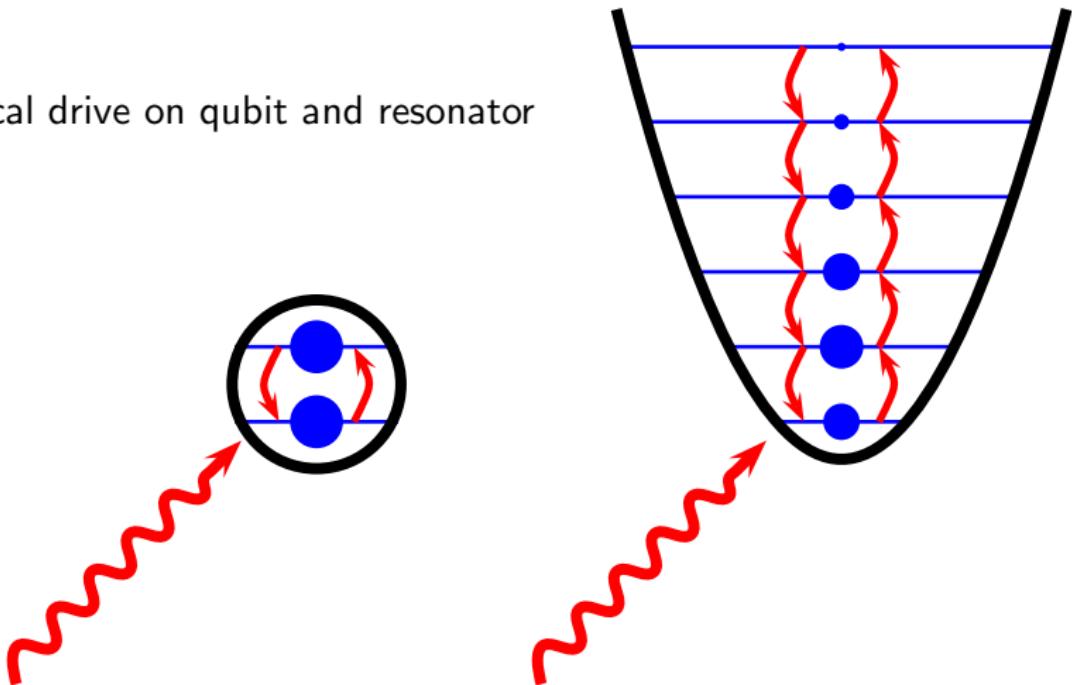
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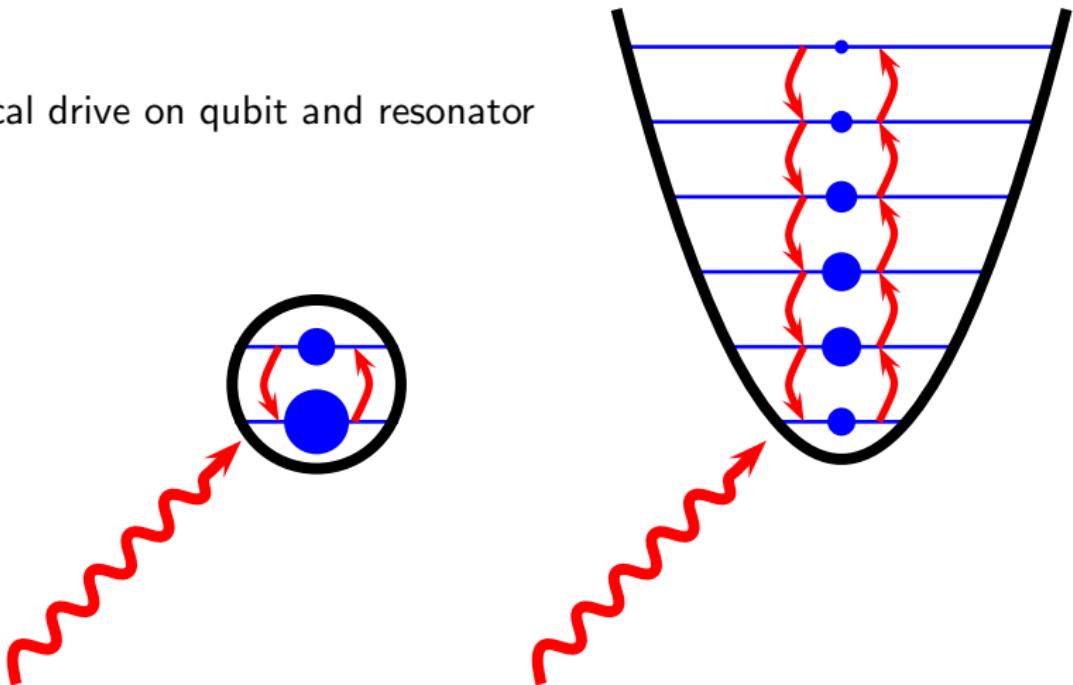
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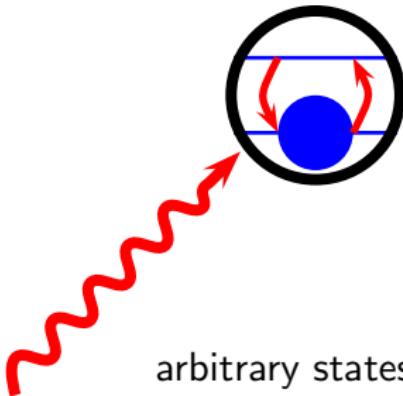
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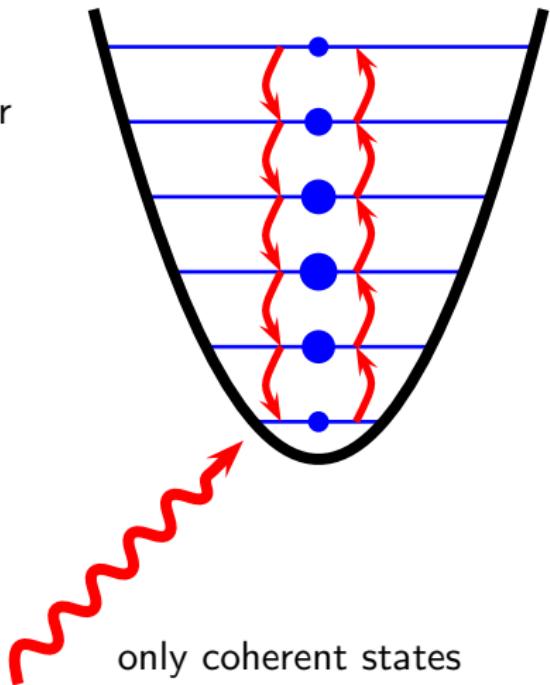


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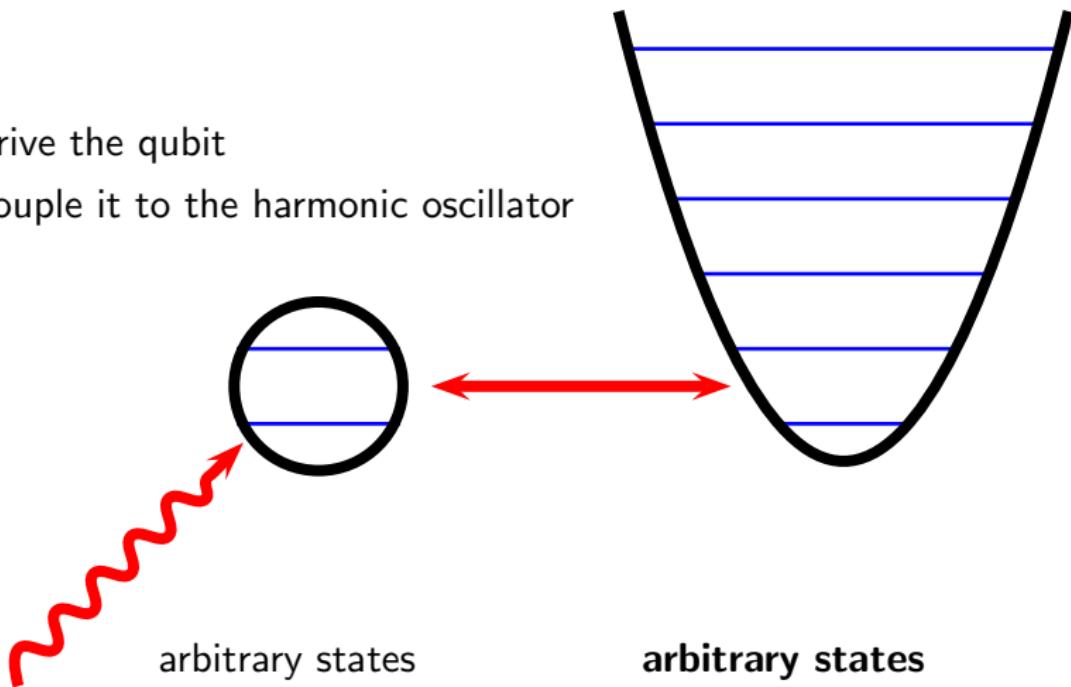
arbitrary states



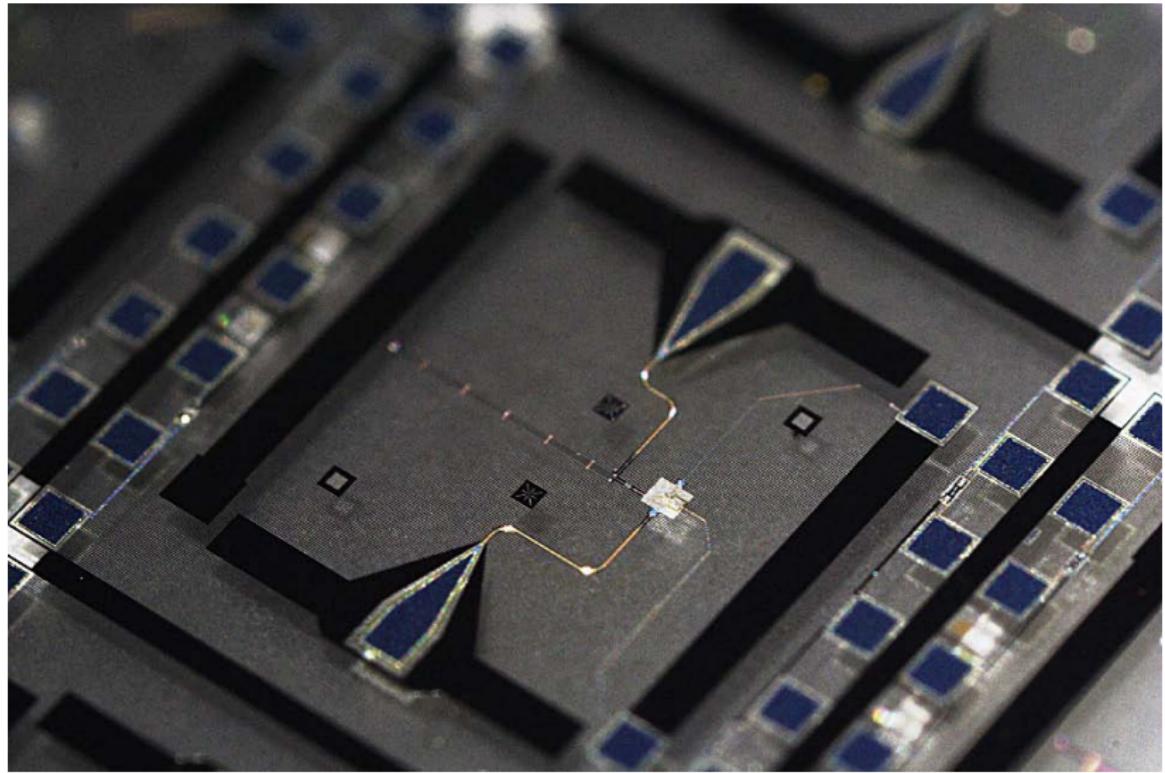
only coherent states

Generation of quantum states in a resonator

- ▶ drive the qubit
- ▶ couple it to the harmonic oscillator

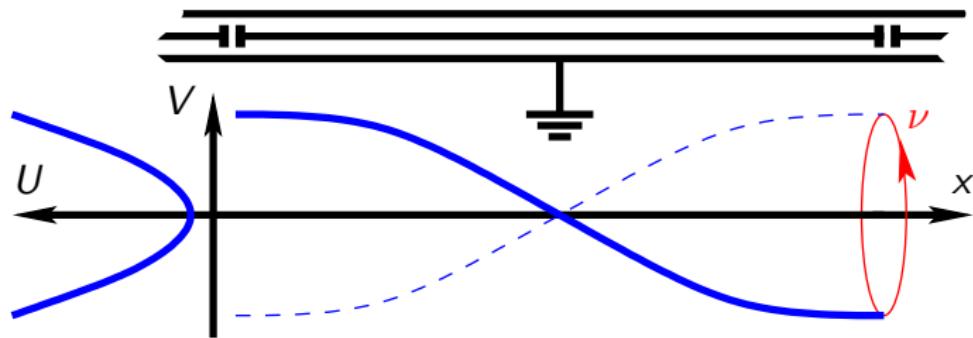


Device



die size: 6×6 mm

Coplanar waveguide resonator

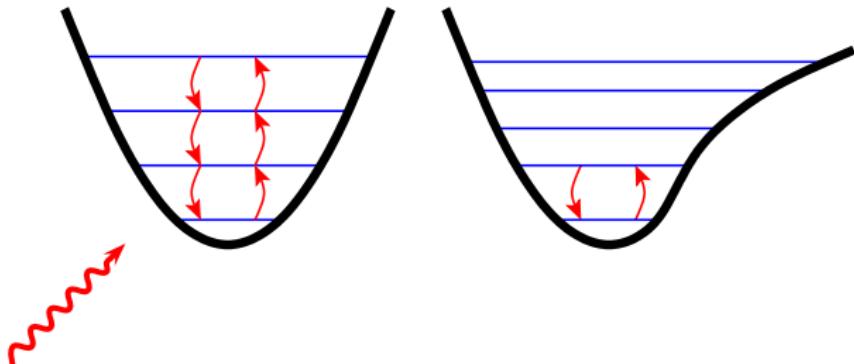


- ▶ very small capacitors are open ends
- ▶ each mode forms a quantum harmonic oscillator
- ▶ we use fundamental mode at frequency $\nu = \frac{c_{\text{eff}}}{2L}$
- ▶ make out of superconductor for high quality factor Q

Performance:

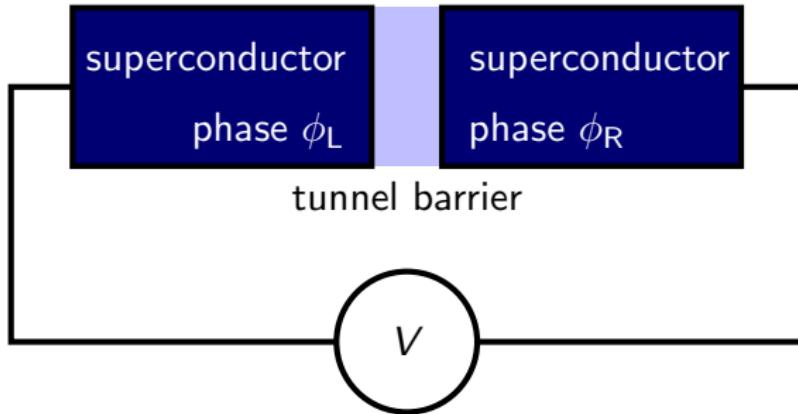
- ▶ $Q \sim 1 \dots 5 \cdot 10^5$ or $T_1 \sim 2.5 \dots 10 \mu\text{s}$

Superconducting qubits



- ▶ 2 level system hard to make in a circuit
- ▶ very strong nonlinearity is enough
- ▶ allows to address 0—1 transition exclusively
- ▶ we use Josephson junction as nonlinear circuit element

Josephson junction



Dynamics are determined by:

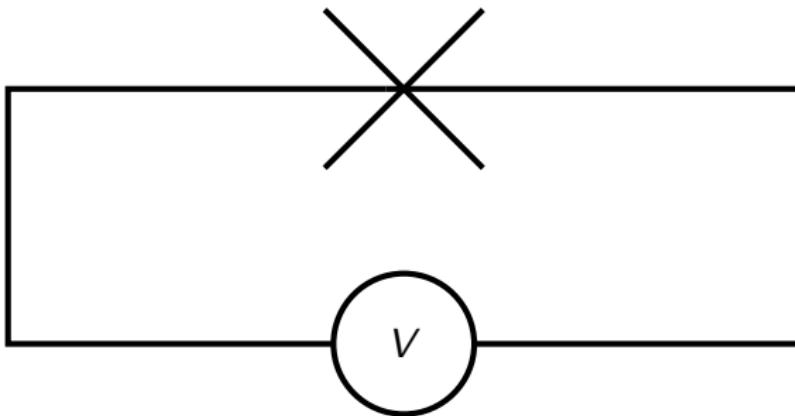
$$I = I_0 \sin \delta \quad \text{DC Josephson relation}$$

$$V = \frac{\hbar}{2e} \frac{\partial \delta}{\partial t} \quad \text{AC Josephson relation}$$

$\delta = \phi_L - \phi_R$ phase difference

I_0 critical current of junction

Josephson junction



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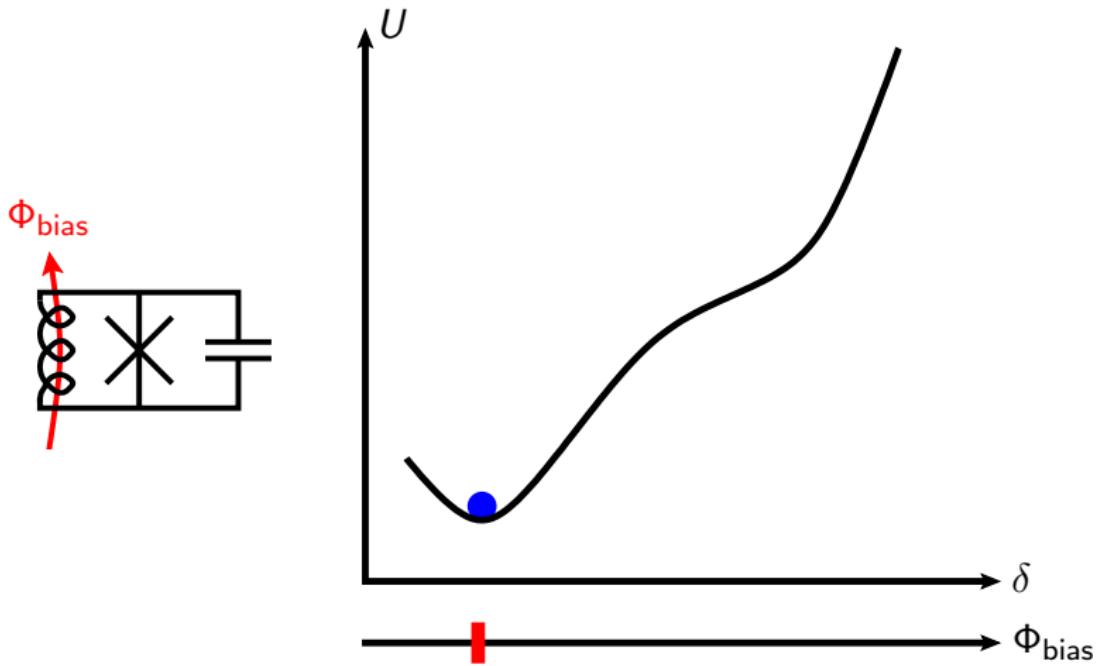
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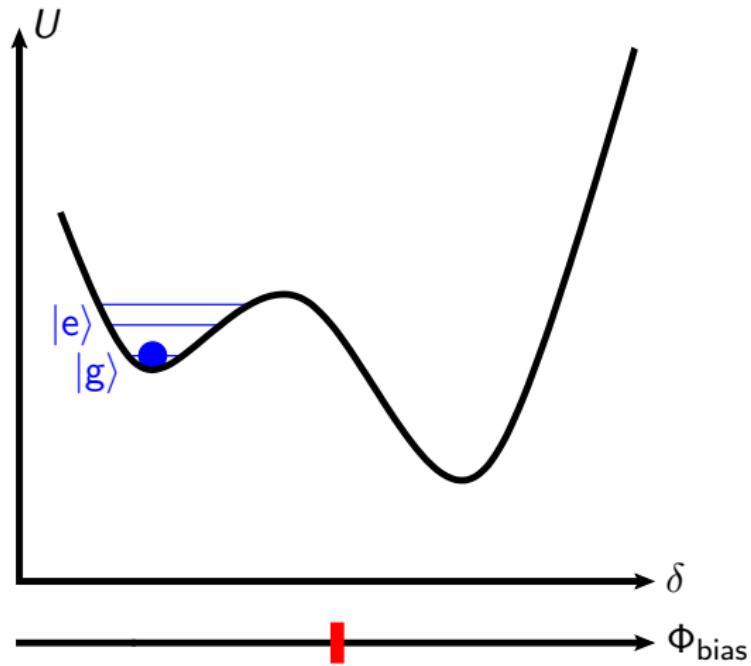
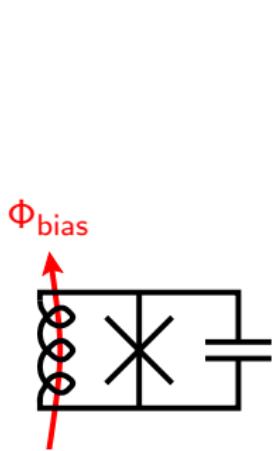
Superconducting phase qubit



$$H = \underbrace{\frac{\left(\frac{\hbar}{2e}\delta - \Phi_{\text{bias}}\right)^2}{2L} - \frac{\hbar}{2e}I_0 \cos \delta}_{U} + \frac{Q^2}{2C}$$

$$[\delta, Q] = 2ei$$

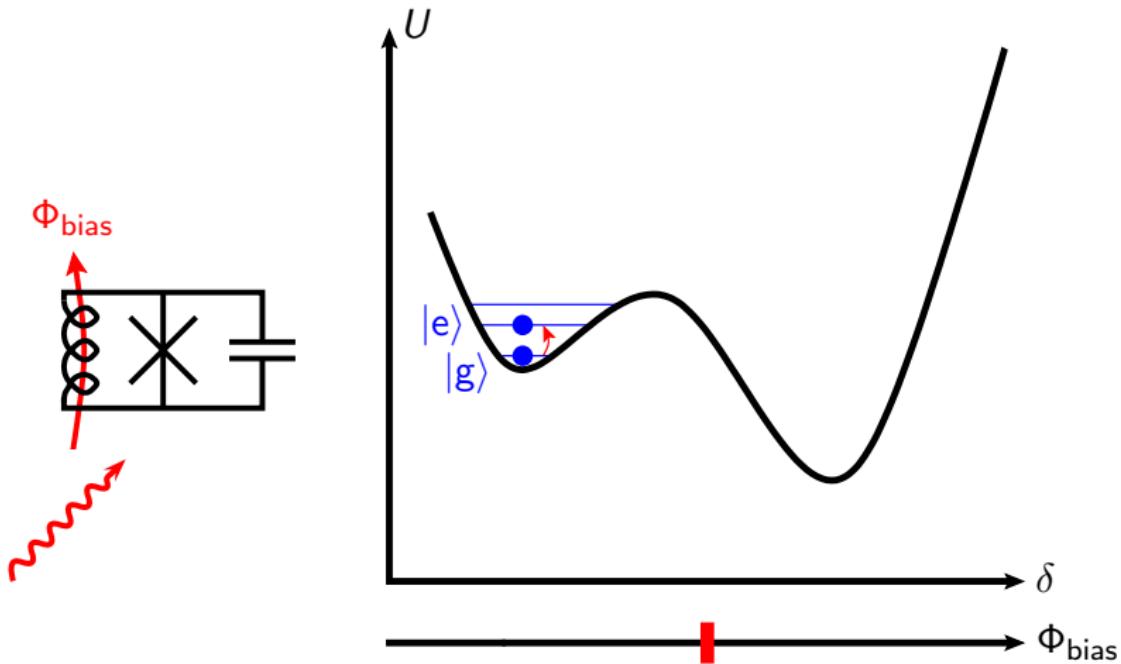
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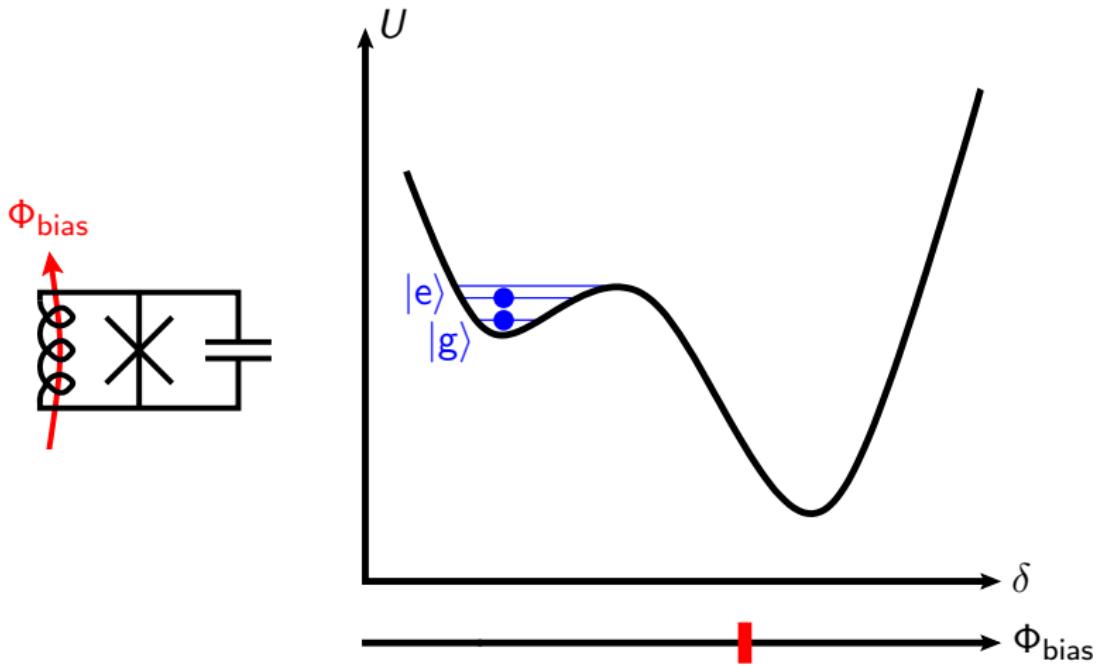
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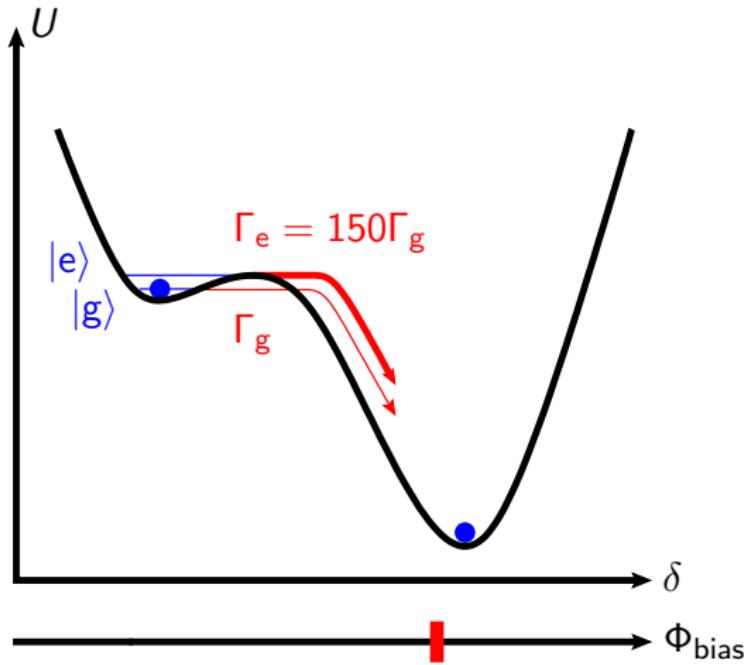
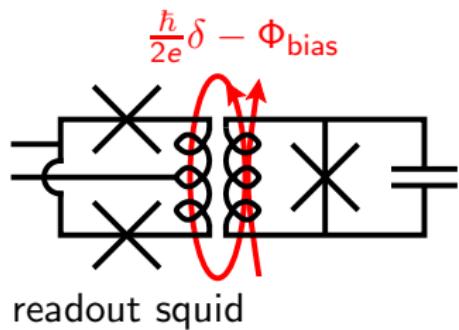
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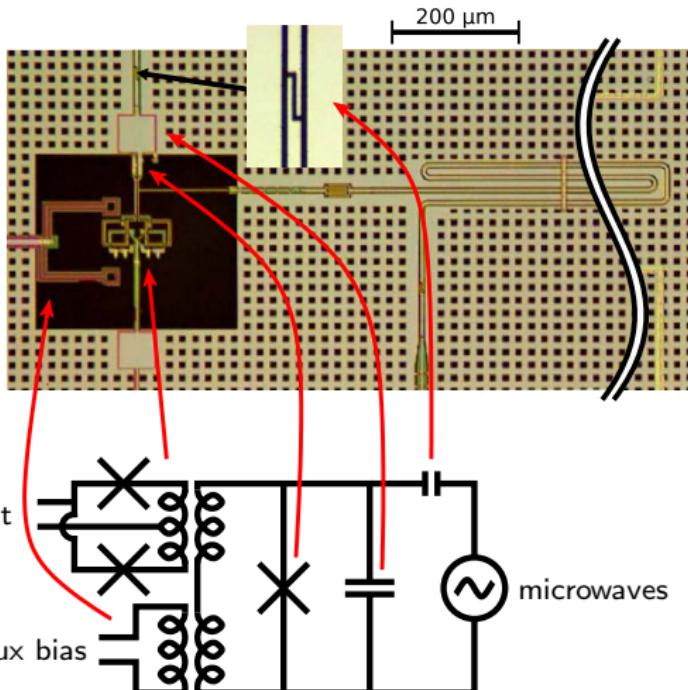
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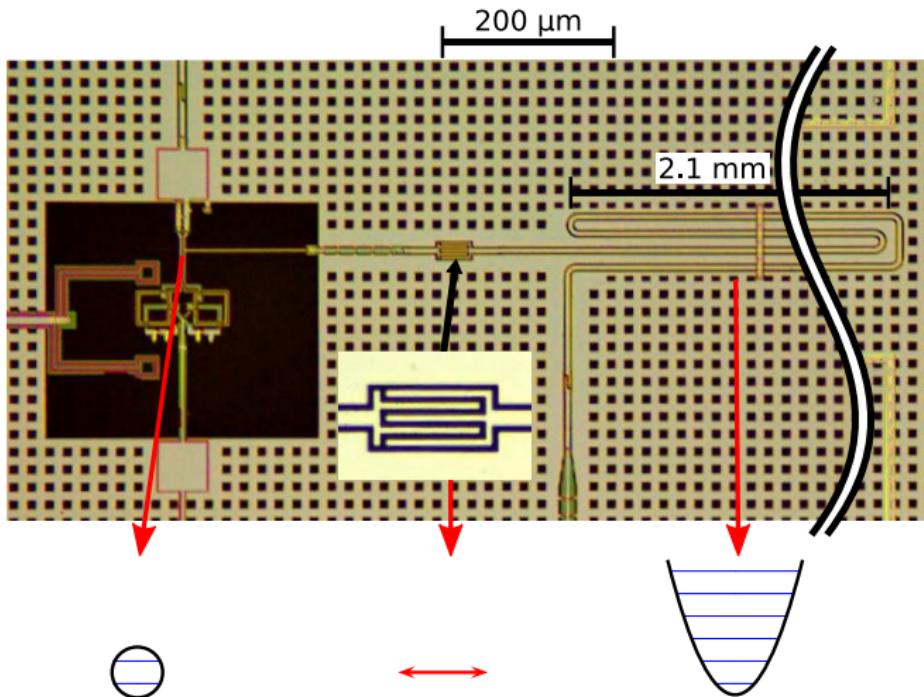
$$[\delta, Q] = 2ei$$

UCSB phase qubit



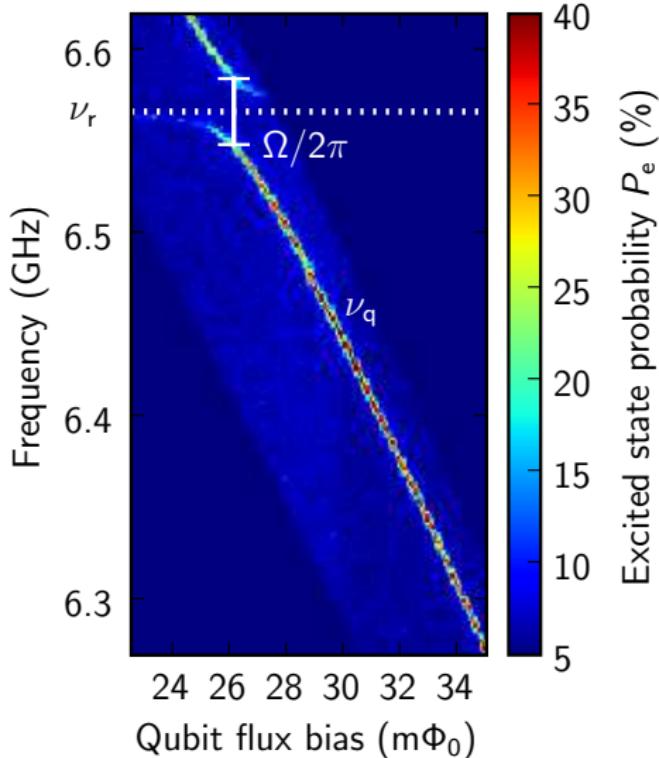
- ▶ energy relaxation time
 $T_1 \approx 600 \text{ ns}$
- ▶ quality factor
 $Q \approx 2 \cdot 10^4$
- ▶ phase coherence time
 $T_2 \approx 150 \text{ ns}$
- ▶ nonlinearity
 $\Delta \approx 200 \text{ MHz}$
limits gate time $\gtrapprox 10 \text{ ns}$
- ▶ readout visibility
 $1 - P_{e'|g} - P_{g'|e} \approx 0.90$

Coupling qubit and resonator



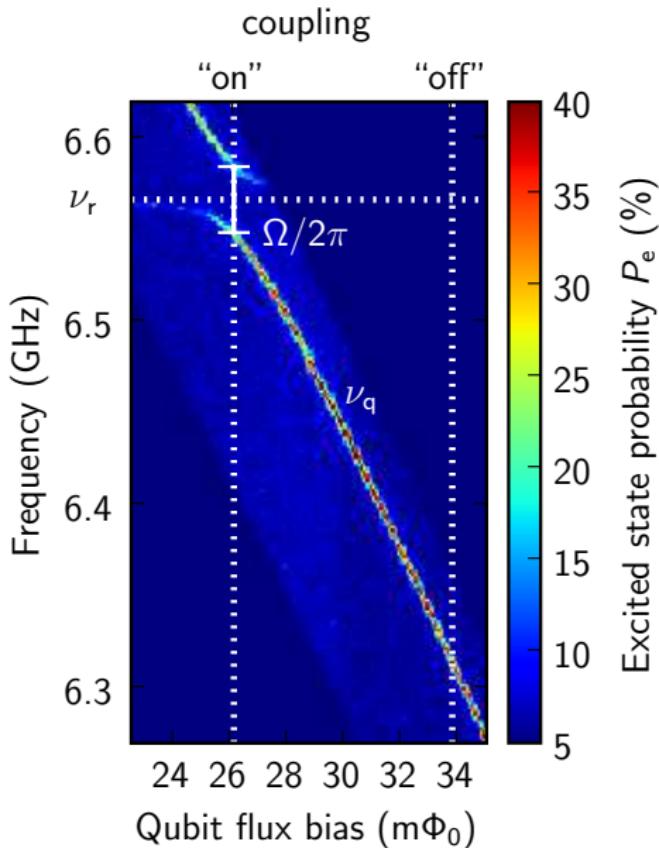
$$H = h\nu_q \sigma^+ \sigma^- + \frac{\hbar\Omega}{2} (a\sigma_+ + a^\dagger\sigma_-) + h\nu_r \left(a^\dagger a + \frac{1}{2} \right)$$

Spectroscopy of qubit-resonator coupling



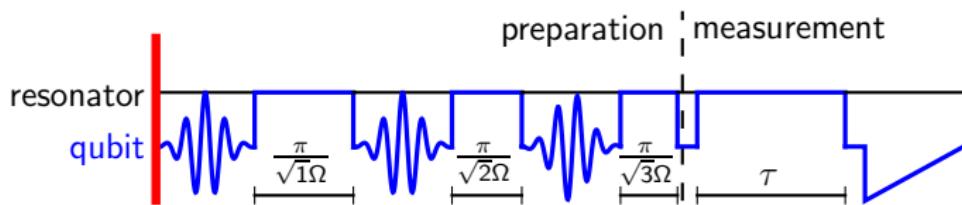
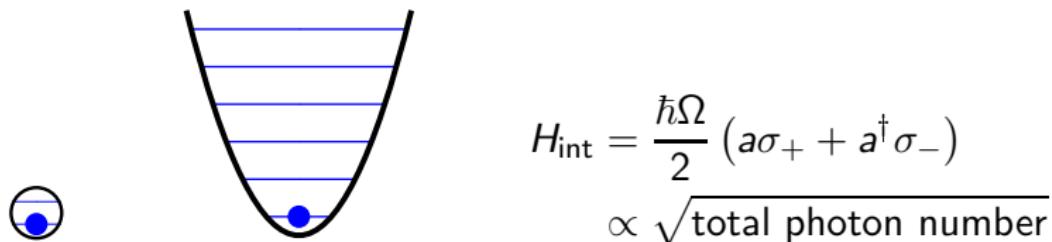
- ▶ high probability indicates eigenfrequencies of qubit-resonator system

Spectroscopy of qubit-resonator coupling



- ▶ high probability indicates eigenfrequencies of qubit-resonator system
- ▶ Qubit and resonator only interact when $\nu_q \approx \nu_r$
- ▶ Coupling can be turned on and off

Generating Fock states: Pumping photons one by one



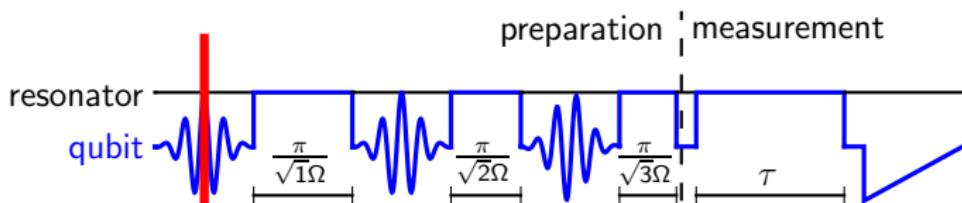
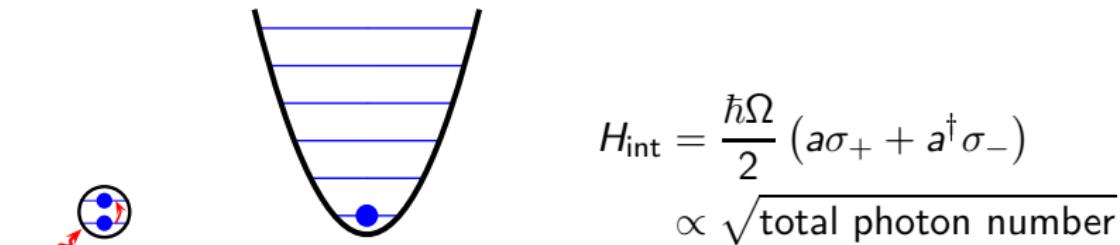
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π -pulse: > 98 % fidelity

E. Lucero *et al.* PRL **100**, 247001 (2008)

swap-pulse: > 98 % fidelity

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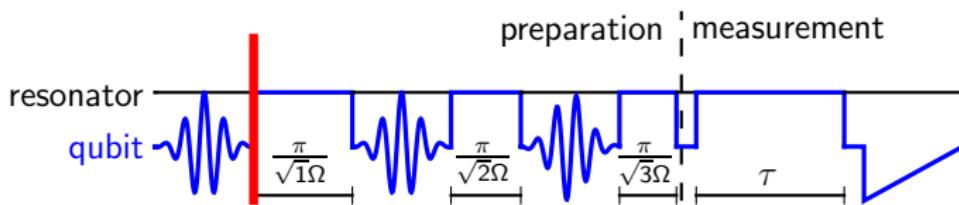
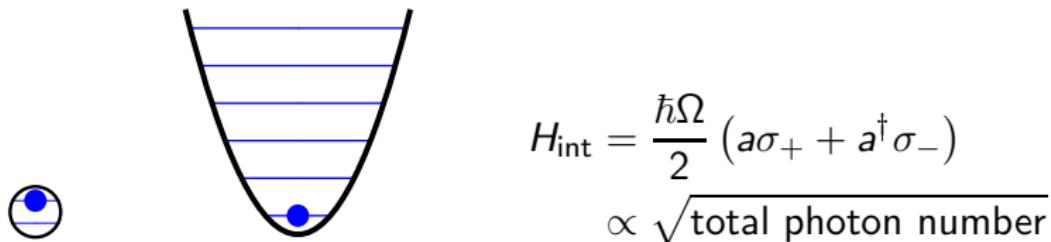
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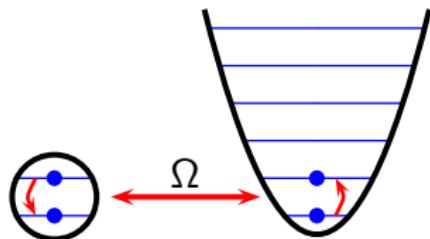
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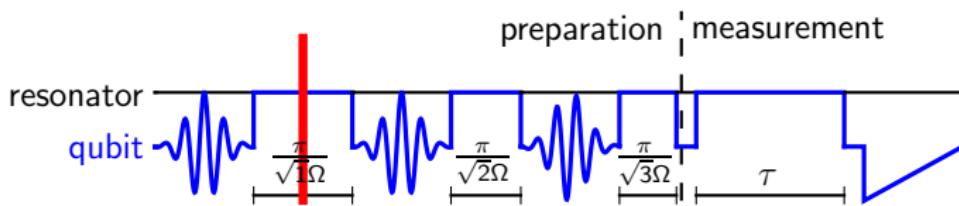
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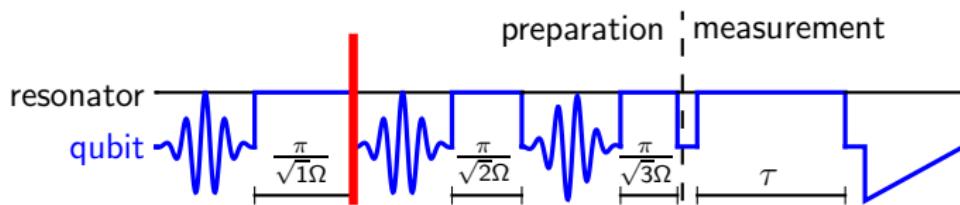
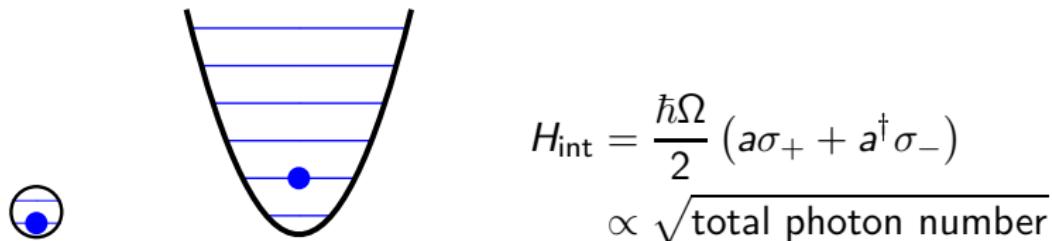
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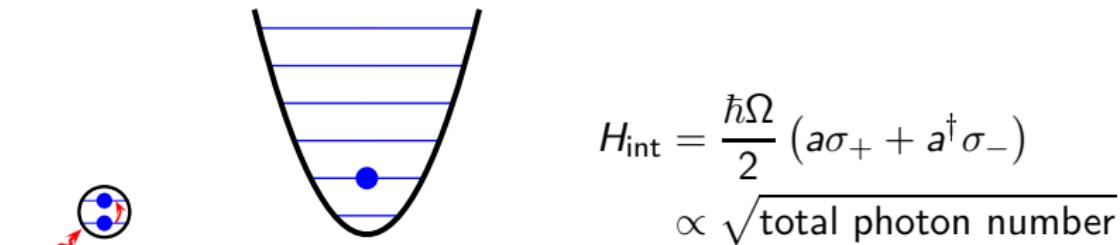
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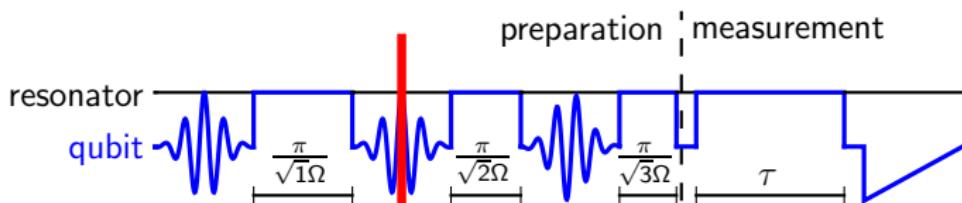
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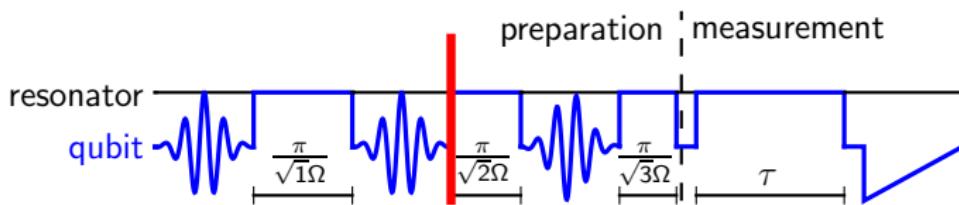
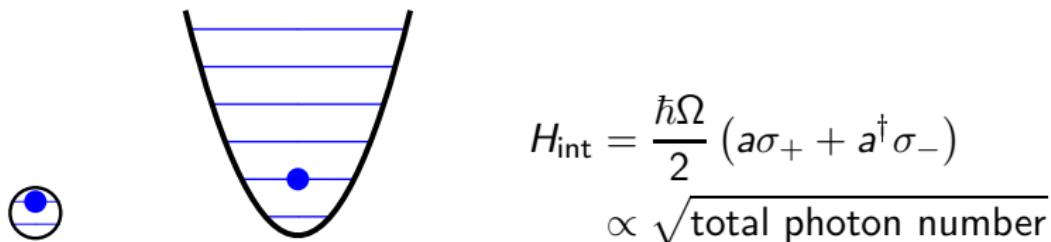
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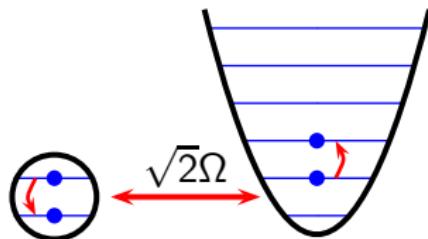
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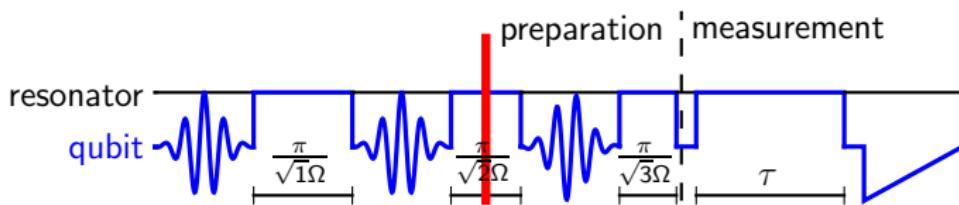
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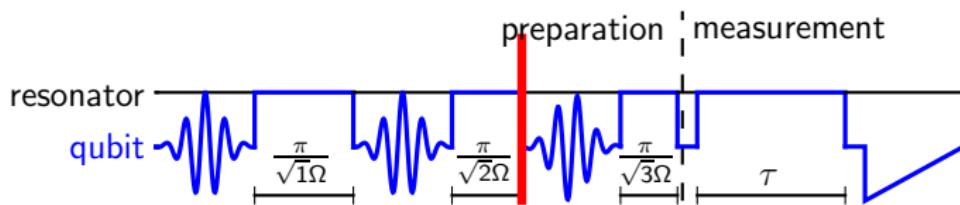
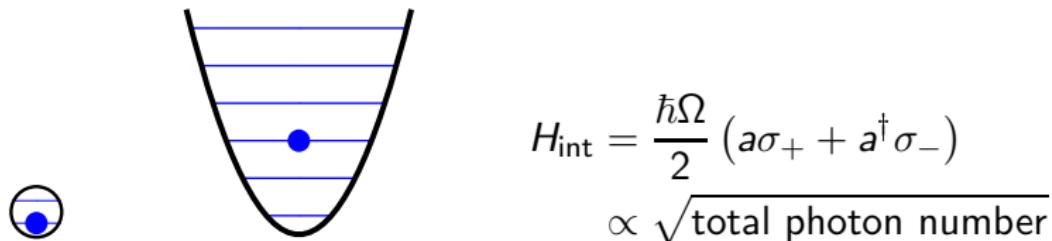
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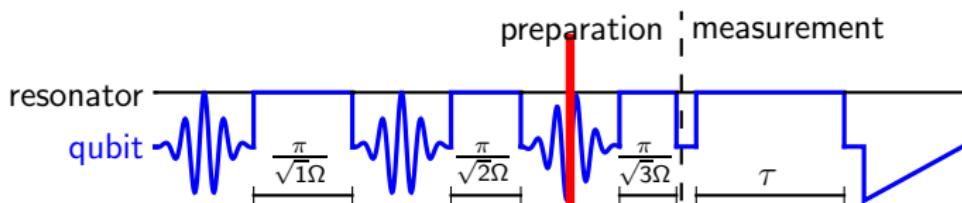
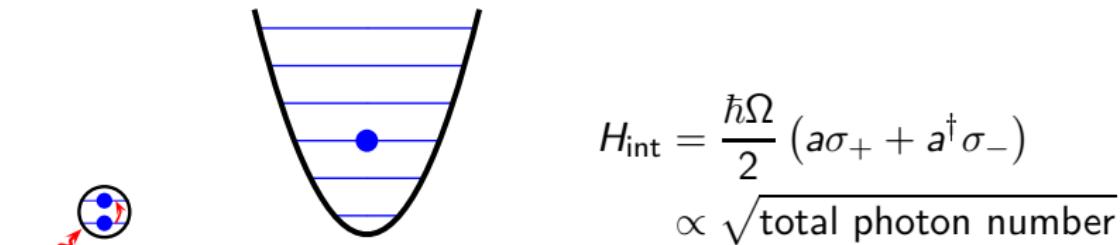
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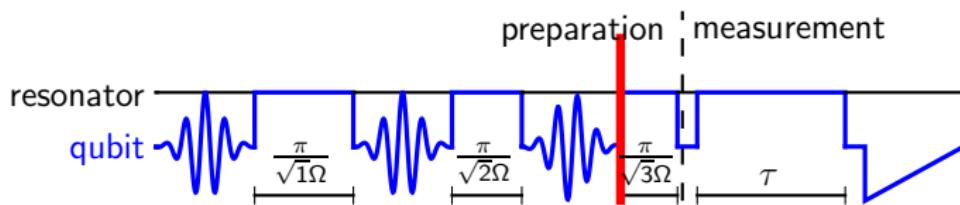
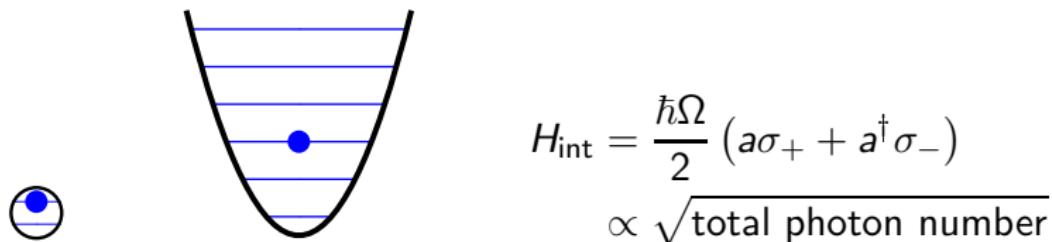
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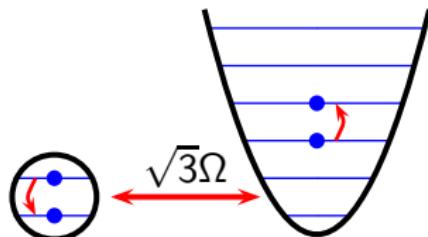
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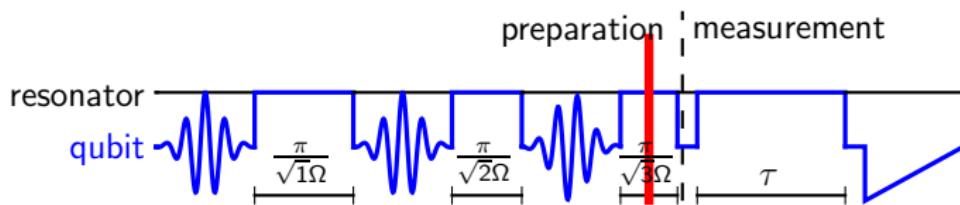
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Generating Fock states: Pumping photons one by one



$$H_{\text{int}} = \frac{\hbar\Omega}{2} (a\sigma_+ + a^\dagger\sigma_-)$$
$$\propto \sqrt{\text{total photon number}}$$



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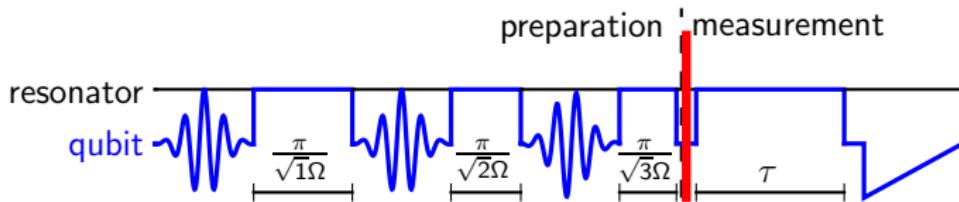
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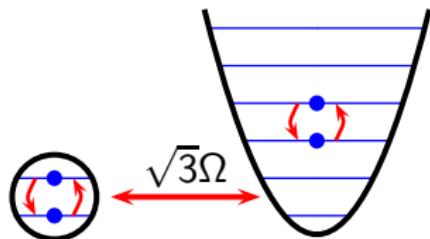
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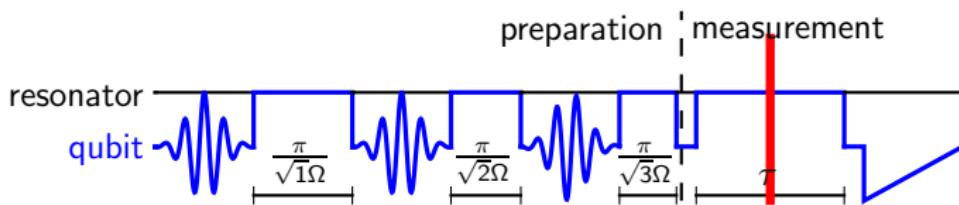
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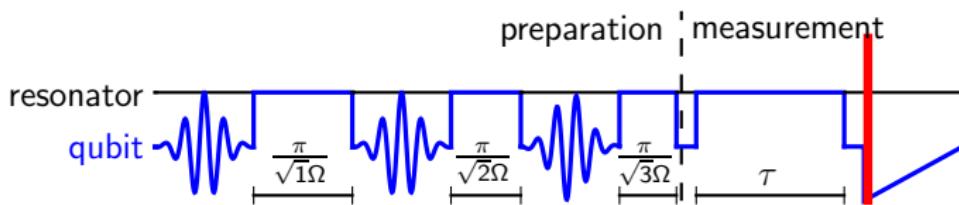
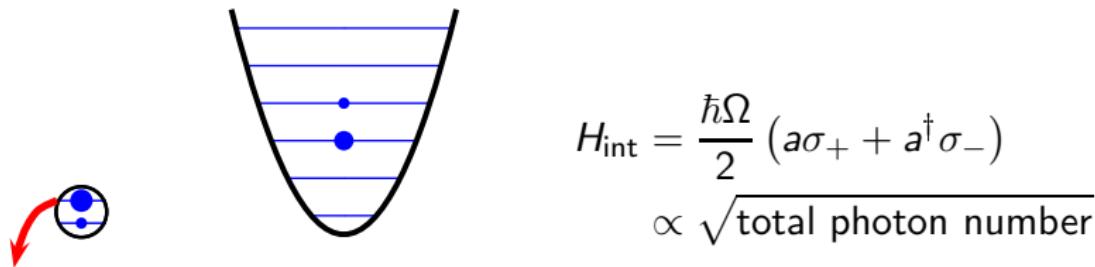
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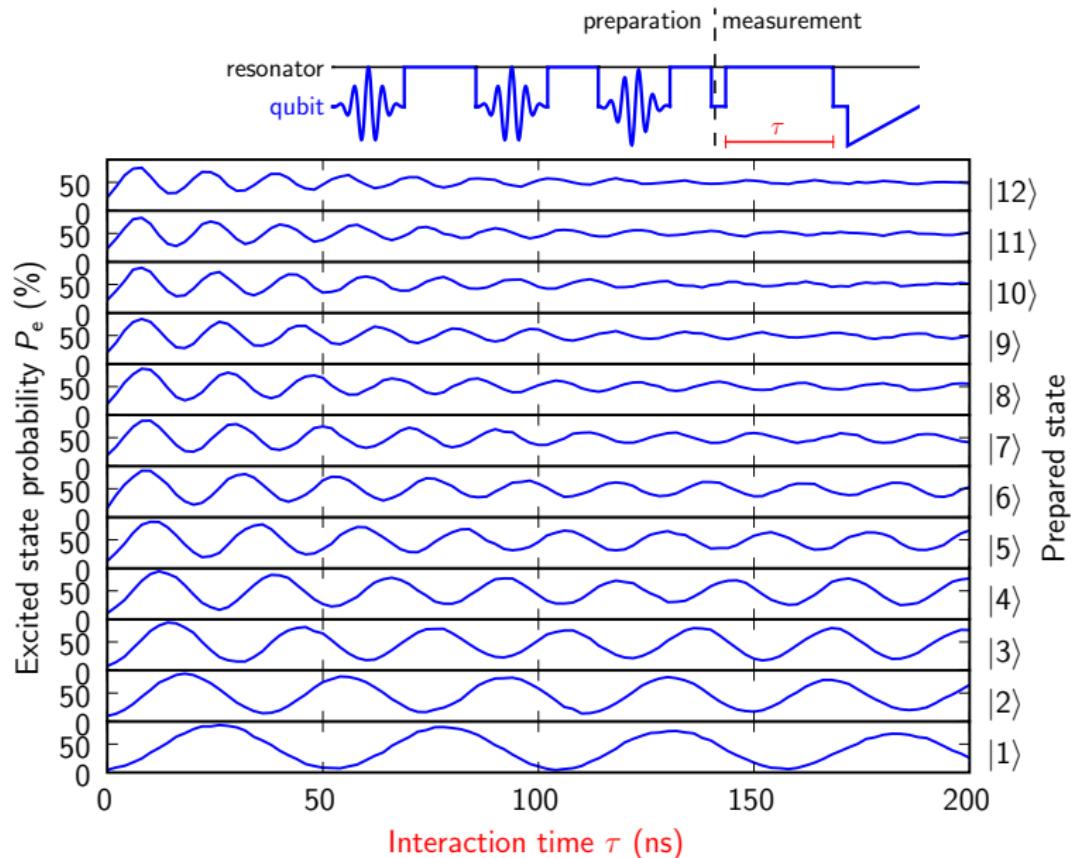
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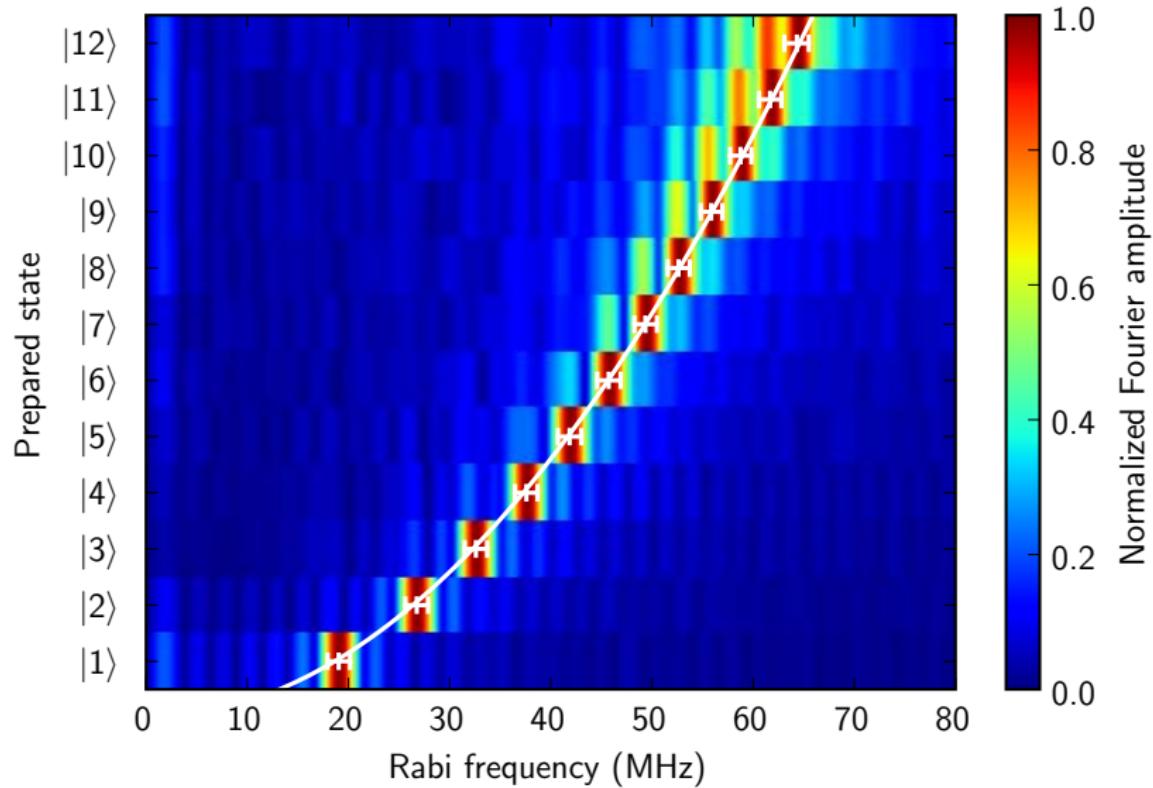
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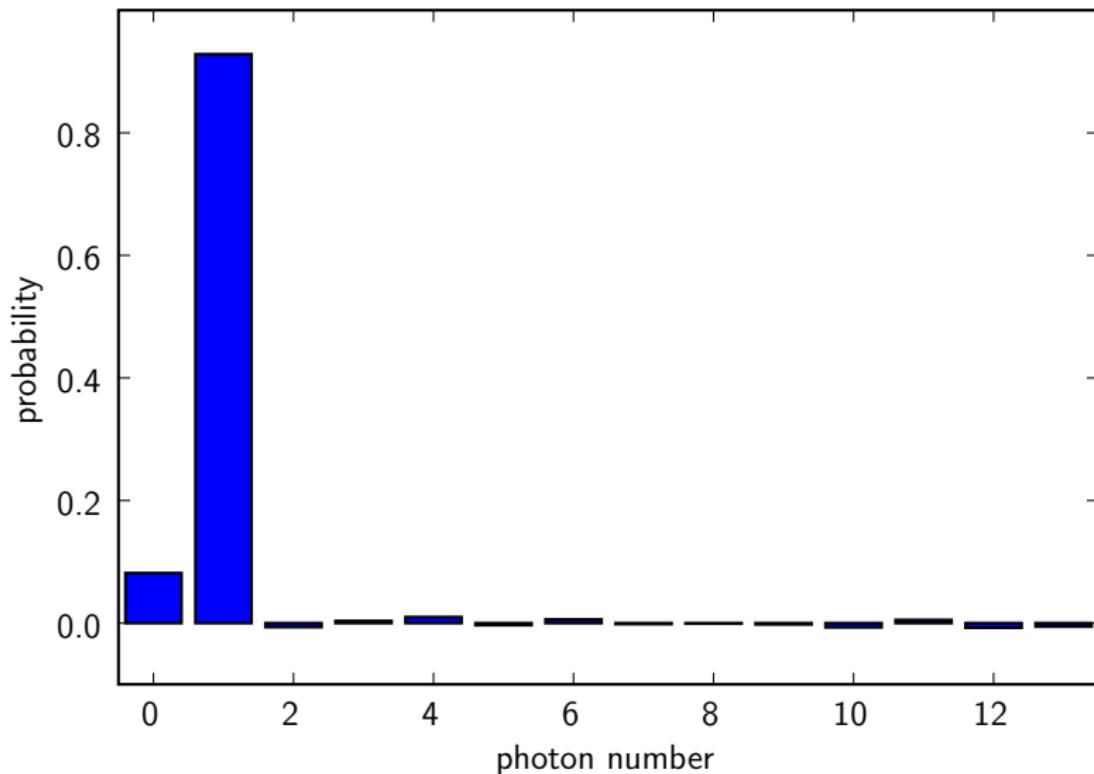
Fock states: Clear Rabi oscillations



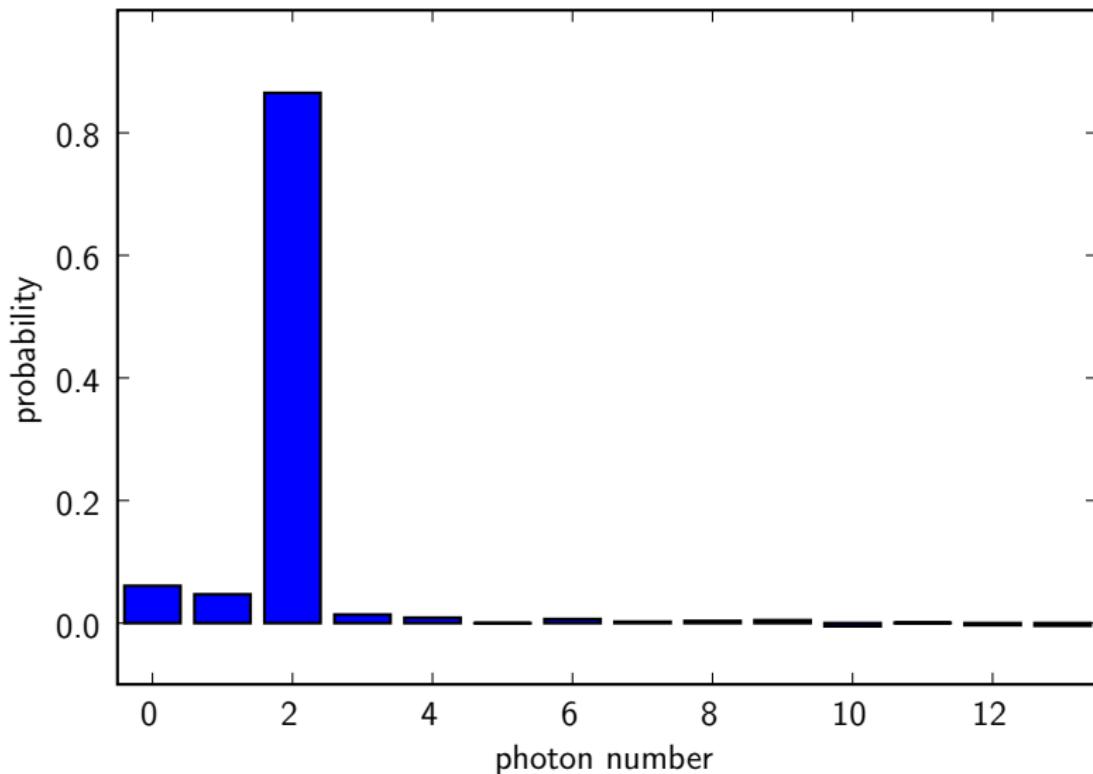
Fock states: Rabi frequencies scale as \sqrt{n}



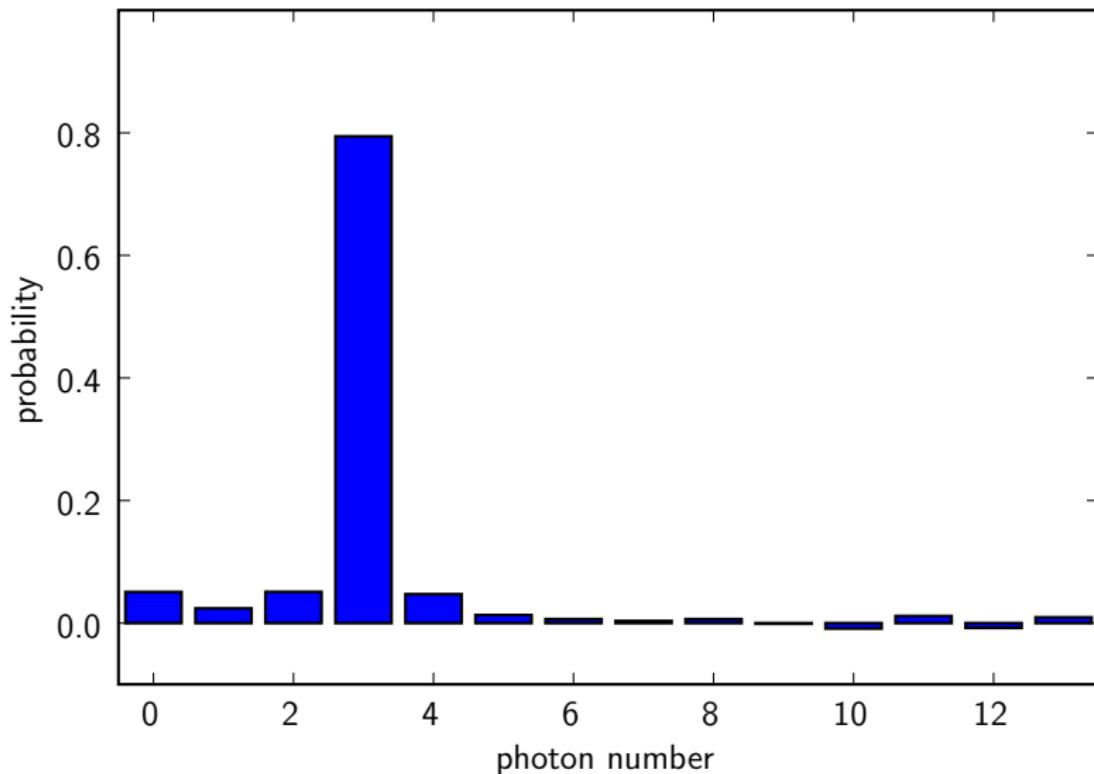
Fock states: Photon number distribution for $|1\rangle$



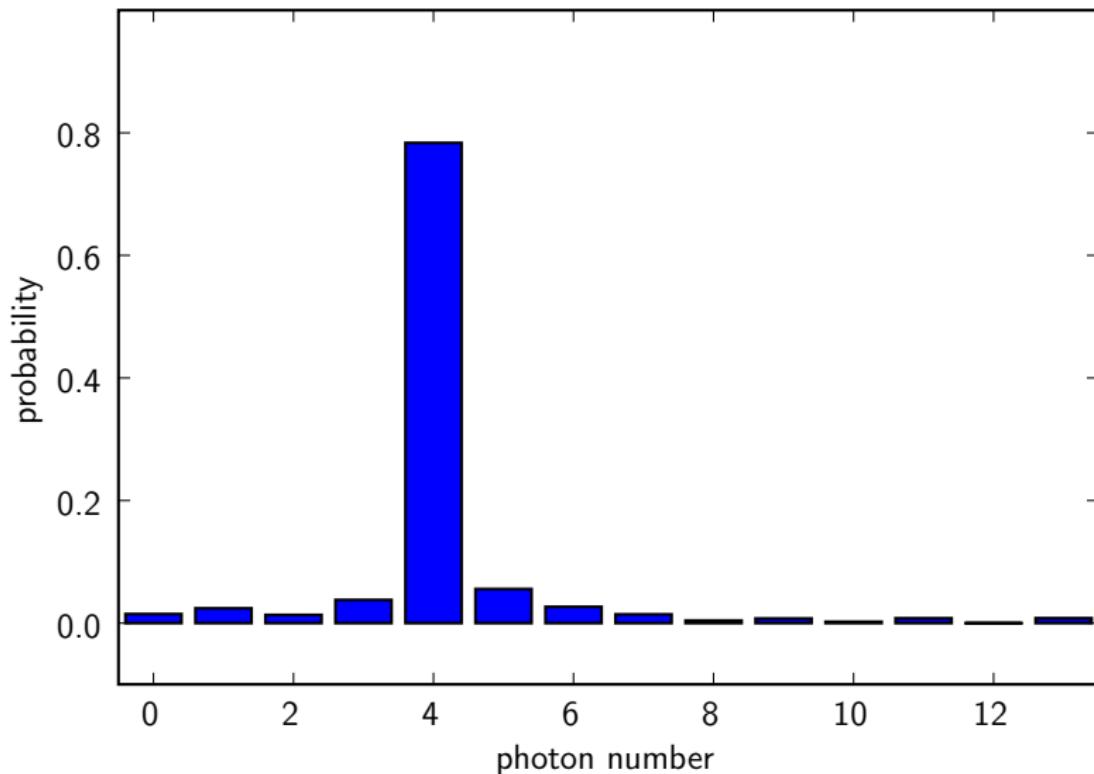
Fock states: Photon number distribution for $|2\rangle$



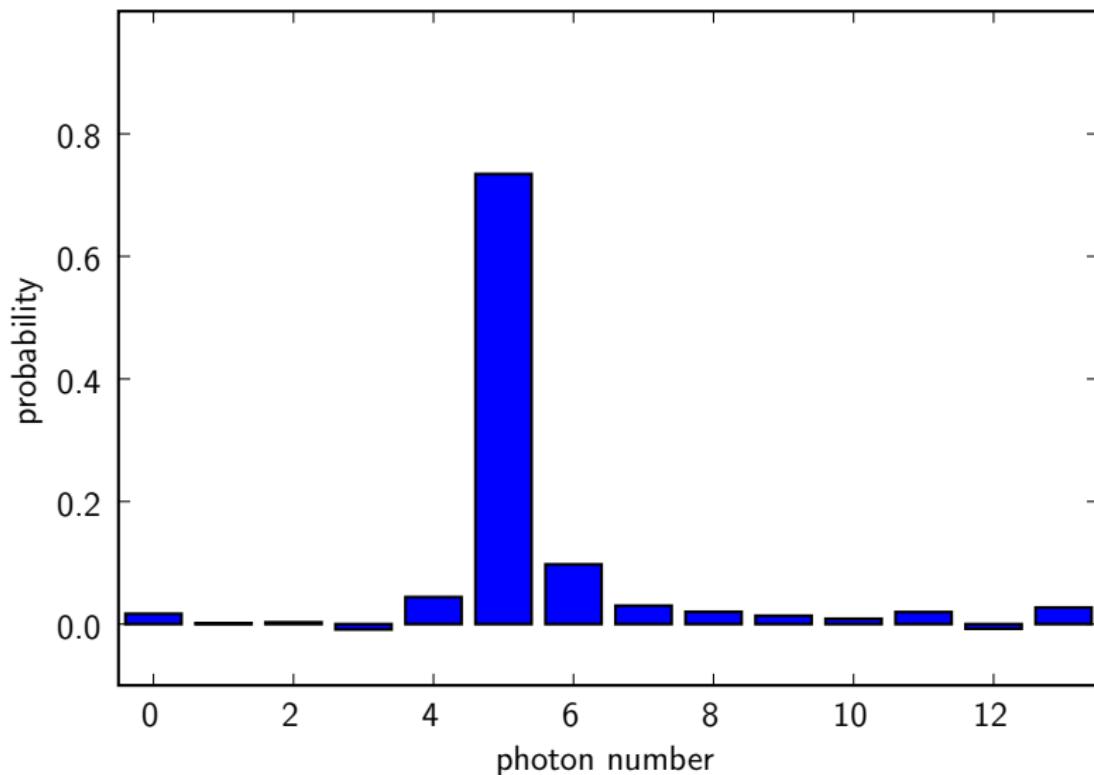
Fock states: Photon number distribution for $|3\rangle$



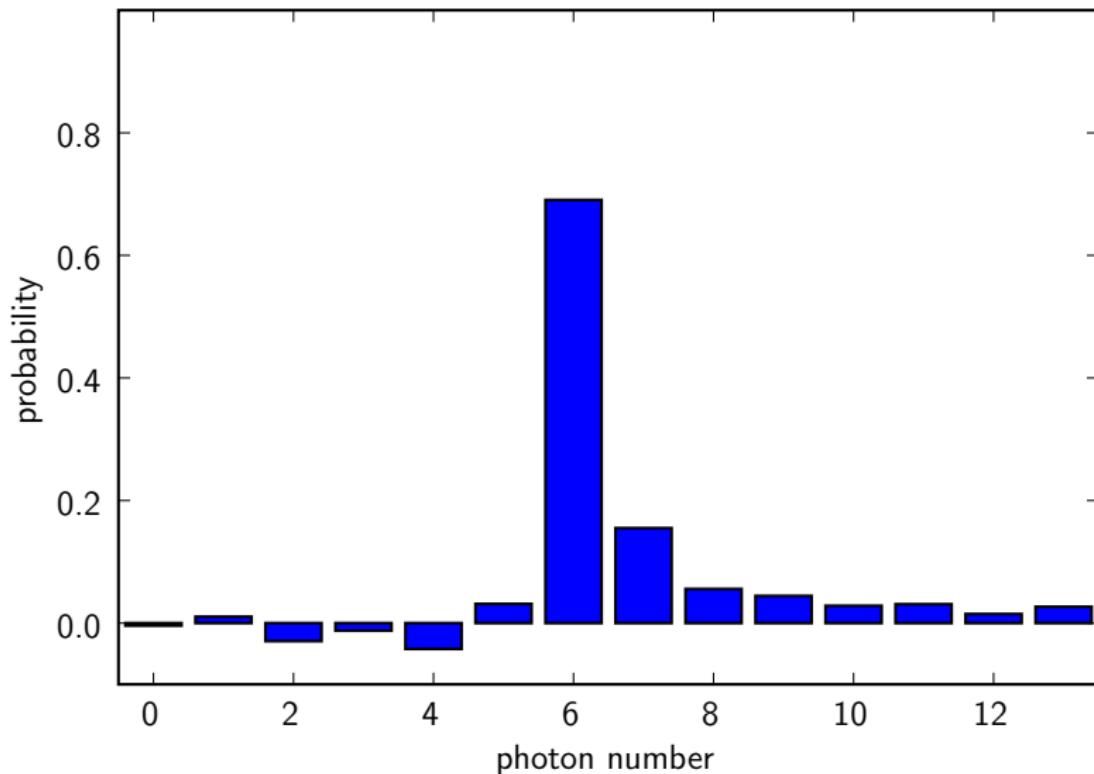
Fock states: Photon number distribution for $|4\rangle$



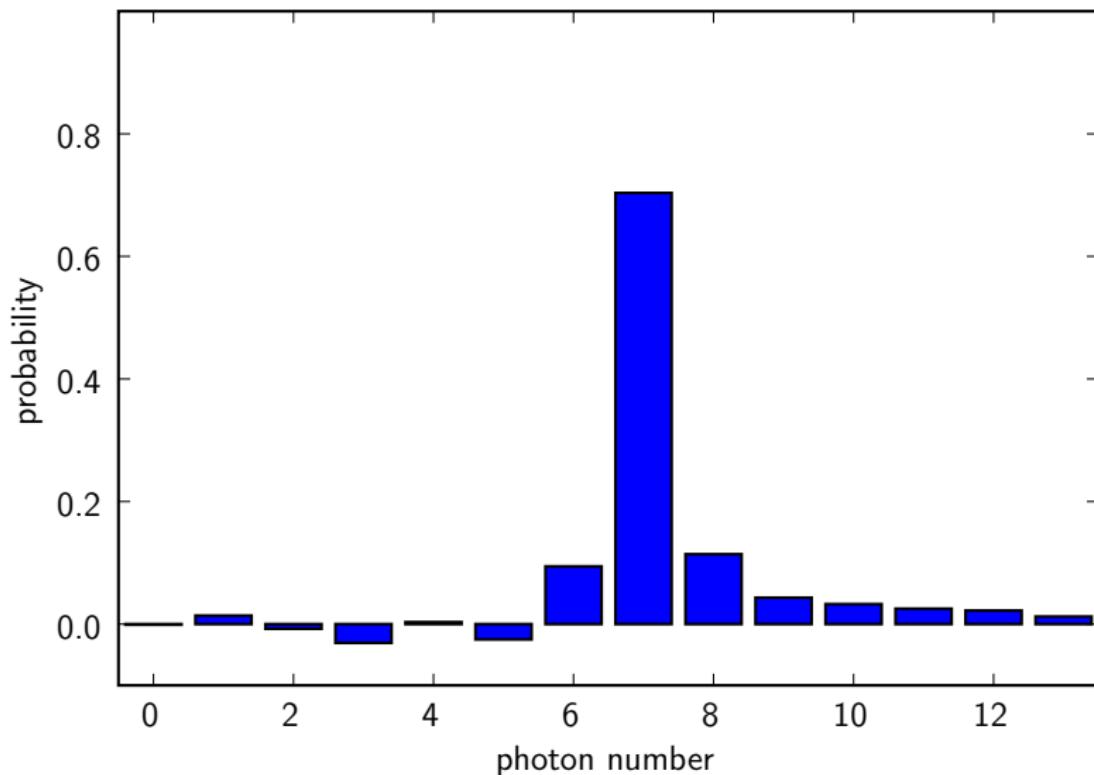
Fock states: Photon number distribution for $|5\rangle$



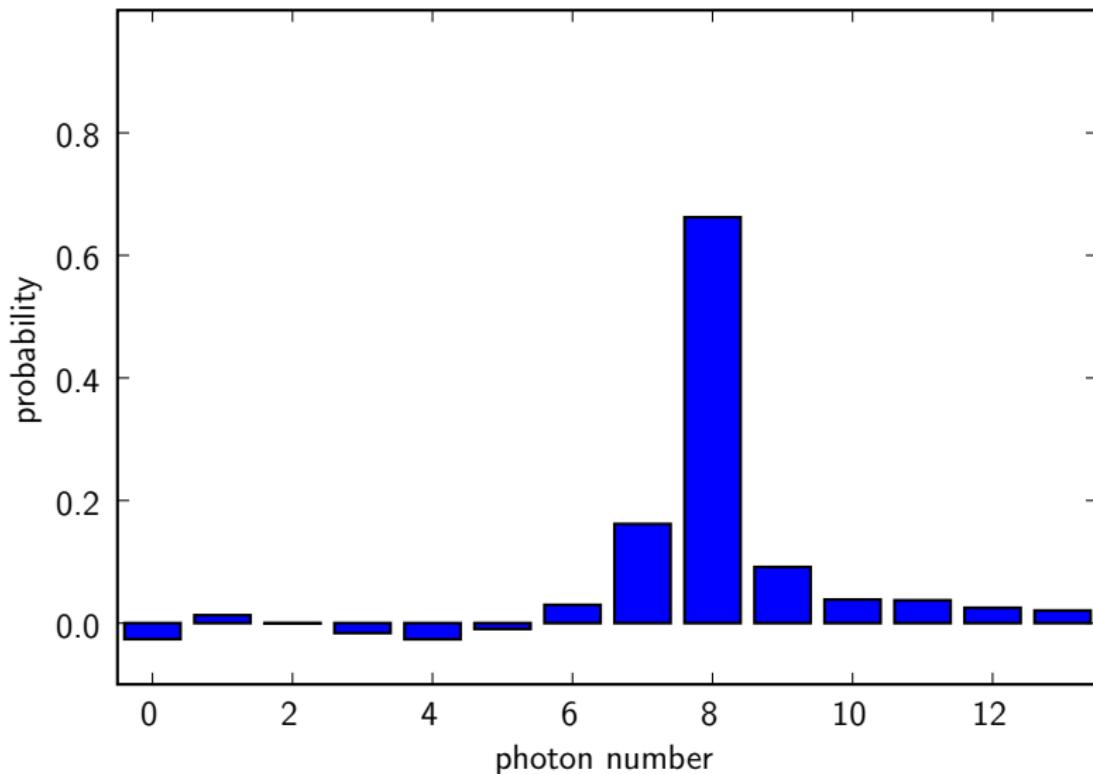
Fock states: Photon number distribution for $|6\rangle$



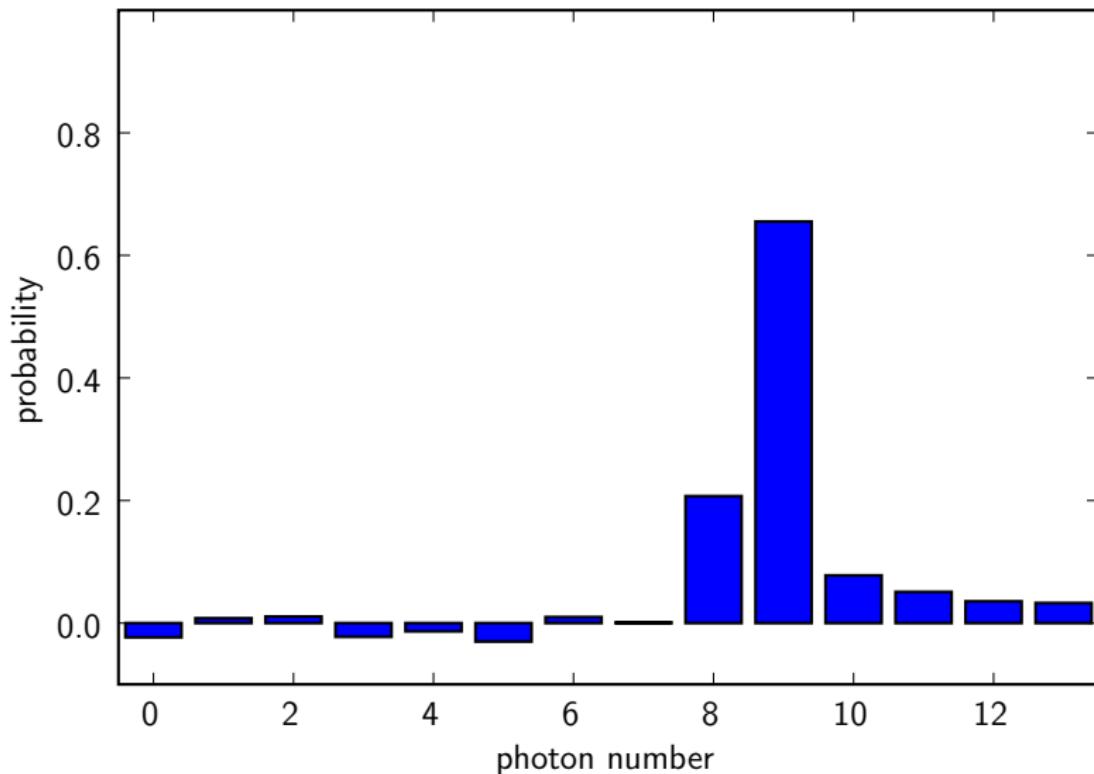
Fock states: Photon number distribution for $|7\rangle$



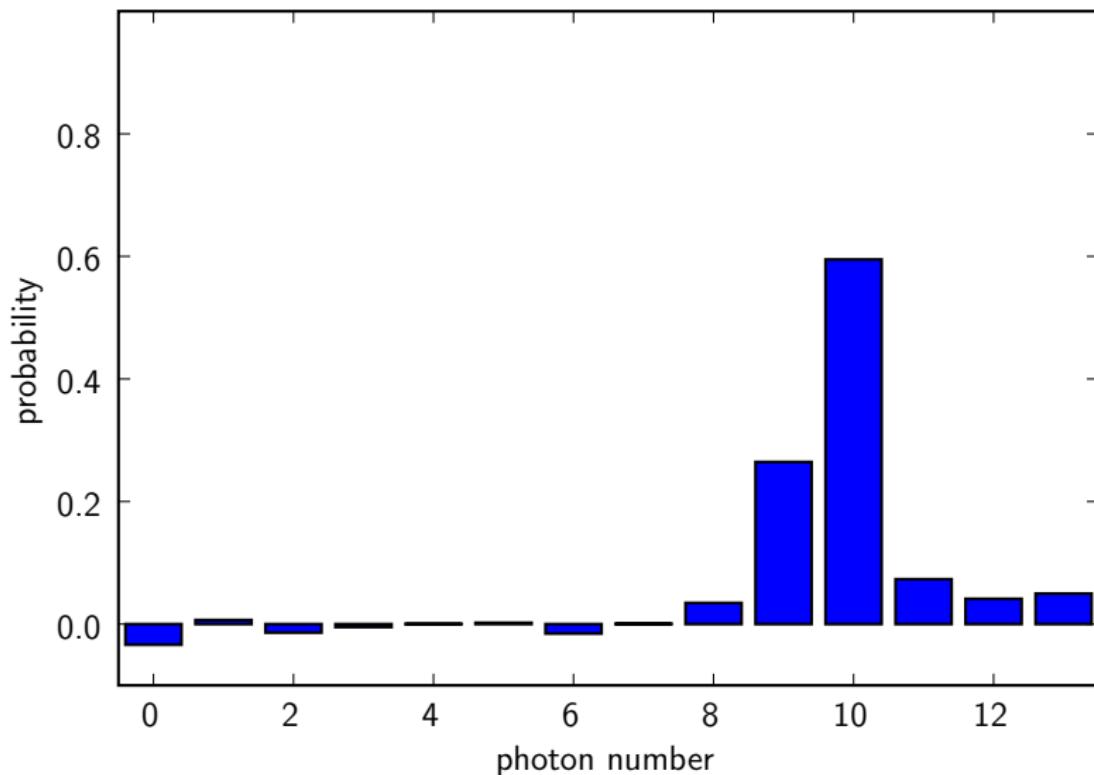
Fock states: Photon number distribution for $|8\rangle$



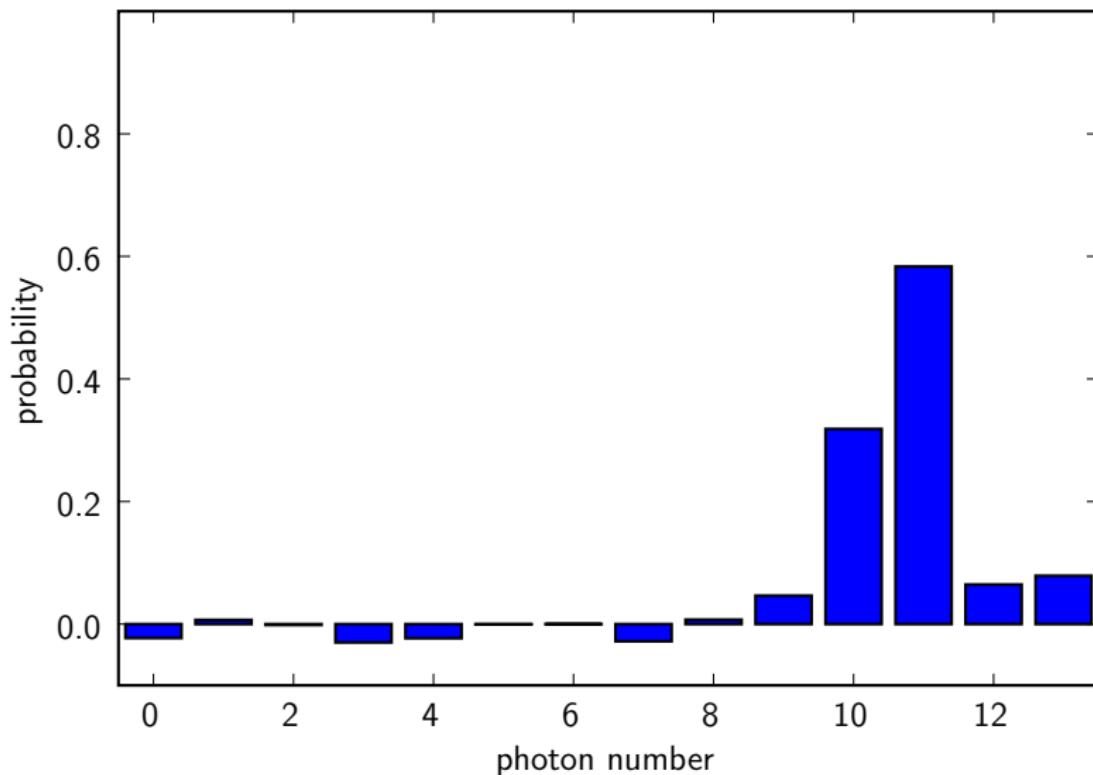
Fock states: Photon number distribution for $|9\rangle$



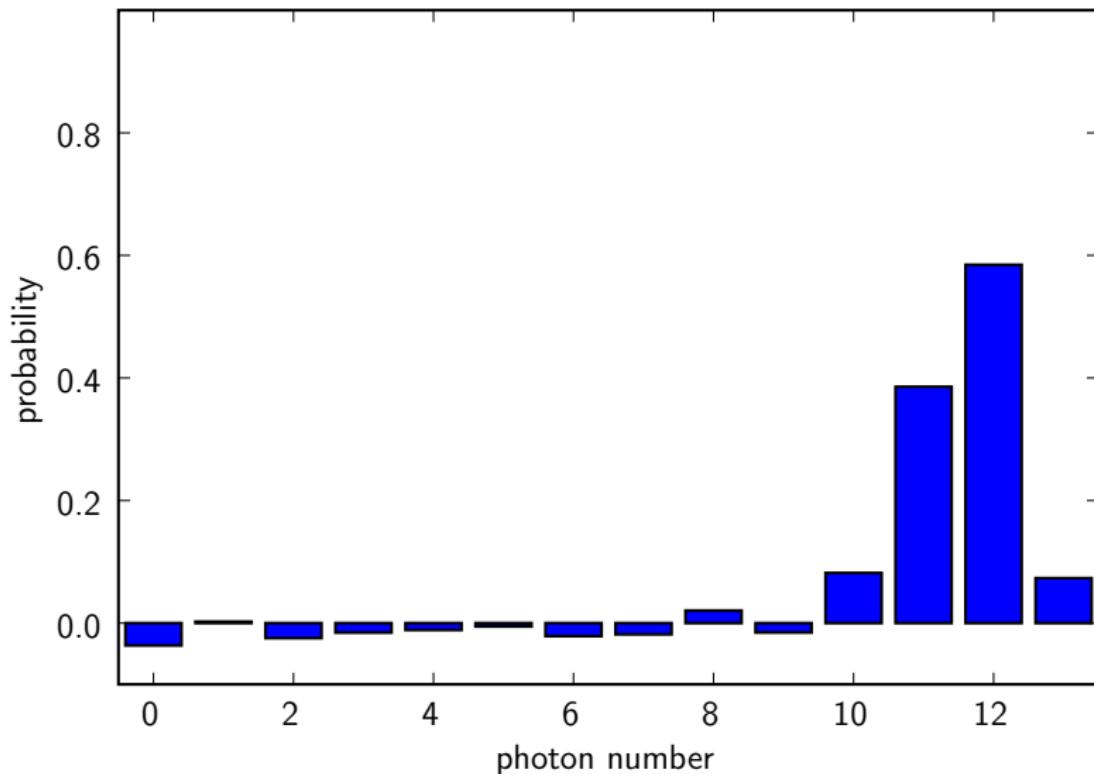
Fock states: Photon number distribution for $|10\rangle$



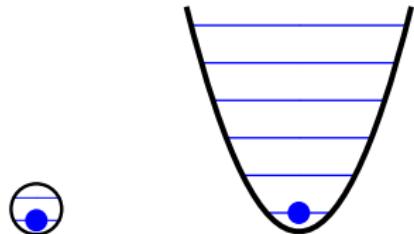
Fock states: Photon number distribution for $|11\rangle$



Fock states: Photon number distribution for $|12\rangle$

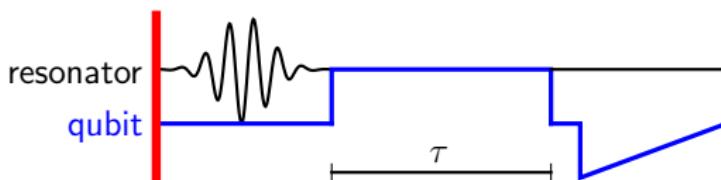


Directly driving the resonator creates a coherent state

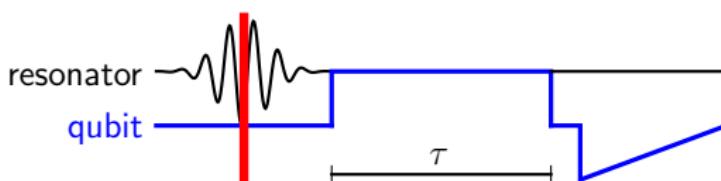
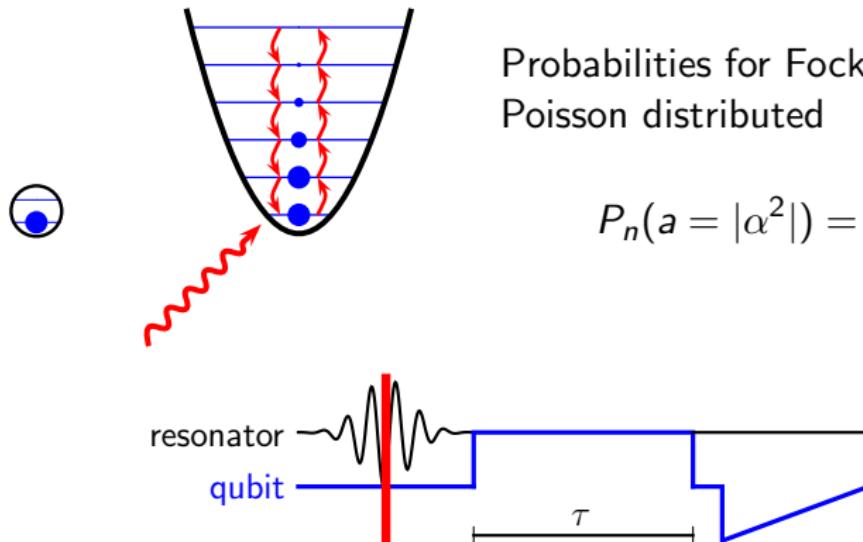


Probabilities for Fock states are
Poisson distributed

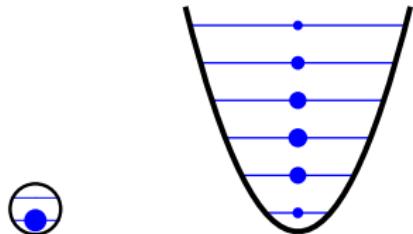
$$P_n(a = |\alpha^2|) = \frac{a^n e^{-a}}{n!}$$



Directly driving the resonator creates a coherent state

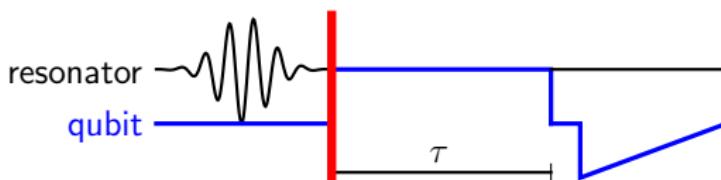


Directly driving the resonator creates a coherent state

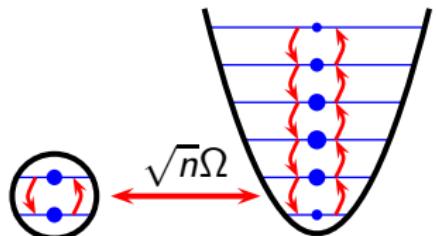


Probabilities for Fock states are
Poisson distributed

$$P_n(a = |\alpha^2|) = \frac{a^n e^{-a}}{n!}$$

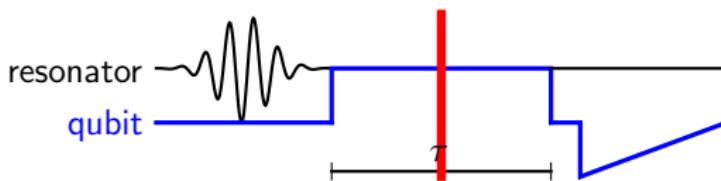


Directly driving the resonator creates a coherent state

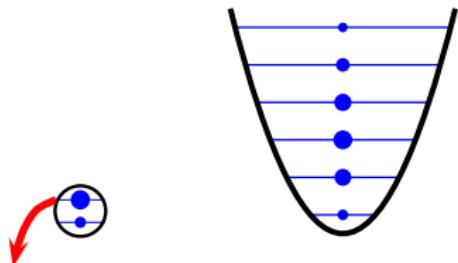


Probabilities for Fock states are Poisson distributed

$$P_n(a = |\alpha^2|) = \frac{a^n e^{-a}}{n!}$$

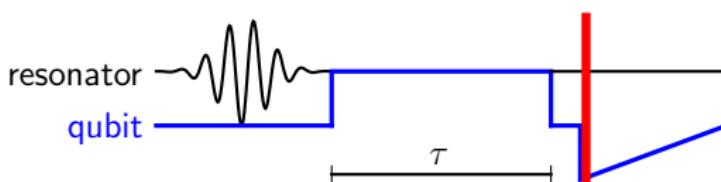


Directly driving the resonator creates a coherent state

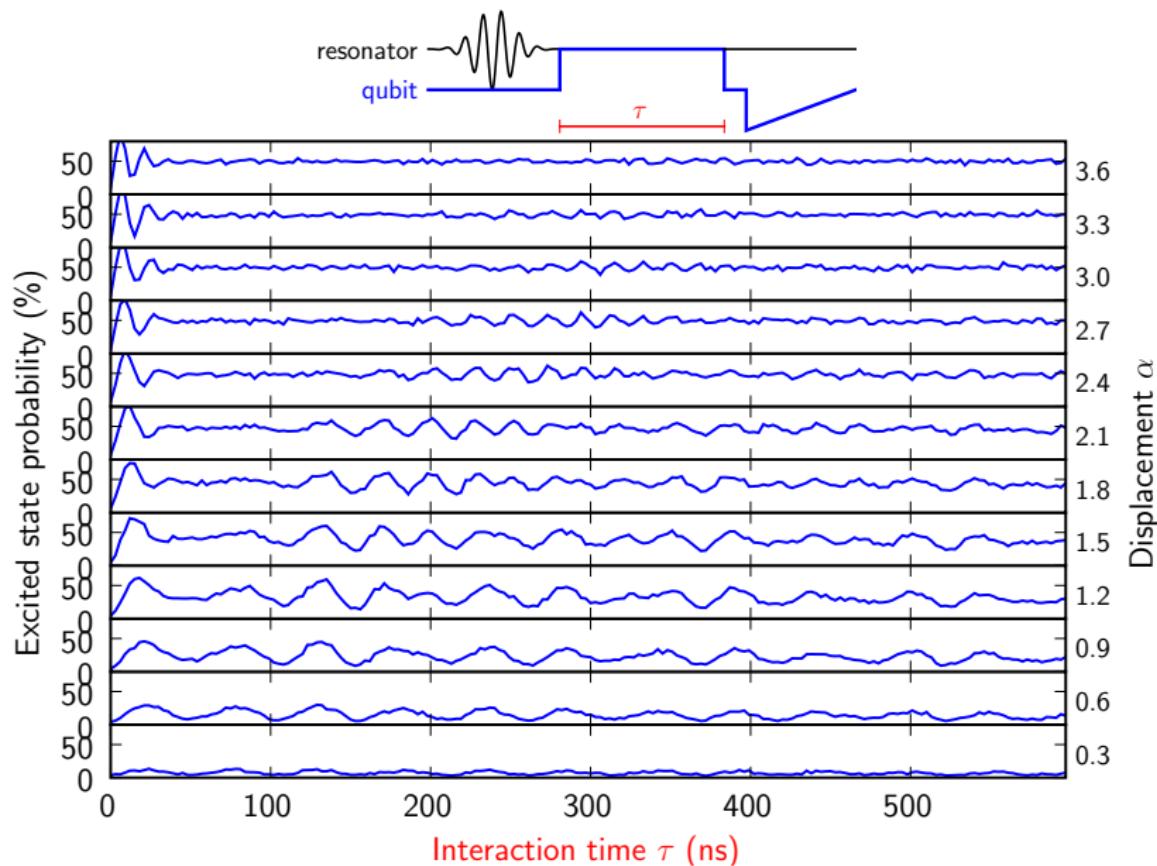


Probabilities for Fock states are
Poisson distributed

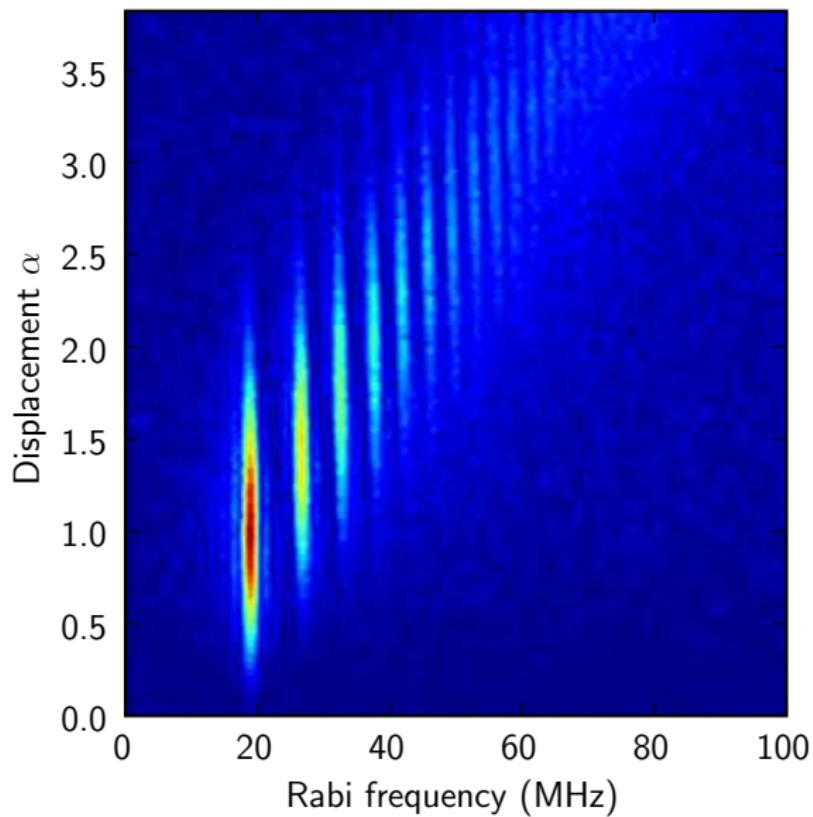
$$P_n(a = |\alpha^2|) = \frac{a^n e^{-a}}{n!}$$



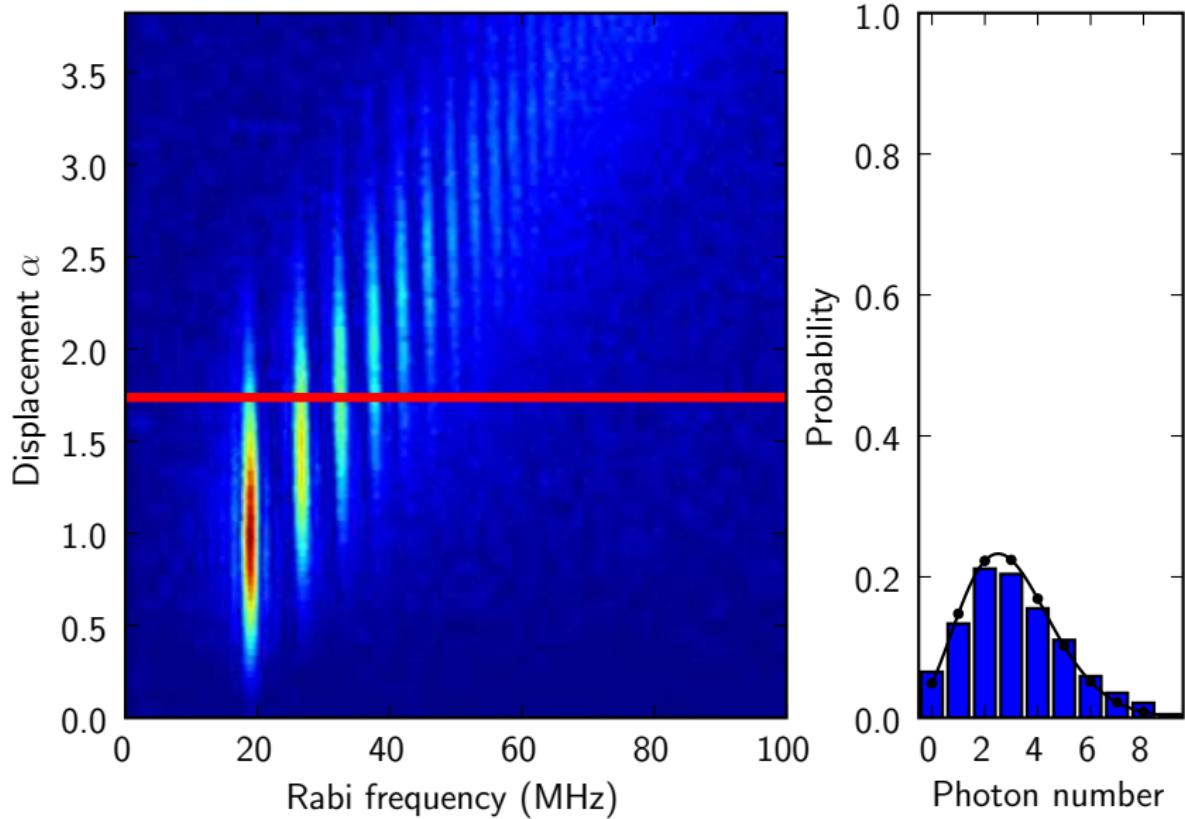
Coherent state: Non-periodic time traces



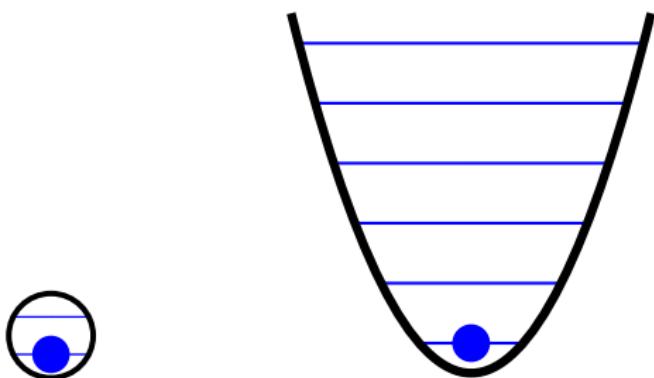
Coherent state: Superposition of Fock states



Coherent state: Superposition of Fock states

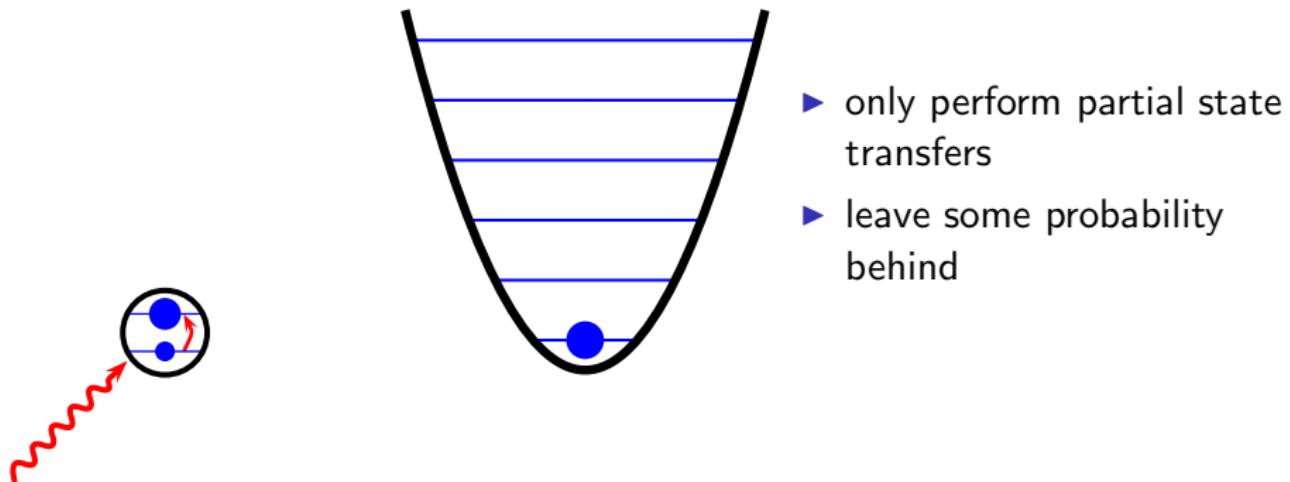


Generalizing the pump sequence to arbitrary states

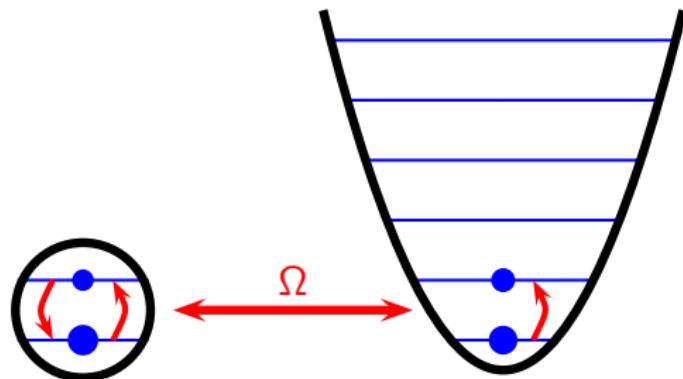


- ▶ only perform partial state transfers
- ▶ leave some probability behind

Generalizing the pump sequence to arbitrary states

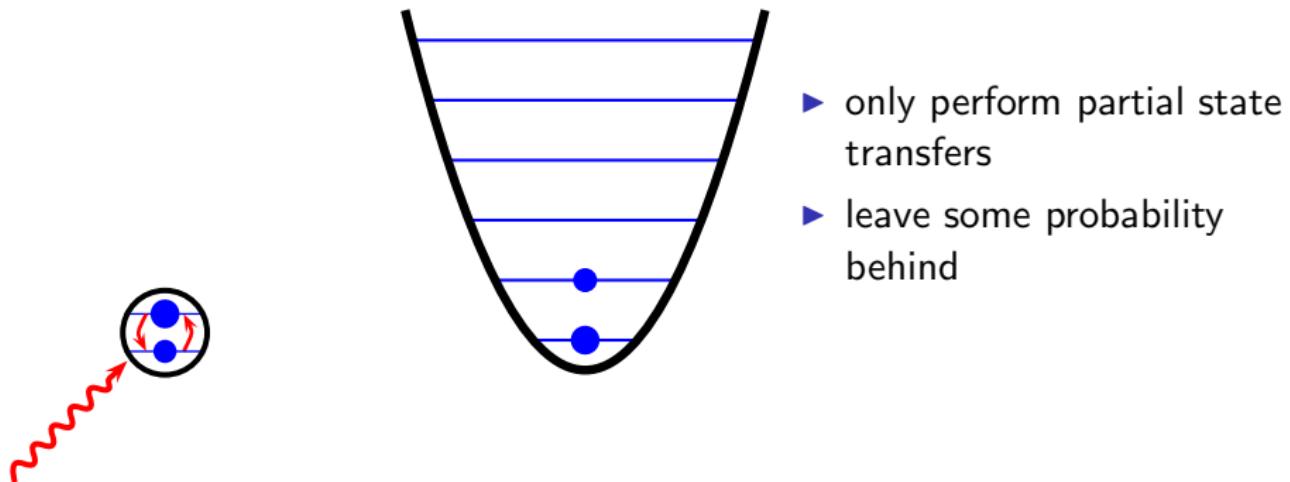


Generalizing the pump sequence to arbitrary states

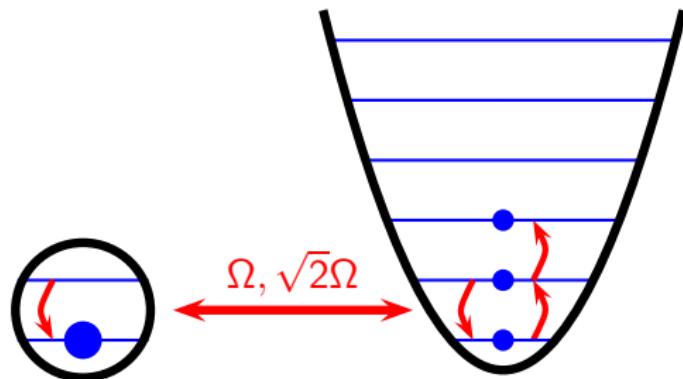


- ▶ only perform partial state transfers
- ▶ leave some probability behind

Generalizing the pump sequence to arbitrary states

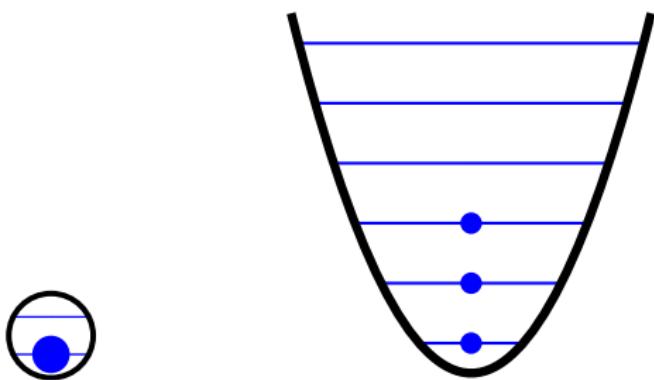


Generalizing the pump sequence to arbitrary states



- ▶ only perform partial state transfers
- ▶ leave some probability behind

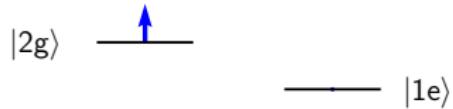
Generalizing the pump sequence to arbitrary states



- ▶ qubit and resonator entangled during sequence
- ▶ phases matter
- ▶ swaps act on several photon numbers

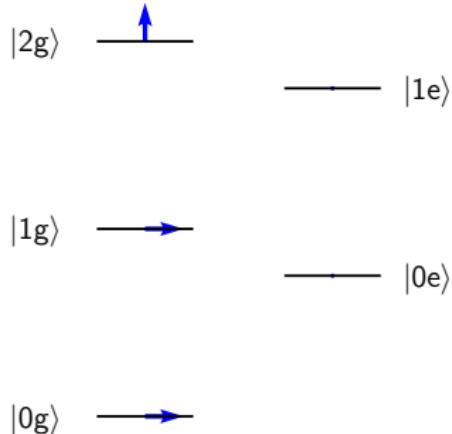
Algorithm for arbitrary states: “Reverse Engineering”

$$|\psi\rangle = \frac{1}{\sqrt{3}}(|0\rangle + |1\rangle + i|2\rangle)|g\rangle$$



Algorithm for arbitrary states: “Reverse Engineering”

$$|\psi\rangle = \frac{1}{\sqrt{3}}(|0\rangle + |1\rangle + i|2\rangle)|g\rangle$$



Time-reversed problem easier:

- ▶ start with the desired state
- ▶ eliminate states from top to bottom

Law, Eberly, PRL **76**, 1055 (1996)

resonator —
qubit —

Algorithm for arbitrary states: “Reverse Engineering”



resonator
qubit



Algorithm for arbitrary states: “Reverse Engineering”

$$|\psi\rangle = (0.577|0\rangle + 0.256|1\rangle)|g\rangle + (0.517i|0\rangle - 0.577|1\rangle)|e\rangle$$

$|2g\rangle$ ——

← |1e\rangle

$|1g\rangle$ —→

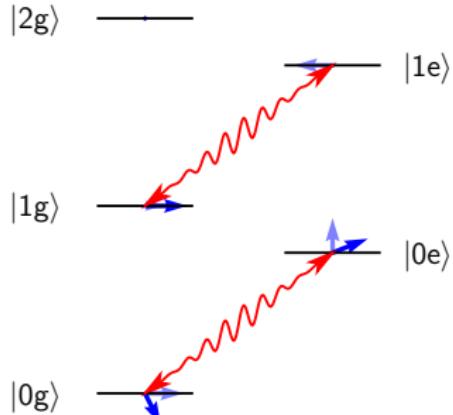
↑ |0e\rangle

$|0g\rangle$ —→

resonator
qubit



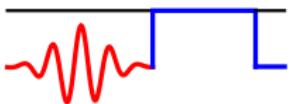
Algorithm for arbitrary states: “Reverse Engineering”



2 \times 2 parameters per photon step:

- ▶ qubit drive pulse amplitude
- ▶ qubit drive pulse phase

resonator
qubit



Algorithm for arbitrary states: “Reverse Engineering”

$$|\psi\rangle = ((0.234 - 0.473i)|0\rangle + 0.632|1\rangle)|g\rangle + (0.528 + 0.210i)|0\rangle|e\rangle$$

$|2g\rangle$ ——

—— $|1e\rangle$

$|1g\rangle$ —→

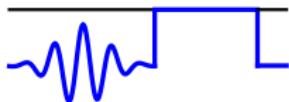
—→ $|0e\rangle$

$|0g\rangle$ —↓

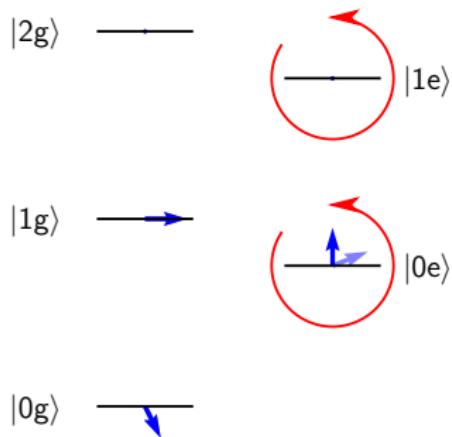
2×2 parameters per photon step:

- ▶ qubit drive pulse amplitude
- ▶ qubit drive pulse phase

resonator
qubit



Algorithm for arbitrary states: “Reverse Engineering”



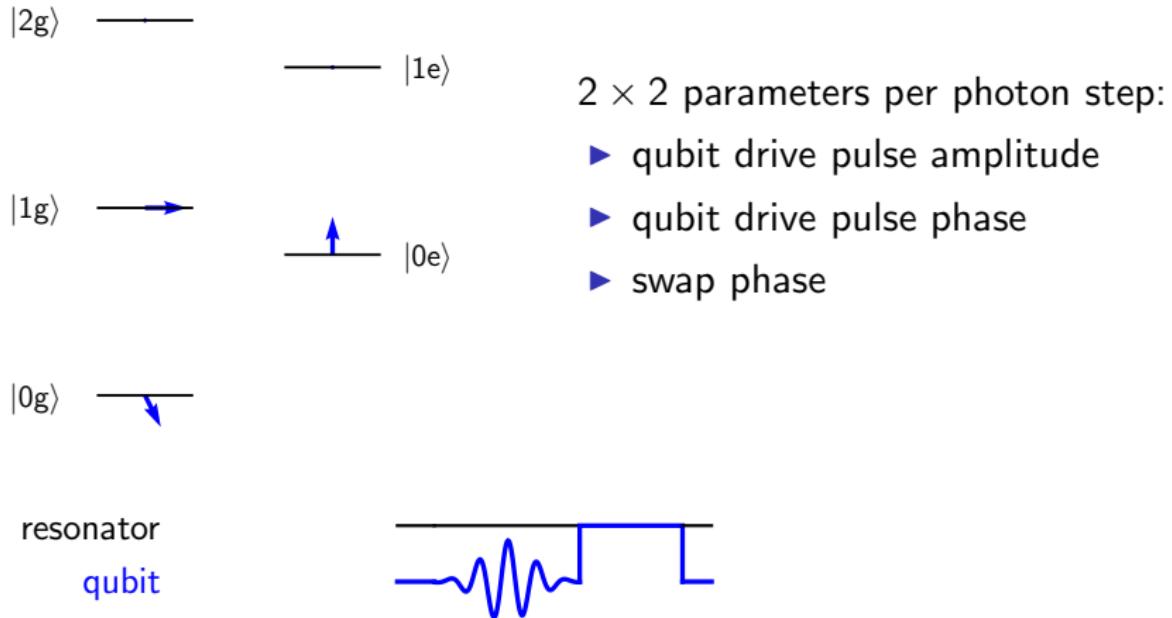
2×2 parameters per photon step:

- ▶ qubit drive pulse amplitude
- ▶ qubit drive pulse phase
- ▶ swap phase

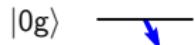


Algorithm for arbitrary states: “Reverse Engineering”

$$|\psi\rangle = ((0.234 - 0.473i)|0\rangle + 0.632|1\rangle)|g\rangle + (0.568i|0\rangle)|e\rangle$$



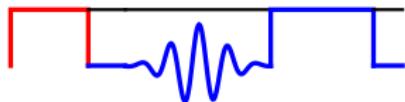
Algorithm for arbitrary states: “Reverse Engineering”



2×2 parameters per photon step:

- ▶ qubit drive pulse amplitude
- ▶ qubit drive pulse phase
- ▶ swap phase
- ▶ swap amplitude

resonator
qubit



Algorithm for arbitrary states: “Reverse Engineering”

$$|\psi\rangle = ((0.234 - 0.473i)|0\rangle)|g\rangle + (0.849i|0\rangle)|e\rangle$$

$|2g\rangle$ ——

—— $|1e\rangle$

$|1g\rangle$ ——

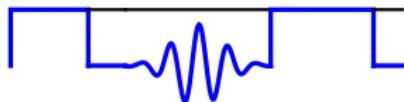
—— ↑ $|0e\rangle$

$|0g\rangle$ —— ↓

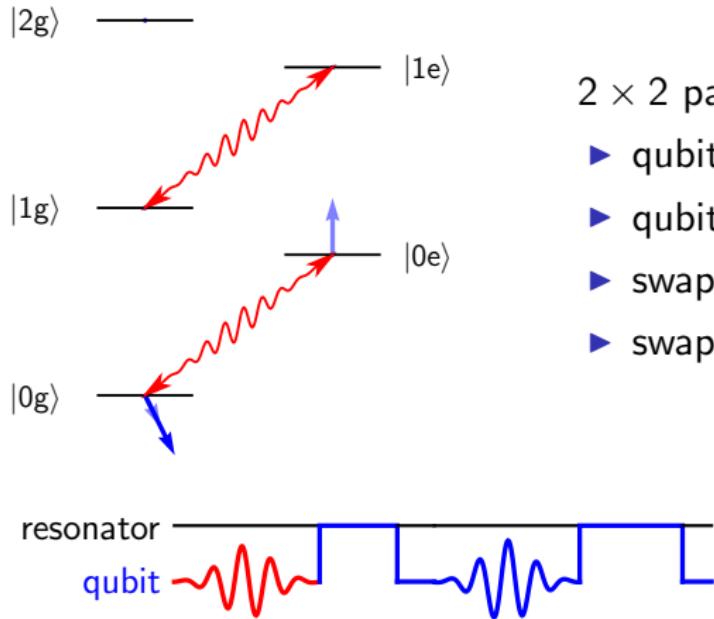
2×2 parameters per photon step:

- ▶ qubit drive pulse amplitude
- ▶ qubit drive pulse phase
- ▶ swap phase
- ▶ swap amplitude

resonator
qubit



Algorithm for arbitrary states: “Reverse Engineering”

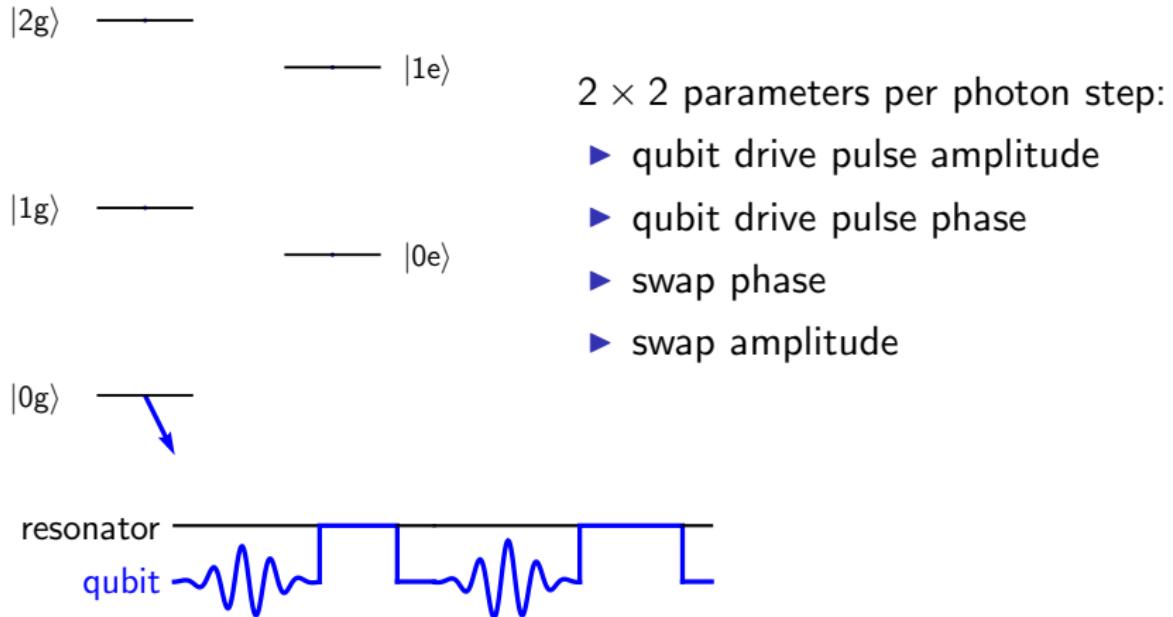


2×2 parameters per photon step:

- ▶ qubit drive pulse amplitude
- ▶ qubit drive pulse phase
- ▶ swap phase
- ▶ swap amplitude

Algorithm for arbitrary states: “Reverse Engineering”

$$|\psi\rangle = (0.444 - 0.896i)|0\rangle|g\rangle$$



Calibrations

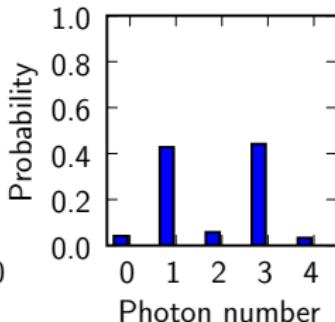
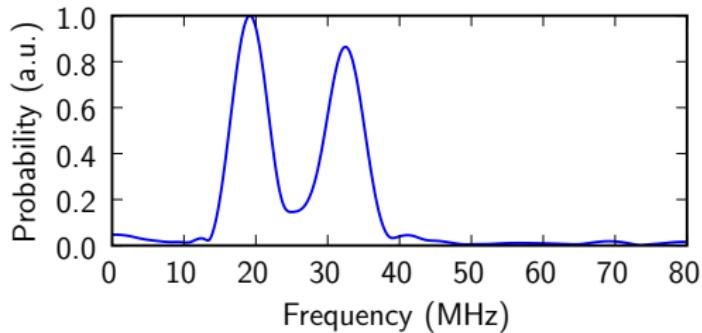
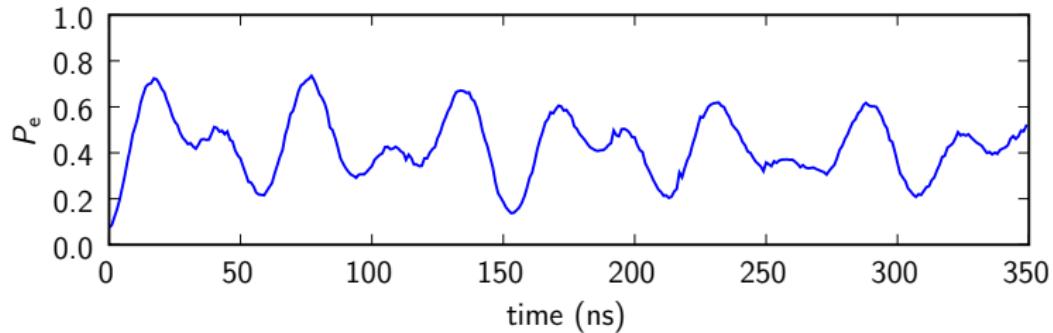
- ▶ impulse response (sampling scope)
- ▶ imperfections of IQ-mixer (spectrum analyzer)
- ▶ cable response (qubit as sampling scope)
- ▶ various parameters for readout SQUID
- ▶ π pulses (amplitude and frequency)
- ▶ measure pulse amplitude
- ▶ swap pulse amplitude
- ▶ swap pulse time for each photon number
- ▶ first order correction for finite rise time of swap pulses
- ▶ qubit/resonator dephasing rate when off resonance
- ▶ resonator displacement/drive amplitude ratio
- ▶ resonator drive phase
- ▶ readout visibility

BUT:

- ▶ **no per-state calibrations; just run the calculated sequence**

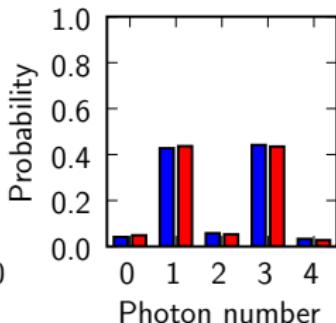
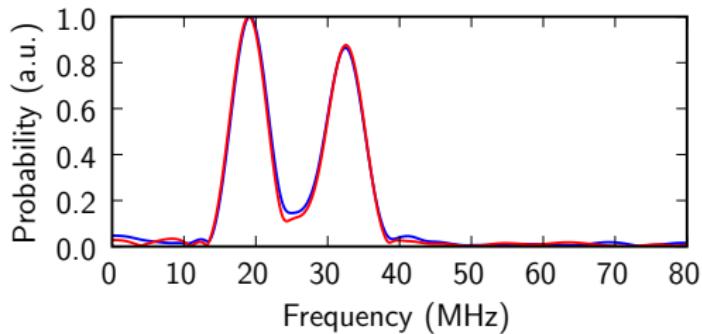
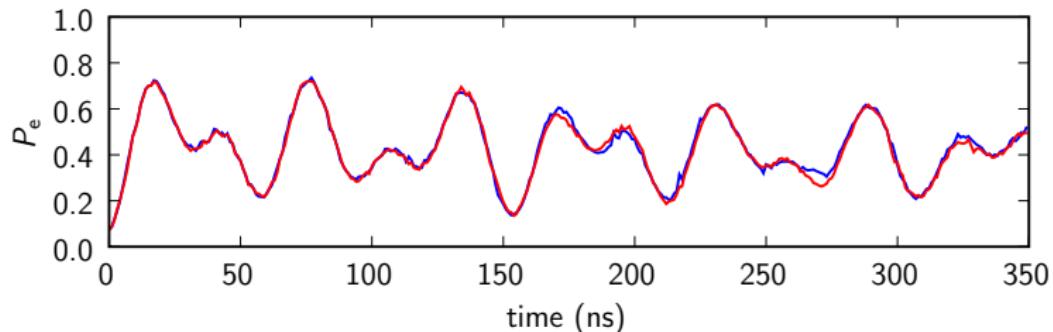
A superposition of $|1\rangle$ and $|3\rangle$

$$|\psi\rangle = |1\rangle + |3\rangle$$



A superposition of $|1\rangle$ and $|3\rangle$

$|\psi\rangle = |1\rangle + |3\rangle$ and $|\phi\rangle = |1\rangle + i|3\rangle$ look the same!

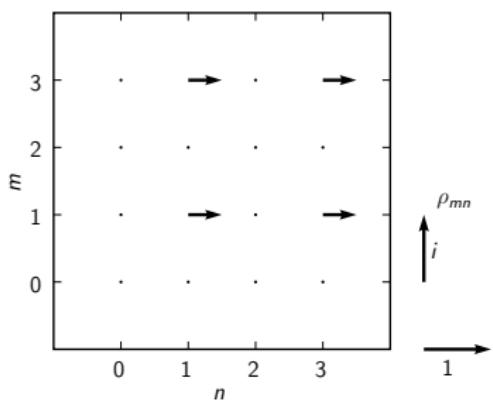


Full description of a resonator state

Example: $|\psi\rangle = |1\rangle + |3\rangle$

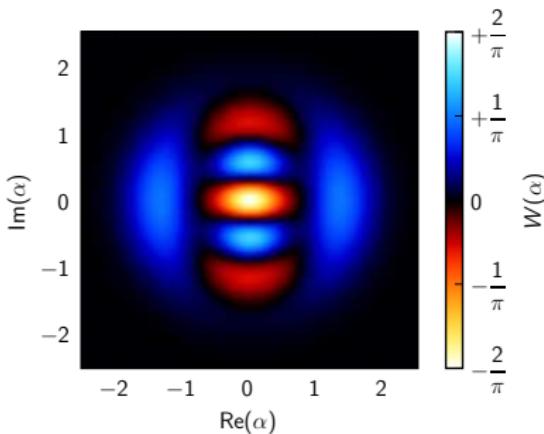
Density matrix:

$$\rho = |\psi\rangle\langle\psi|$$



Wigner function:

$$W(\alpha) = \frac{2}{\pi} \langle \psi | D^\dagger(-\alpha) \Pi D(-\alpha) | \psi \rangle$$

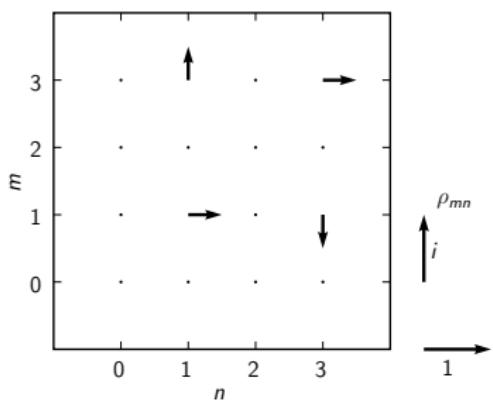


Full description of a resonator state

Example: $|\psi\rangle = |1\rangle + i|3\rangle$

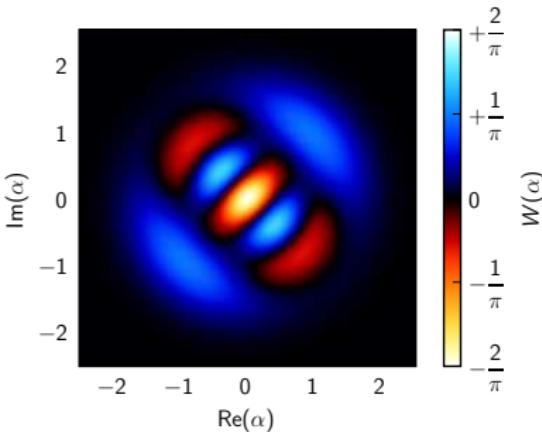
Density matrix:

$$\rho = |\psi\rangle\langle\psi|$$



Wigner function:

$$W(\alpha) = \frac{2}{\pi} \langle \psi | D^\dagger(-\alpha) \Pi D(-\alpha) | \psi \rangle$$



Wigner tomography: measurement protocol

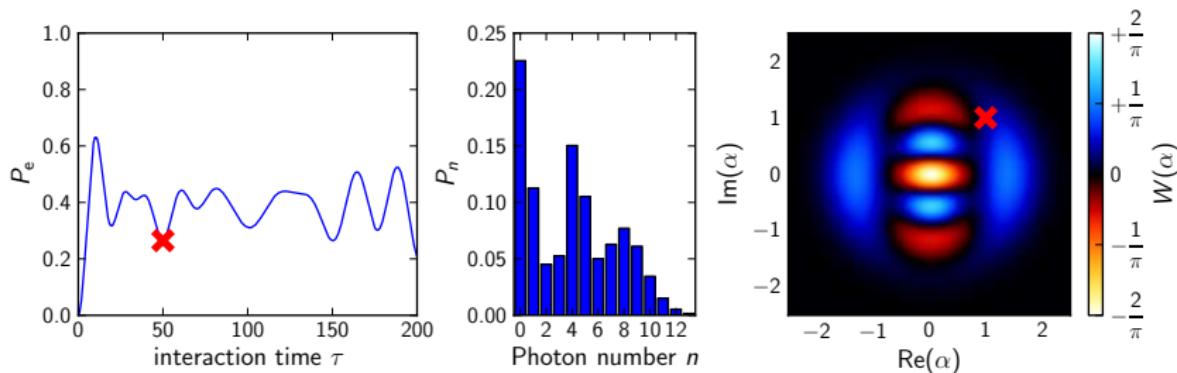


Preparation of state $|\psi\rangle$
(Law & Eberly protocol)

$$\rho = F^2 |\psi\rangle\langle\psi| + \dots$$

State reconstruction

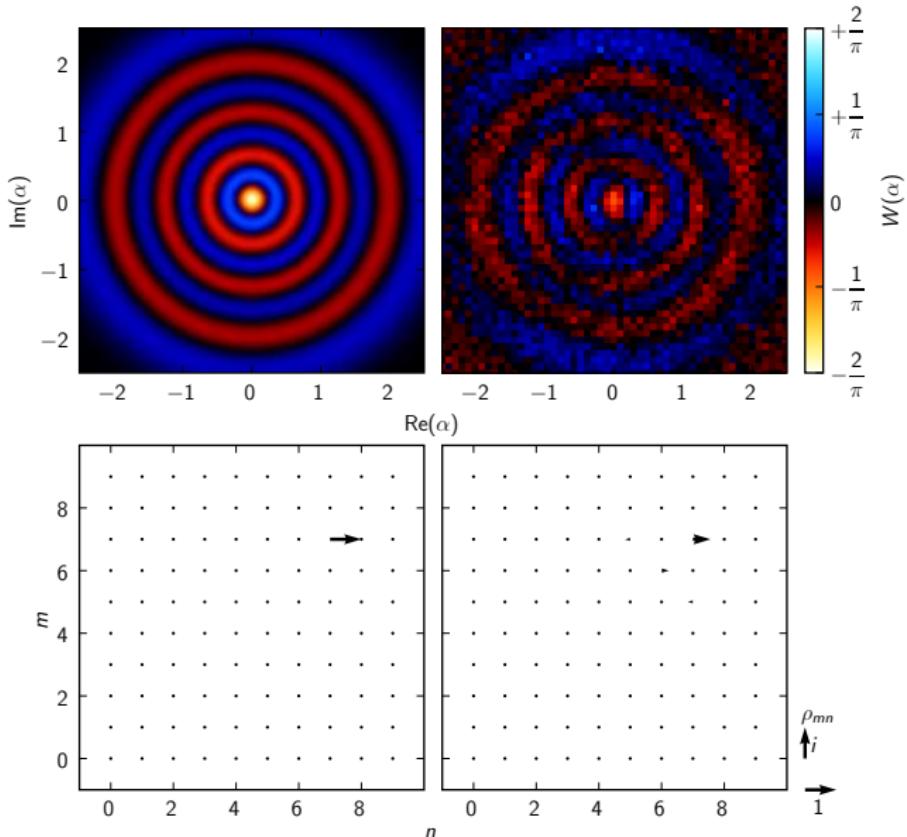
- ▶ displacement: $\rho \rightarrow \rho_\alpha = D(-\alpha)\rho D(\alpha)$
- ▶ qubit–resonator interaction for time τ
- ▶ qubit measurement



$$W(\alpha) = \frac{2}{\pi} \text{Tr} (\rho_\alpha \Pi) = \frac{2}{\pi} \sum_{n=0}^{\infty} (-1)^n P_n$$

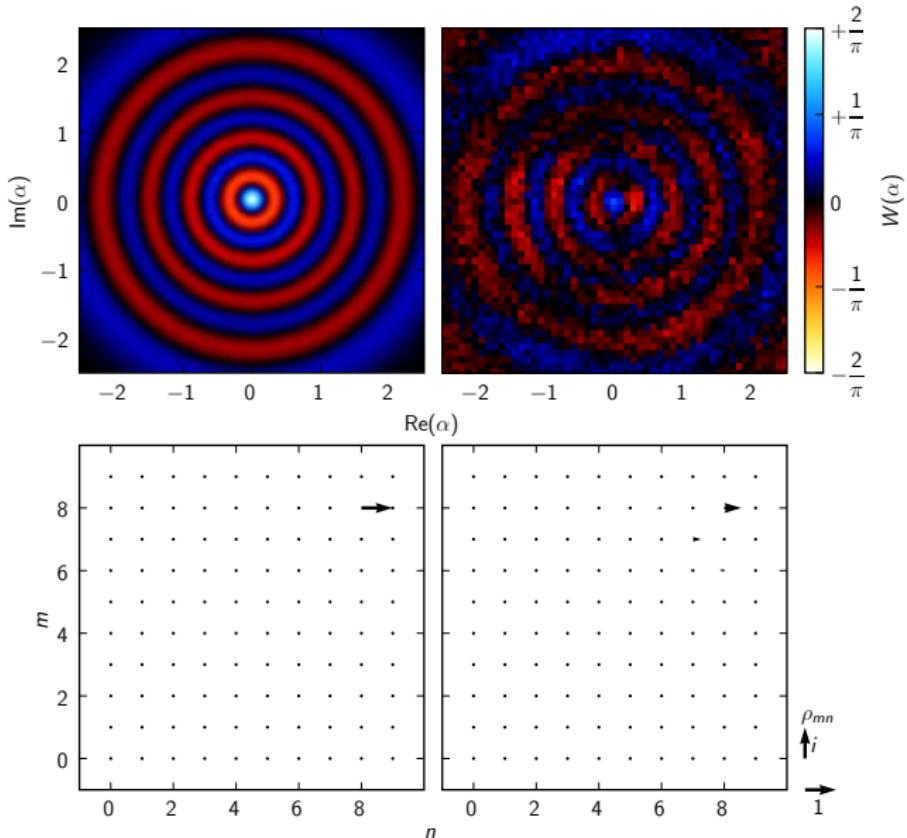
Controlling photon number

$$|\psi\rangle = |7\rangle$$



Controlling photon number

$$|\psi\rangle = |8\rangle$$

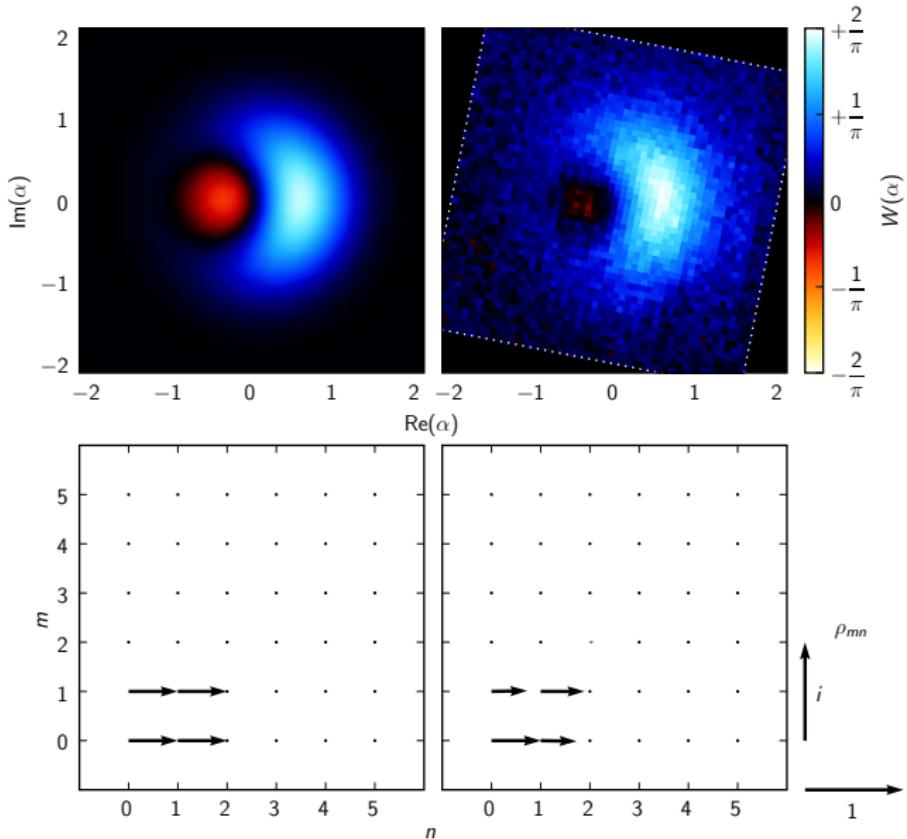


Fidelity:

$$F = \sqrt{\langle \psi | \rho | \psi \rangle} = 0.72$$

Controlling superpositions

$$|\psi\rangle = |0\rangle + |1\rangle$$

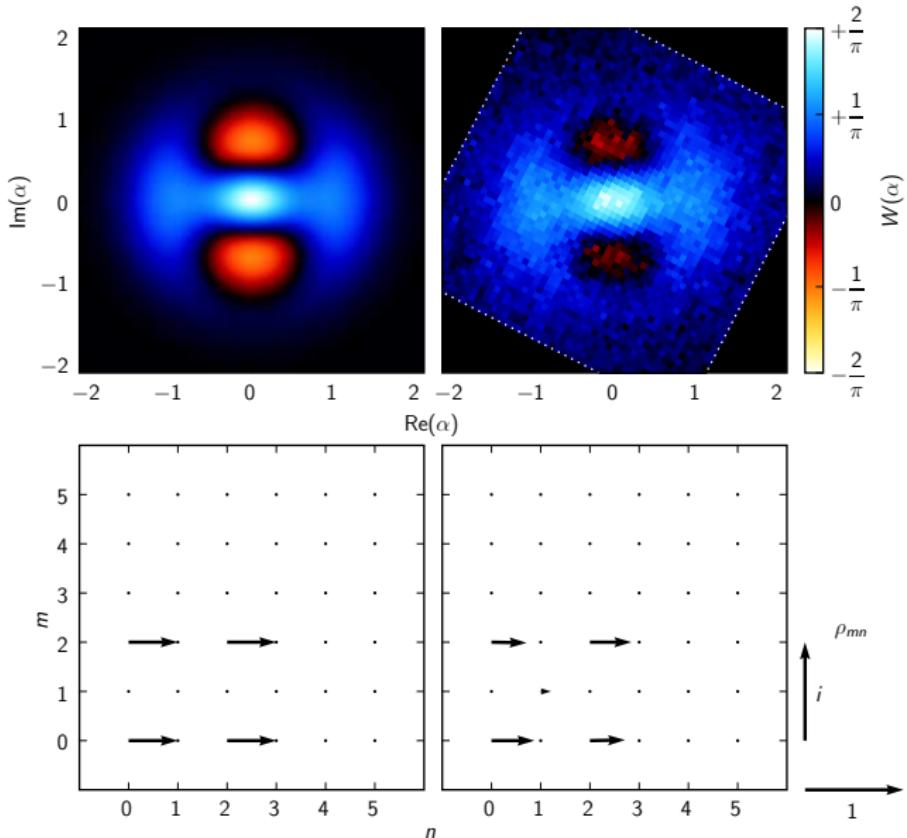


Fidelity:

$$F = \sqrt{\langle \psi | \rho | \psi \rangle} = 0.92$$

Controlling superpositions

$$|\psi\rangle = |0\rangle + |2\rangle$$

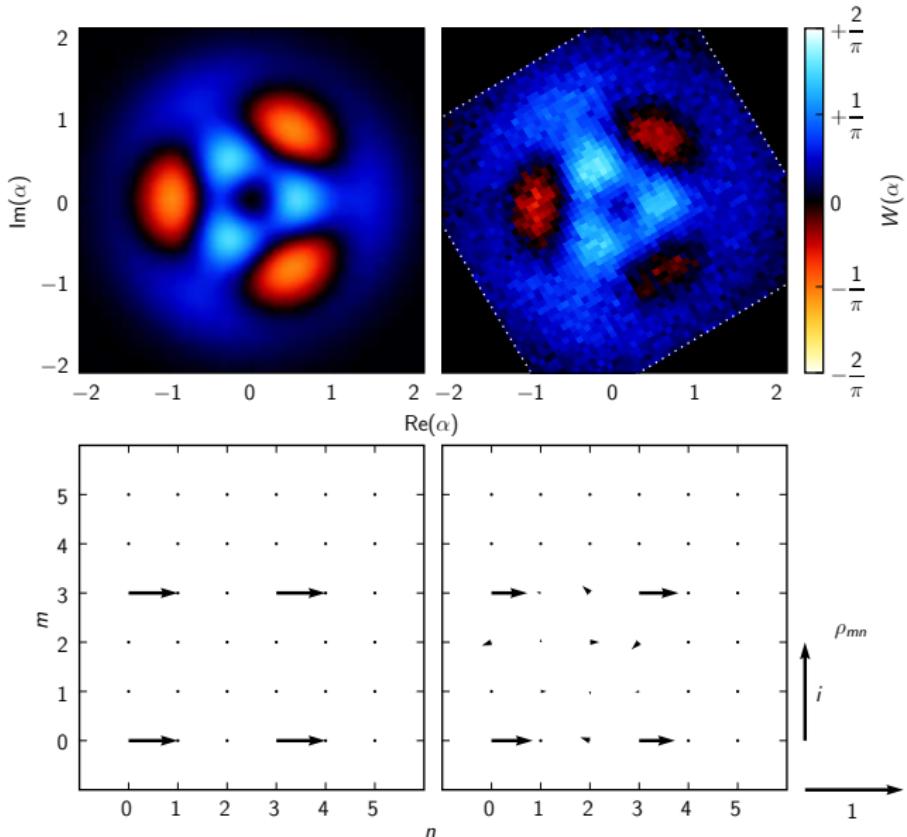


Fidelity:

$$F = \sqrt{\langle \psi | \rho | \psi \rangle} = 0.89$$

Controlling superpositions

$$|\psi\rangle = |0\rangle + |3\rangle$$

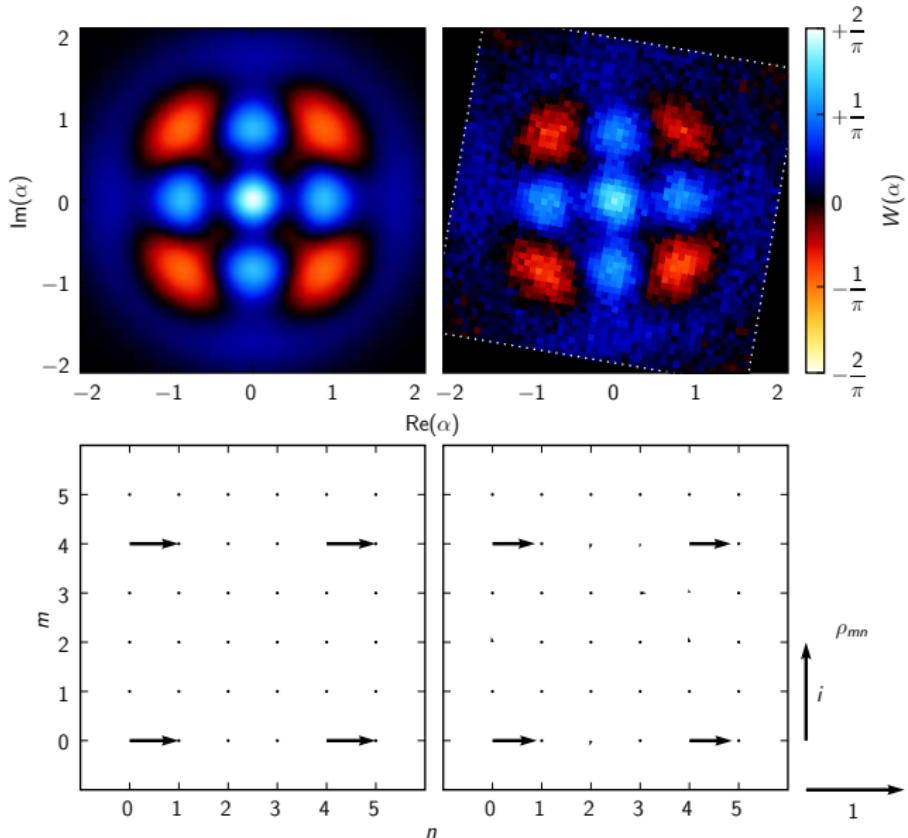


Fidelity:

$$F = \sqrt{\langle \psi | \rho | \psi \rangle} = 0.88$$

Controlling superpositions

$$|\psi\rangle = |0\rangle + |4\rangle$$

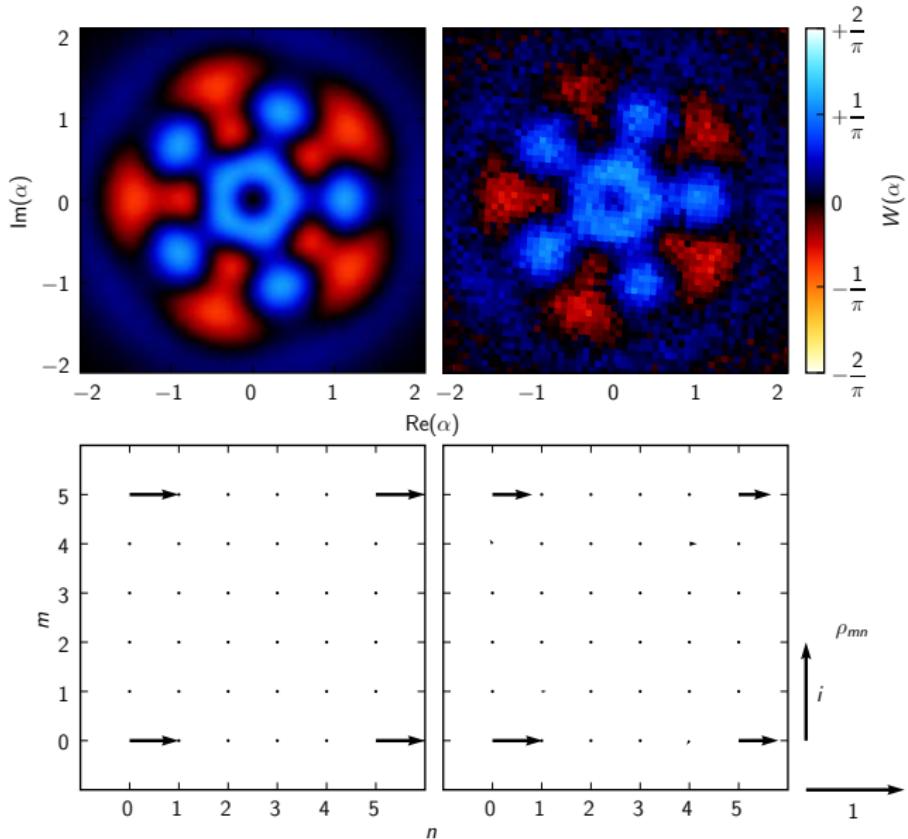


Fidelity:

$$F = \sqrt{\langle \psi | \rho | \psi \rangle} = 0.94$$

Controlling superpositions

$$|\psi\rangle = |0\rangle + |5\rangle$$

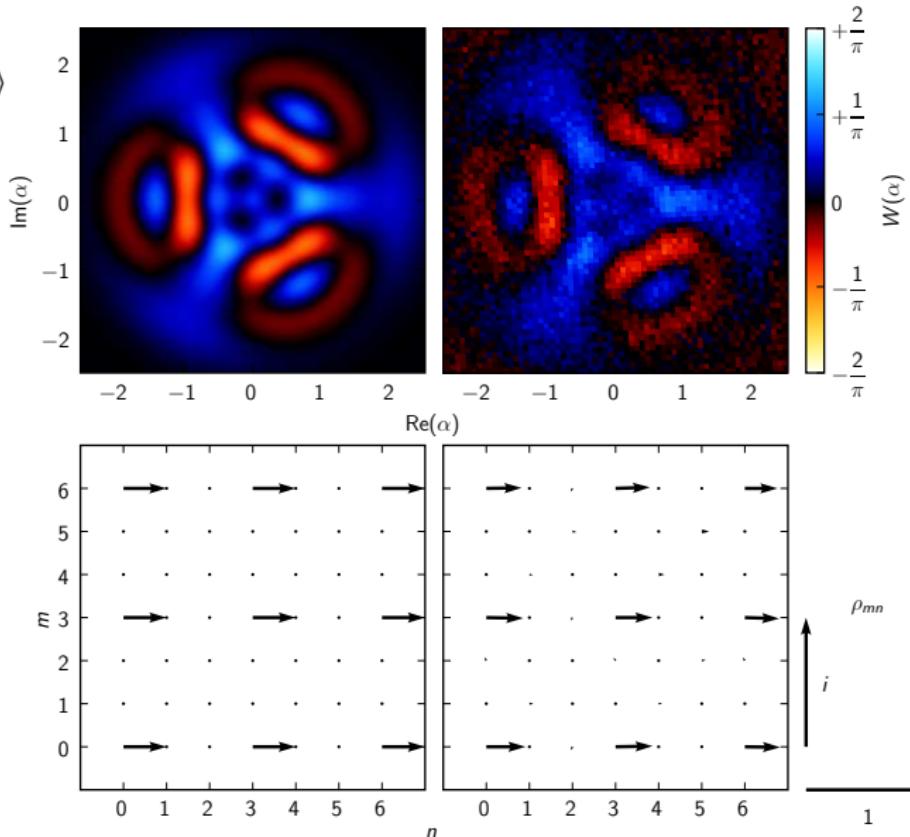


Fidelity:

$$F = \sqrt{\langle \psi | \rho | \psi \rangle} = 0.91$$

Controlling phase

$$|\psi\rangle = |0\rangle + e^{\frac{0}{8}i\pi}|3\rangle + |6\rangle$$

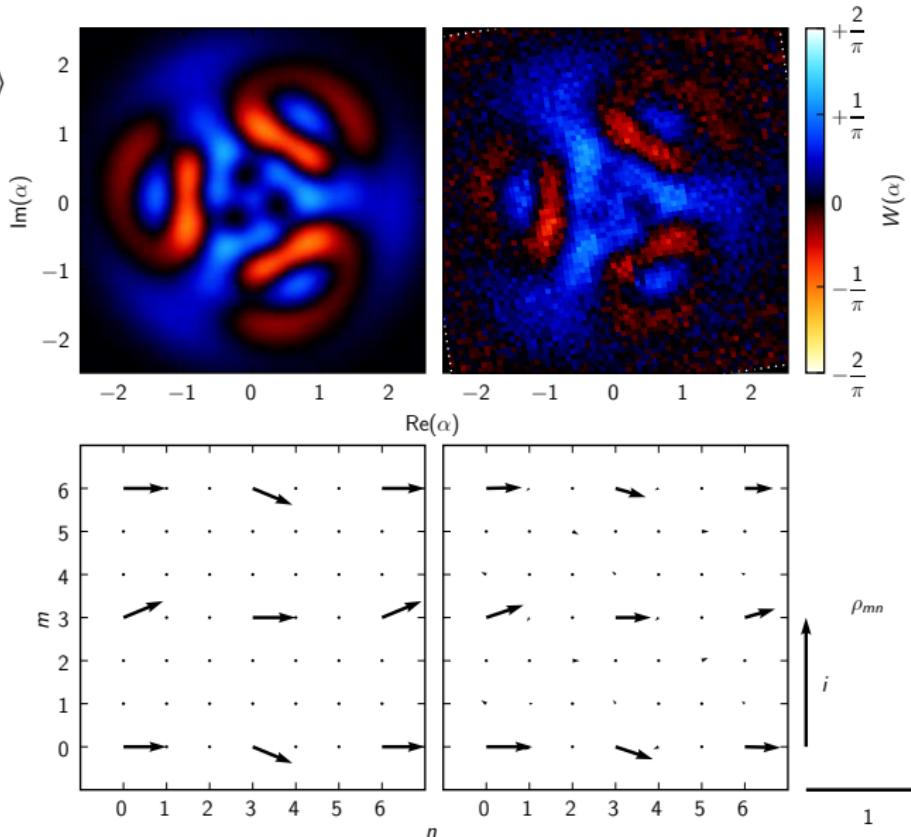


Fidelity:

$$F = \sqrt{\langle \psi | \rho | \psi \rangle} = 0.91$$

Controlling phase

$$|\psi\rangle = |0\rangle + e^{\frac{1}{8}i\pi}|3\rangle + |6\rangle$$

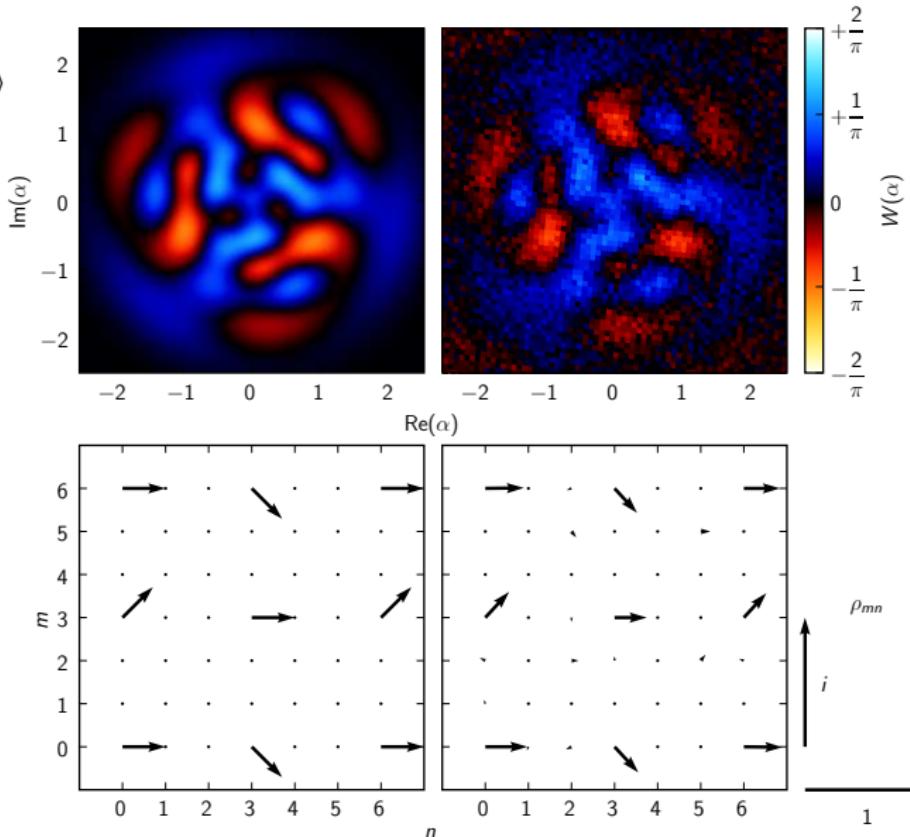


Fidelity:

$$F = \sqrt{\langle \psi | \rho | \psi \rangle} = 0.91$$

Controlling phase

$$|\psi\rangle = |0\rangle + e^{\frac{2}{8}i\pi}|3\rangle + |6\rangle$$

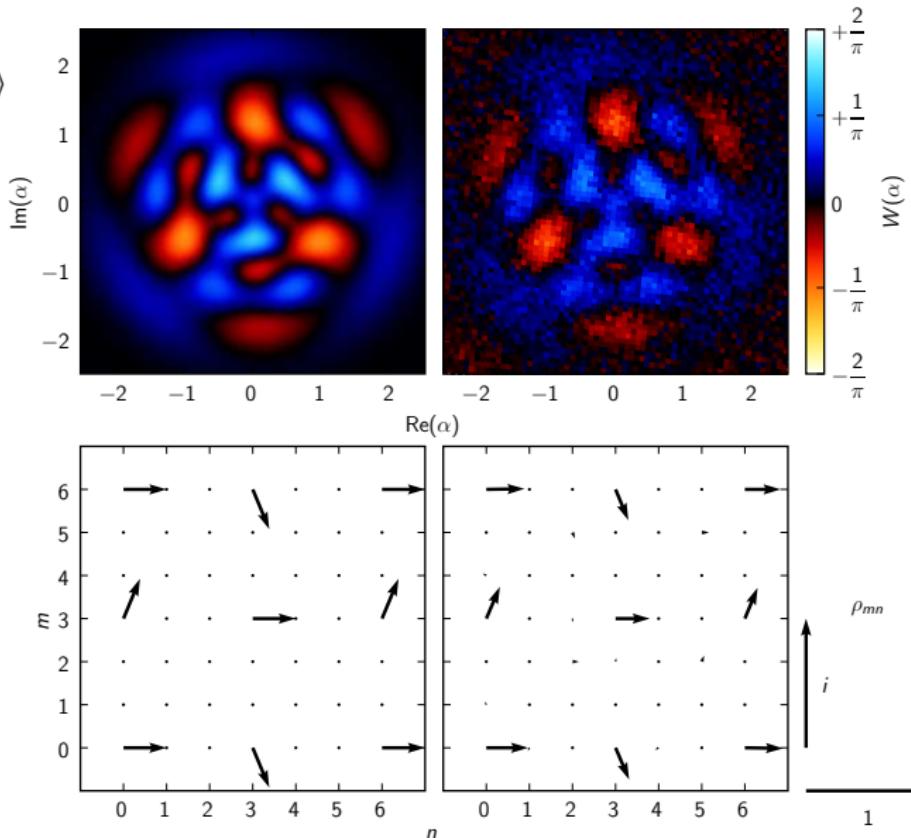


Fidelity:

$$F = \sqrt{\langle \psi | \rho | \psi \rangle} = 0.91$$

Controlling phase

$$|\psi\rangle = |0\rangle + e^{\frac{3}{8}i\pi}|3\rangle + |6\rangle$$

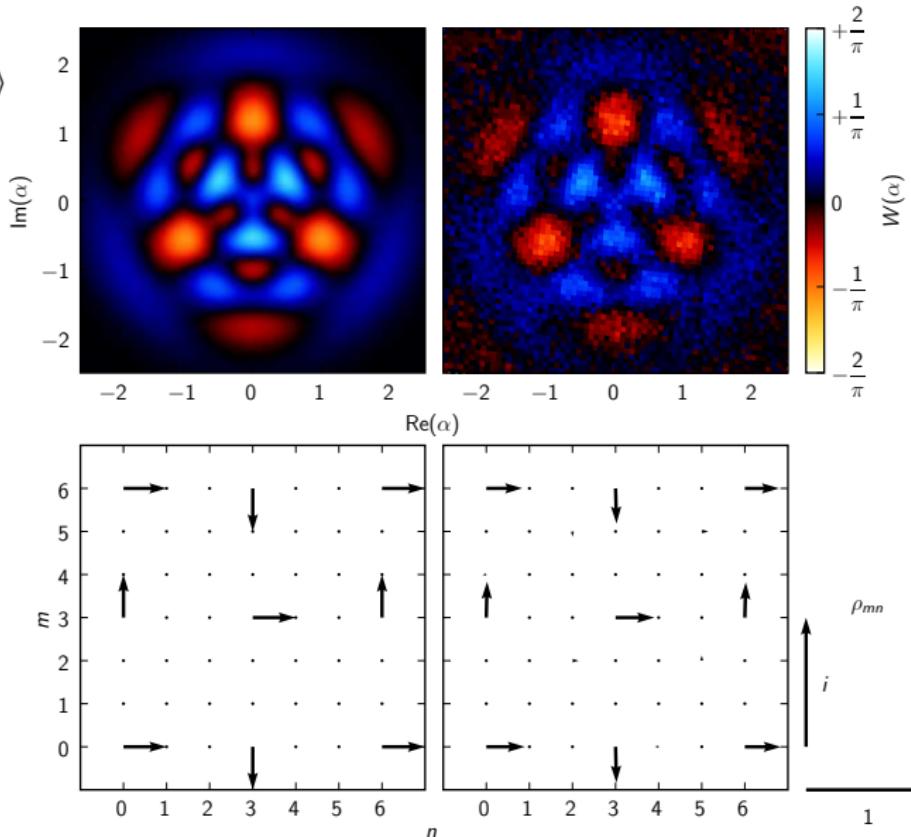


Fidelity:

$$F = \sqrt{\langle \psi | \rho | \psi \rangle} = 0.92$$

Controlling phase

$$|\psi\rangle = |0\rangle + e^{\frac{4}{8}i\pi}|3\rangle + |6\rangle$$

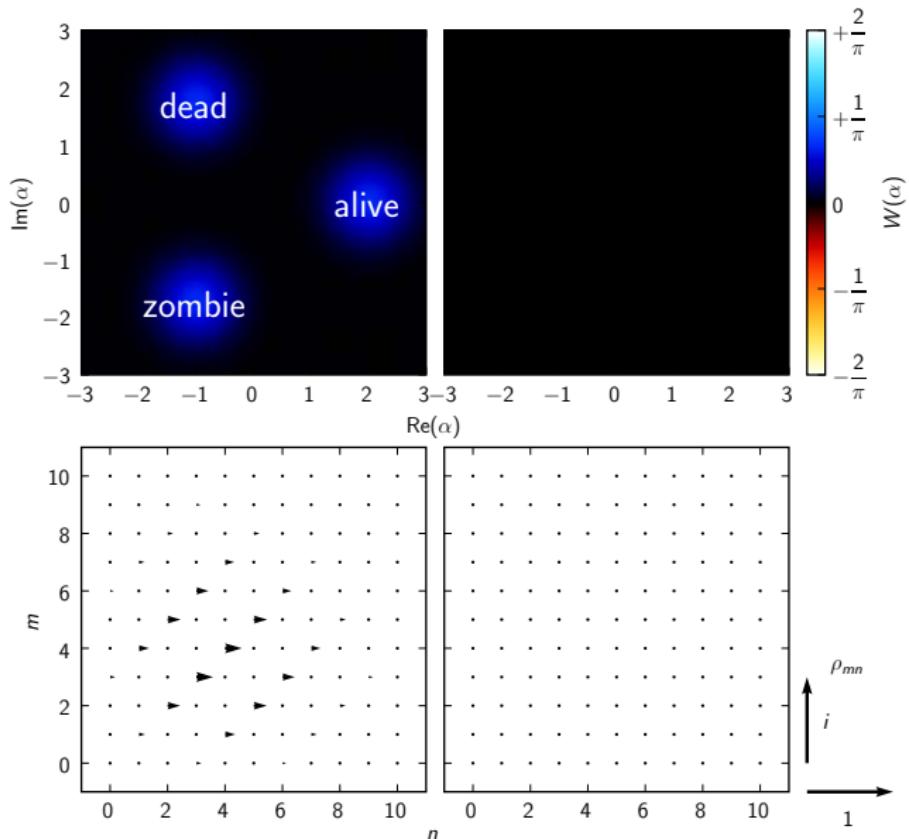


Fidelity:

$$F = \sqrt{\langle \psi | \rho | \psi \rangle} = 0.92$$

Voodoo cat

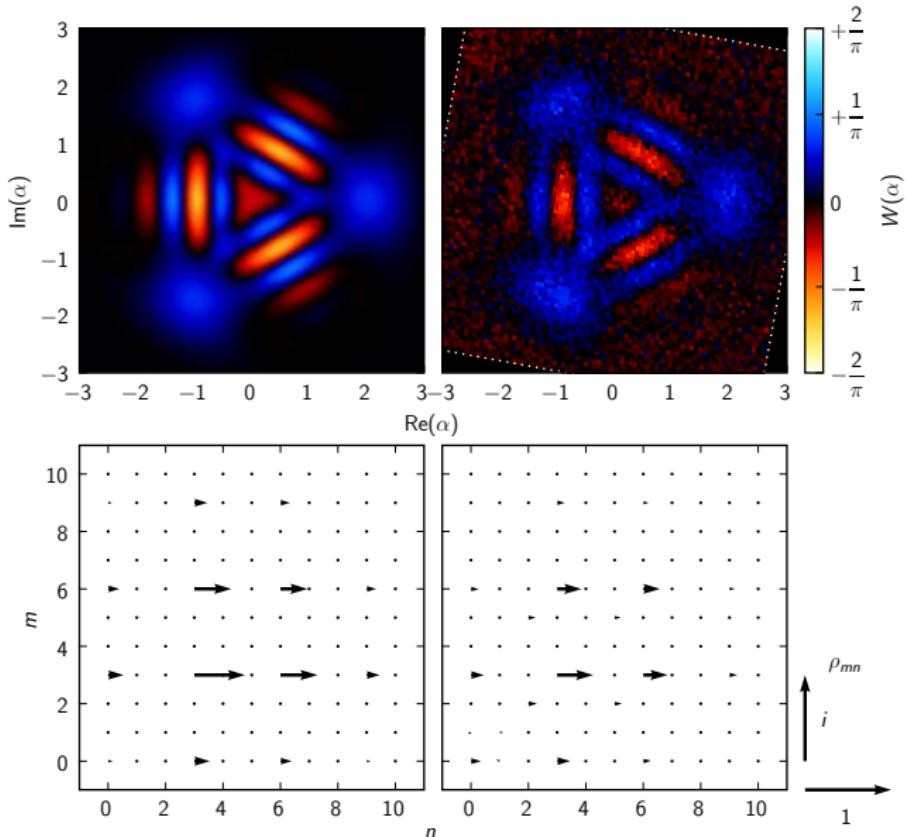
$$|\psi_k\rangle = |2e^{i\frac{2\pi}{3}k}\rangle$$



Voodoo cat

$$|\psi\rangle = \sum_{k=0,1,2} |2e^{i\frac{2\pi}{3}k}\rangle$$

$$\propto \sum_{n=0,3,6,9\dots} \frac{2^n}{\sqrt{n!}} |n\rangle$$

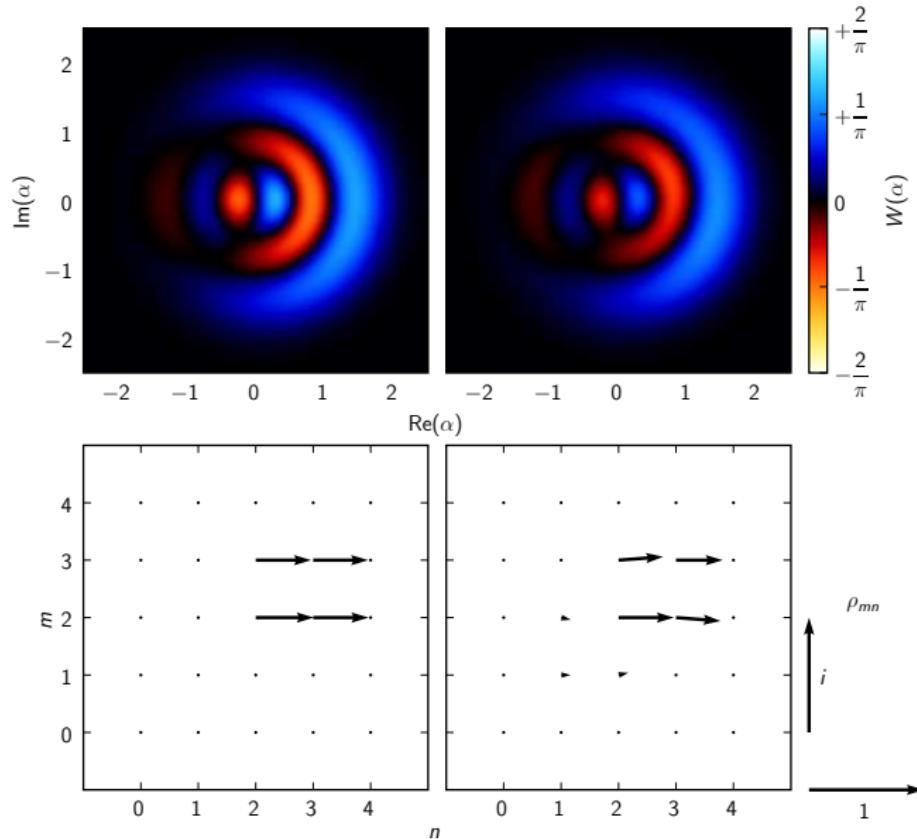


Fidelity:

$$F = \sqrt{\langle\psi|\rho|\psi\rangle} = 0.83$$

Decay of resonator states (Haohua Wang)

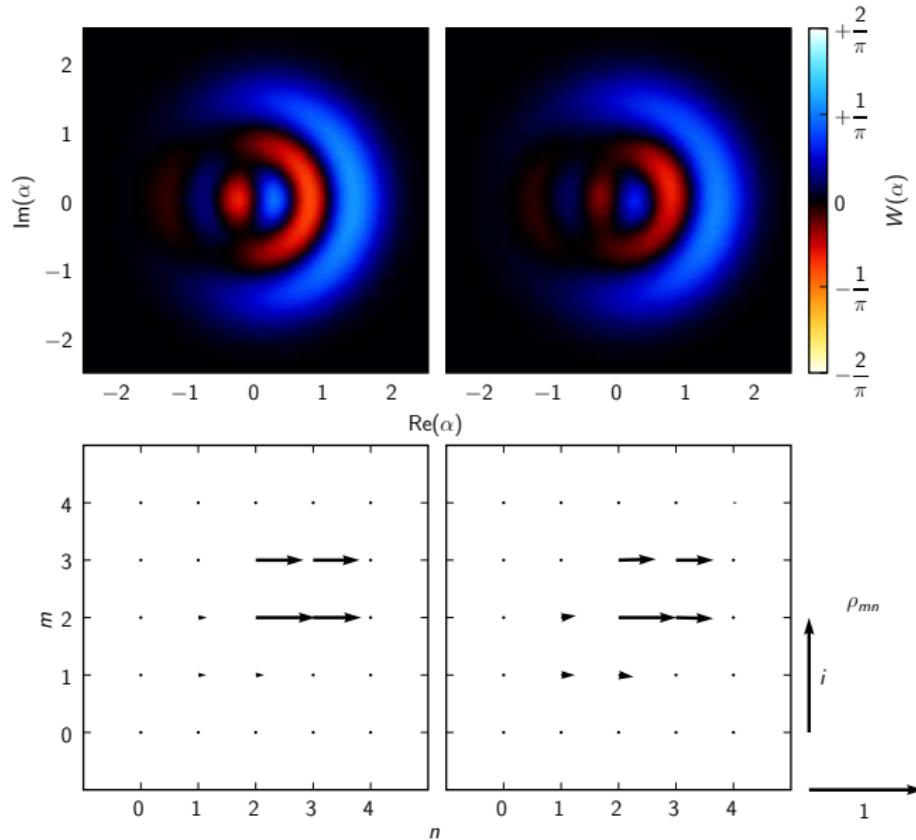
$$|\psi\rangle = |2\rangle + |3\rangle$$



Delay:
 $\tau = 0.05 \mu\text{s}$

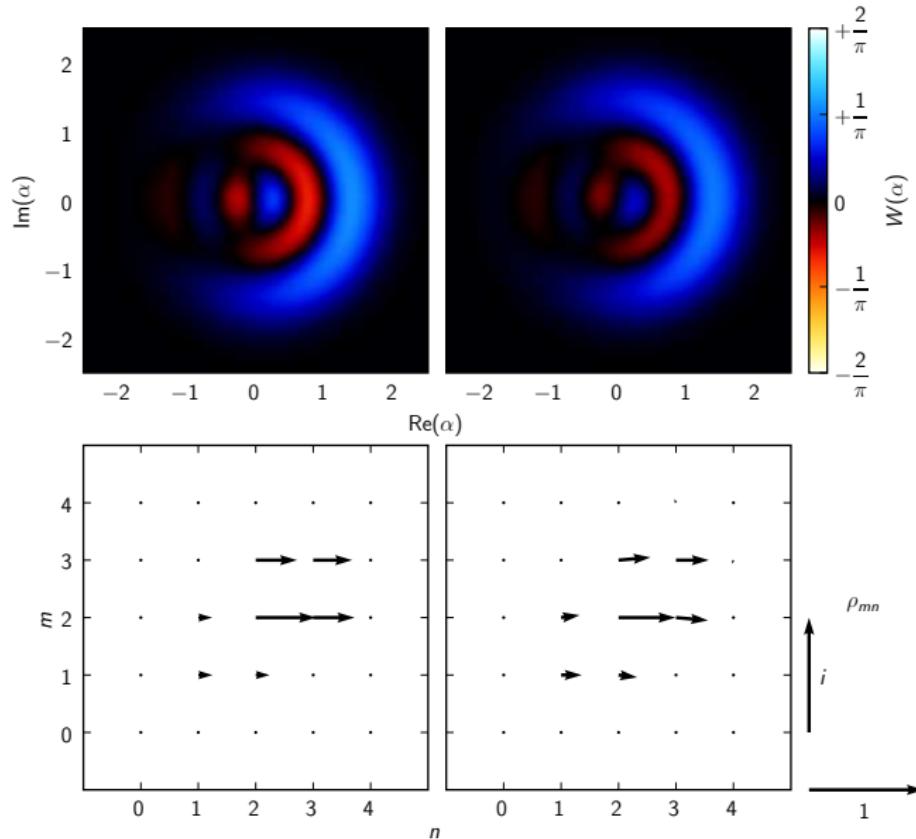
Decay of resonator states (Haohua Wang)

$$|\psi\rangle = |2\rangle + |3\rangle$$



Decay of resonator states (Haohua Wang)

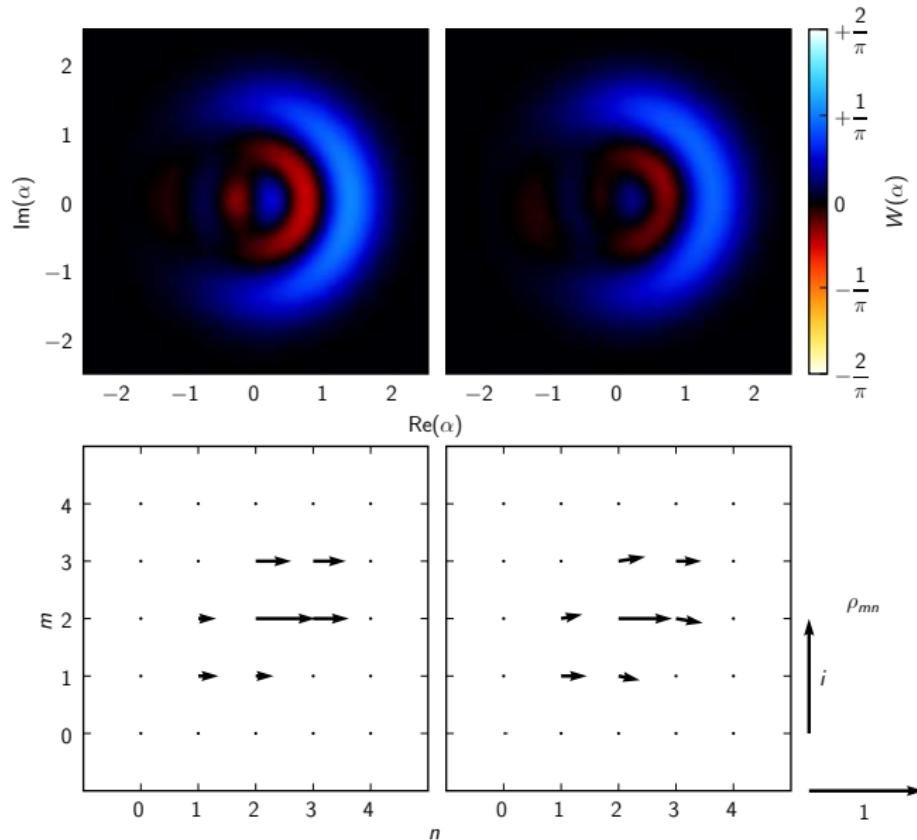
$$|\psi\rangle = |2\rangle + |3\rangle$$



Delay:
 $\tau = 0.35 \mu\text{s}$

Decay of resonator states (Haohua Wang)

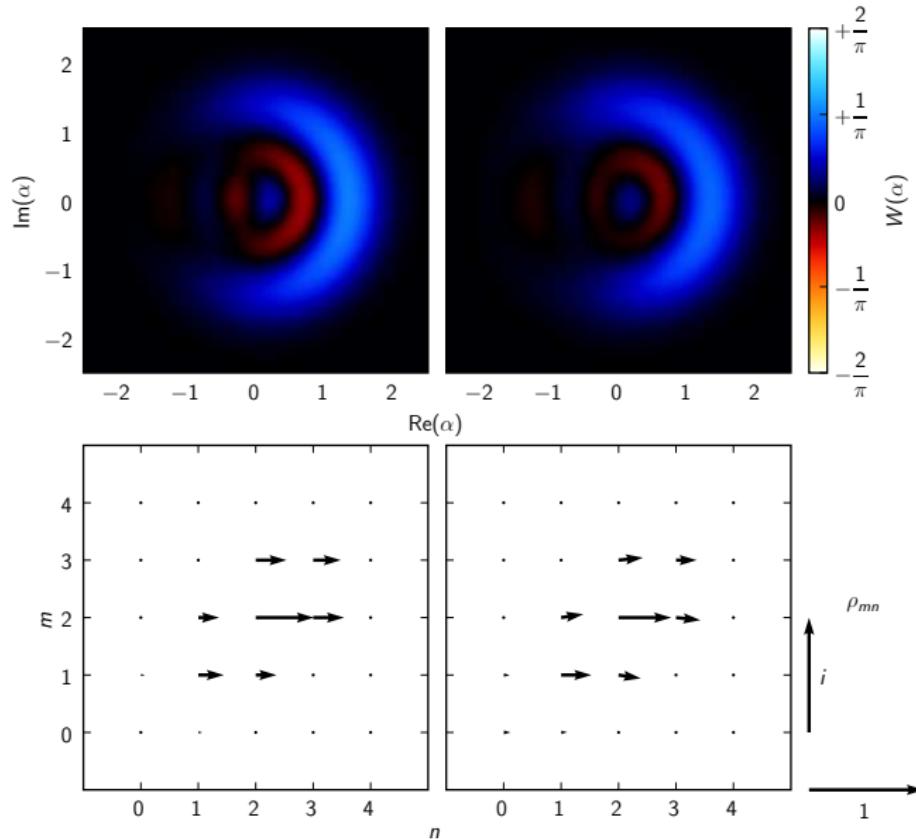
$$|\psi\rangle = |2\rangle + |3\rangle$$



Delay:
 $\tau = 0.50 \mu\text{s}$

Decay of resonator states (Haohua Wang)

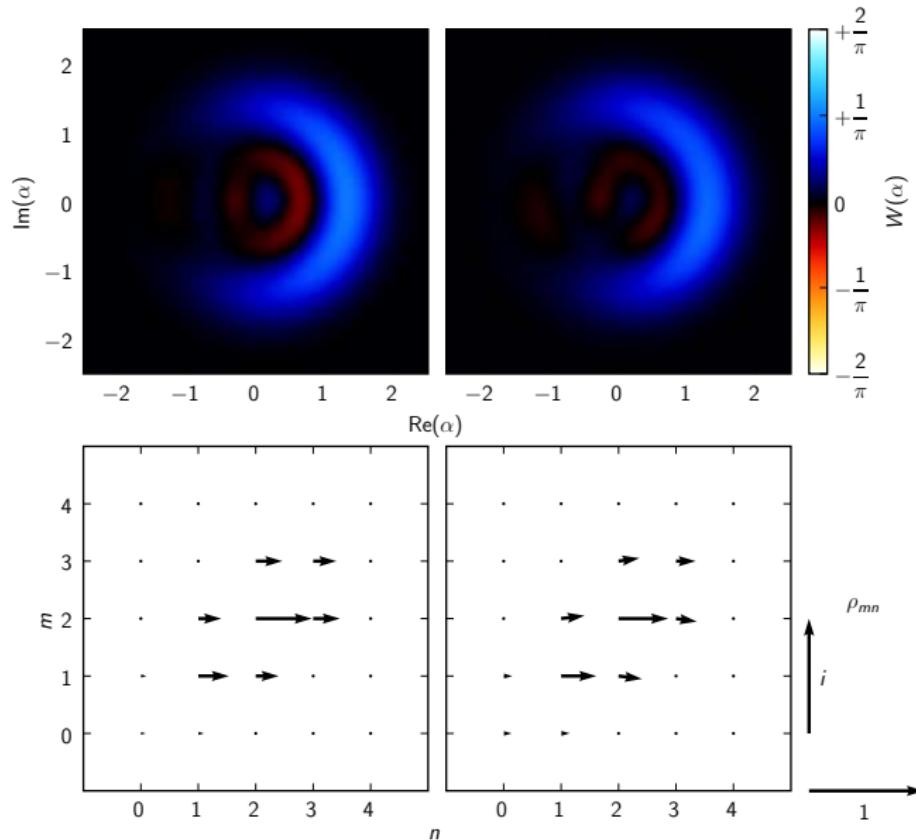
$$|\psi\rangle = |2\rangle + |3\rangle$$



Delay:
 $\tau = 0.65 \mu\text{s}$

Decay of resonator states (Haohua Wang)

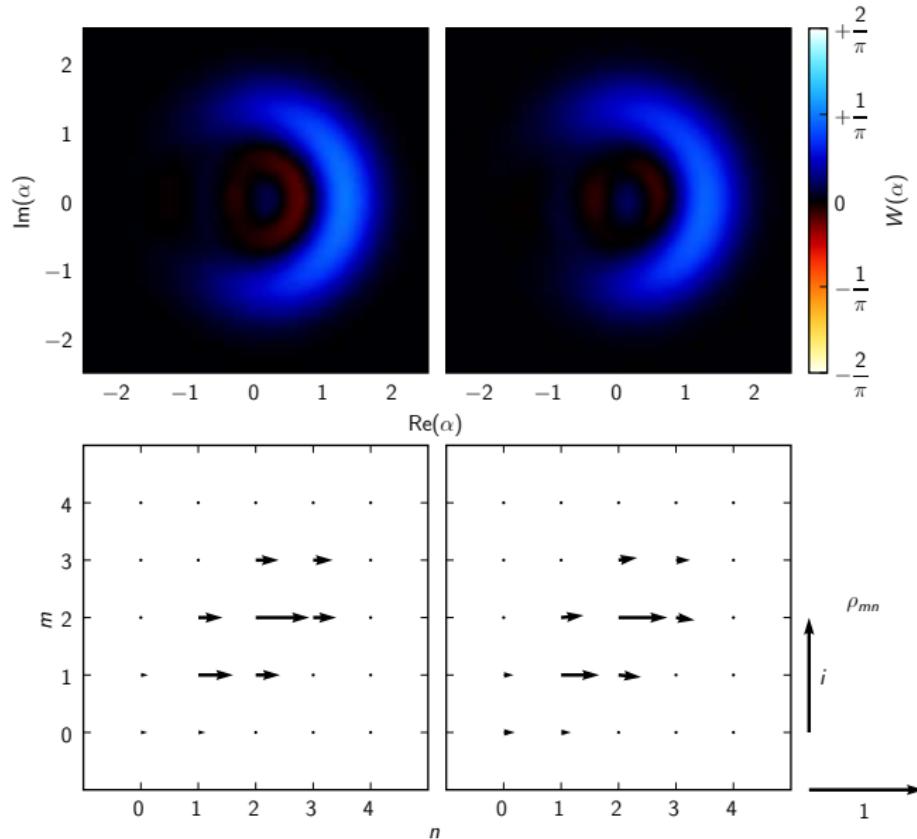
$$|\psi\rangle = |2\rangle + |3\rangle$$



Delay:
 $\tau = 0.80 \mu\text{s}$

Decay of resonator states (Haohua Wang)

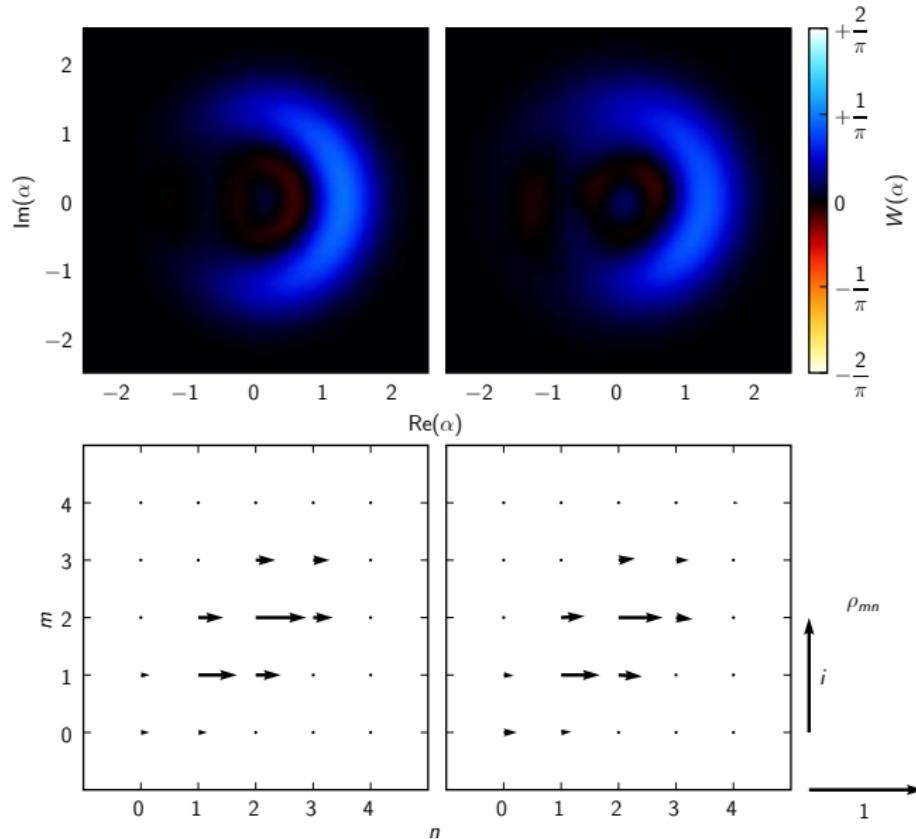
$$|\psi\rangle = |2\rangle + |3\rangle$$



Delay:
 $\tau = 0.95 \mu\text{s}$

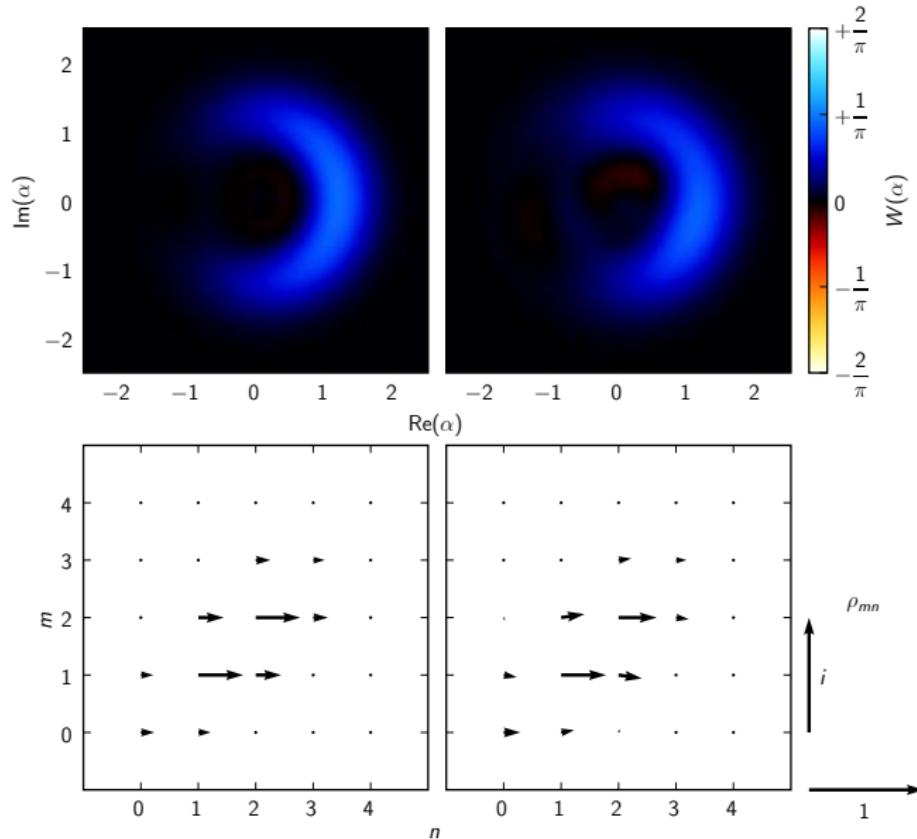
Decay of resonator states (Haohua Wang)

$$|\psi\rangle = |2\rangle + |3\rangle$$



Decay of resonator states (Haohua Wang)

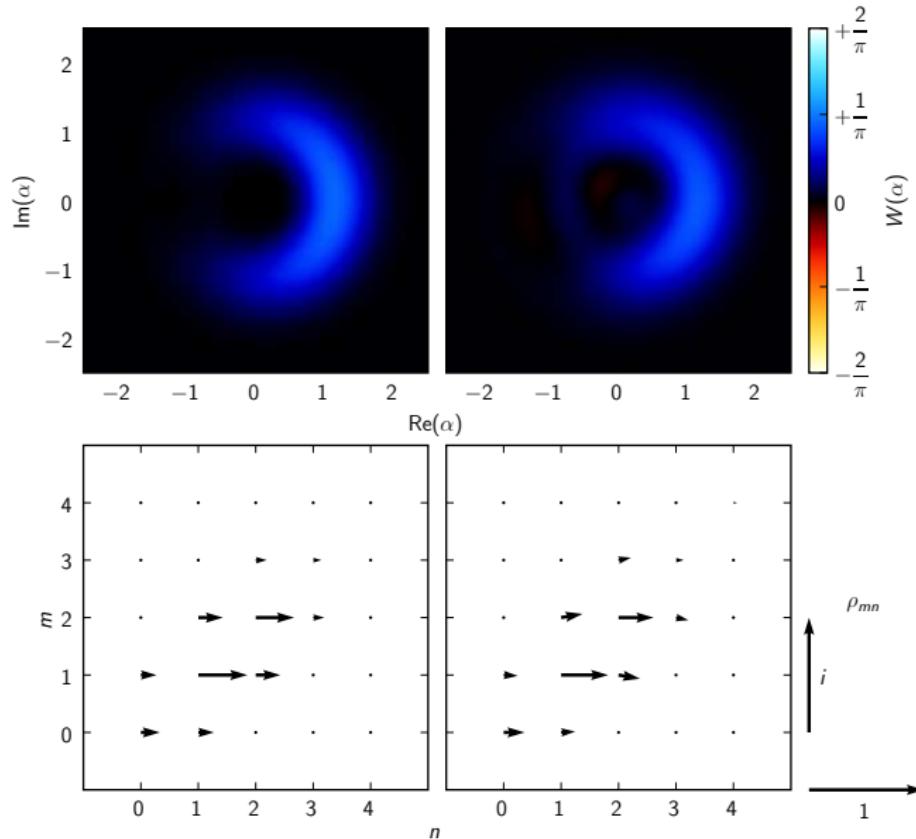
$$|\psi\rangle = |2\rangle + |3\rangle$$



Delay:
 $\tau = 1.4 \mu\text{s}$

Decay of resonator states (Haohua Wang)

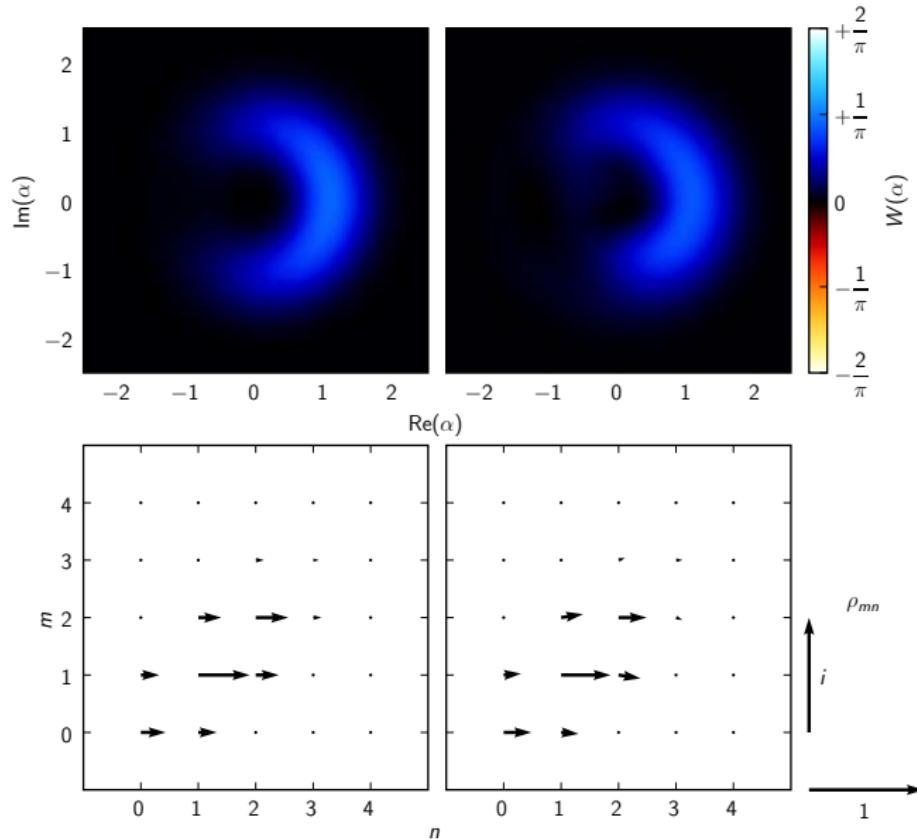
$$|\psi\rangle = |2\rangle + |3\rangle$$



Delay:
 $\tau = 1.7 \mu\text{s}$

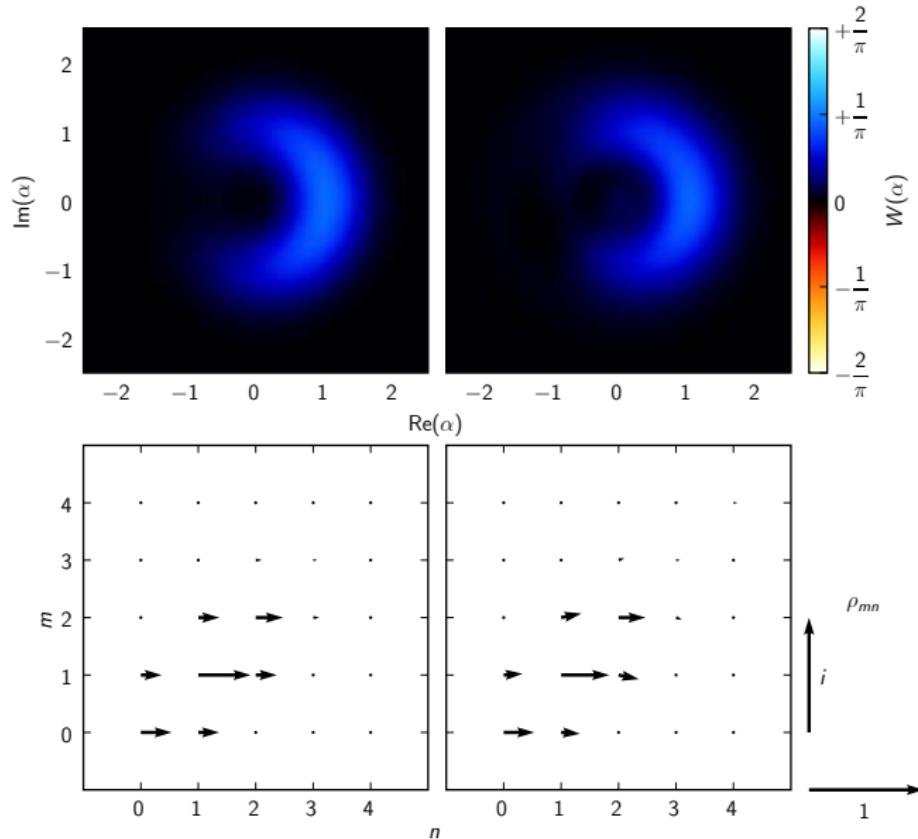
Decay of resonator states (Haohua Wang)

$$|\psi\rangle = |2\rangle + |3\rangle$$



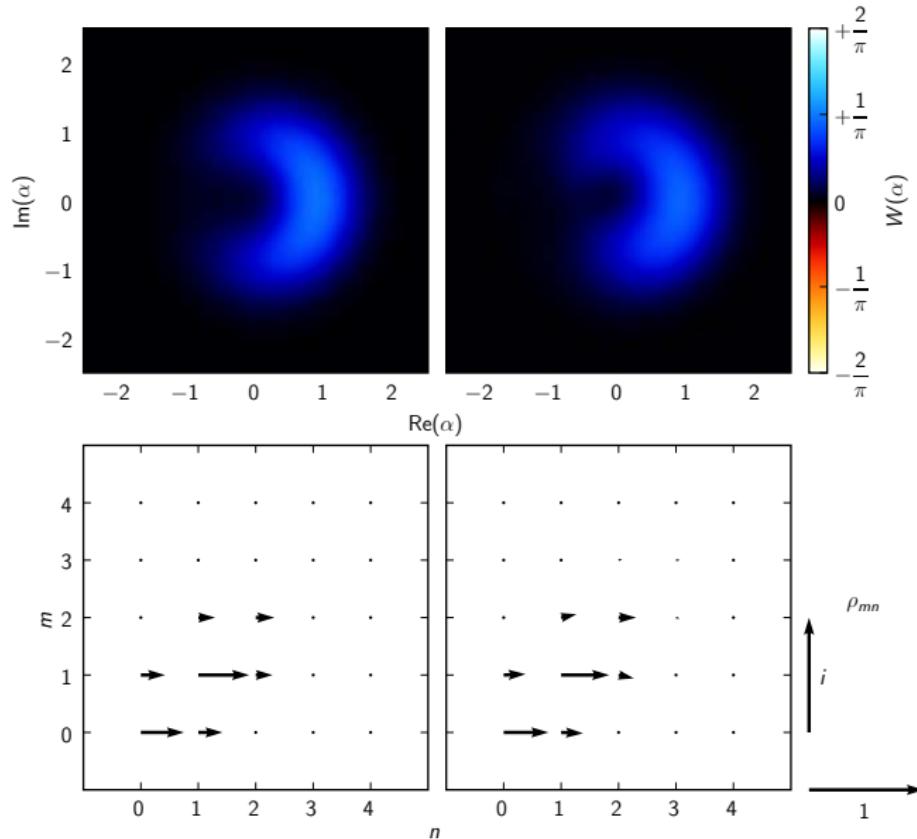
Decay of resonator states (Haohua Wang)

$$|\psi\rangle = |2\rangle + |3\rangle$$



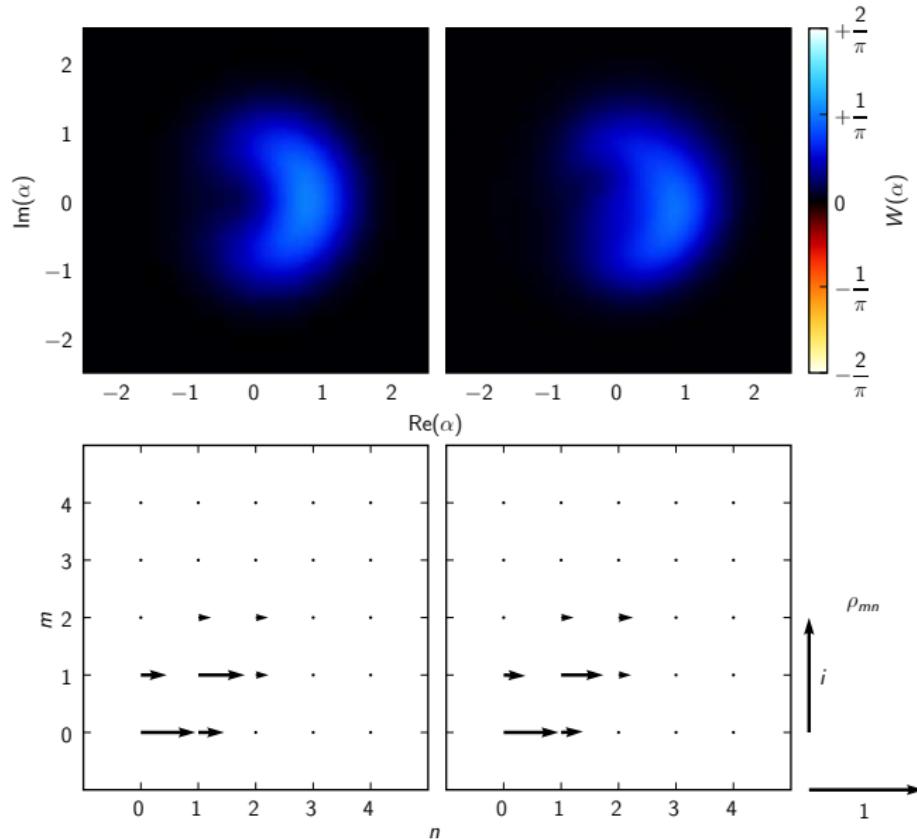
Decay of resonator states (Haohua Wang)

$$|\psi\rangle = |2\rangle + |3\rangle$$



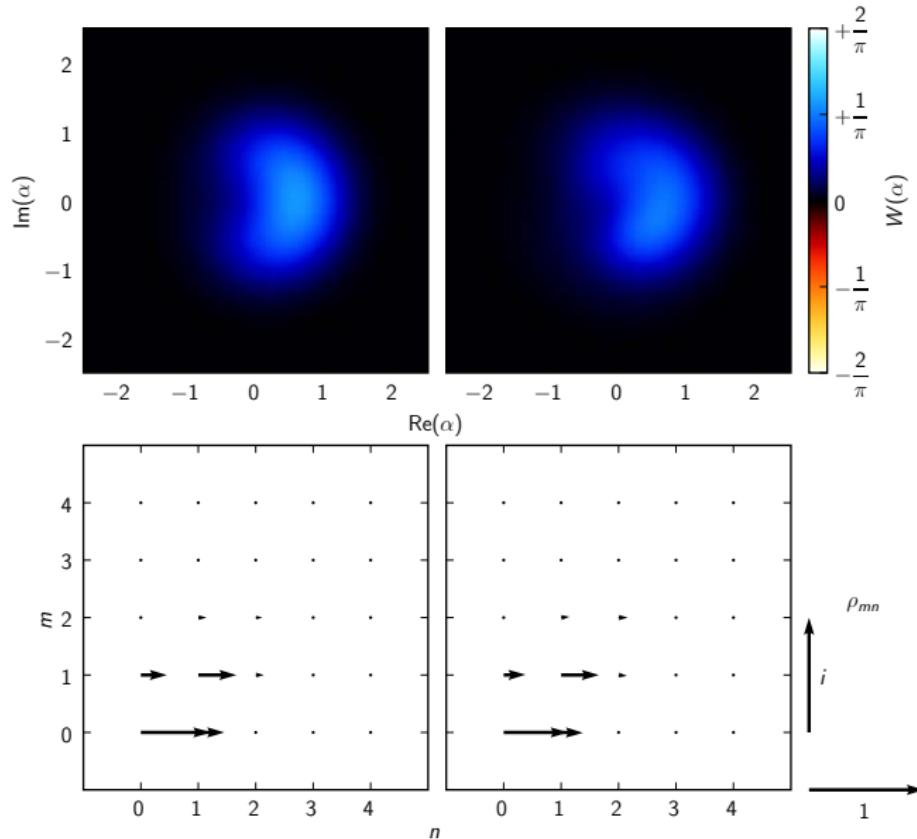
Decay of resonator states (Haohua Wang)

$$|\psi\rangle = |2\rangle + |3\rangle$$



Decay of resonator states (Haohua Wang)

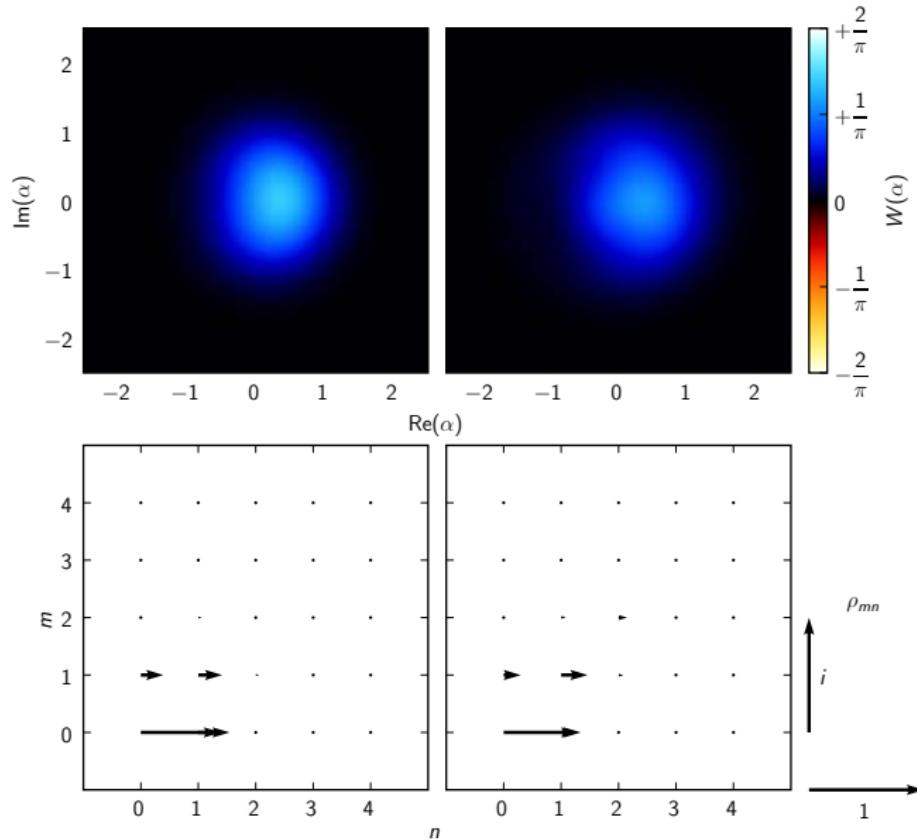
$$|\psi\rangle = |2\rangle + |3\rangle$$



Delay:
 $\tau = 4.4 \mu\text{s}$

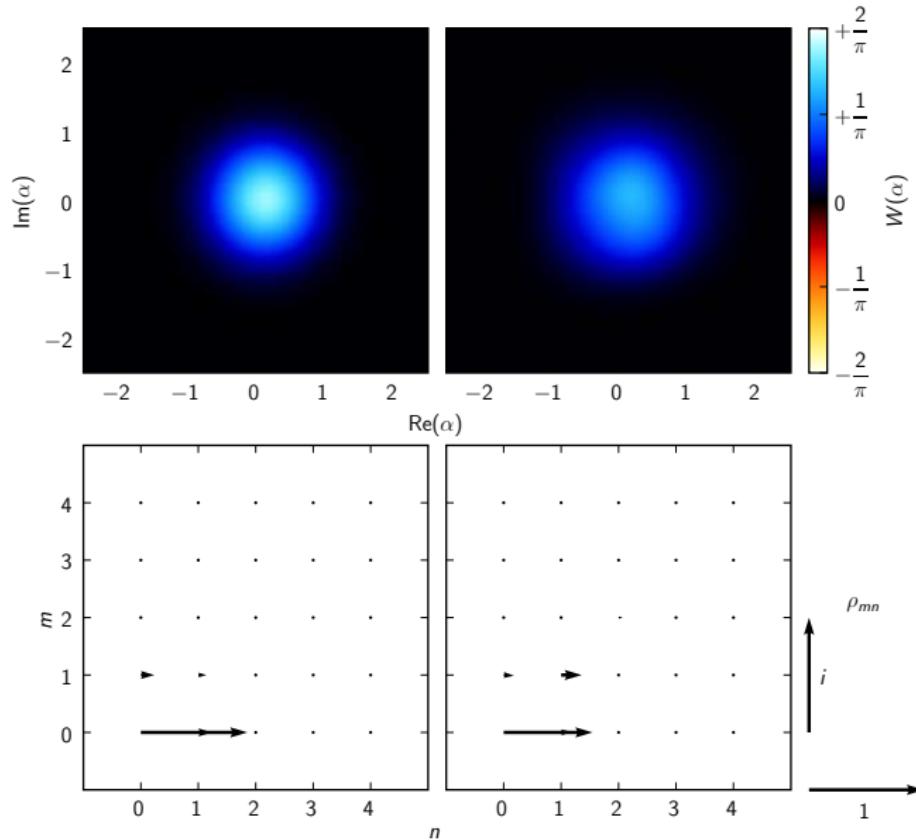
Decay of resonator states (Haohua Wang)

$$|\psi\rangle = |2\rangle + |3\rangle$$



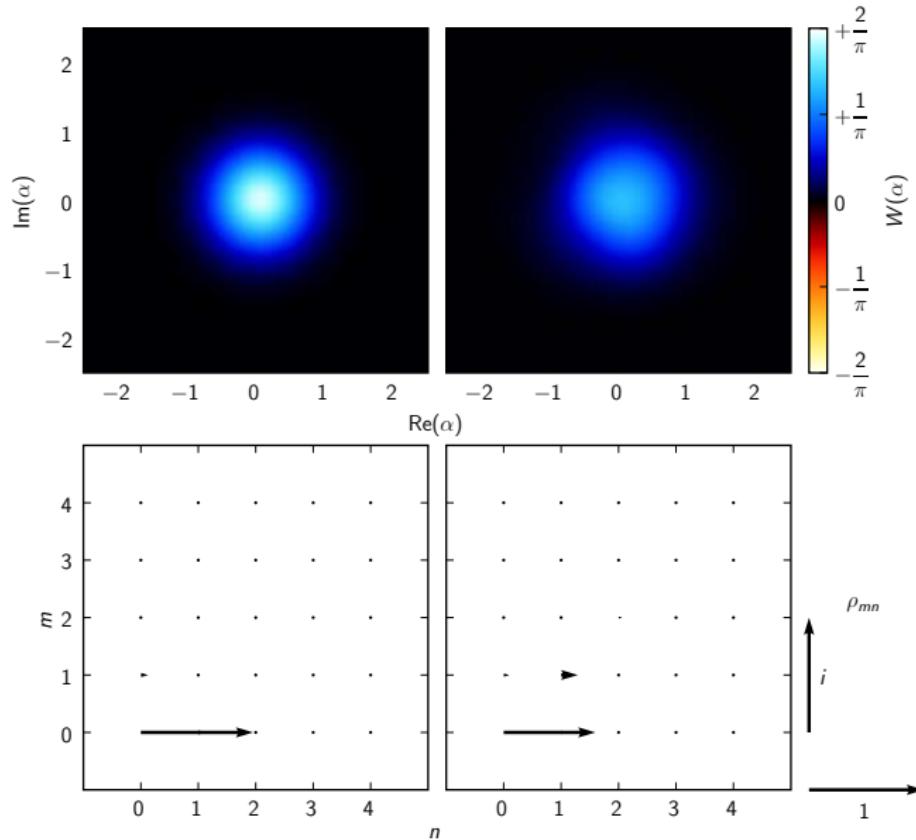
Decay of resonator states (Haohua Wang)

$$|\psi\rangle = |2\rangle + |3\rangle$$



Decay of resonator states (Haohua Wang)

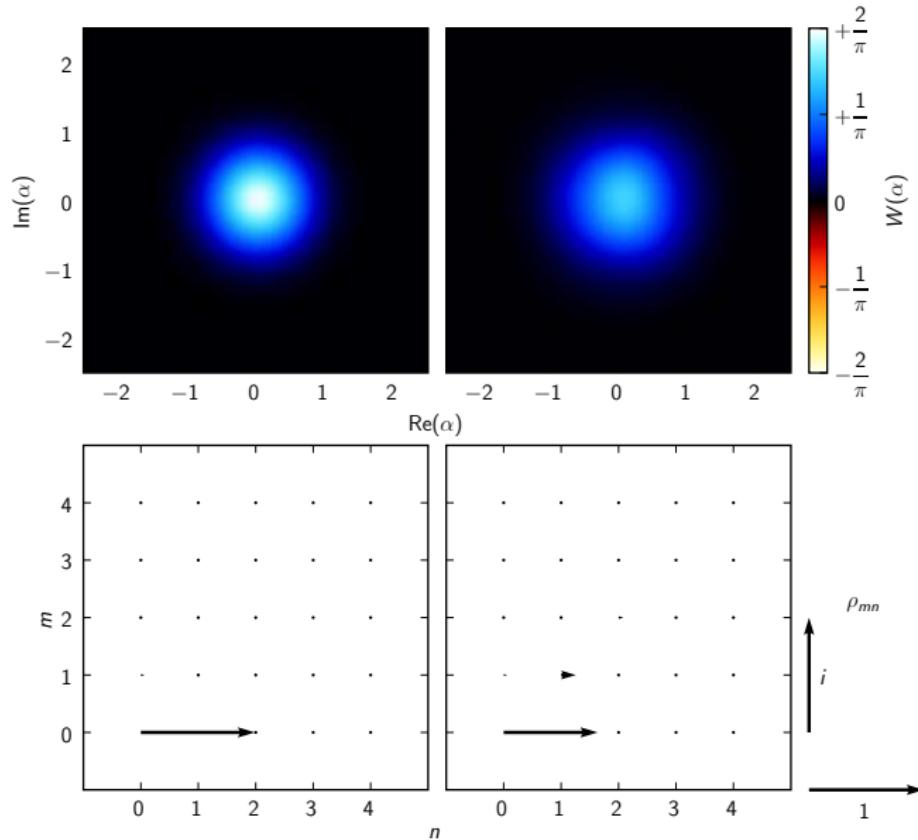
$$|\psi\rangle = |2\rangle + |3\rangle$$



Delay:
 $\tau = 12 \mu\text{s}$

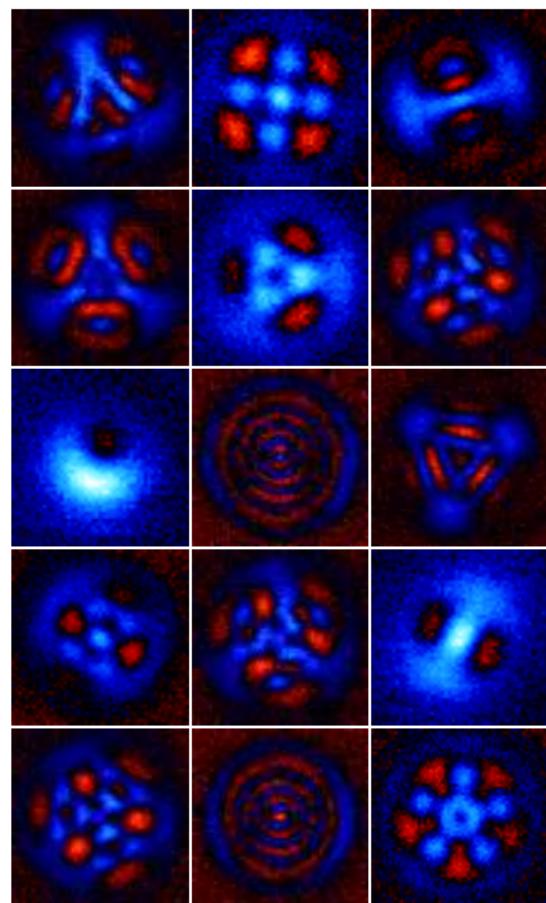
Decay of resonator states (Haohua Wang)

$$|\psi\rangle = |2\rangle + |3\rangle$$



Delay:
 $\tau = 15 \mu\text{s}$

Conclusions



Full control over qubit state extended to resonator:

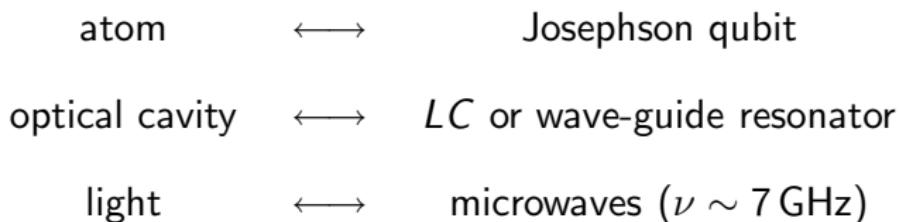
- ▶ arbitrary states
($N < 10$ due to coherence)
- ▶ deterministic generation
- ▶ full characterization

Max Hofheinz, H. Wang, M. Ansmann,
R. Bialczak, E. Lucero, M. Neeley, A. O'Connell,
D. Sank, M. Weides, J. Wenner, J.M. Martinis,
A.N. Cleland

- Fock states: Nature **454**, 310–314 (2008)
~ decoherence: H. Wang, PRL **101**, 240401 (2008)
Arbitrary states: Nature **459**, 546–549 (2009)
~ decoherence: H. Wang, PRL **103**, 200404 (2009)

Funding:  IARPA,  ARO,  NSF

Superconducting circuits: Microwave quantum optics



Advantages:

- ▶ Qubits and cavities can be coupled at will
- ▶ Microwave signal waveforms can be fully controlled

Drawbacks:

- ▶ Low photon energy requires low temperatures $T < 50\text{ mK}$
- ▶ Large solid state system limits coherence times

Calibrations

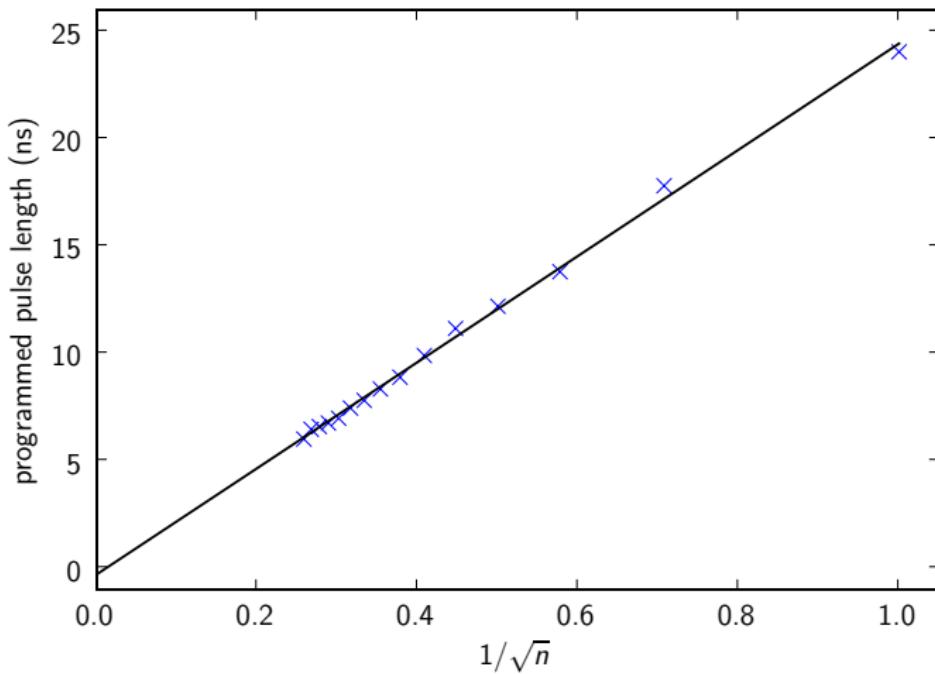
- ▶ impulse response (sampling scope)
- ▶ imperfections of IQ-mixer (spectrum analyzer)
- ▶ cable response (qubit as sampling scope)
- ▶ various parameters for readout SQUID
- ▶ π pulses (amplitude and frequency)
- ▶ measure pulse amplitude
- ▶ swap pulse amplitude
- ▶ swap pulse time for each photon number
- ▶ first order correction for finite rise time of swap pulses
- ▶ qubit/resonator dephasing rate when off resonance
- ▶ resonator displacement/drive amplitude ratio
- ▶ resonator drive phase
- ▶ readout visibility

BUT:

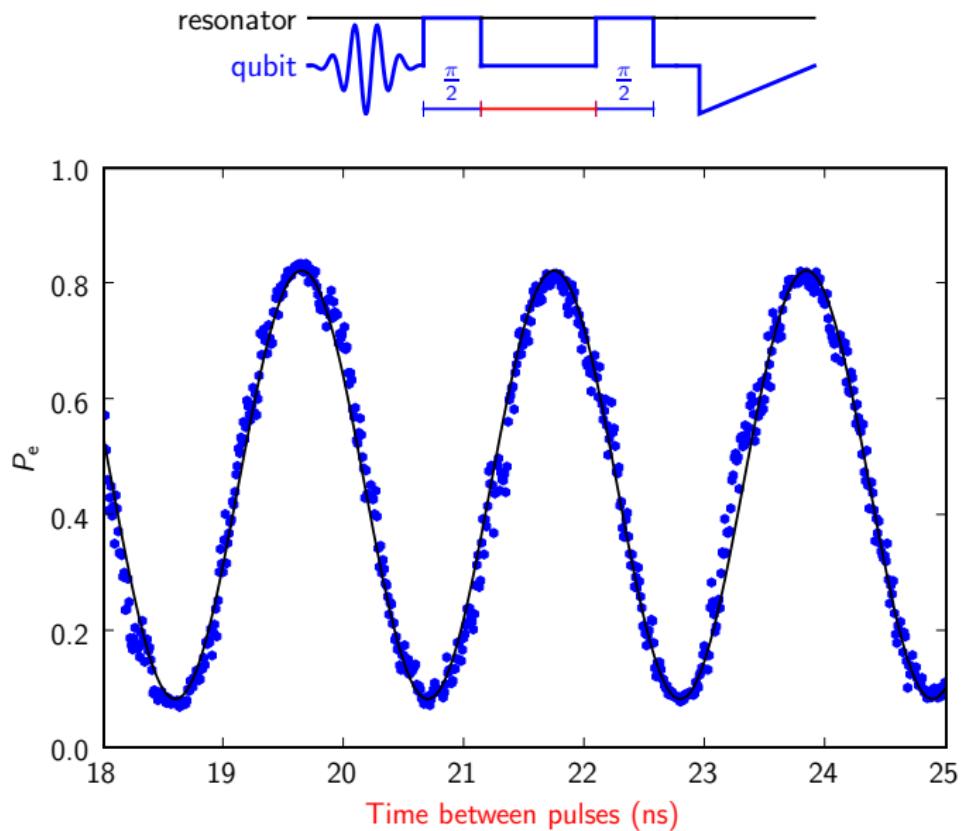
- ▶ no per-state calibrations; just run the calculated sequence

Calibrating the swap pulse length

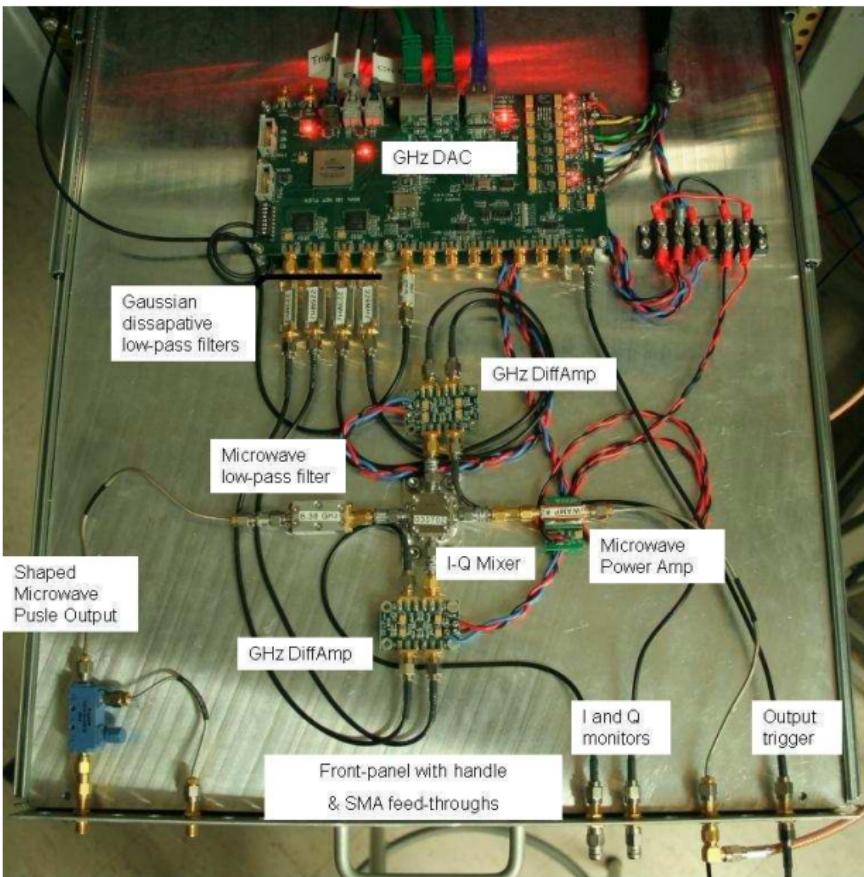
Calibrate with the pulse lengths from the Fock state tune-up.



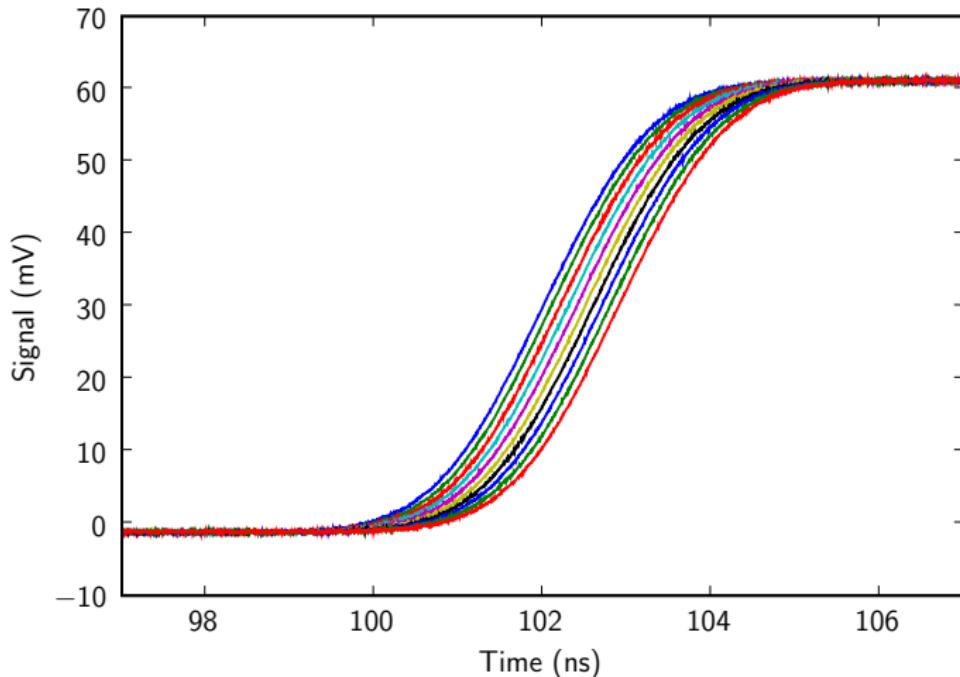
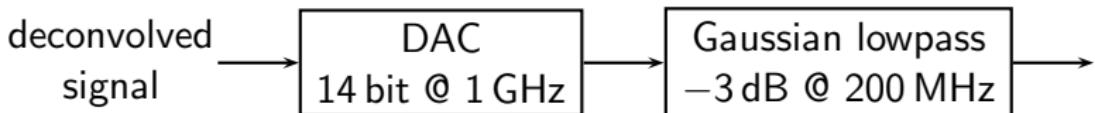
Detuning time calibration: Qubit–resonator Ramsey fringes



Custom microwave electronics



Sub-ns timing resolution with a 1 GHz DAC



Sub-ns timing resolution with a 1 GHz DAC

