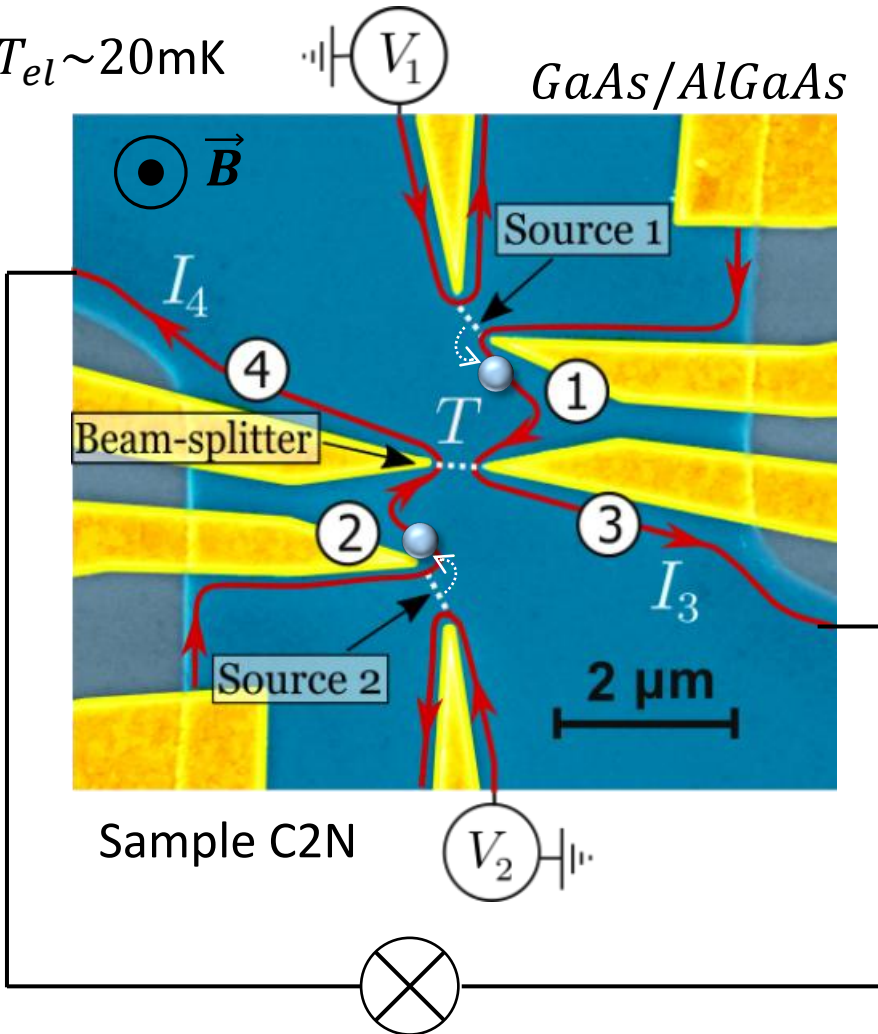


Electron optics experiments in quantum Hall conductors

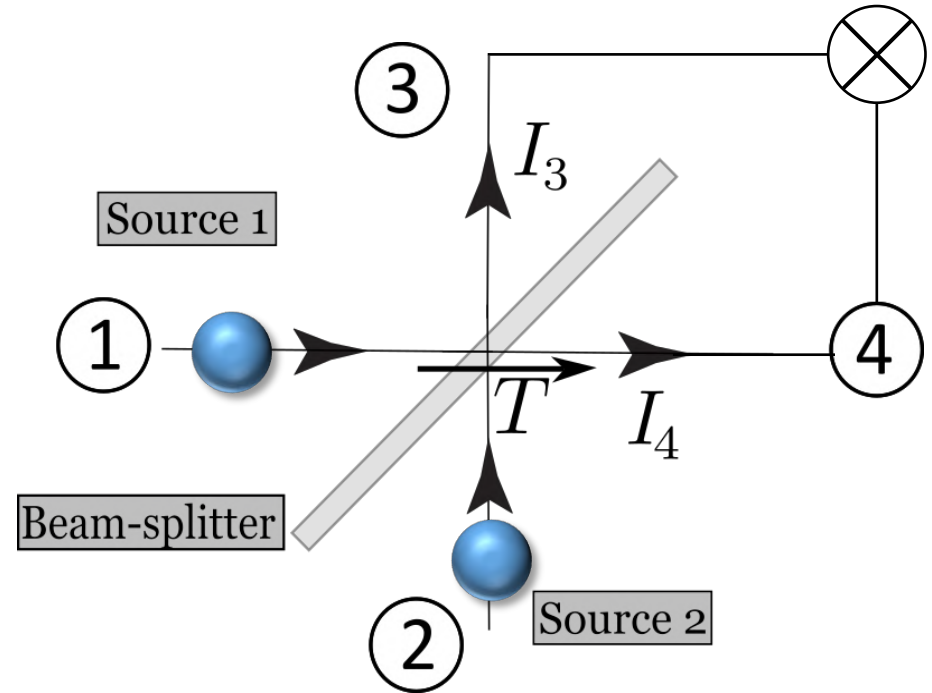
$B \sim \text{few T}$

$T_{el} \sim 20 \text{mK}$



Current correlations $\langle \Delta I_3 \Delta I_4 \rangle$

Current correlations $\langle \Delta I_3 \Delta I_4 \rangle$

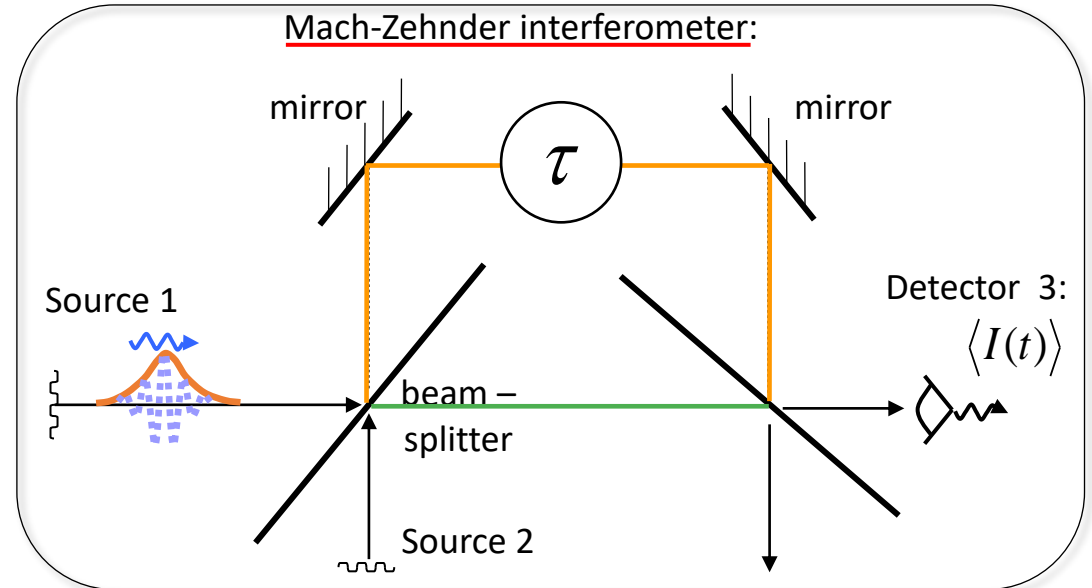


Single-particle vs two-particle interferometry (optics)

Single particle interferometer

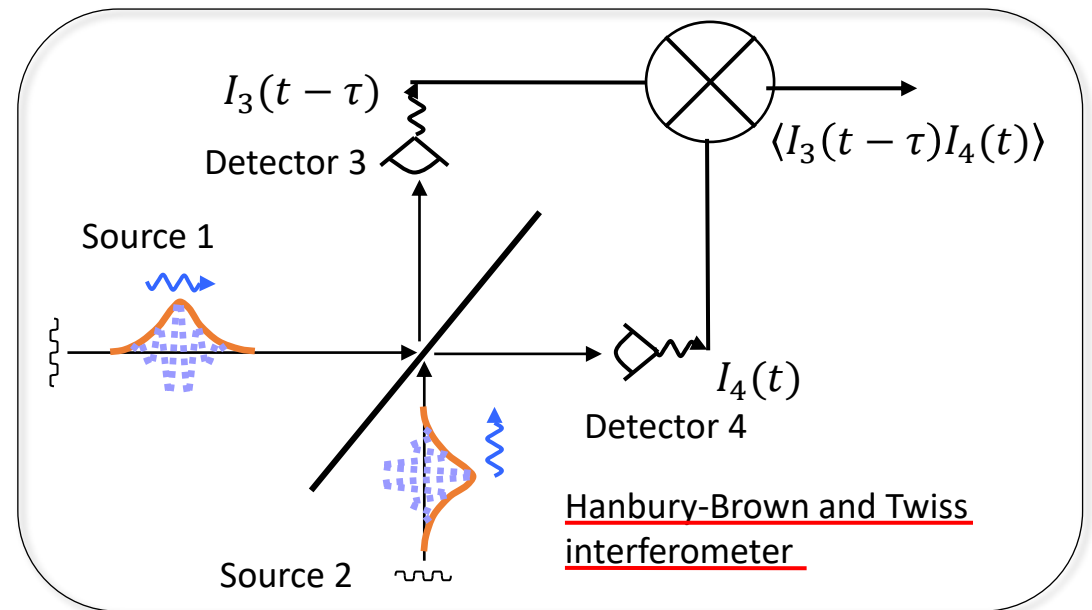
$$G^{(1)}(t + \tau, t) \propto \langle E^*(t - \tau)E(t) \rangle$$

Coherence of electric field



Two-particle interferometer

HBT interferometry $\langle I_3(t - \tau)I_4(t) \rangle$

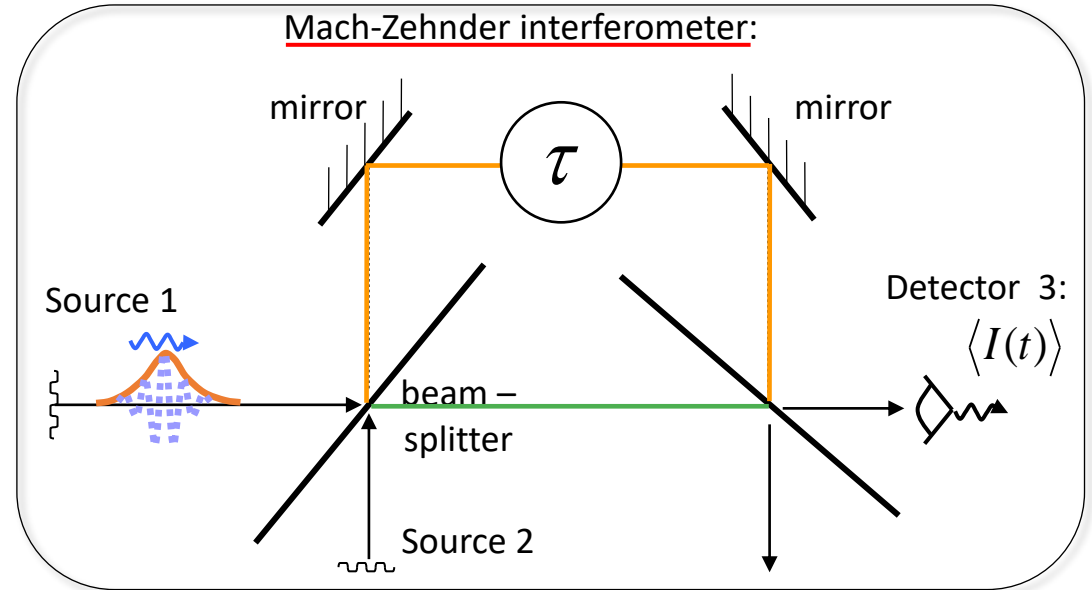


Single-particle vs two-particle interferometry (optics)

Single particle interferometer

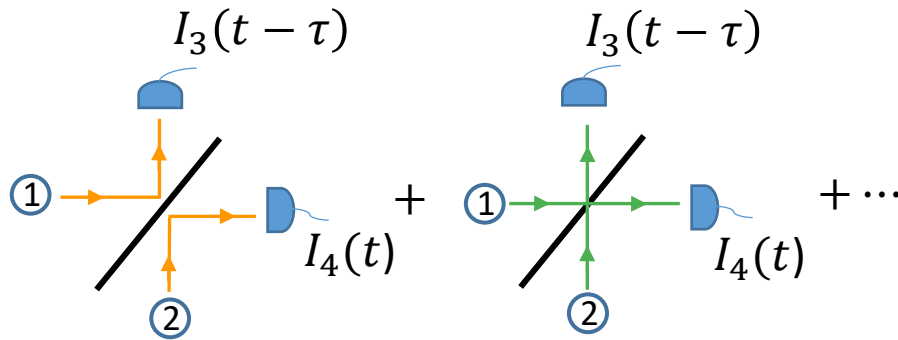
$$G^{(1)}(t + \tau, t) \propto \langle E^*(t - \tau)E(t) \rangle$$

Coherence of electric field

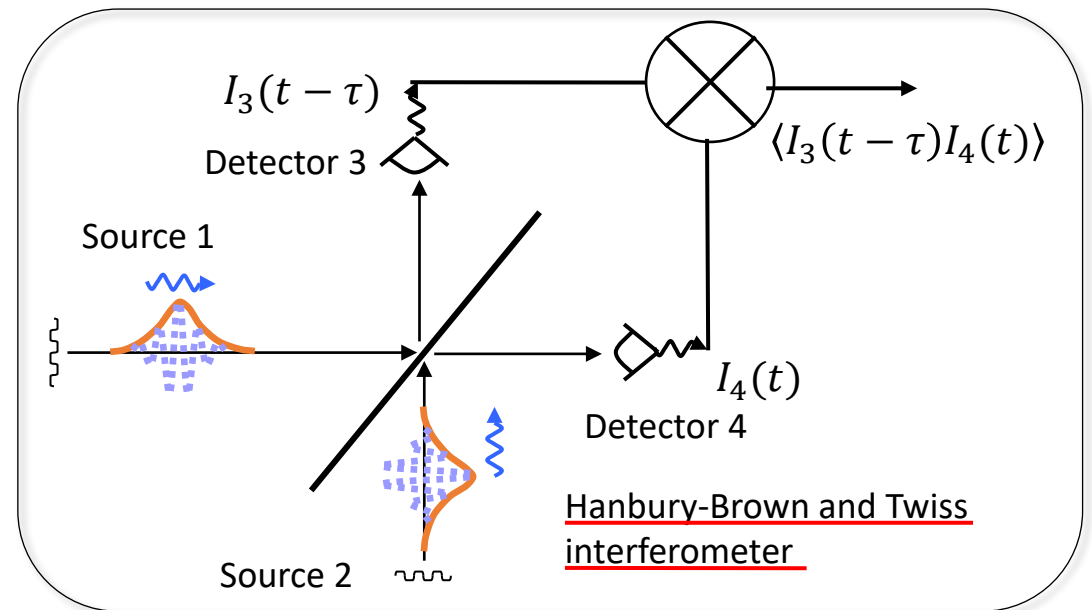


Two-particle interferometer

HBT interferometry $\langle I_3(t - \tau)I_4(t) \rangle$

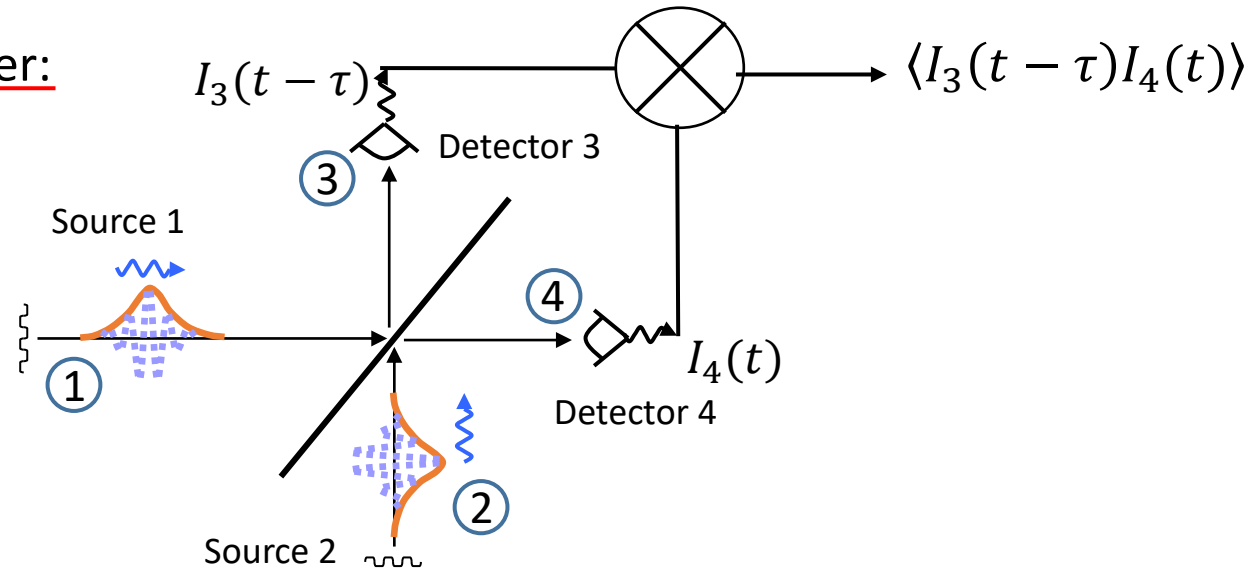


$$\propto \langle E_1^*(t - \tau)E_1(t) \rangle \langle E_2(t - \tau)E_2^*(t) \rangle$$

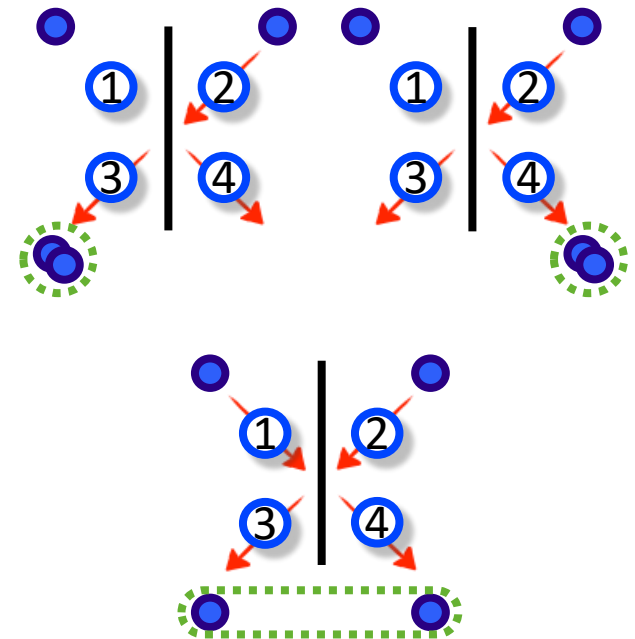
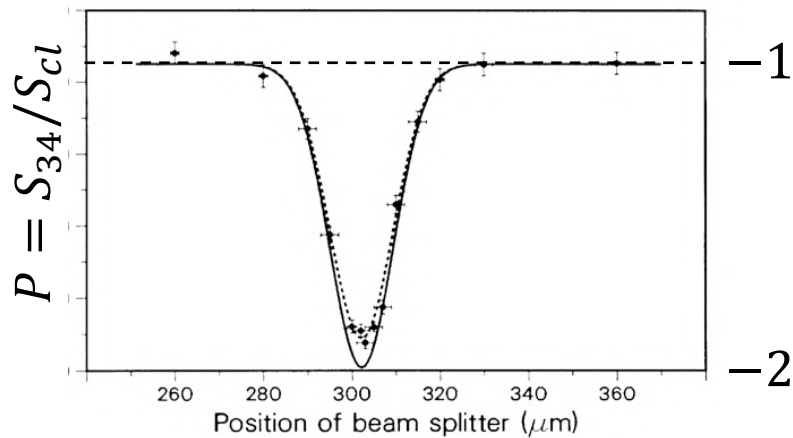


Single-particle vs two-particle interferometry (optics)

Two-particle interferometer:

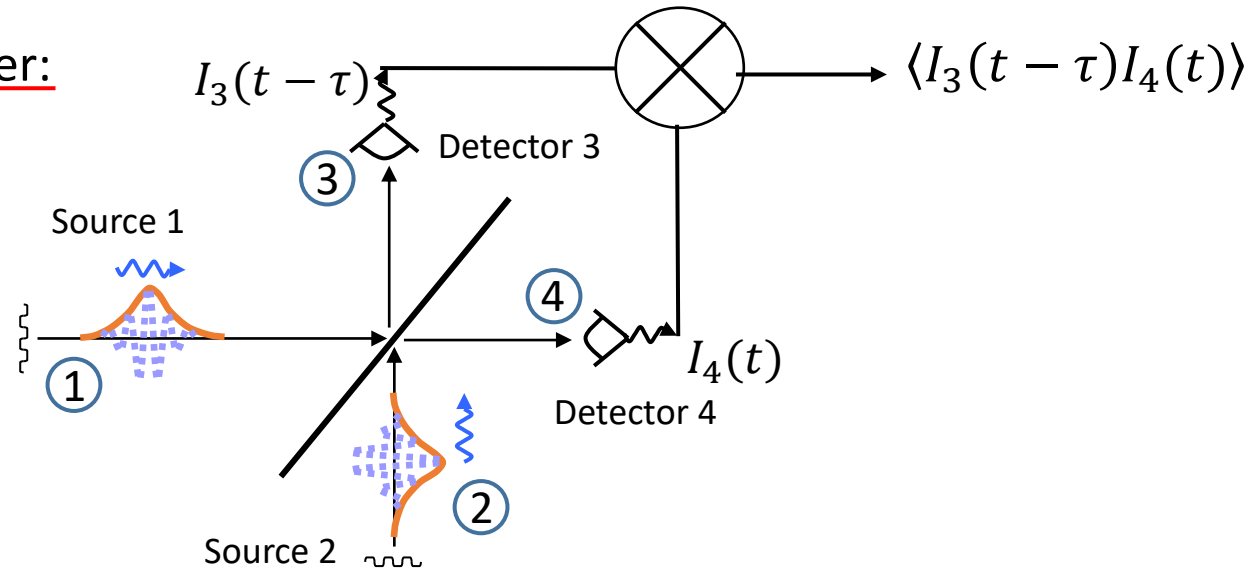


Cross-correlations:
$$S_{34} = 2 \int d\tau \overline{\Delta I_3(t - \tau) \Delta I_4(t)}^t$$

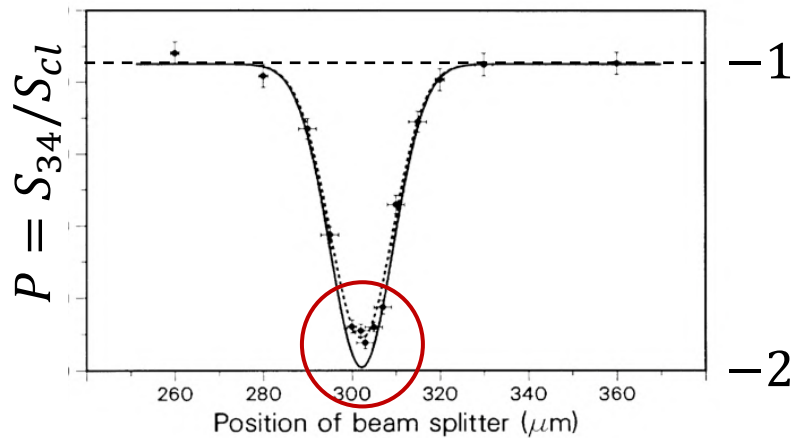


Single-particle vs two-particle interferometry (optics)

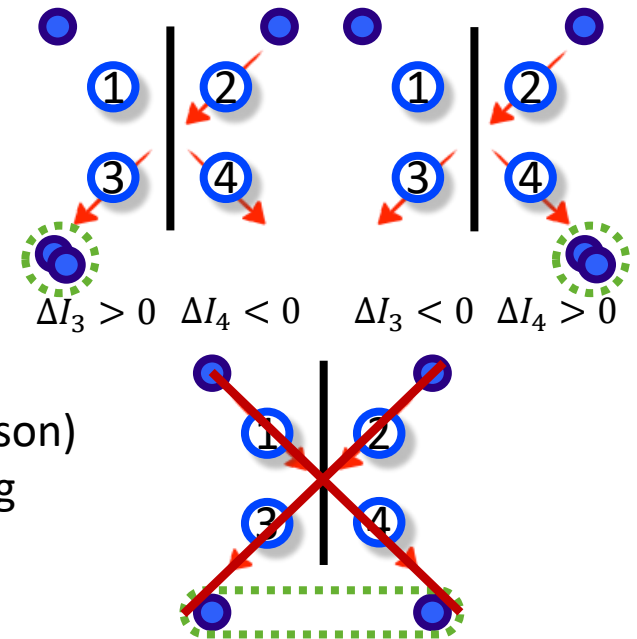
Two-particle interferometer:



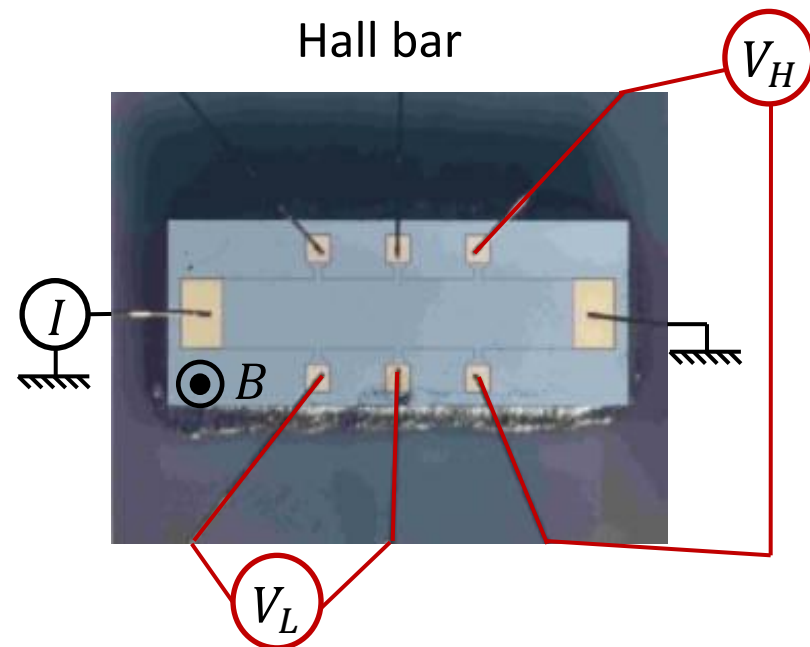
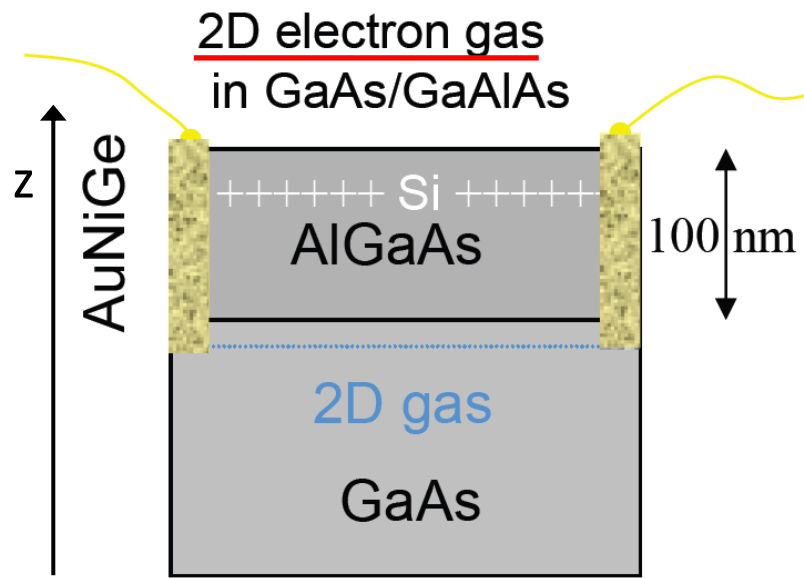
Cross-correlations:
$$S_{34} = 2 \int d\tau \overline{\langle \Delta I_3(t - \tau) \Delta I_4(t) \rangle}^t$$



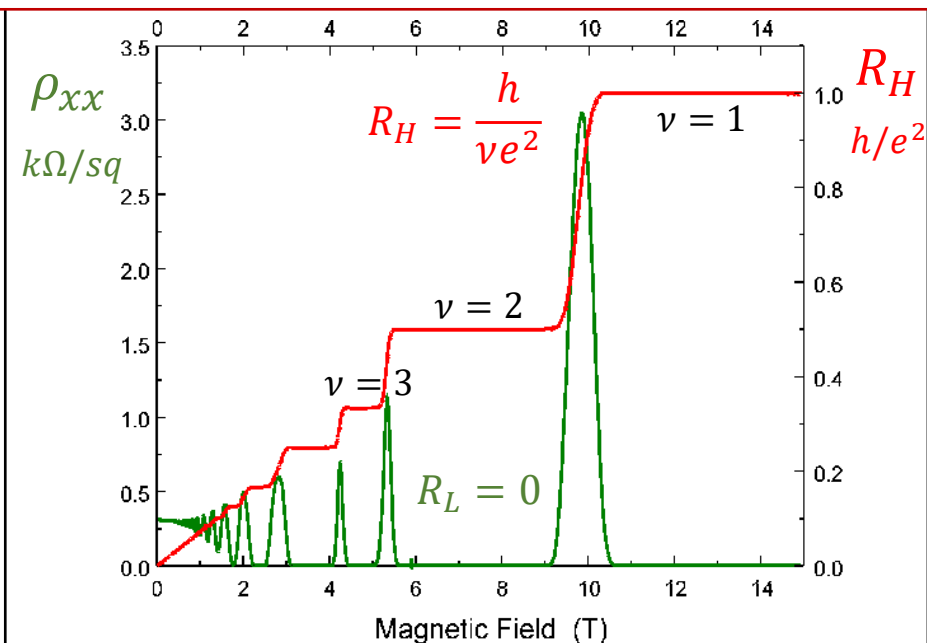
C. Hong *et al.*, PRL **59**(18), 2044 (1987)



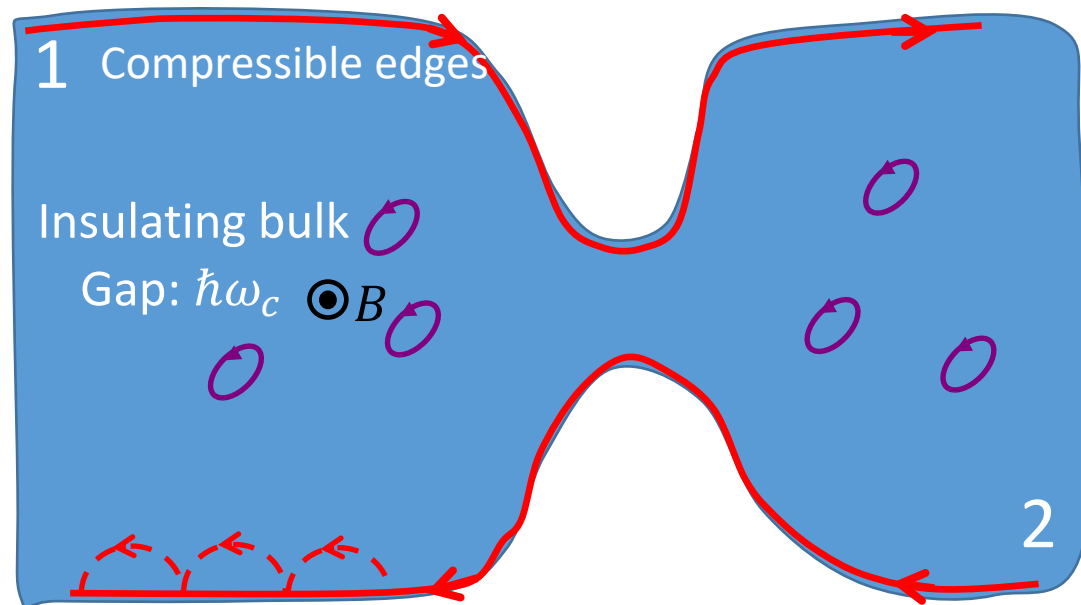
Photon (boson)
bunching



K. v. Klitzing, G. Dorda, and M. Pepper, PRL **45**, 494 (1980).

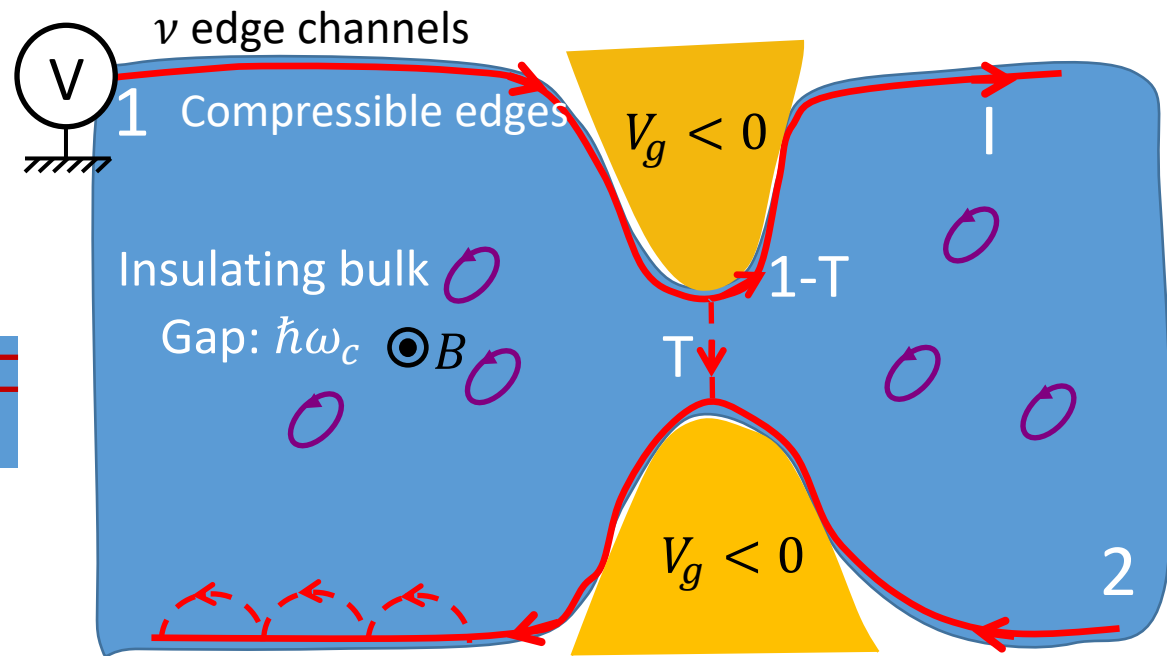
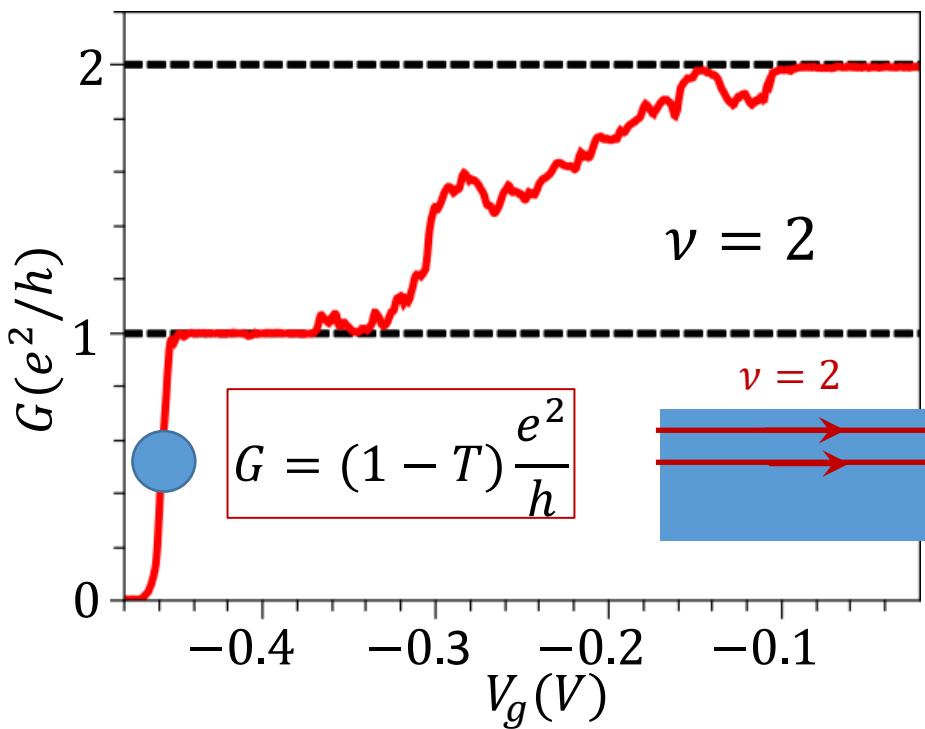
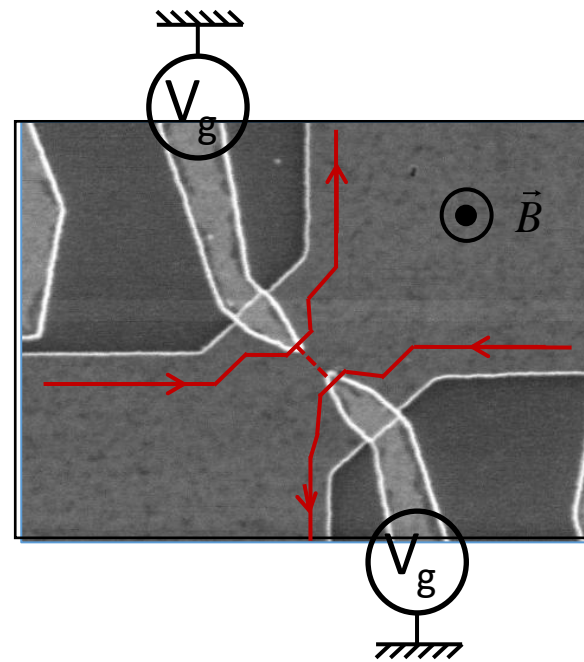
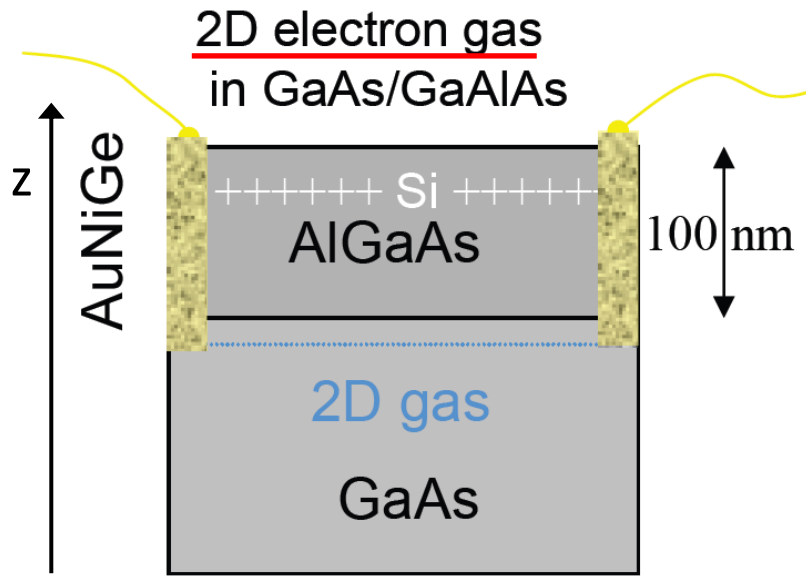


Ballistic propagation along ν edge channels

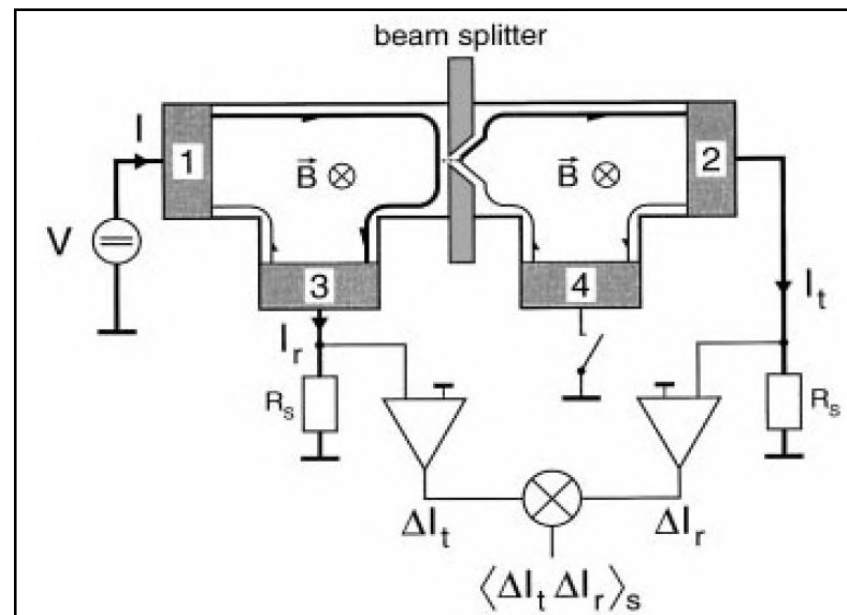
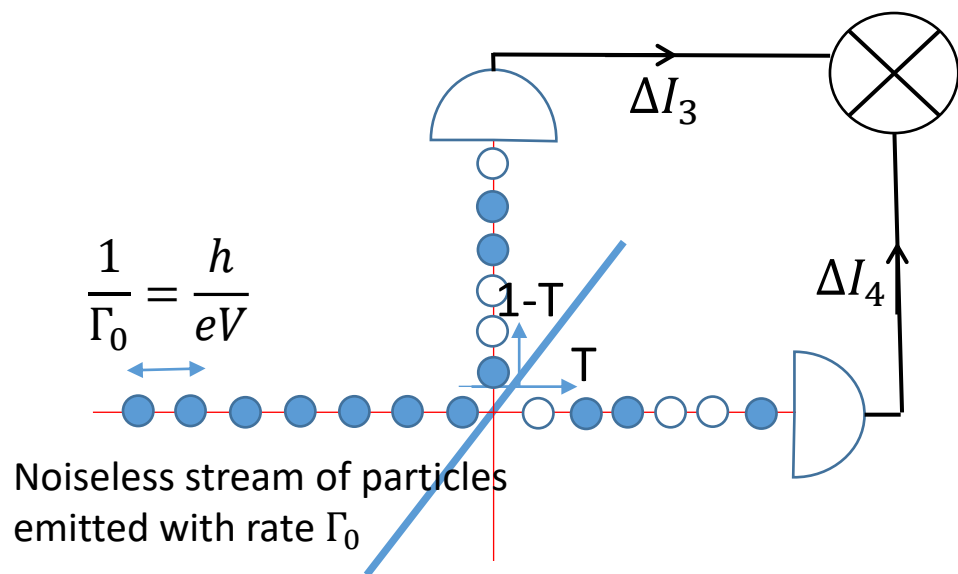


LPENS The quantum point contact: tunable beam-splitter

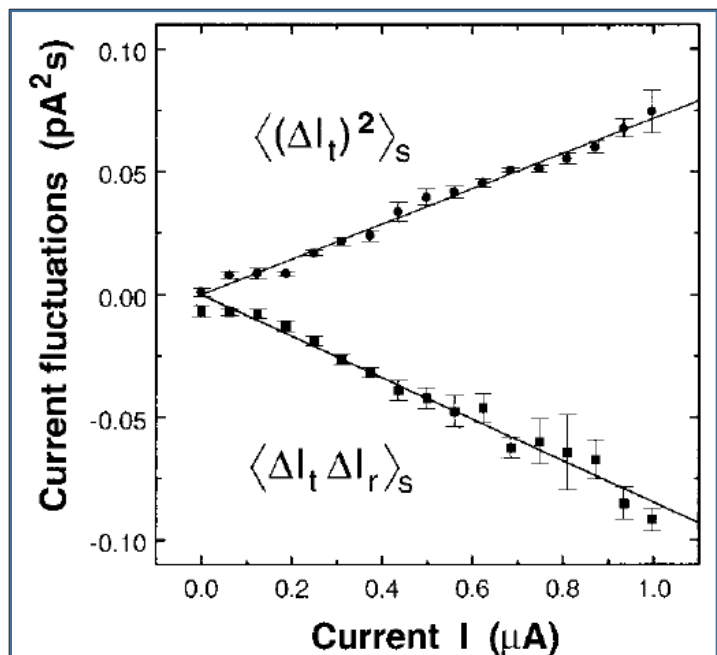
LABORATOIRE DE PHYSIQUE DE L'ÉCOLE NORMALE SUPÉRIEURE



Electron optics experiments: random partitioning of electron beams



M. Henny et al., Science **284** 296 (1999)
W.D. Oliver et al., Science **284** 299 (1999)
A. Kumar et al., PRL **76** 2778 (1996)



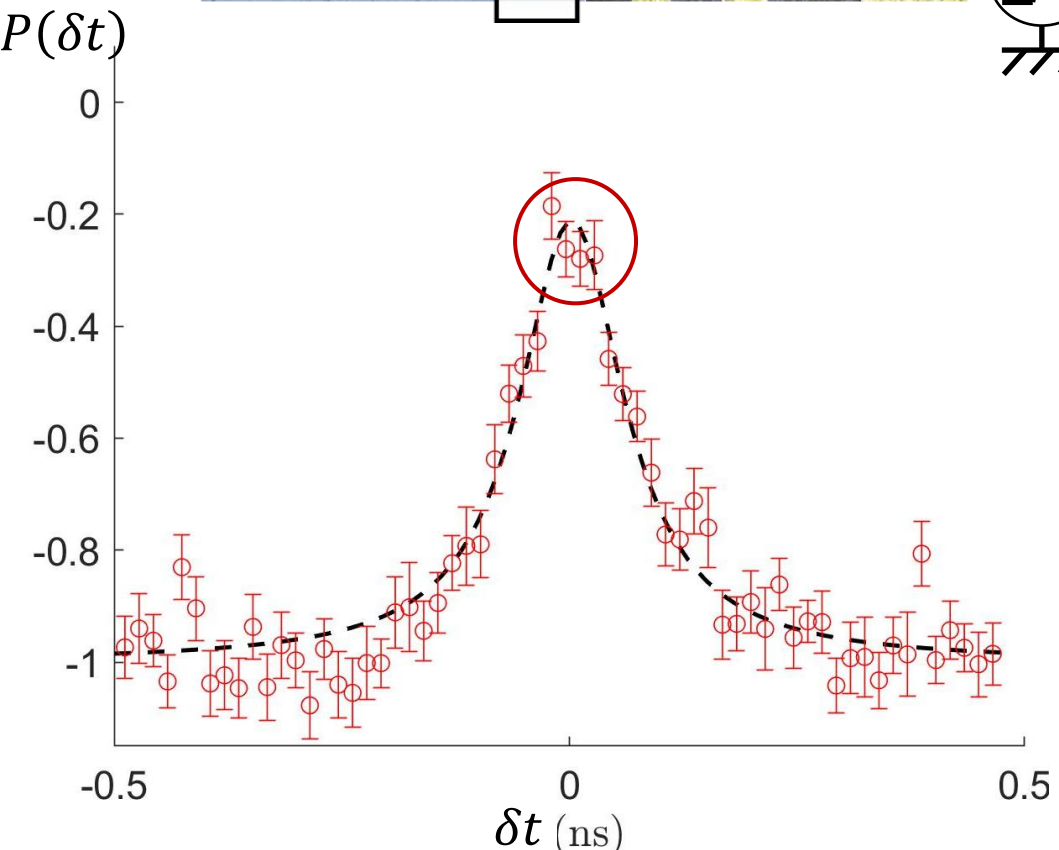
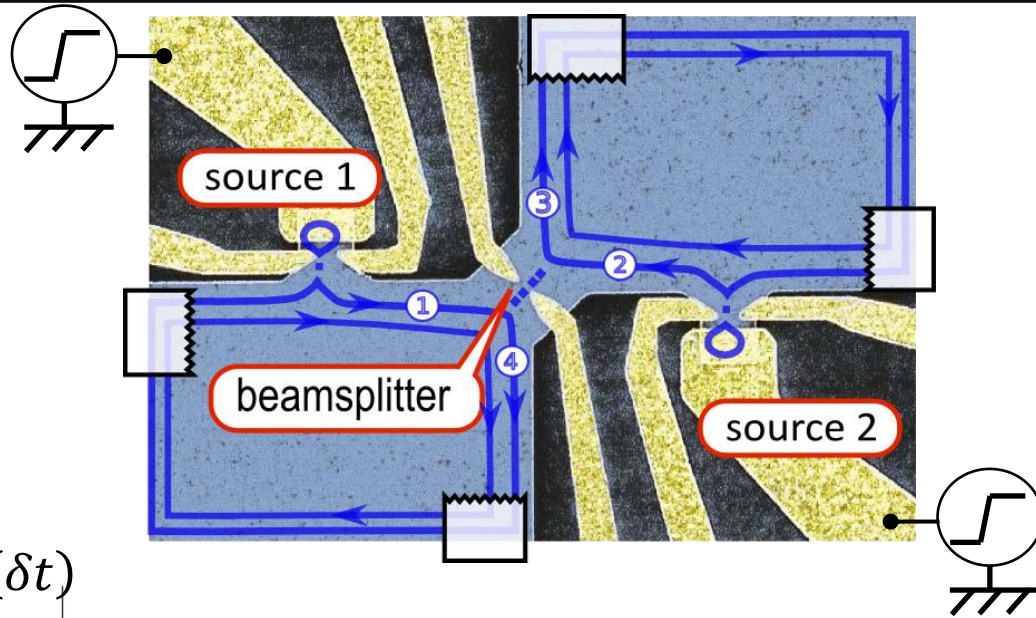
Binomial law: $\langle \Delta N_4^2 \rangle = T(1 - T) N_0$

$$S_{44} = \frac{2q^2}{T_{meas}} \langle \Delta N_4^2 \rangle = 2qT(1 - T)I_0 \equiv S_{cl}$$

$$S_{34} = -S_{44} = -S_{cl}$$

$$P = S_{34}/S_{cl} = -1$$

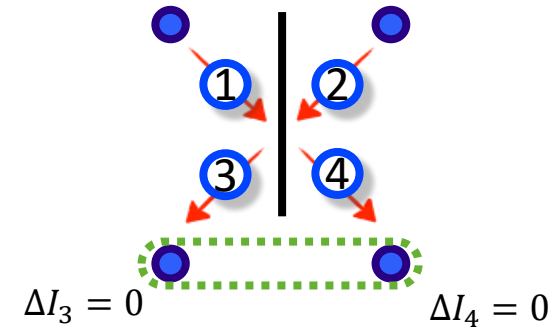
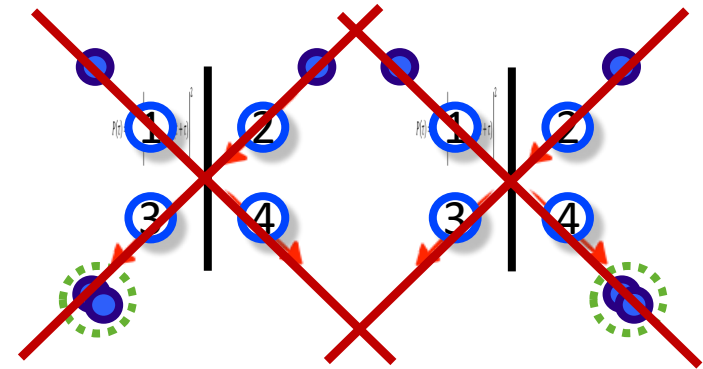
$\nu = 2$: Two-particle interferometry with single electrons



$$\langle \Delta I_3(t - \tau) \Delta I_4(t) \rangle$$

$$\propto \langle \psi_1^+(t - \tau) \psi_1(t) \rangle \langle \psi_2(t - \tau) \psi_2^+(t) \rangle$$

Electron interferometry

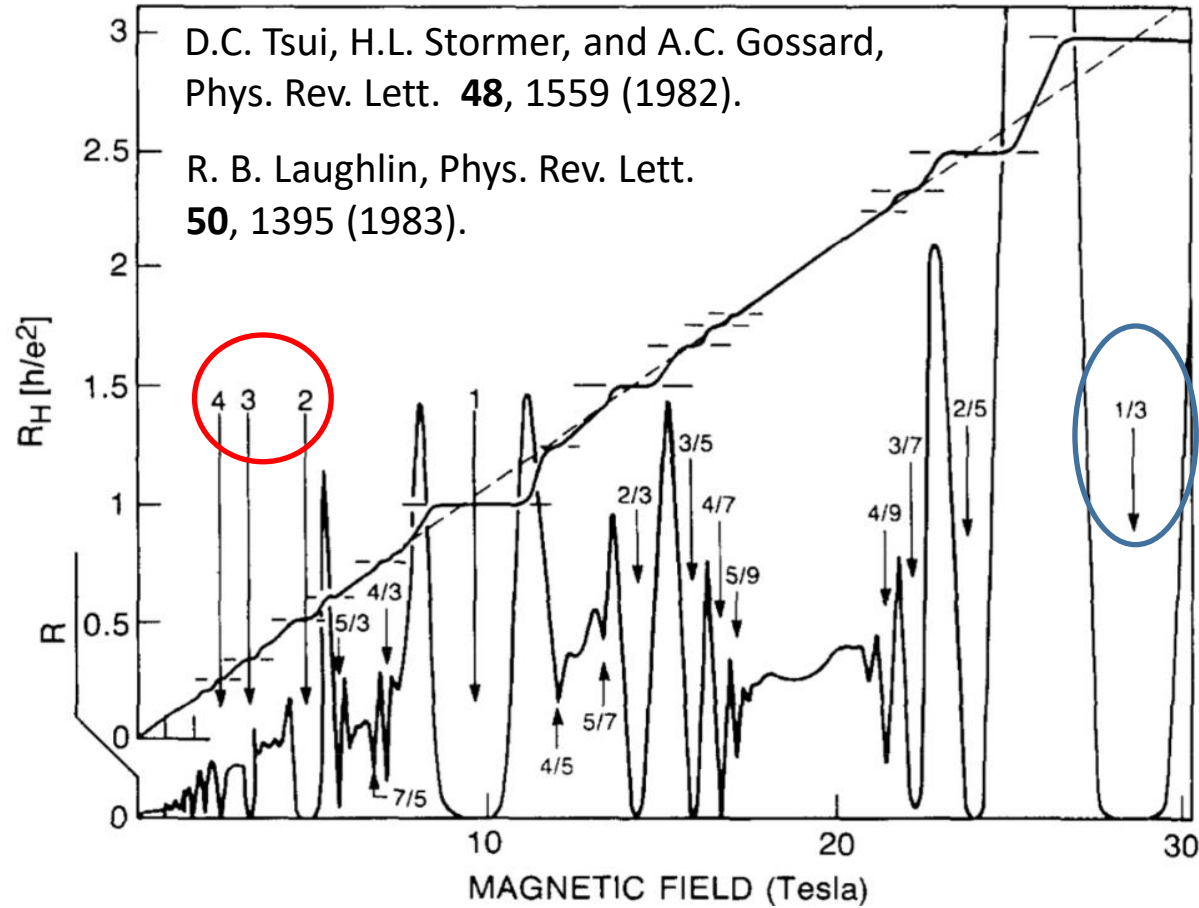


Fermion antibunching

S. Ol'khovskaya et al., PRL **101**, 166802, (2008).

E. Bocquillon et al., Science **339**, 1054 (2013).

$$\begin{aligned} \nu &= N \\ q &= e \\ \varphi &= \pi \\ \theta &= 2\varphi = 2\pi \end{aligned}$$



$$G = \frac{1}{3} \frac{e^2}{h}$$

$$\begin{aligned} \nu &= 1/3 \\ q &= e/3 \\ \varphi &= \pi/3 \\ \theta &= 2\varphi = 2\pi/3 \end{aligned}$$

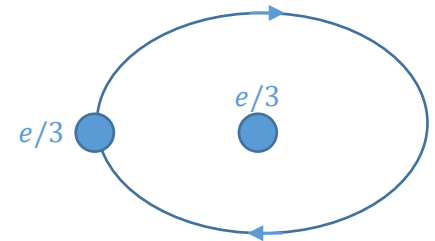
Each FQHE phase hosts a specific variety of anyons characterized by their fractional charge q and their fractional statistics φ

Halperin, PRL **52** 1583 (1984)

Arovas, Schrieffer, Wilczek PRL **53** 722 (1984)

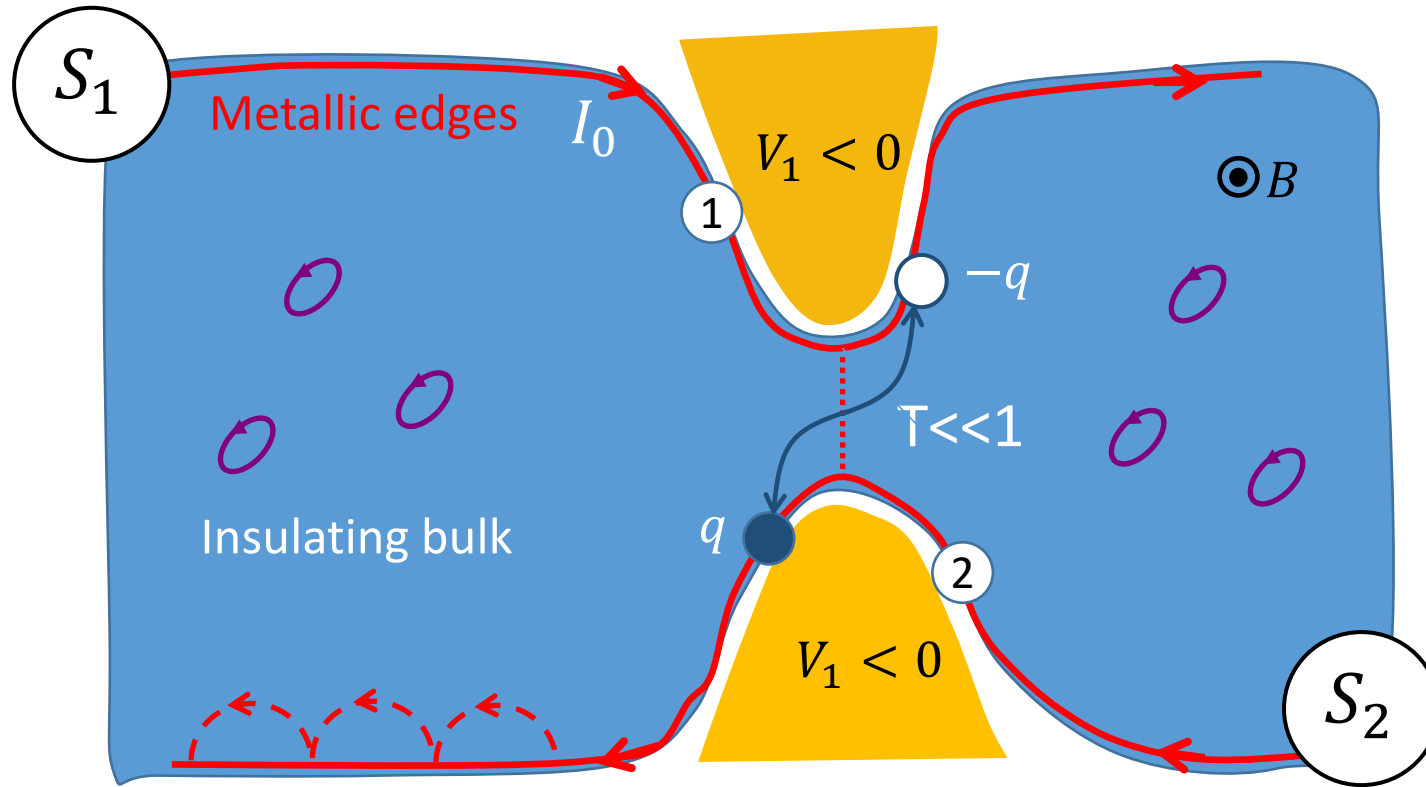
J. K. Jain, PRL **63**, 199 (1989).

Anyons have non-trivial braiding



$$e^{i\theta} = e^{i2\pi/3} \quad (\nu = 1/3)$$

Electron/anyons tunneling at quantum point contact



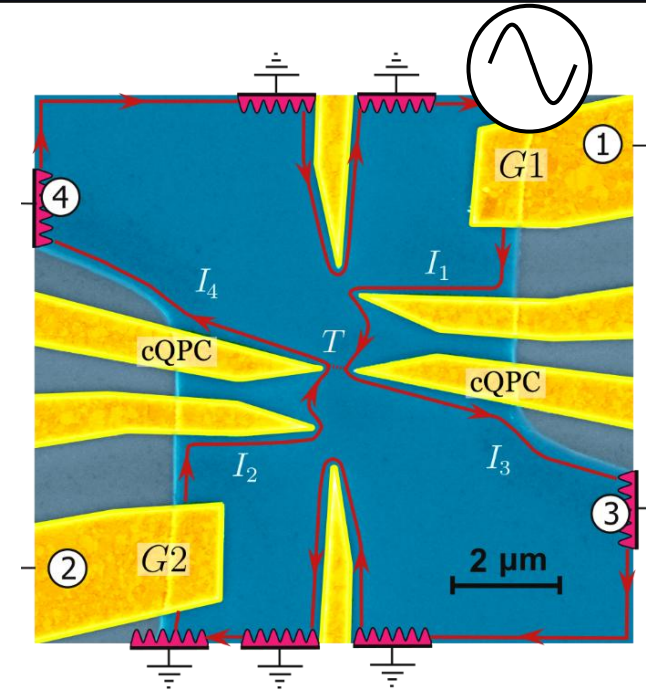
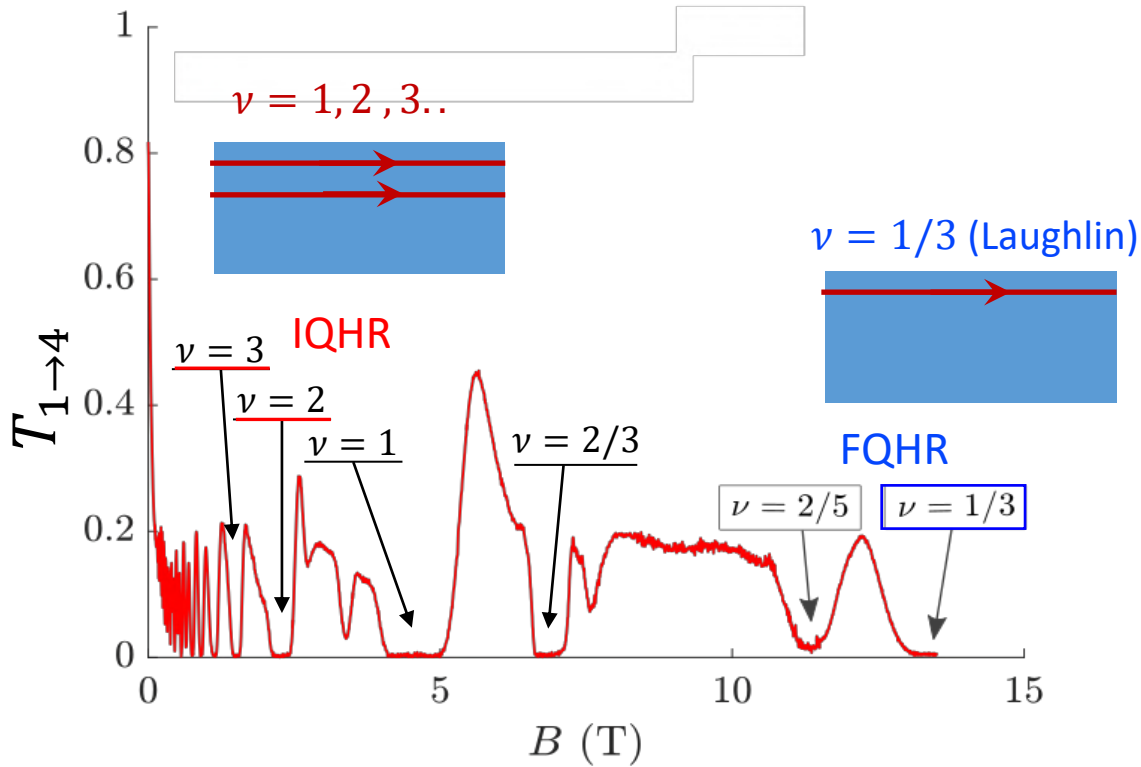
Random tunneling of electrons in the integer quantum Hall regime: $H_T \propto \psi_{2,e}^+ \psi_{1,e} + h.c$

$$\nu = 2, \nu = 3 \quad q = e \quad \varphi = \pi$$

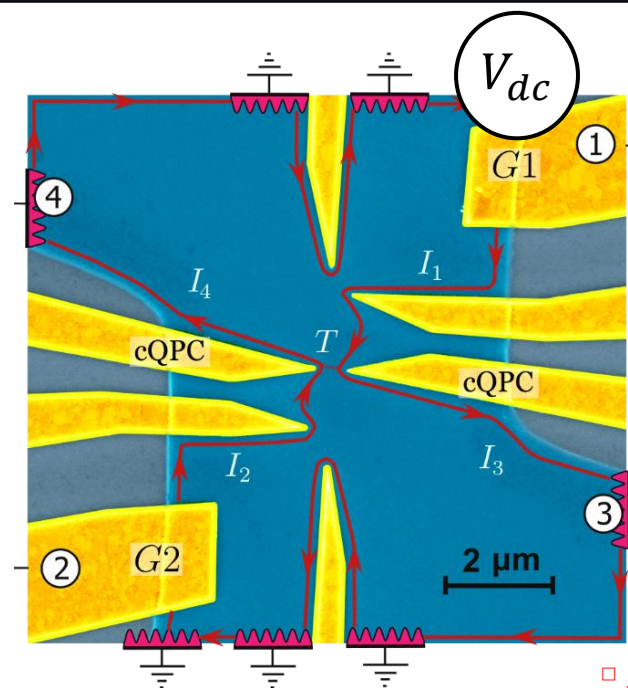
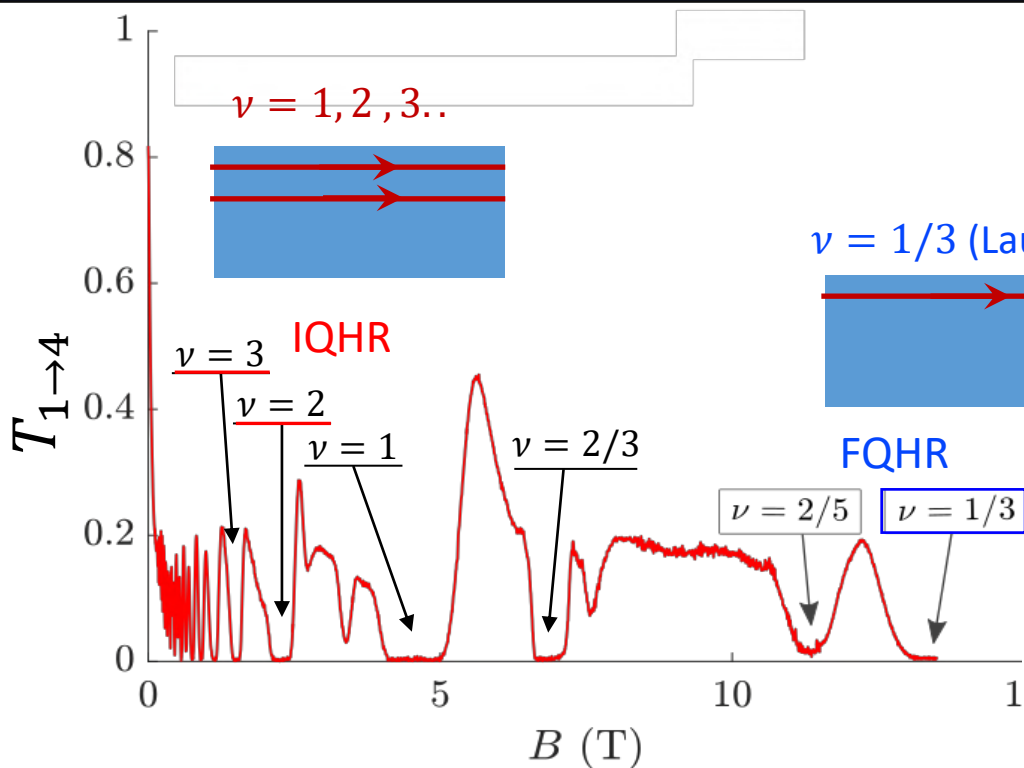
Random tunneling of anyons in the fractional quantum Hall regime: $H_T \propto \psi_{2,a}^+ \psi_{1,a} + h.c$

$$\nu = 1/3 \quad q = e/3 \quad \varphi = \pi/3 \quad \theta = 2\varphi = \frac{2\pi}{3}$$

Electron/anyon beam splitters: random partition noise and charge measurement

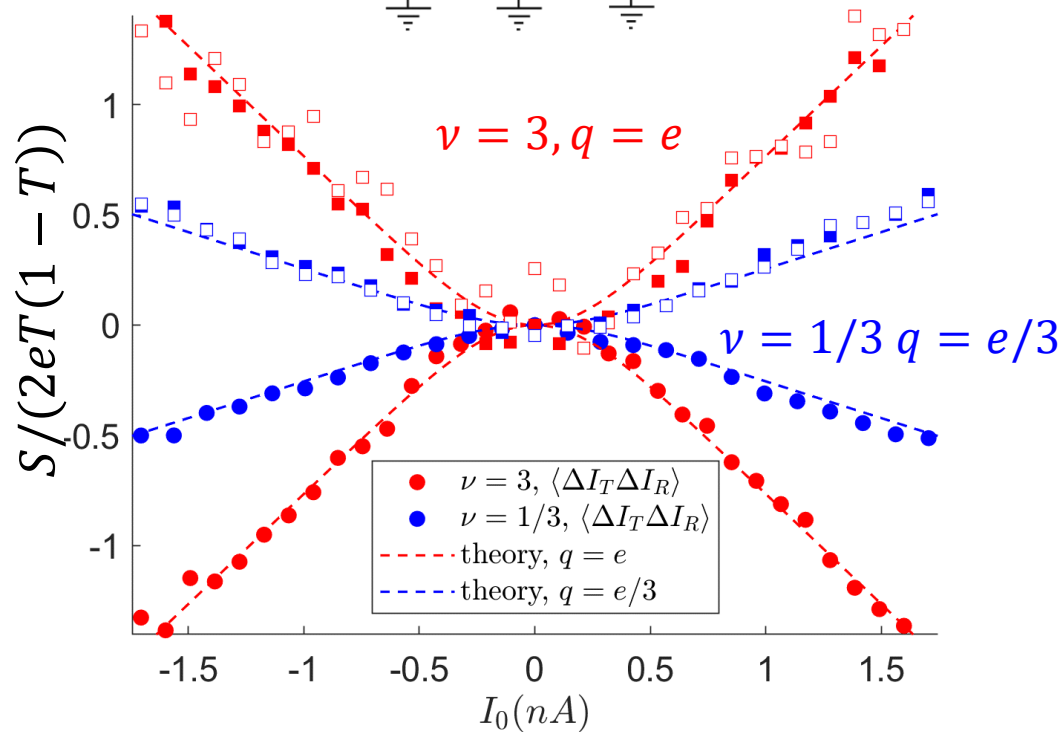


Electron/anyon beam splitters: random partition noise and charge measurement



$$S_{33} = 2qT(1 - T)I_0 \equiv S_{cl}$$

$$S_{34} = -S_{cl}$$



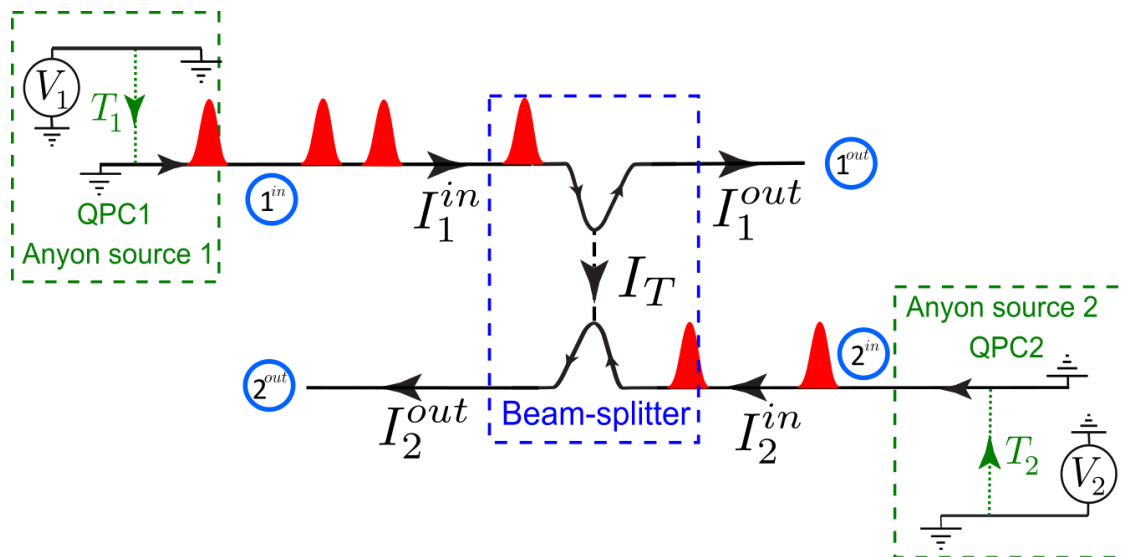
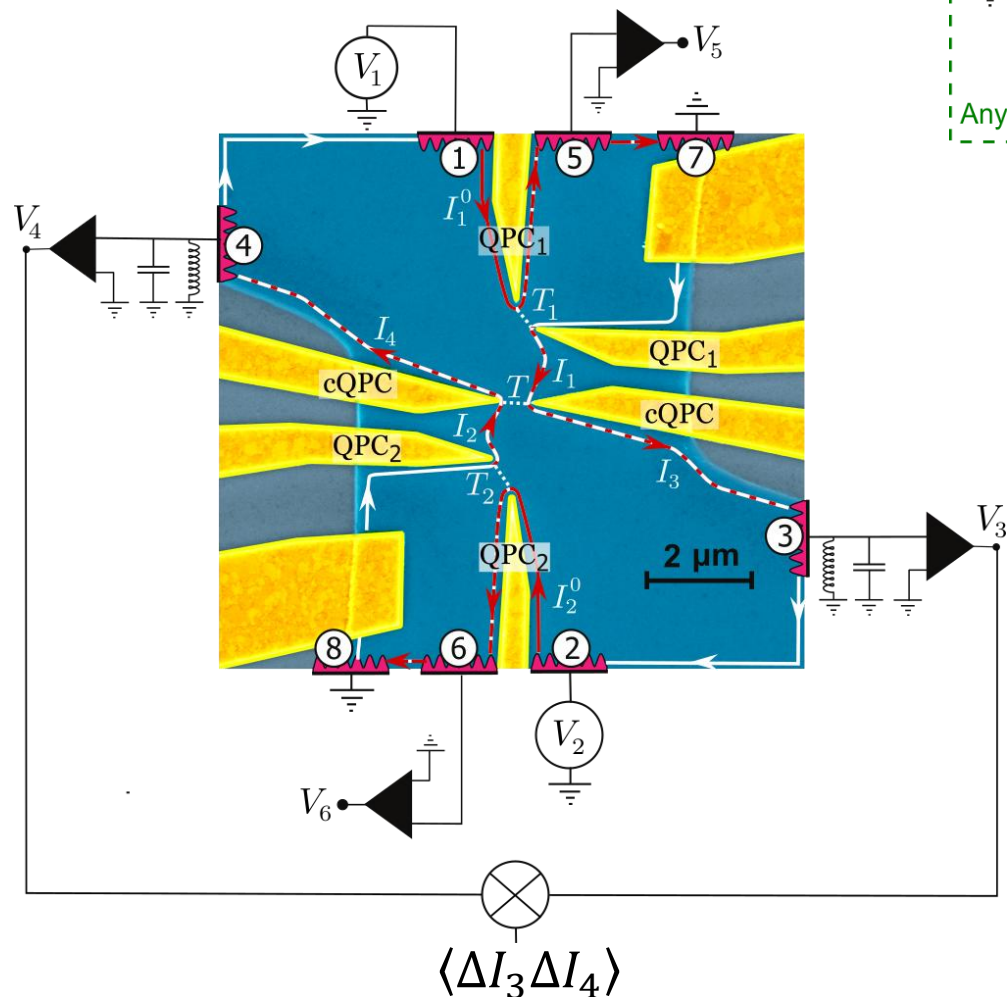
R. de Picciotto et al., Nature **389**, 162 (1997).

L. Saminadayar et al., Phys. Rev. Lett. **79**, 2526 (1997).

C. L. Kane and M. P. A. Fisher, Phys. Rev. Lett. **72**, 724 (1994)

B. Rosenow, I.P. Levkivskiy, B. Halperin PRL **116**, 156802 (2016)

H. Bartolomei et al., Science **368**, 173 (2020)



Random emission of particles:
probabilities $T_1 = T_2 = T_S$

Poissonian limit, $T_S \ll 1$

Fano factor:
$$P = \frac{S_{34}}{S_{cl}}$$

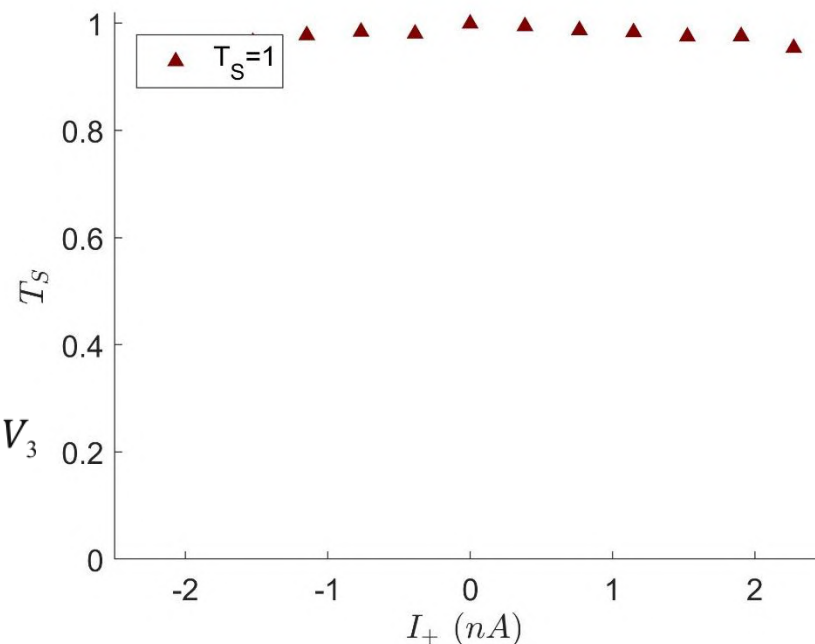
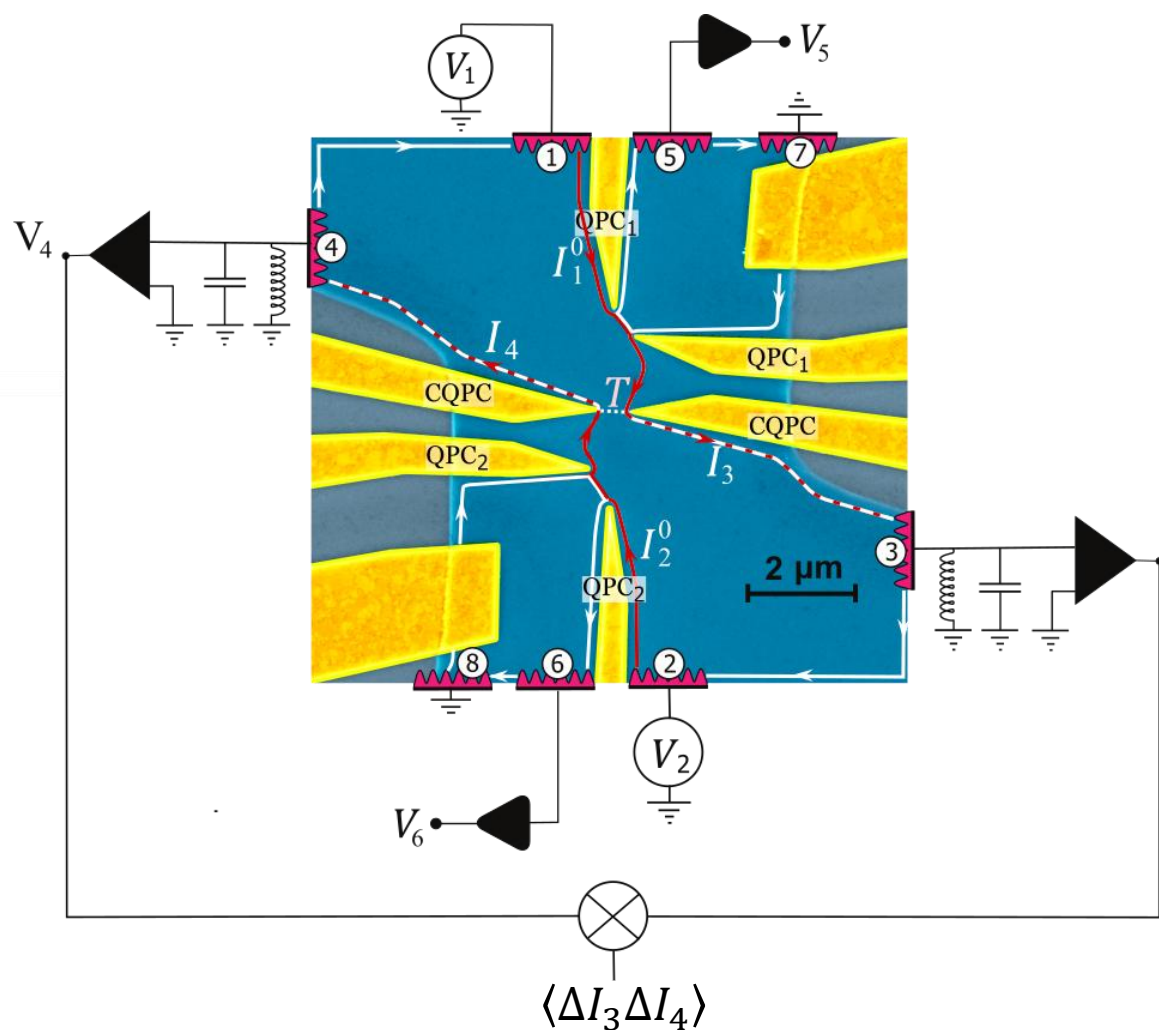
$$S_{cl} = 2qT(1 - T)I_+$$

Total input current:
$$I_+ = I_1^{in} + I_2^{in}$$

Balanced case: $I_1^{in} = I_2^{in}$

Integer case: $q = e$, fermions

$\nu = 2, T = 0.4, T_S = 1$



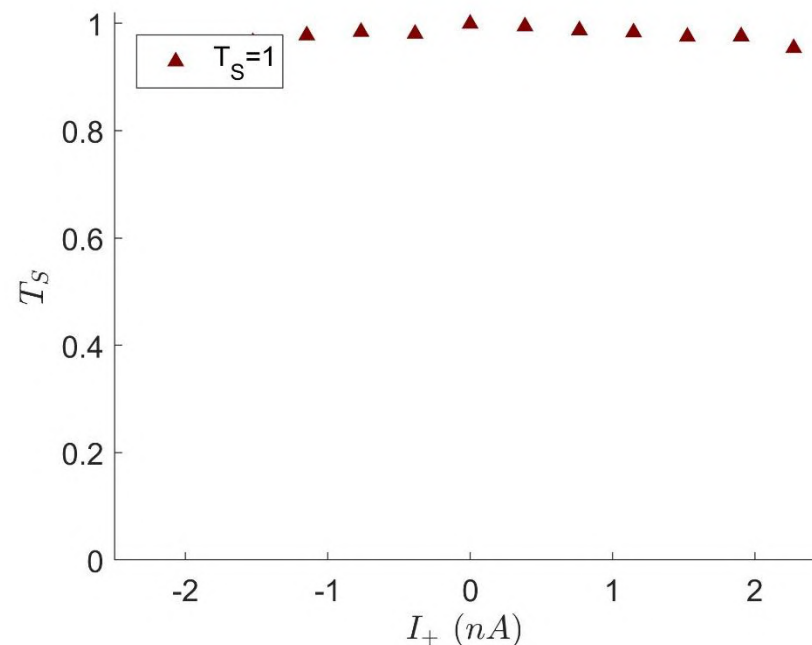
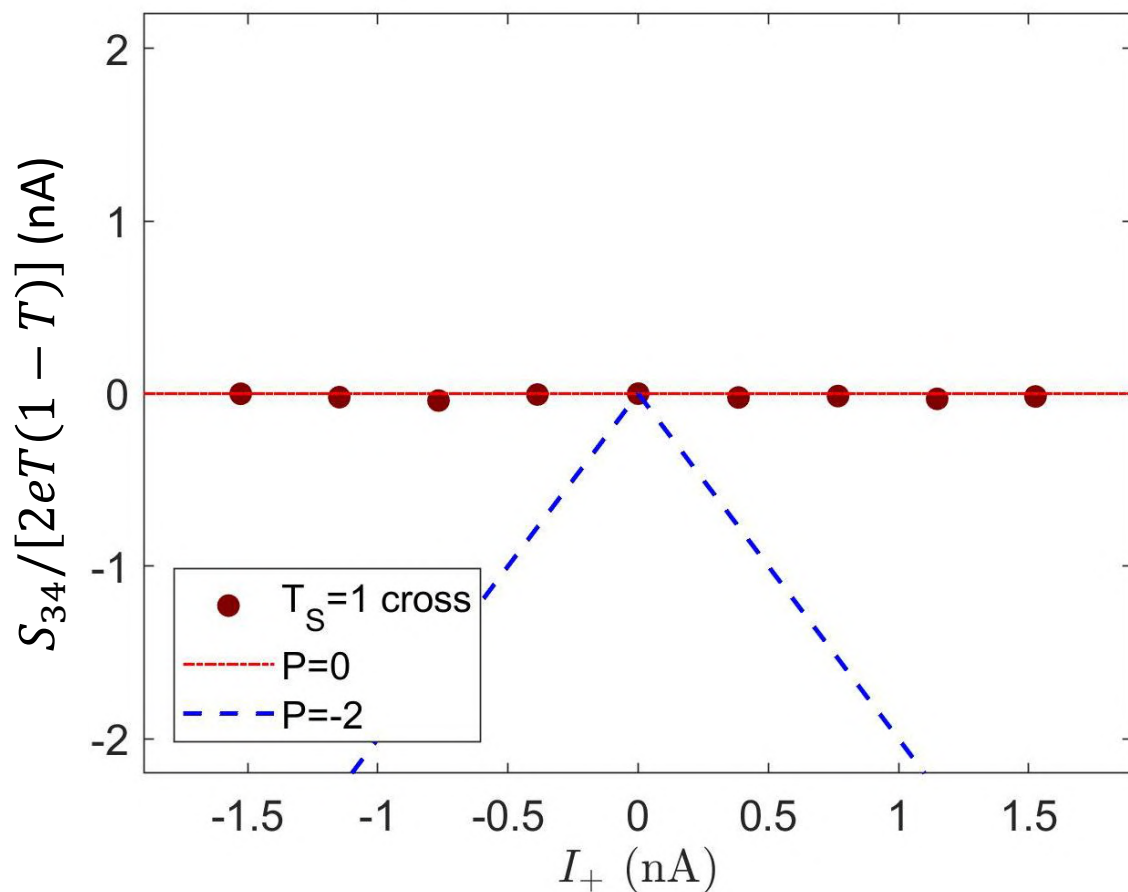
H. Bartolomei, M. Kumar et al., Science **368** 173 (2020)

M. Ruelle et al., PRX **13**, 011031 (2023)

Balanced case: $I_1^{in} = I_2^{in}$

Integer case: $q = e$, fermions

$\nu = 2, T = 0.4, T_S = 1$



H. Bartolomei, M. Kumar et al., Science **368** 173 (2020)

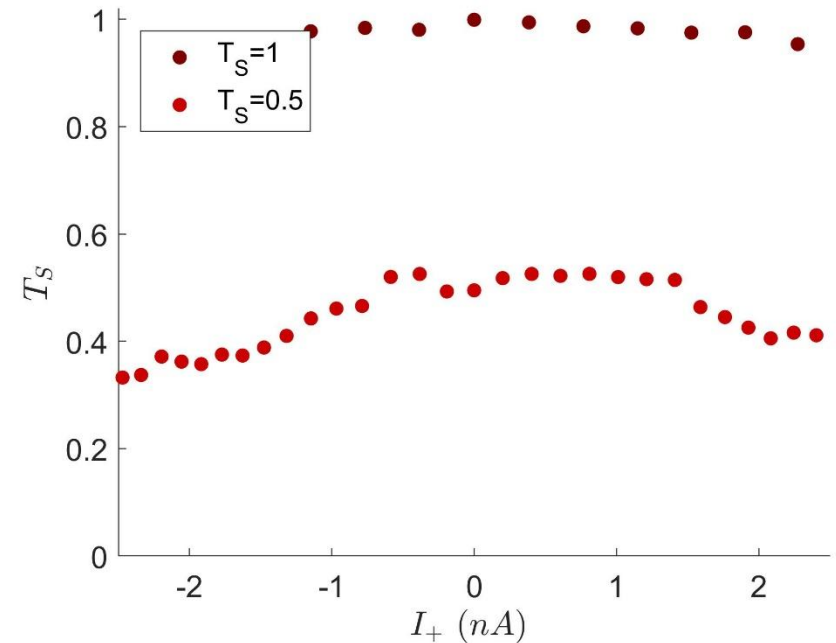
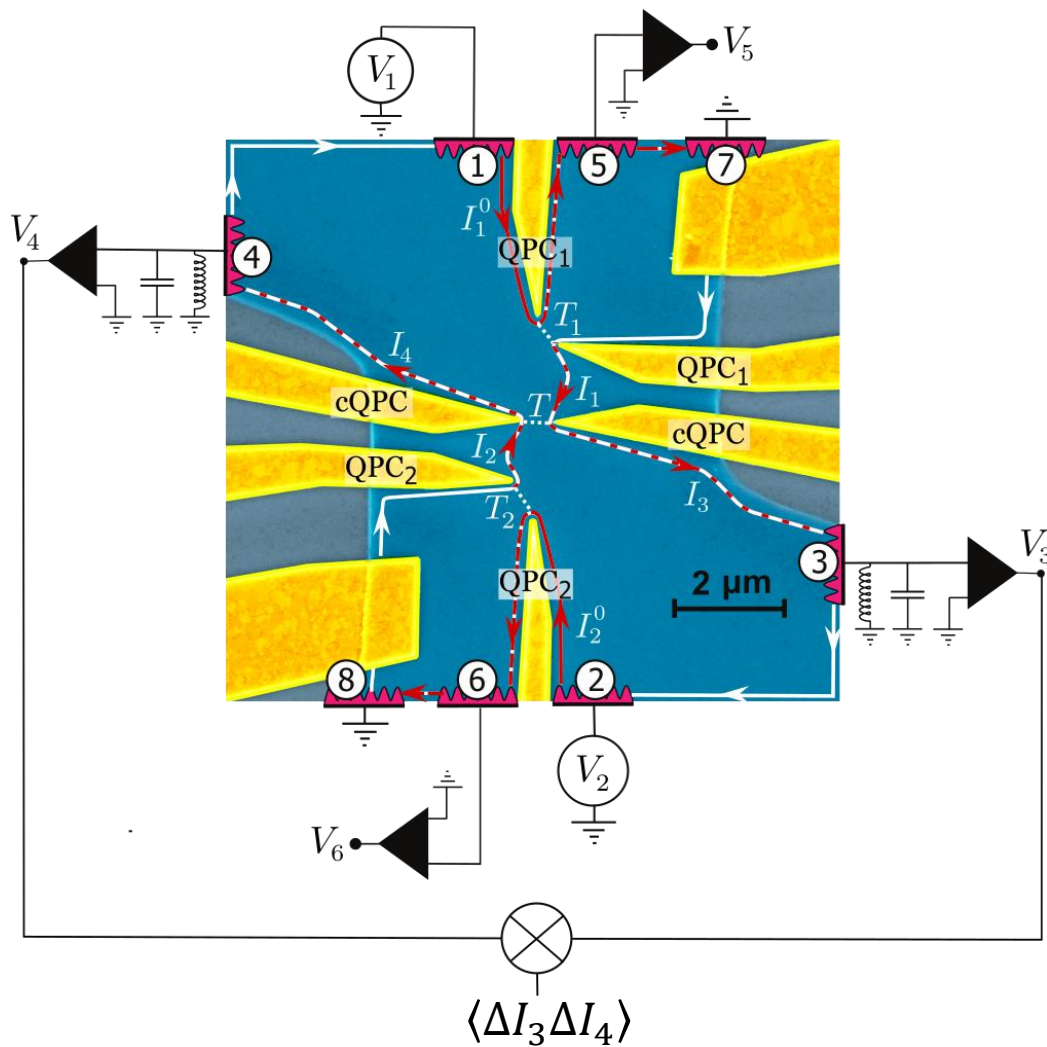
M. Ruelle et al., PRX **13**, 011031 (2023)

$P \approx 0$ fermions

Balanced case: $I_1^{in} = I_2^{in}$

Integer case: $q = e$, fermions

$\nu = 2, T = 0.4, T_S = 0.5$



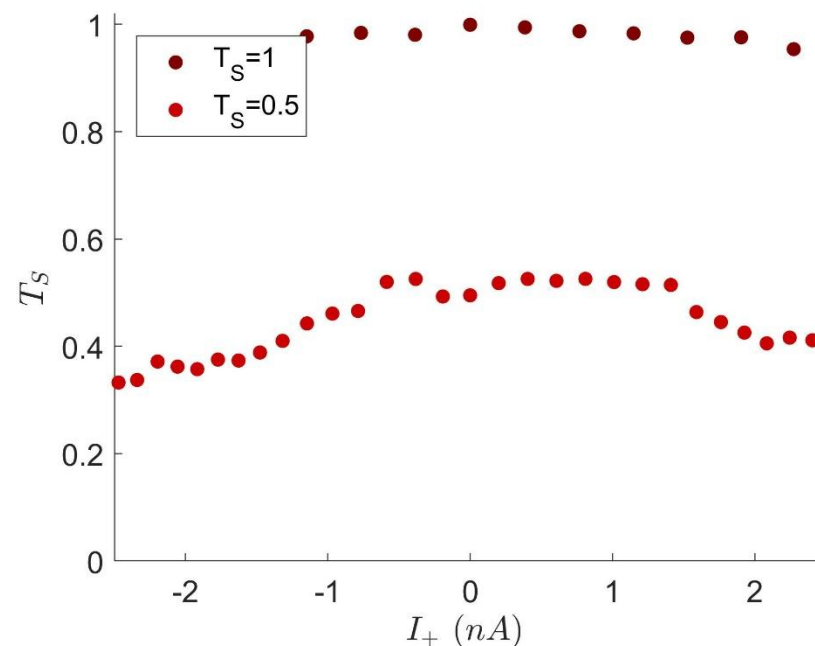
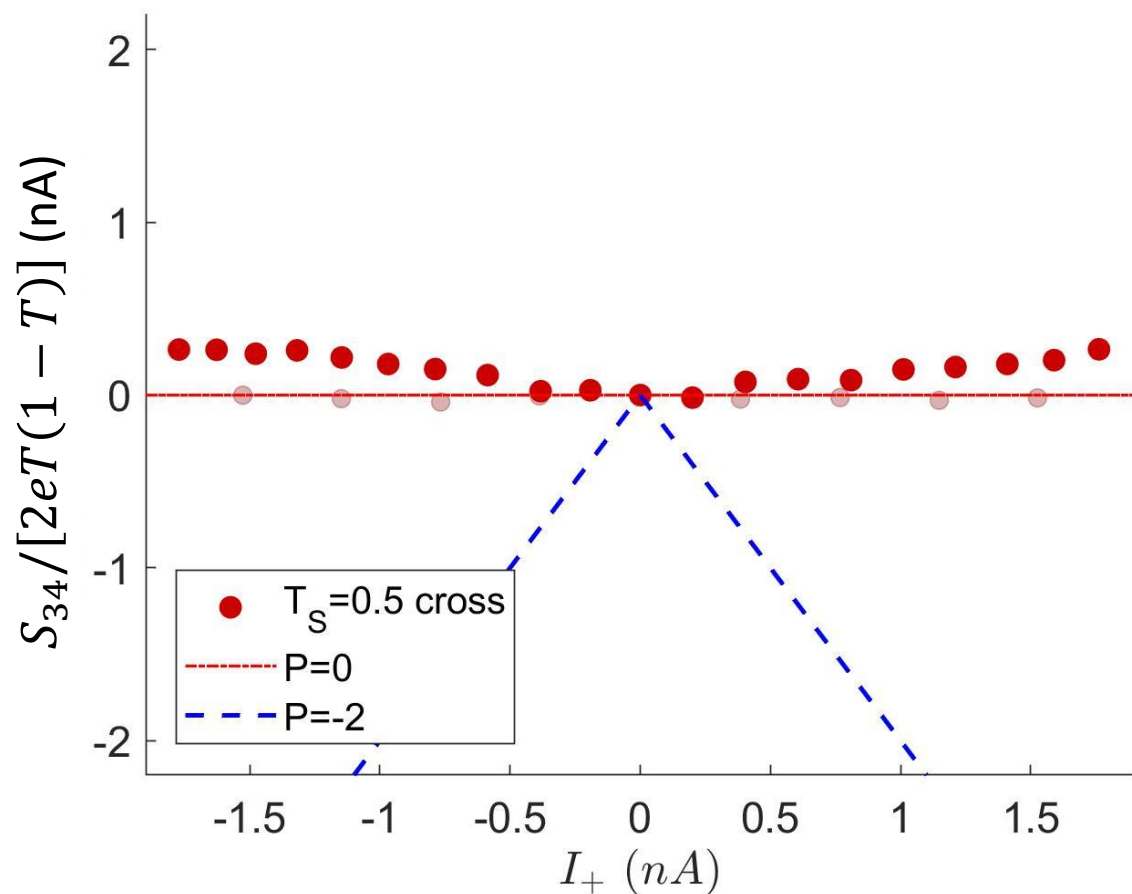
H. Bartolomei, M. Kumar et al., Science **368** 173 (2020)

M. Ruelle et al., PRX **13**, 011031 (2023)

Balanced case: $I_1^{in} = I_2^{in}$

Integer case: $q = e$, fermions

$\nu = 2, T = 0.4, T_S = 0.5$



H. Bartolomei, M. Kumar et al., Science **368** 173 (2020)

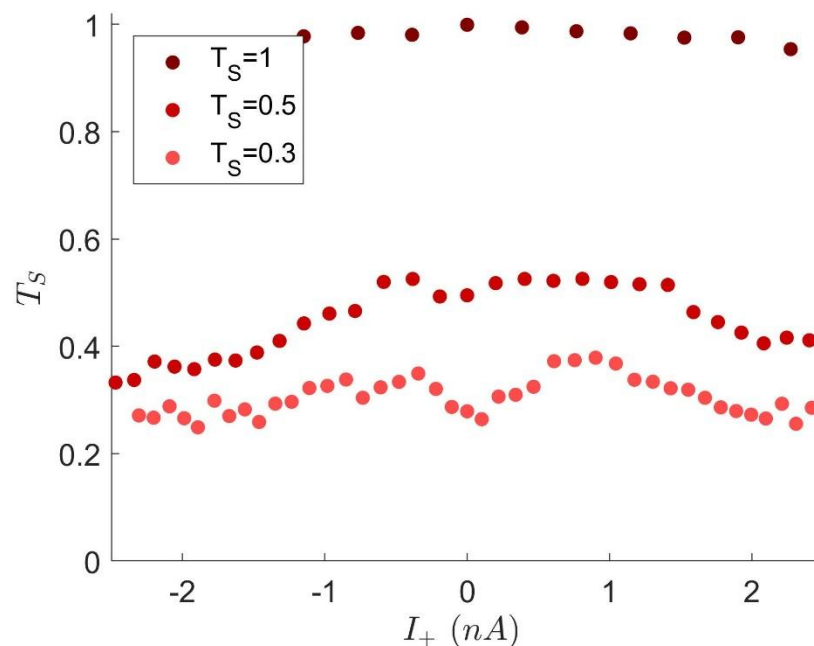
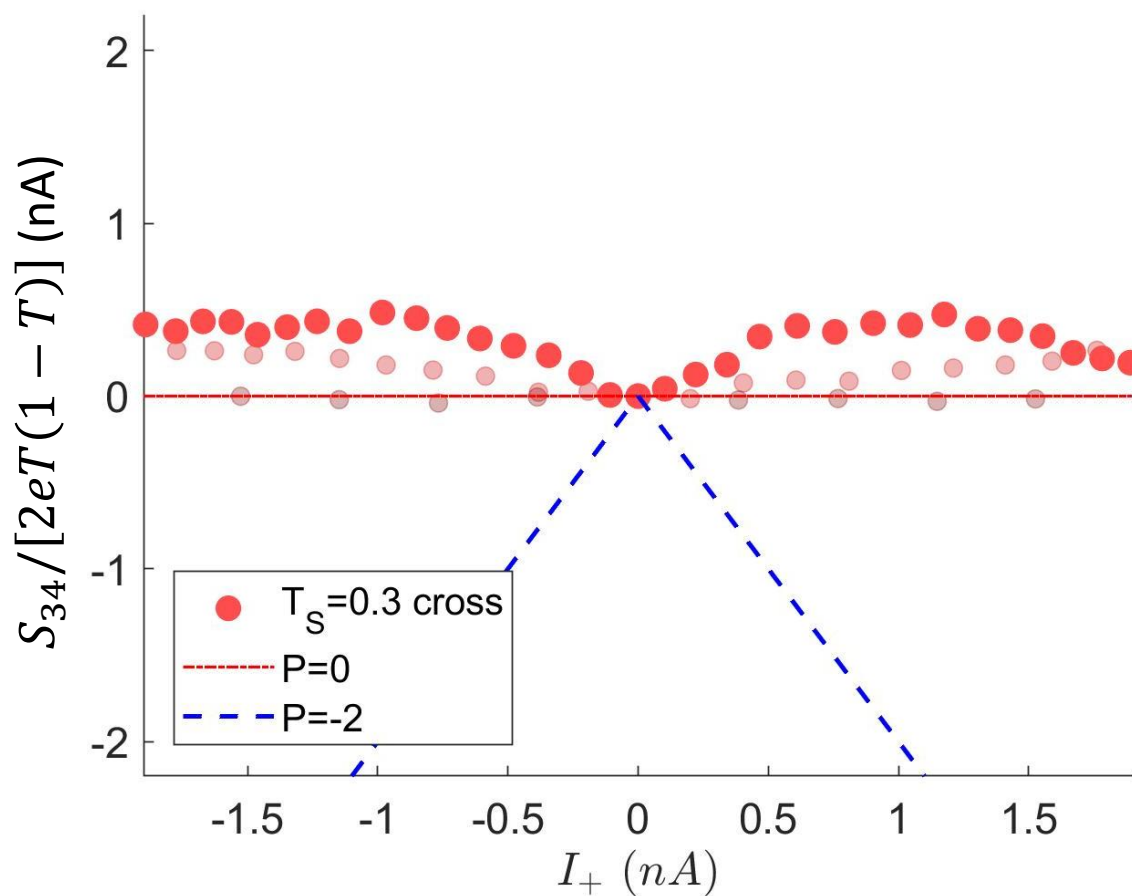
M. Ruelle et al., PRX **13**, 011031 (2023)

$P \approx 0$ fermions

Balanced case: $I_1^{in} = I_2^{in}$

Integer case: $q = e$, fermions

$\nu = 2, T = 0.4, T_S = 0.3$



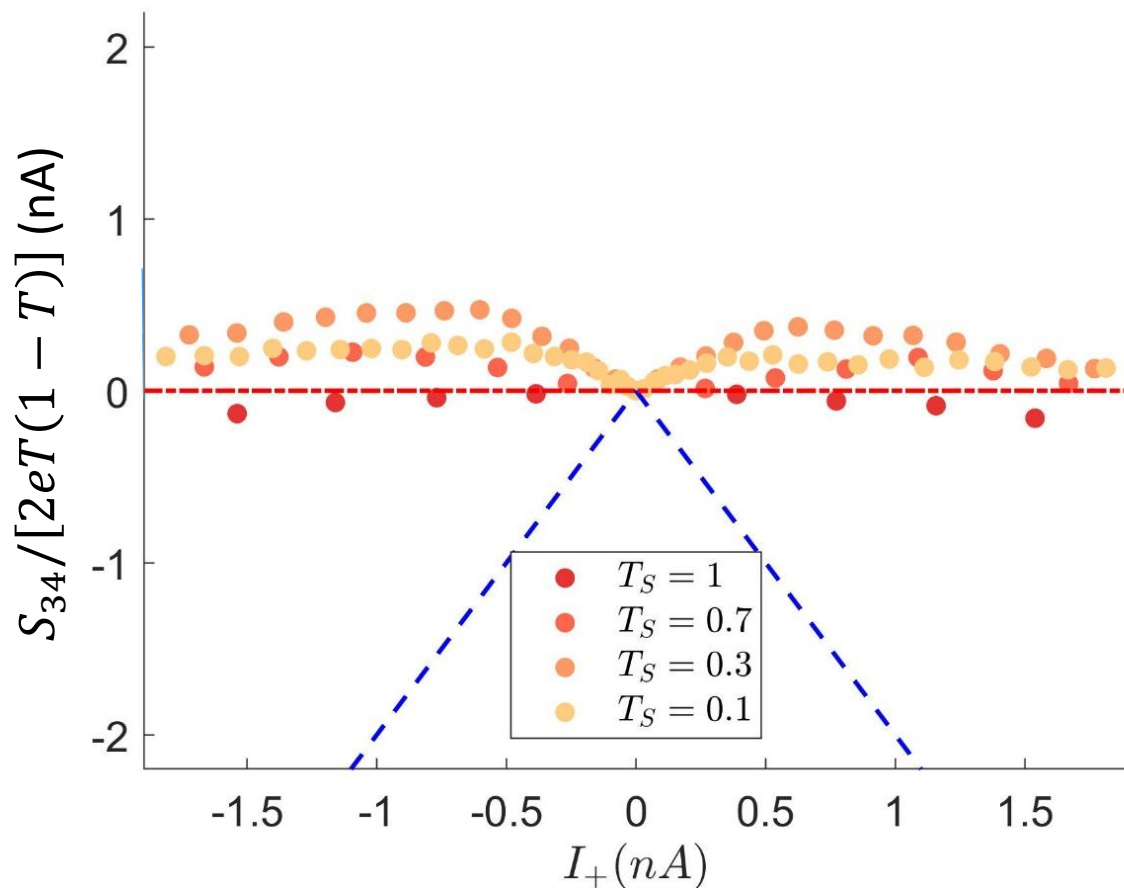
H. Bartolomei, M. Kumar et al., Science **368** 173 (2020)

M. Ruelle et al., PRX **13**, 011031 (2023)

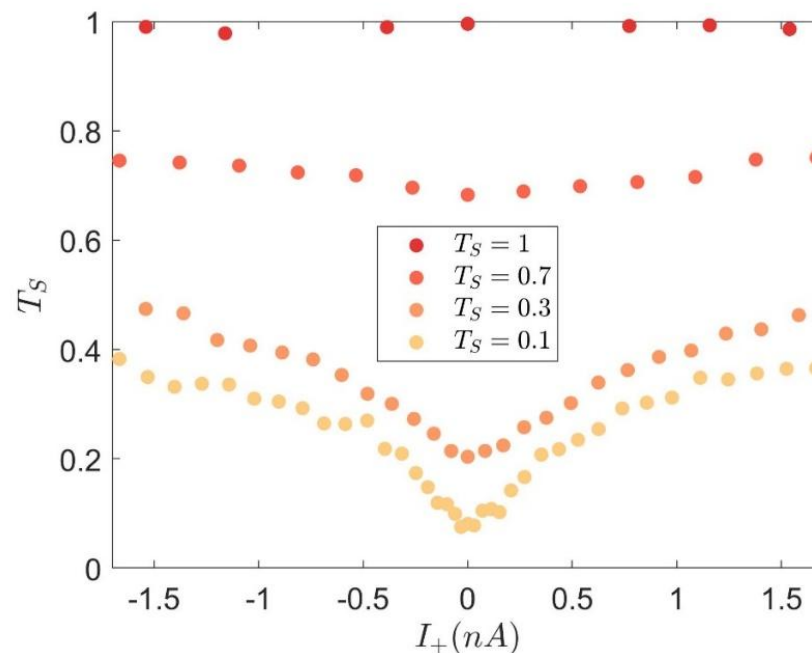
$P \approx 0$ fermions

Balanced case: $I_1^{in} = I_2^{in}$

Integer case: $q = e$, fermions
 $\nu = 3, T = 0.4, T_S = \{1; 0.7; 0.3; 0.1\}$



$P \approx 0$ fermions



H. Bartolomei, M. Kumar et al., Science **368** 173 (2020)

M. Ruelle et al., PRX **13**, 011031 (2023)

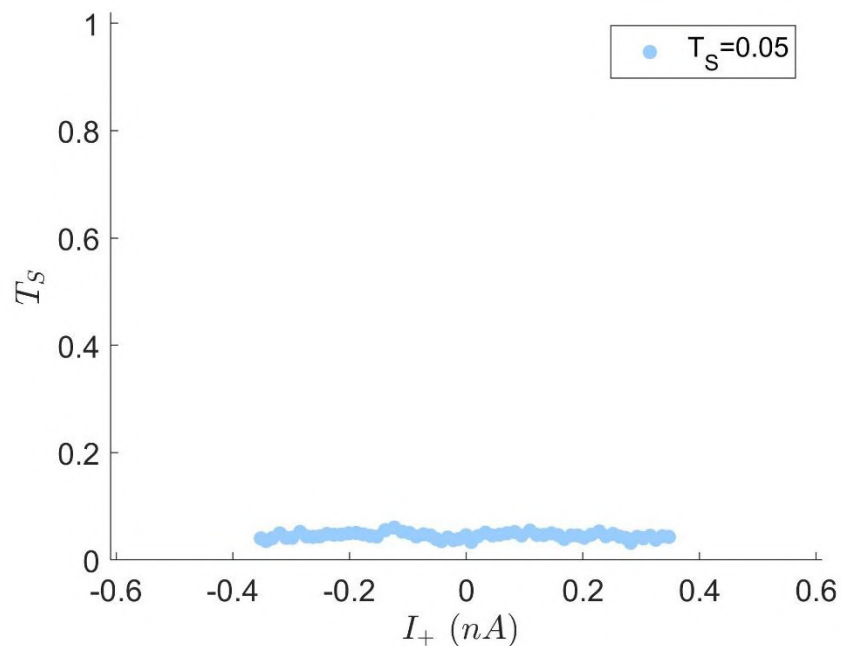
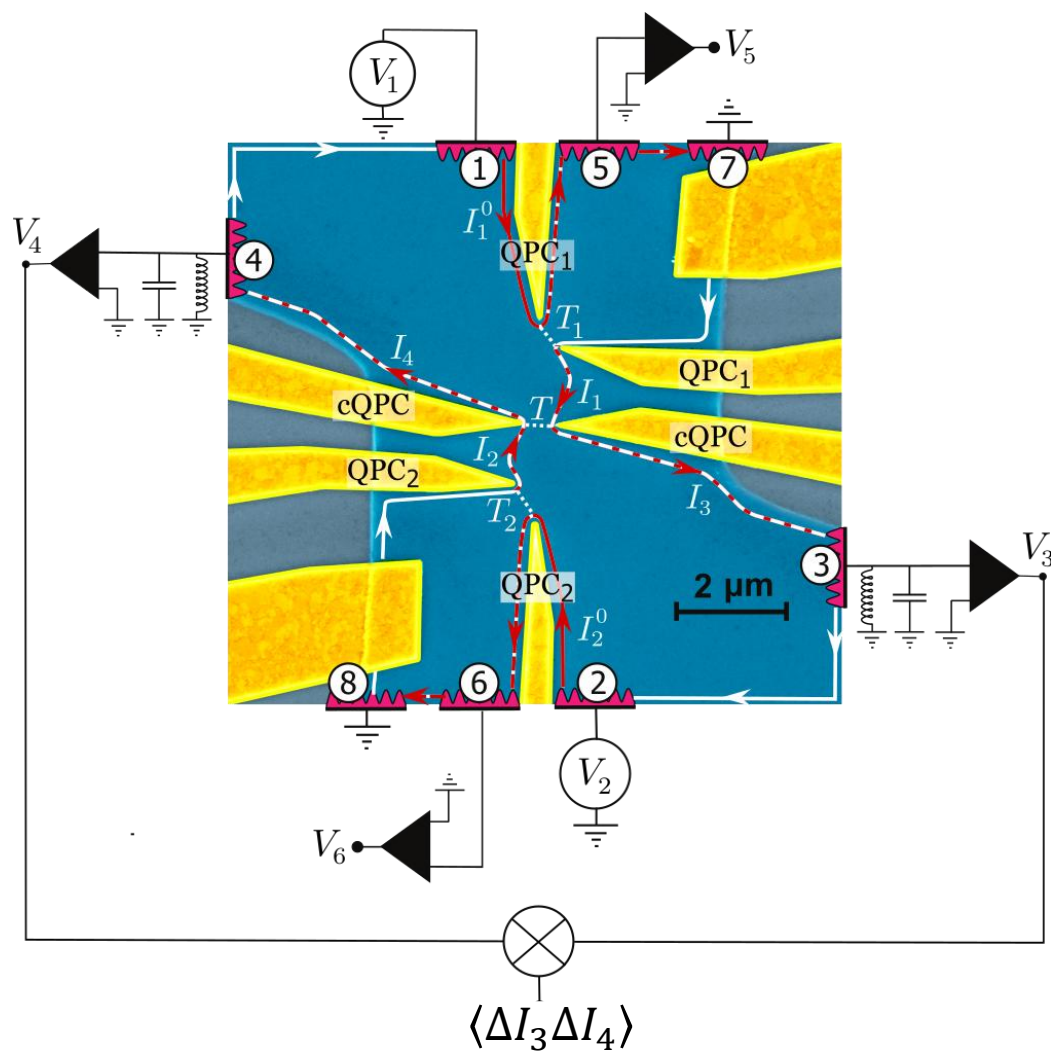
Other experiment in F. Pierre and
A. Anthore group

P. Glidic et al., Phys. Rev. X **13**, 011030 (2023).

Balanced case: $I_1^{in} = I_2^{in}$

Fractional case: $q = e/3$, anyons

$\nu = 1/3, T = 0.3, T_S = 0.05$

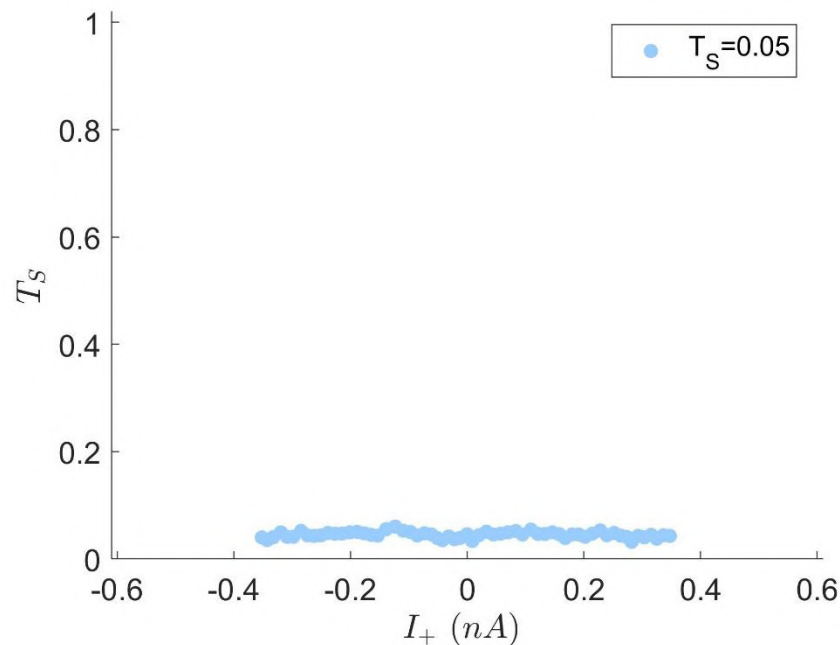
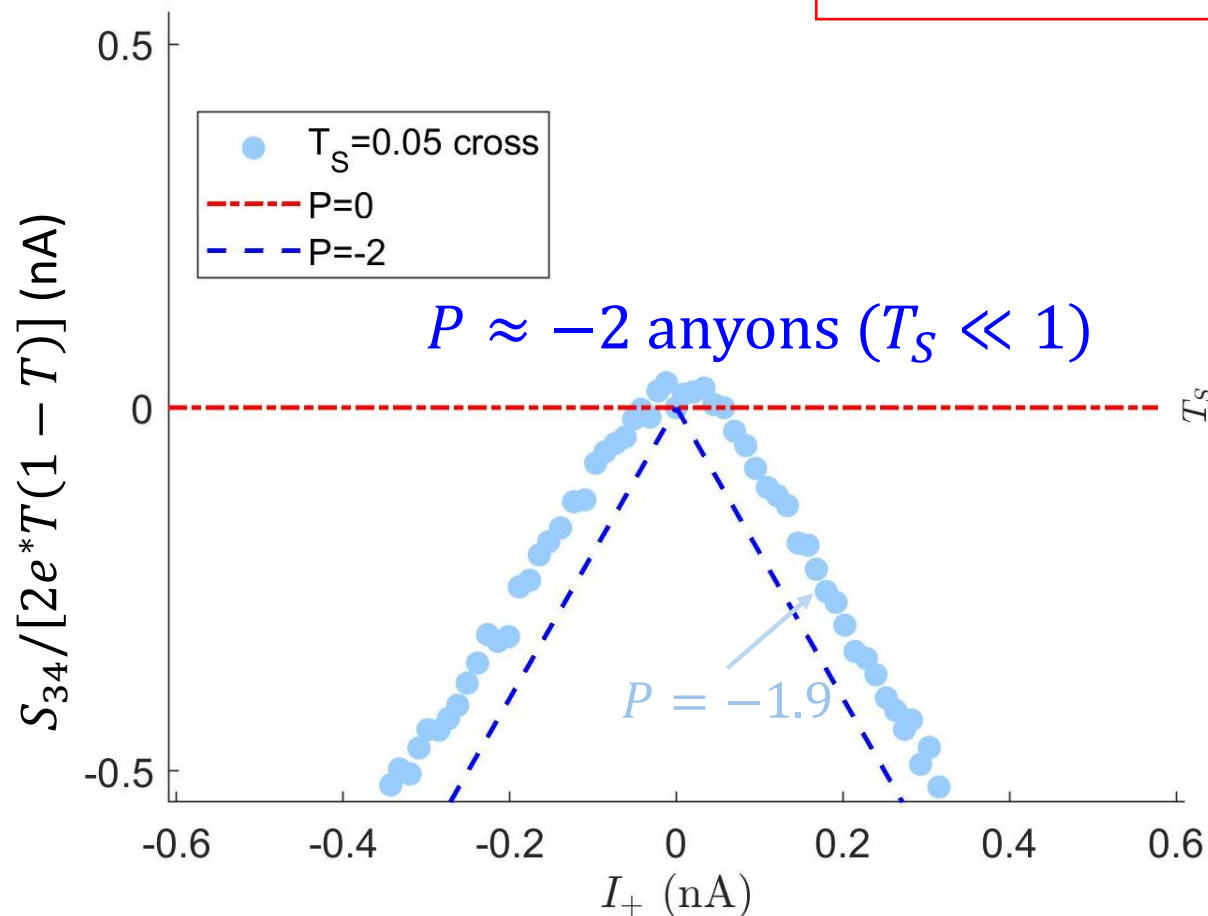


H. Bartolomei, M. Kumar et al., Science **368** 173 (2020)

M. Ruelle et al., PRX **13**, 011031 (2023)

Balanced case: $I_1^{in} = I_2^{in}$

Fractional case: $q = e/3$, anyons
 $\nu = 1/3, T = 0.3, T_S = 0.05$



H. Bartolomei, M. Kumar et al., Science **368** 173 (2020)
M. Ruelle et al., PRX **13**, 011031 (2023)

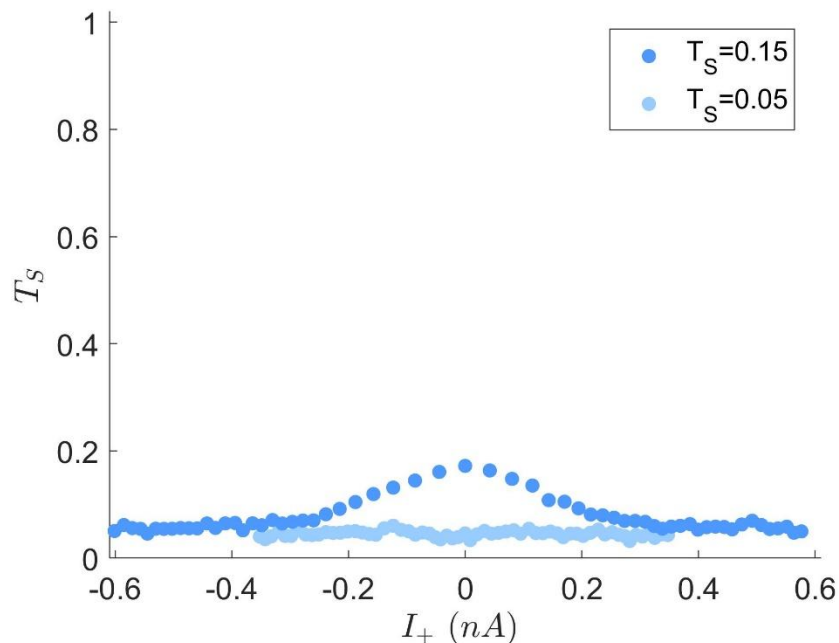
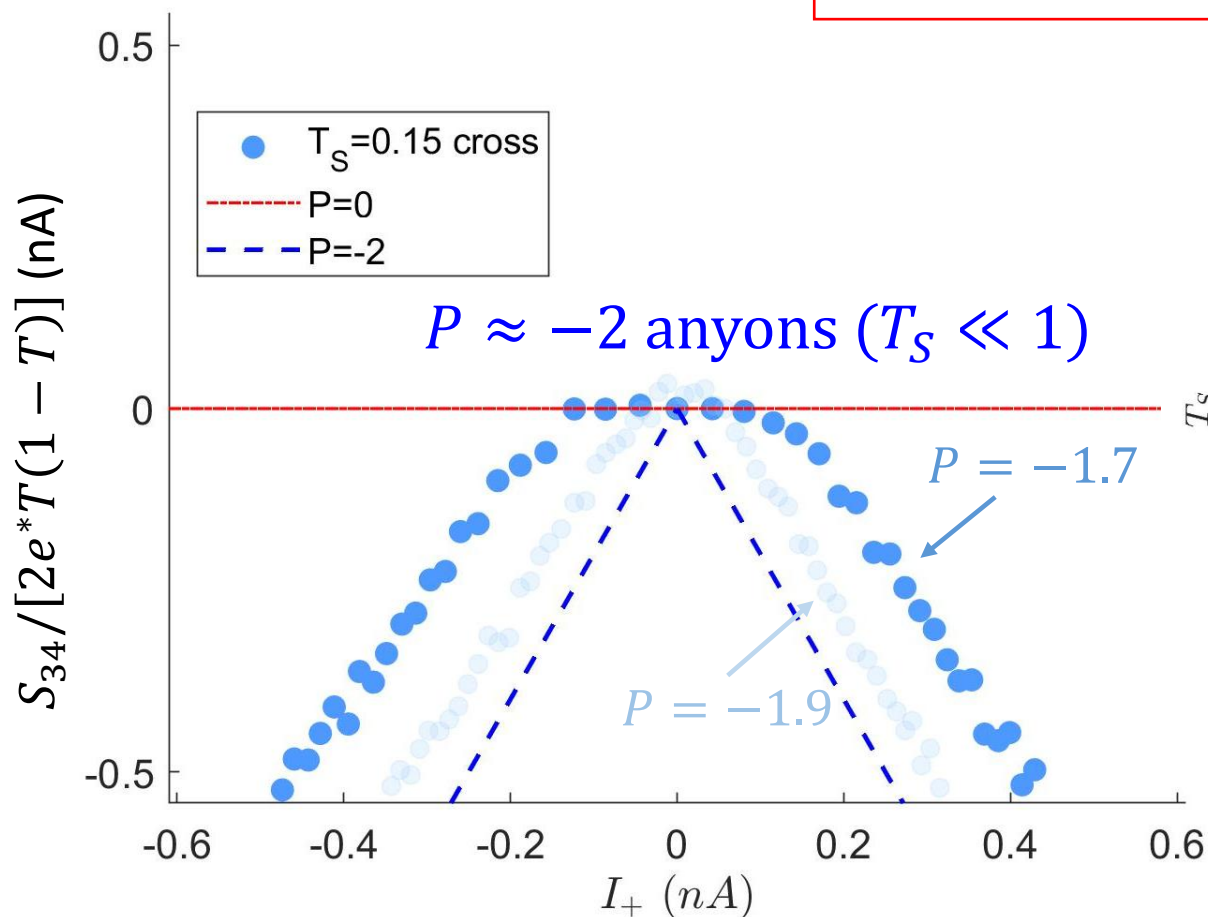
$P(\theta = 2\pi/3) = -2$, Rosenow et al. PRL 2016

$T_S \ll 1$: standard bunching (similar to photons) $\propto T_S^2$ can be excluded!

Fractional case: $q = e/3$, anyons

$\nu = 1/3, T = 0.3, T_S = 0.15$

Balanced case: $I_1^{in} = I_2^{in}$



H. Bartolomei, M. Kumar et al., Science **368** 173 (2020)
M. Ruelle et al., PRX **13**, 011031 (2023)

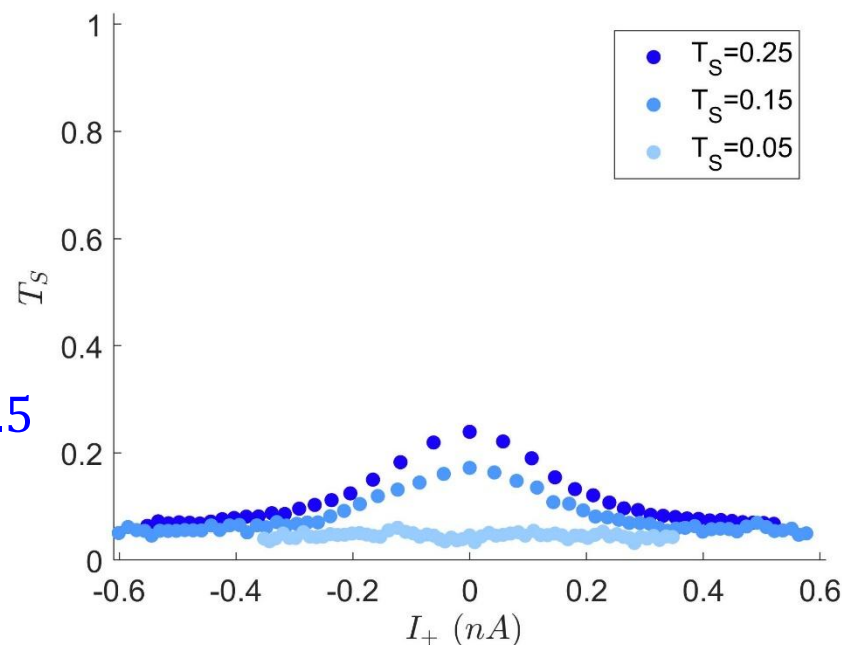
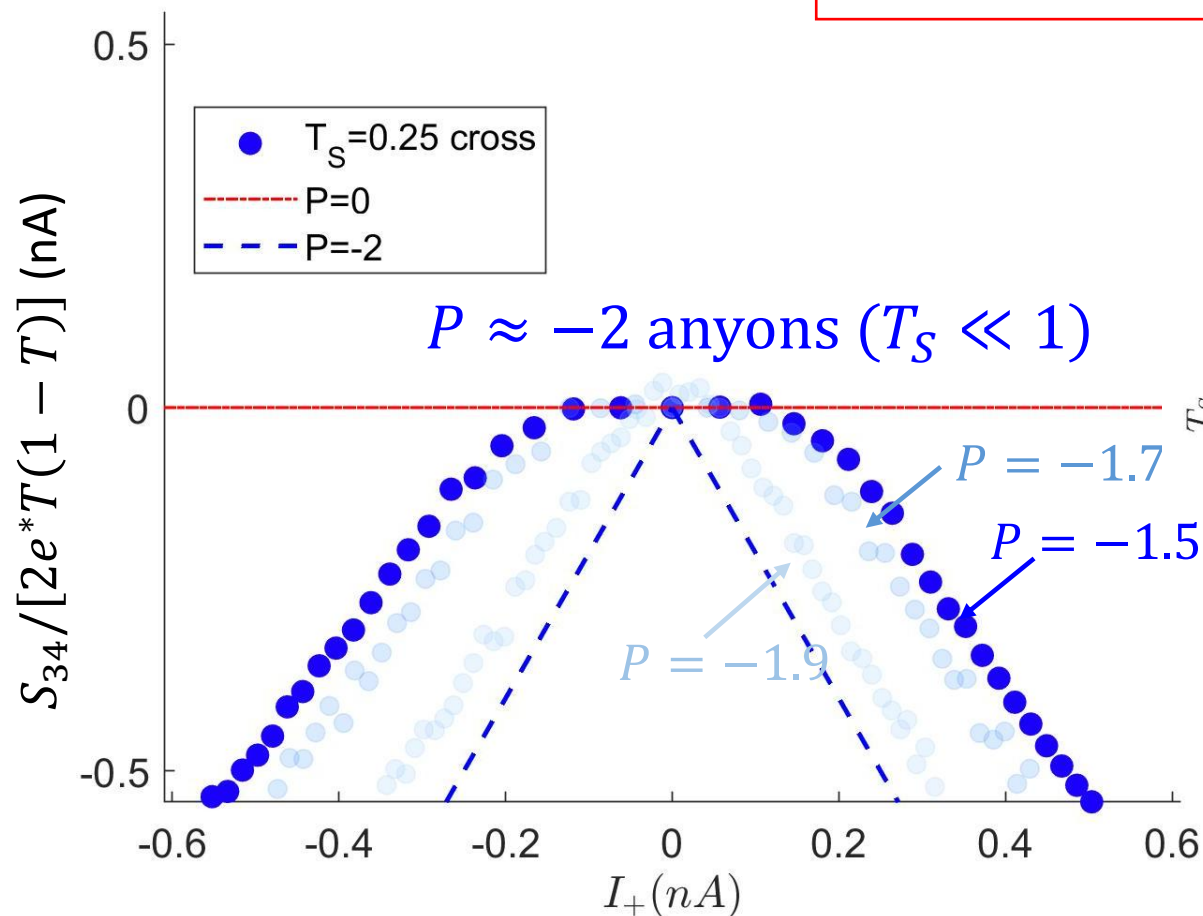
$P(\theta = 2\pi/3) = -2$, Rosenow et al. PRL 2016

$T_S \ll 1$: standard bunching (similar to photons) $\propto T_S^2$ can be excluded!

Balanced case: $I_1^{in} = I_2^{in}$

Fractional case: $q = e/3$, anyons

$\nu = 1/3, T = 0.3, T_S = 0.25$

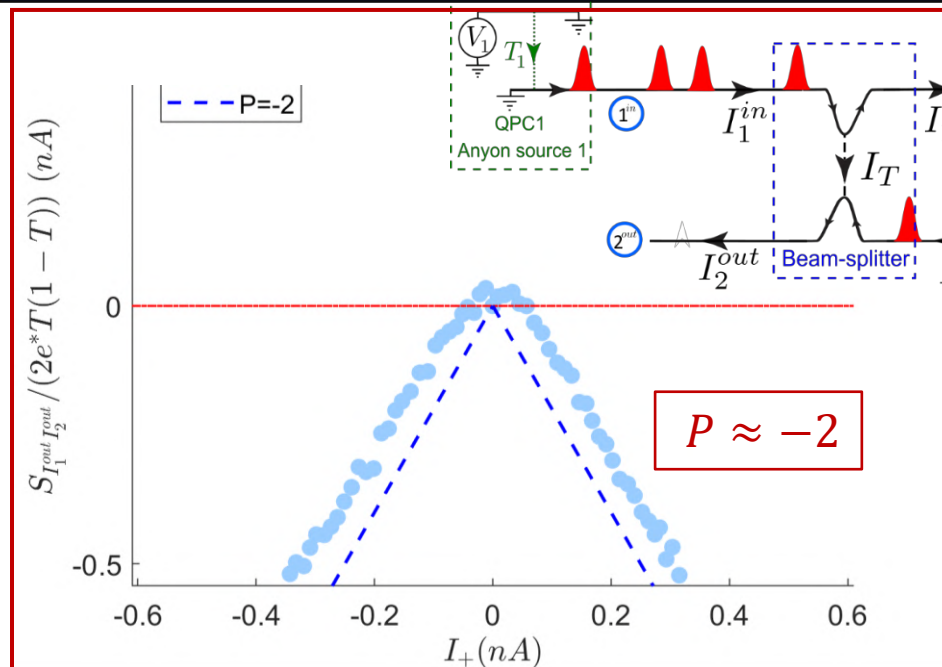


H. Bartolomei, M. Kumar et al., Science **368** 173 (2020)
M. Ruelle et al., PRX **13**, 011031 (2023)

$P(\theta = 2\pi/3) = -2$, Rosenow et al. PRL 2016

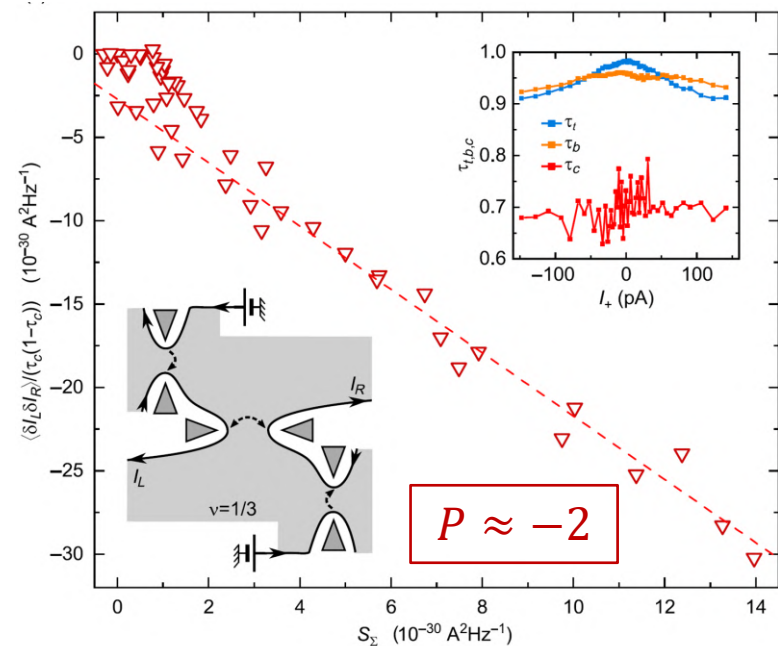
$T_S \ll 1$: standard bunching (similar to photons) $\propto T_S^2$ can be excluded!

Partitioning of diluted anyon beams, $\nu = 1/3$

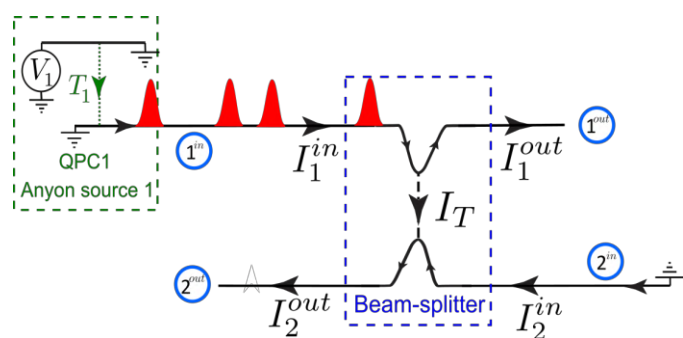


H. Bartolomei et al., Science **368** 173 (2020)

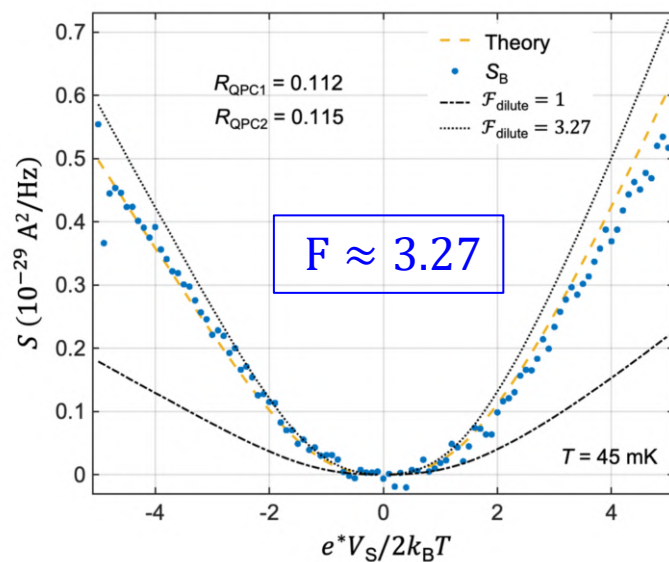
M. Ruelle et al., PRX **13**, 011031 (2023)



P. Glidic et al., PRX **13**, 011030 (2023).



J.Y.M. Lee et al., Nature **617**, 277 (2023).



Time-domain braiding:

H.S Sim KAIST

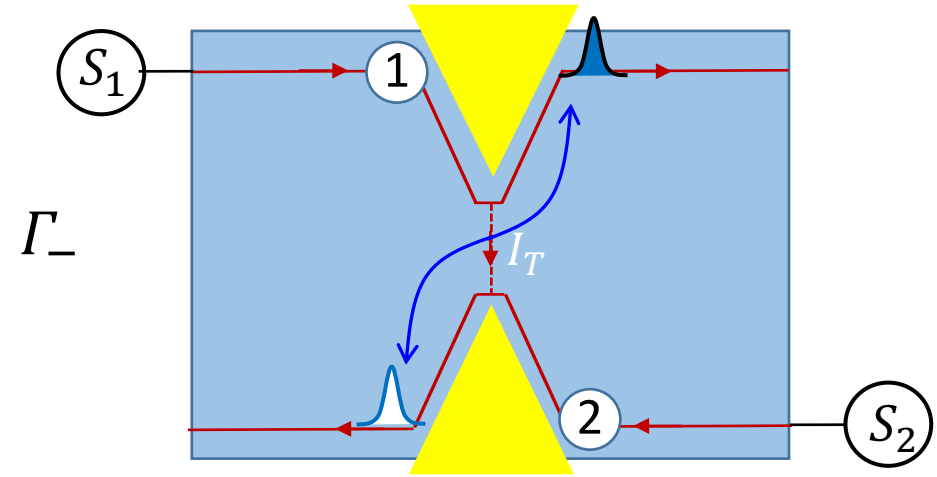
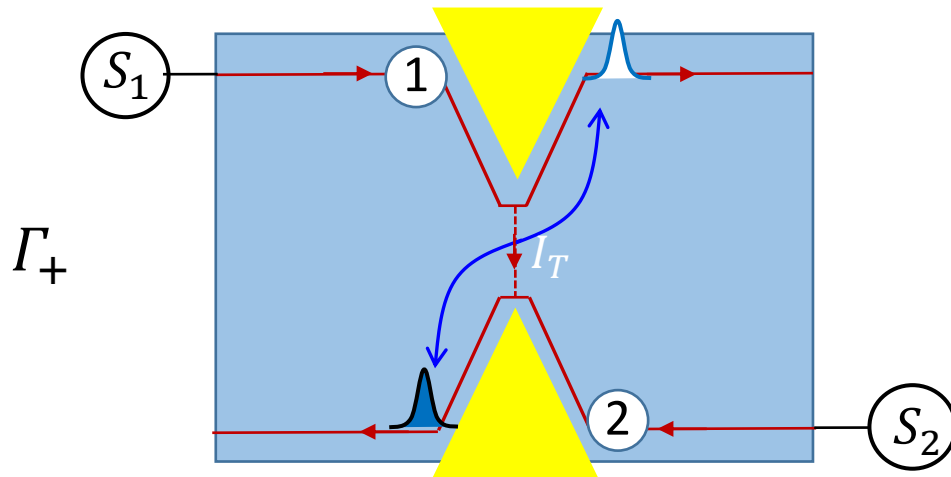
Han et al., Nat. Commun. **7**, 11131 (2016)

Morel et al., PRB **105**, 075433 (2022)

Lee et al., Nat. Commun. **13**, 6660 (2022)

Schiller et al., PRL **131** 186601 (2023)

Weak backscattering regime: lowest order in tunneling $H_T = \zeta \psi_{2,a}^+ \psi_{1,a} + \zeta^* \psi_{1,a}^+ \psi_{2,a}$



$$I_T = q(\Gamma_+ - \Gamma_-)$$

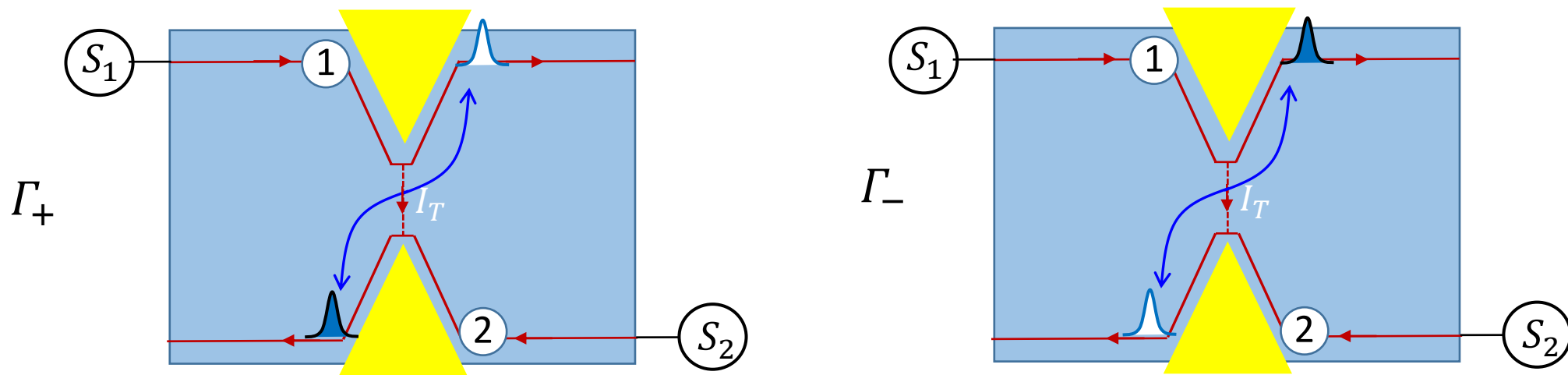
Tunneling current

$$S_{I_T} = 2q^2(\Gamma_+ + \Gamma_-)$$

Noise of tunneling current

Time-domain braiding: Han et al., Nat. Commun. **7**, 11131 (2016) Lee et al., Nat. Commun. **13**, 6660 (2022)
 H.S Sim KAIST Morel et al., PRB **105**, 075433 (2022) Schiller et al., PRL **131** 186601 (2023)

Weak backscattering regime: lowest order in tunneling $H_T = \zeta \psi_{2,a}^+ \psi_{1,a} + \zeta^* \psi_{1,a}^+ \psi_{2,a}$



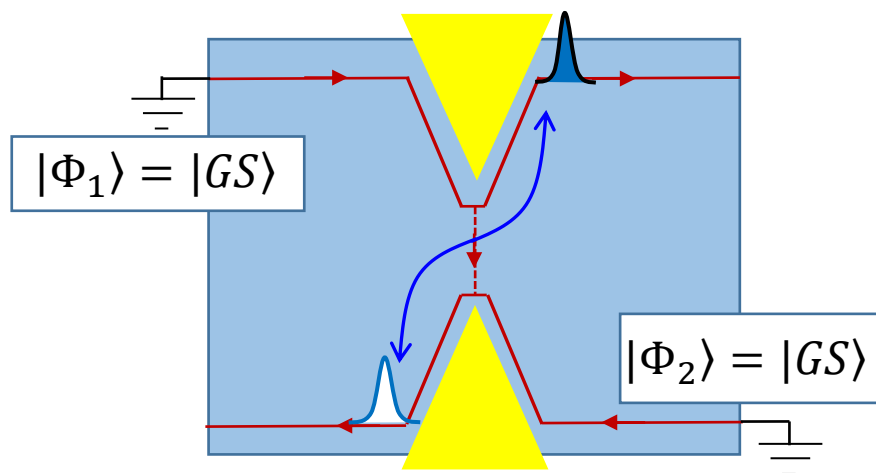
$$I_T = q(\Gamma_+ - \Gamma_-) \quad \text{Tunneling current}$$

$$S_{I_T} = 2q^2(\Gamma_+ + \Gamma_-) \quad \text{Noise of tunneling current}$$

$$\Gamma_-(t) \propto \int_0^{+\infty} d\tau \underbrace{\langle \psi_{1,a}(t) \psi_{1,a}^+(t-\tau) \rangle \langle \psi_{2,a}^+(t) \psi_{2,a}(t-\tau) \rangle}_{\text{Two particle interferometry}} + h.c.$$

$$\Gamma_+(t) \propto \dots$$

- Rate Γ_-



Han et al., Nat. Commun. **7**, 11131 (2016)

Morel et al., PRB **105**, 075433 (2022)

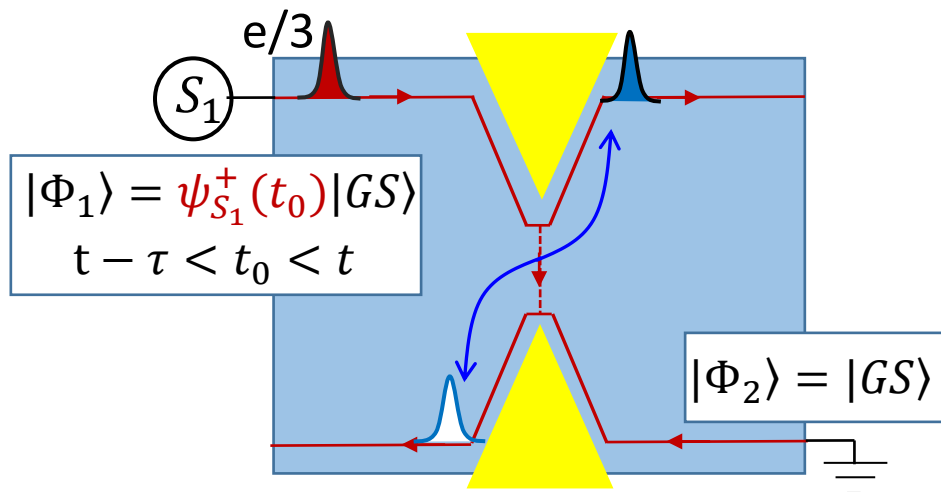
Lee et al., Nat. Commun. **13**, 6660 (2022)

Schiller et al., PRL **131** 186601 (2023)

$$\langle GS | \psi_1(t) \psi_1^\dagger(t - \tau) | GS \rangle = \langle GS | \psi_2^\dagger(t) \psi_2(t - \tau) | GS \rangle = G_{eq}(\tau)$$

Equilibrium correlations

- Rate Γ_-



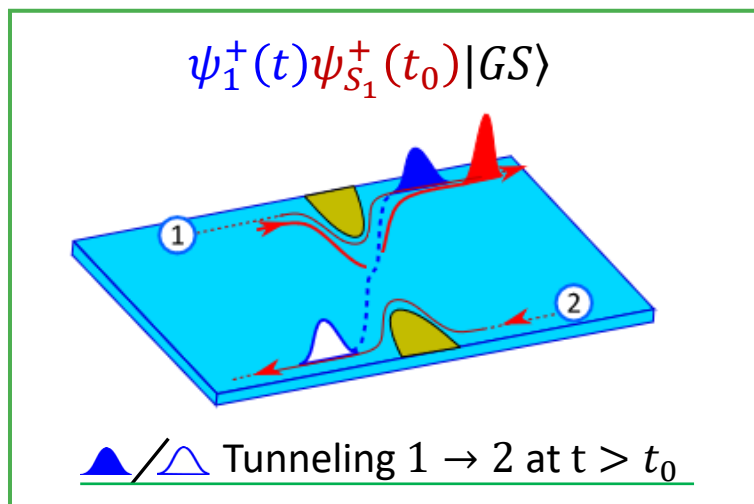
Han et al., Nat. Commun. **7**, 11131 (2016)

Morel et al., PRB **105**, 075433 (2022)

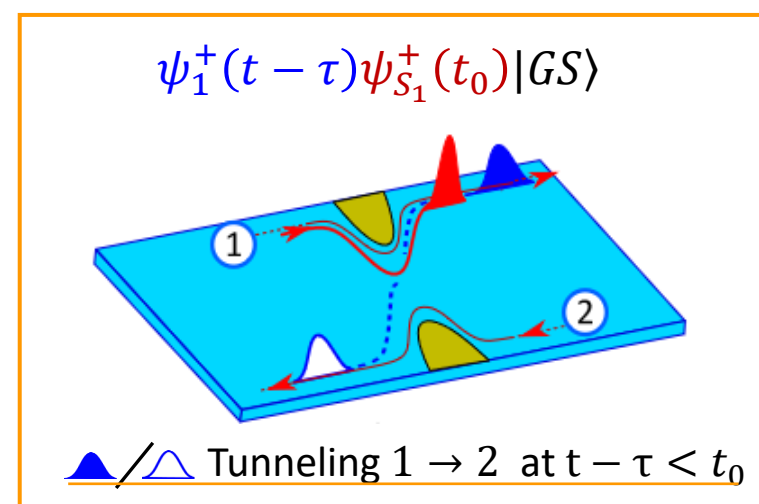
Lee et al., Nat. Commun. **13**, 6660 (2022)

Schiller et al., PRL **131** 186601 (2023)

$$\langle \Phi_1 | \psi_1(t) \psi_1^+(t - \tau) | \Phi_1 \rangle = \langle GS | \underbrace{\psi_{S_1}(t_0) \psi_1(t)}_{\text{green}} \underbrace{\psi_1^+(t - \tau) \psi_{S_1}^+(t_0)}_{\text{yellow}} | GS \rangle$$

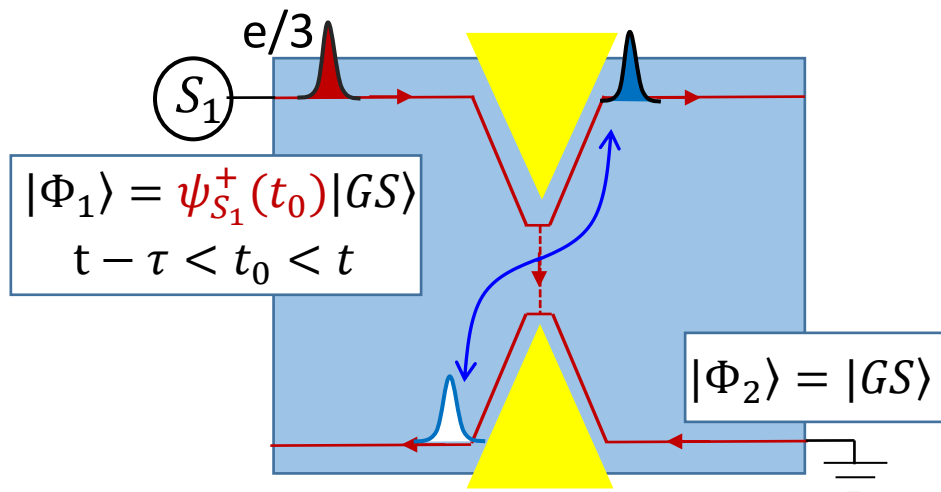


+



Interferences between tunneling of blue anyon at times t and $t - \tau$

- Rate Γ_-



Han et al., Nat. Commun. **7**, 11131 (2016)

Morel et al., PRB **105**, 075433 (2022)

Lee et al., Nat. Commun. **13**, 6660 (2022)

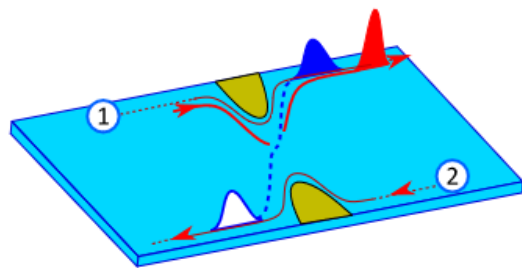
Schiller et al., PRL **131** 186601 (2023)

$$\langle \Phi_1 | \psi_1(t) \psi_1^+(t-\tau) | \Phi_1 \rangle = \langle GS | \underbrace{\psi_{S_1}(t_0) \psi_1(t)}_{\text{green}} \underbrace{\psi_1^+(t-\tau) \psi_{S_1}^+(t_0)}_{\text{yellow}} | GS \rangle$$

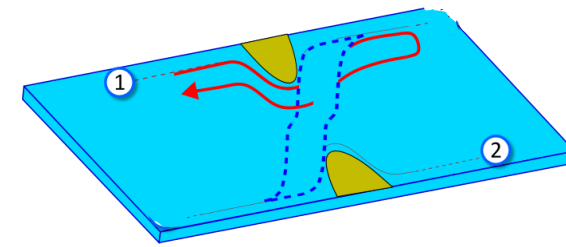
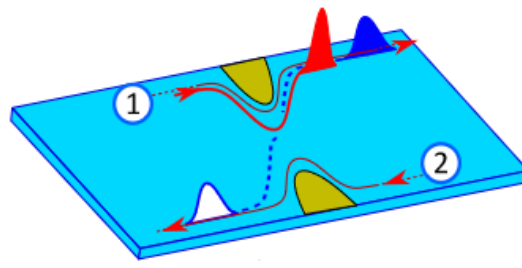
$$\psi_1^+(t) \psi_{S_1}^+(t_0) | GS \rangle$$

$$\psi_1^+(t-\tau) \psi_{S_1}^+(t_0) | GS \rangle$$

$$\langle \psi_1(t) \psi_1^+(t-\tau) \rangle = e^{-i\theta_{rb}} G_{eq}(\tau)$$



+

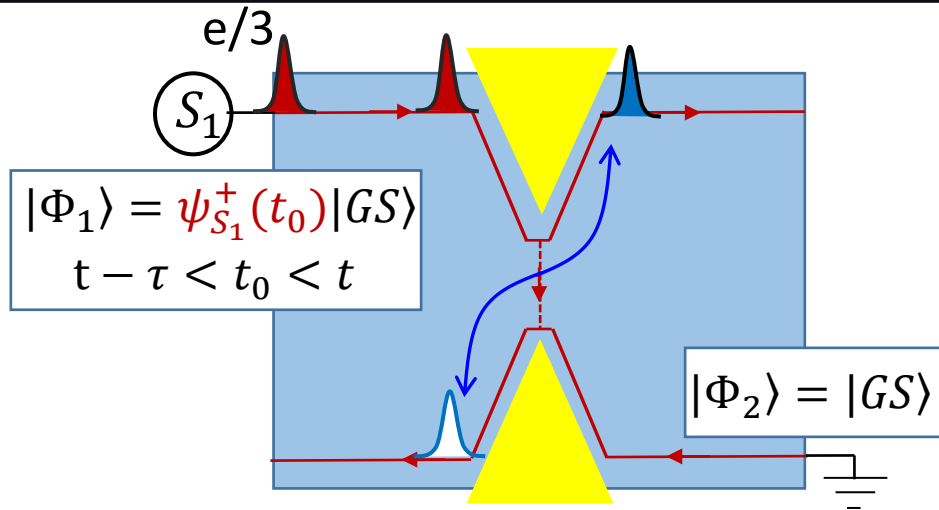


/ Tunneling 1 \rightarrow 2 at $t > t_0$

/ Tunneling 1 \rightarrow 2 at $t - \tau < t_0$

θ_{rb} : braiding phase between red and blue anyons: $\nu = 1/3$, $\theta_{rb} = 2\pi/3$

- Rate Γ_{\pm}



Han et al., Nat. Commun. **7**, 11131 (2016)

Morel et al., PRB **105**, 075433 (2022)

Lee et al., Nat. Commun. **13**, 6660 (2022)

Schiller et al., PRL **131** 186601 (2023)

$$\Gamma_{\pm}(t) \propto \text{Re} \left[\int_{-\infty}^t e^{\pm i\theta N_1(t,t')} G_{eq,\delta}(t-t')^2 dt' \right]$$

$N_1(t, t')$ set by the nature of the source S_1 (stochastic in «collider»)

θ : braiding phase (one expects $\theta = 2\pi/3$ at $\nu = 1/3$)

δ : scaling dimension, time-correlations at the edge $G_{eq,\delta}(t-t') \sim \frac{1}{|t-t'|^\delta}$

one expects $\delta = 1/3$ at $\nu = 1/3$

X. G Wen, Advances in Physics, 44, 405 (1995)

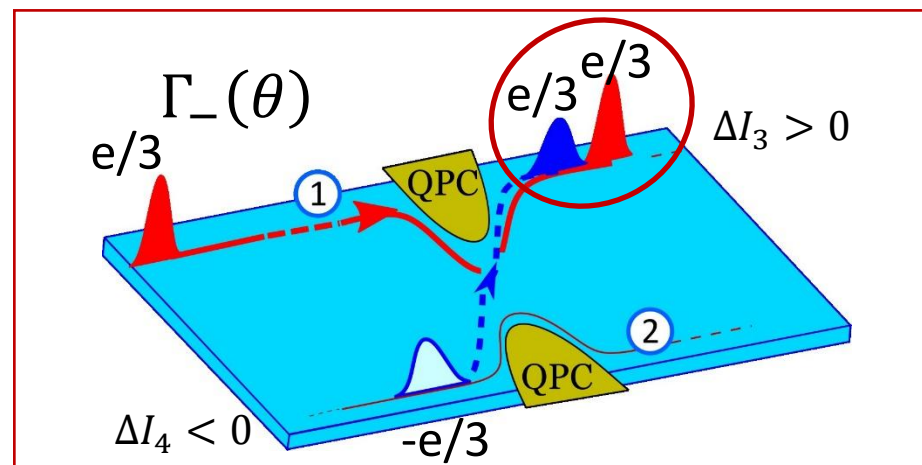
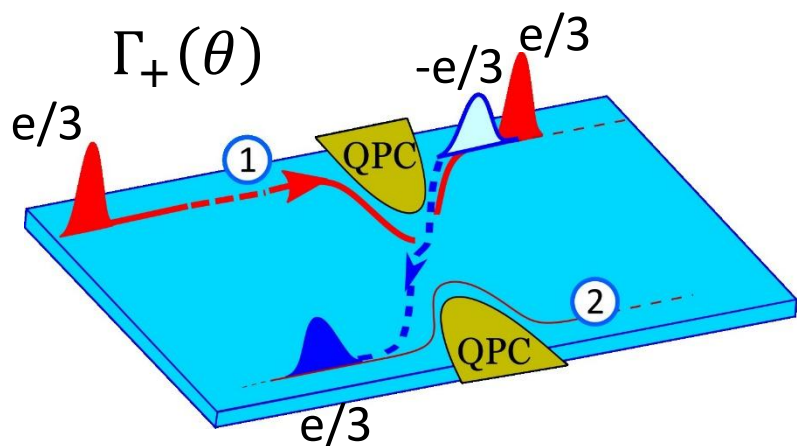
I. Safi, arXiv:2510.10525

A. Latyshev and I. Safi, arXiv:2510.2059

$\delta = 1/3$ at $\nu = 1/3$ observed in

A. Veillon et al., Nature **632**, 517–521 (2024).

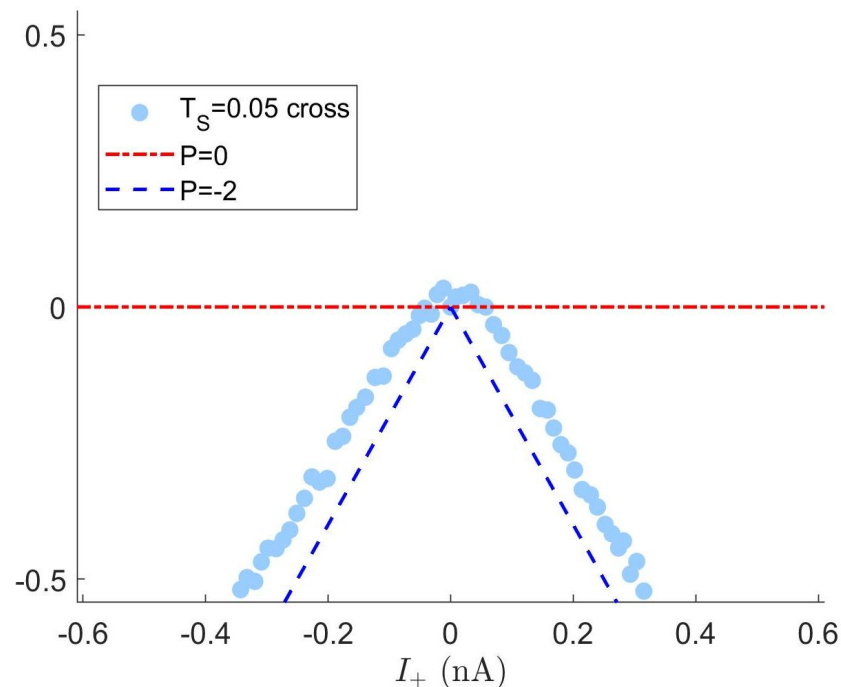
R. Guerrero-Suarez et al, Nat. Phys. **21**, 1787 (2025)

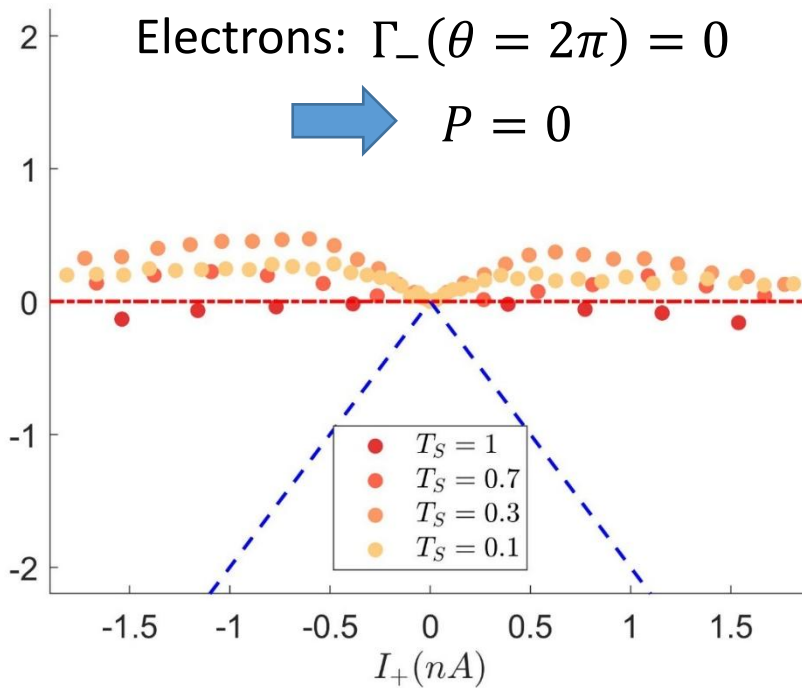
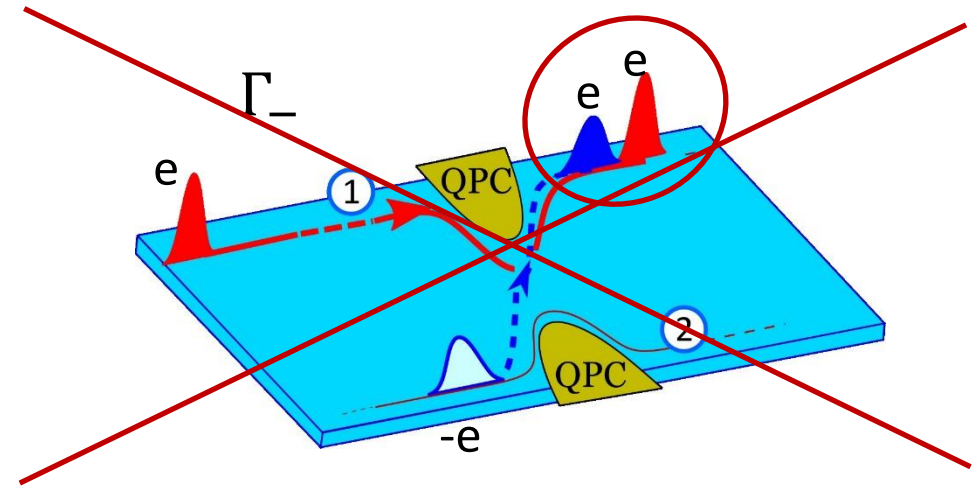
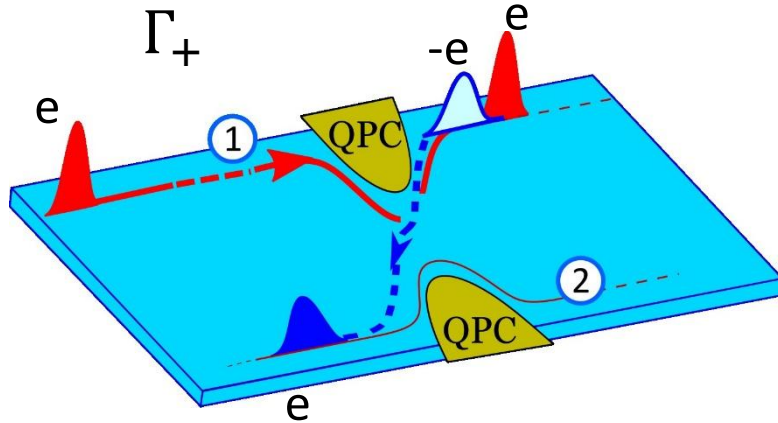


Anyons: $\Gamma_-(\theta \neq 2\pi) \neq 0$: bunching

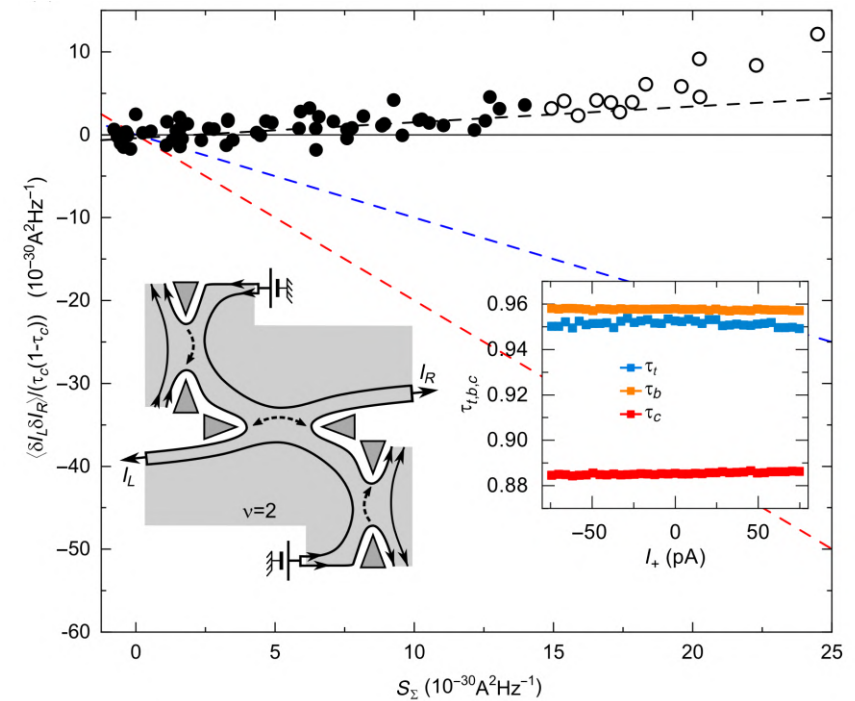
➔ $P < 0$ and $P(\theta = 2\pi/3) = -2$

One anyon in 1 is enough to produce a pair of anyons in 1 together with an anyon hole in 2



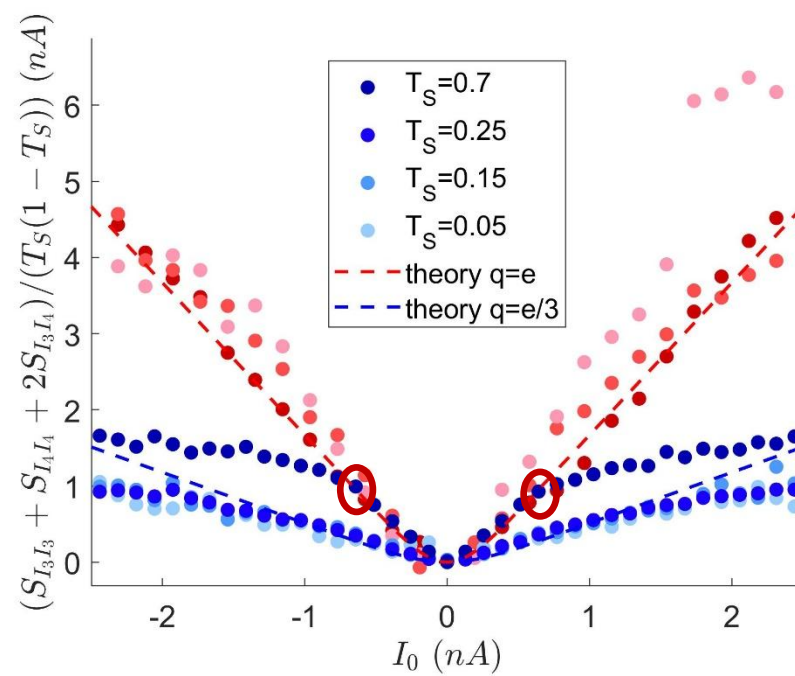
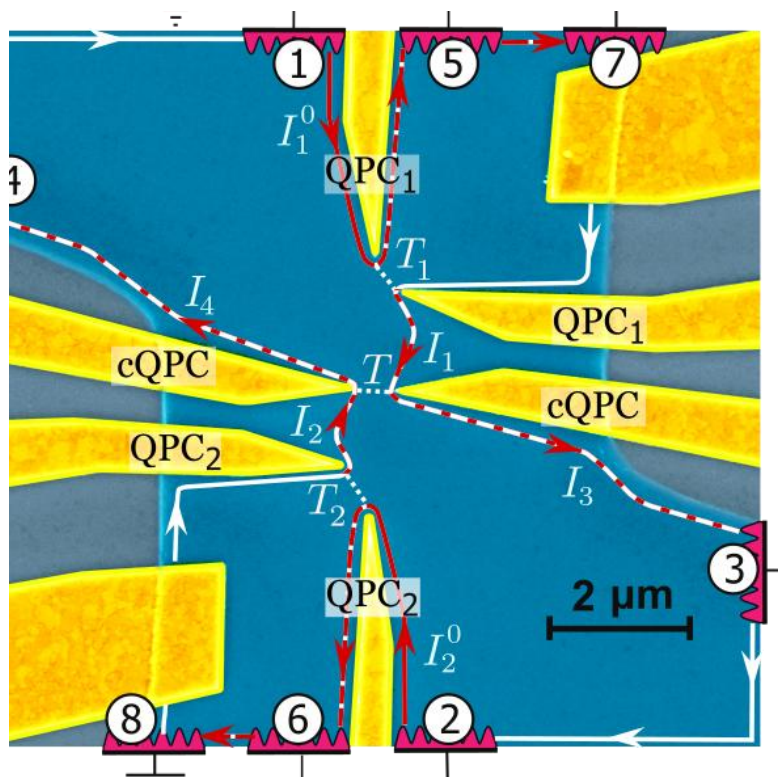
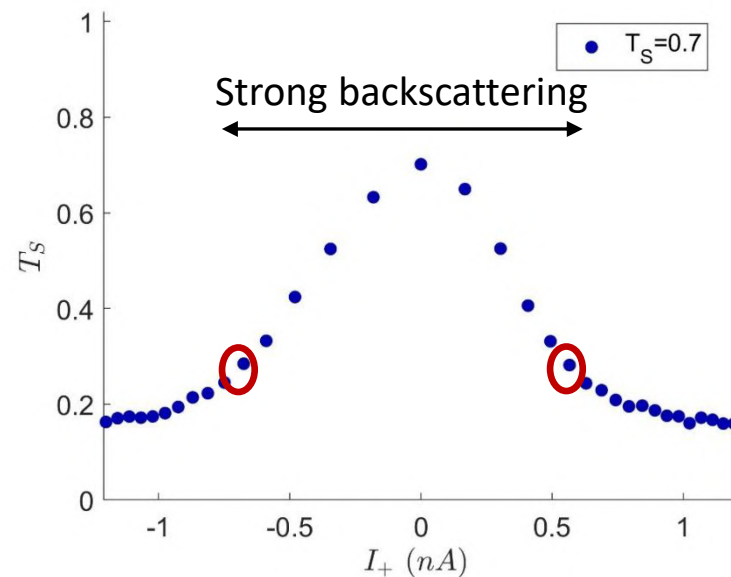
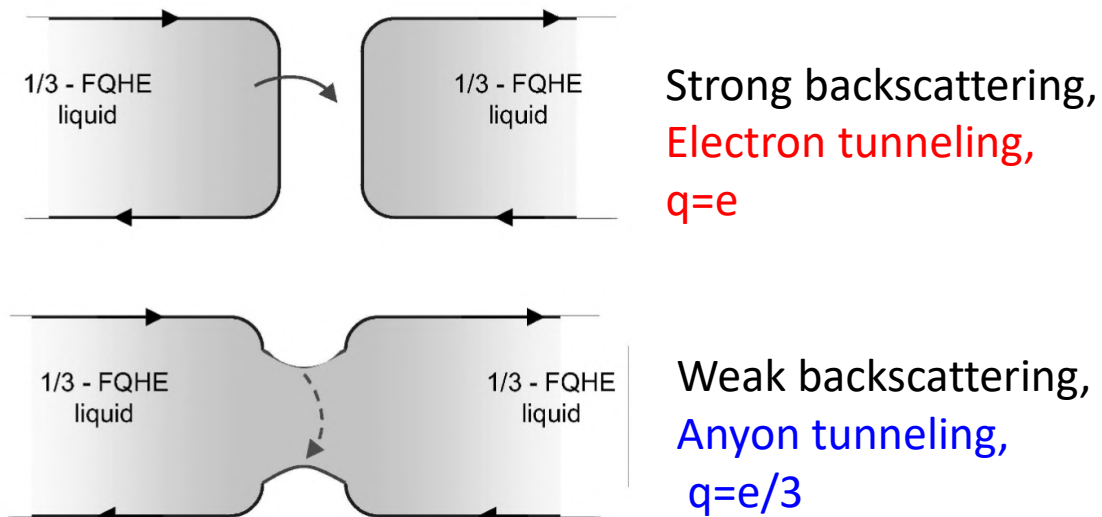


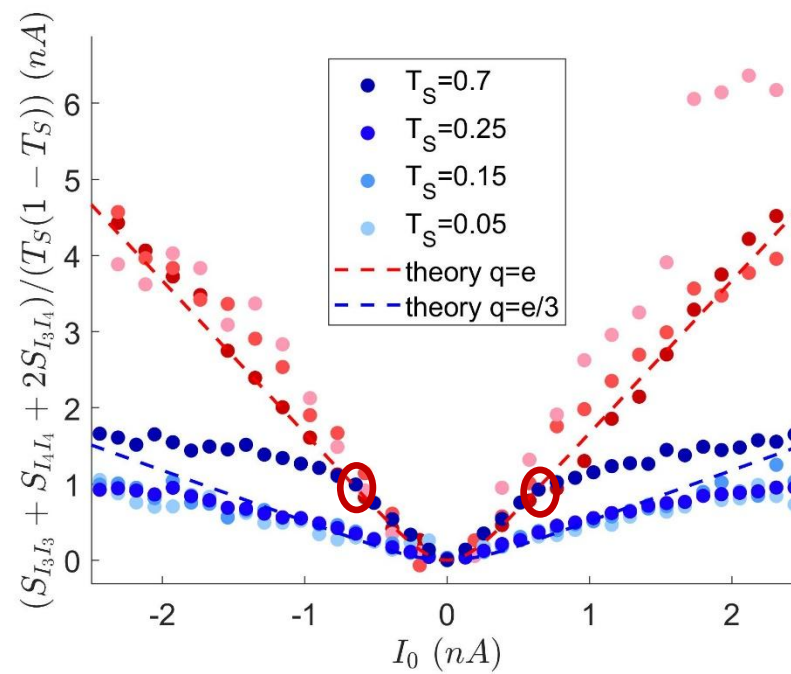
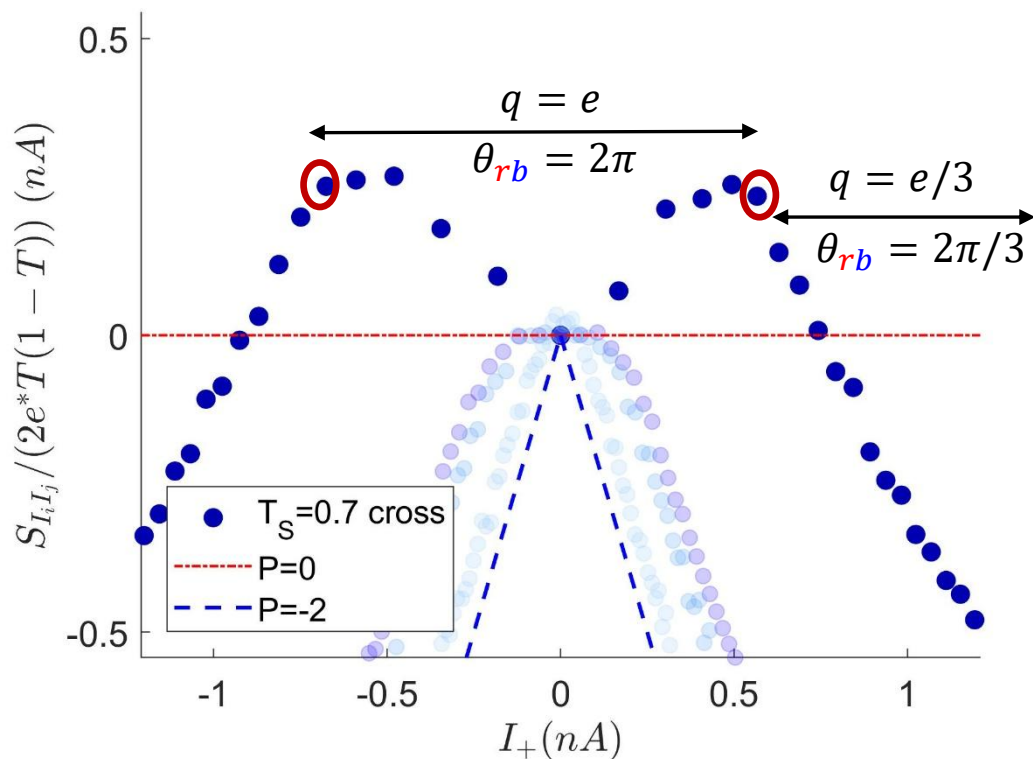
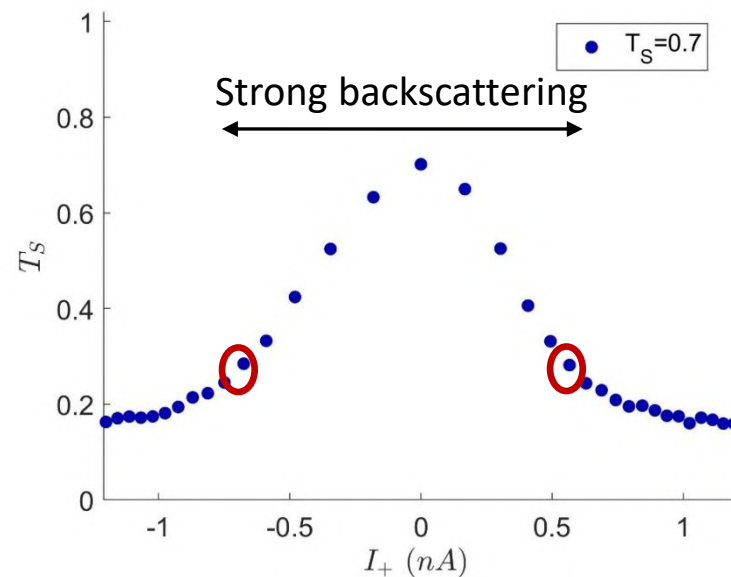
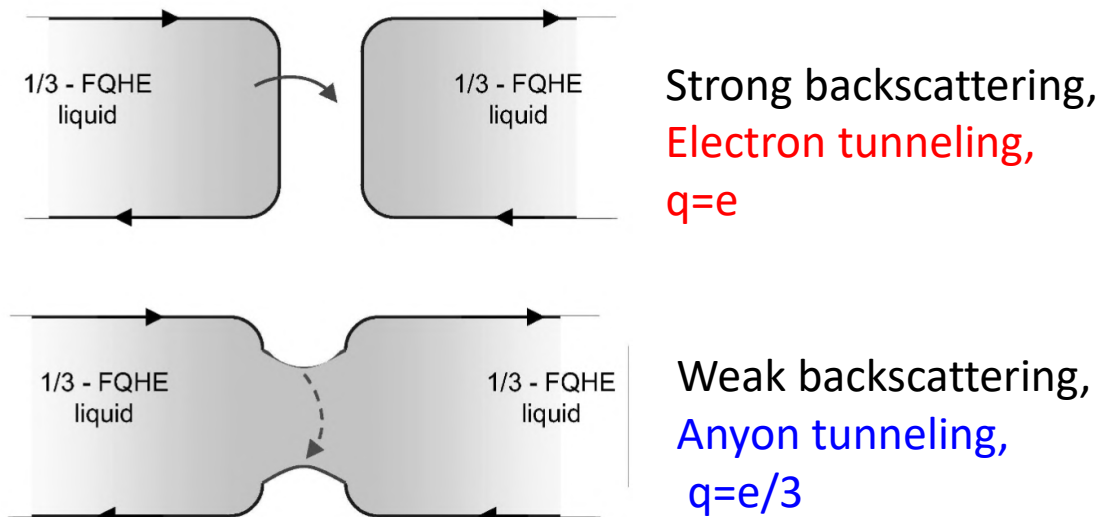
H. Bartolomei, M. Kumar et al., Science **368** 173 (2020)



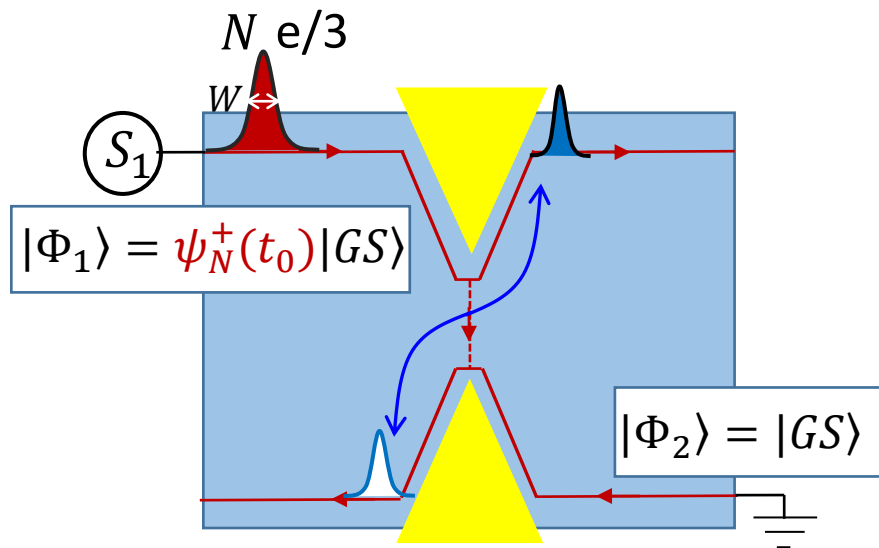
P. Glidic et al., Phys. Rev. X **13**, 011030 (2023).

Electron and anyon tunneling, $\nu = 1/3$





T. Jonckheere, et al., PRL **130**, 186203 (2023).

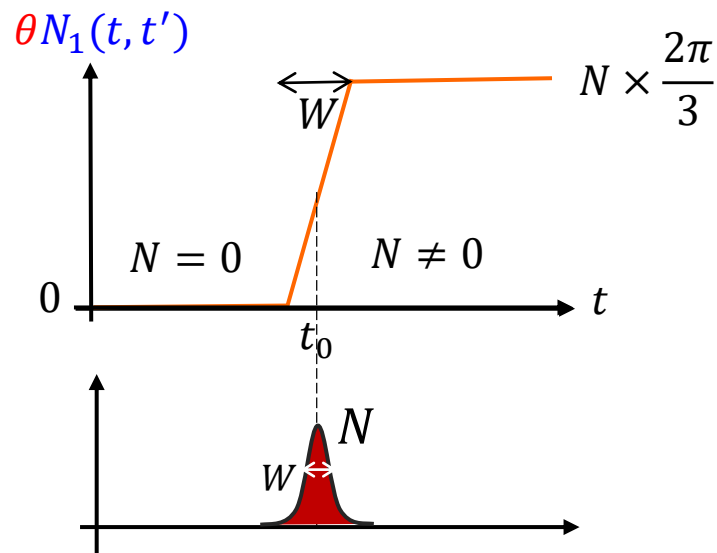


$$\Gamma_{\pm}(t) \propto \text{Re} \left[\int_{-\infty}^t e^{\pm i\theta N(t,t')} G_{eq}(t-t')^2 dt' \right]$$

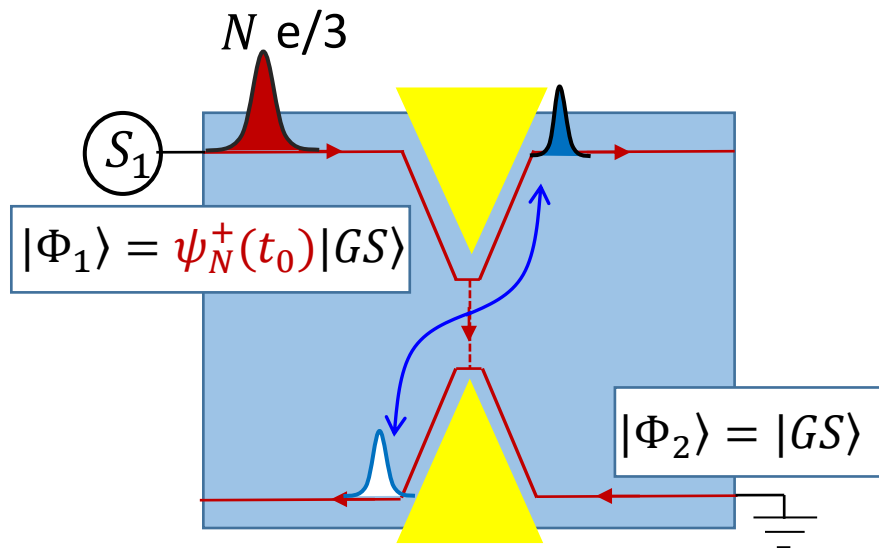
- $N = 1$: single anyon emission at $\nu = 1/3$
- $N = 3$: single electron emission at $\nu = 1/3$

Stochastic source replaced by deterministic anyon source:

 N anyons cross the QPC at time t_0



T. Jonckheere, et al., PRL **130**, 186203 (2023).



$$I_T = q(\Gamma_+ - \Gamma_-)$$

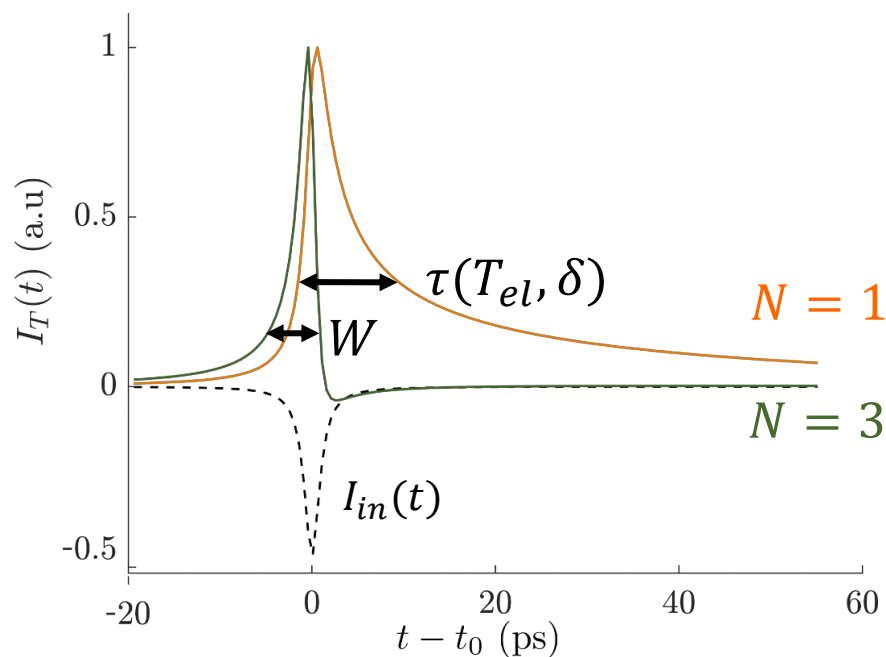
Braiding phase factor,
varies on timescale W

Memory or correlation time
decays on $\tau(T_{el}, \delta)$

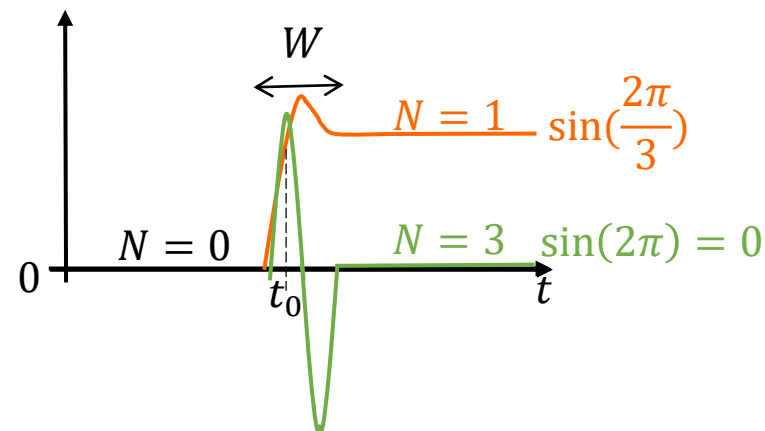
$$I_T(t) \propto \int_{-\infty}^t dt' \sin(\theta N_1(t, t')) \text{Im}[G_\delta(t - t')^2]$$

$N = 1$: single anyon emission at $\nu = 1/3$

$N = 3$: single electron emission at $\nu = 1/3$

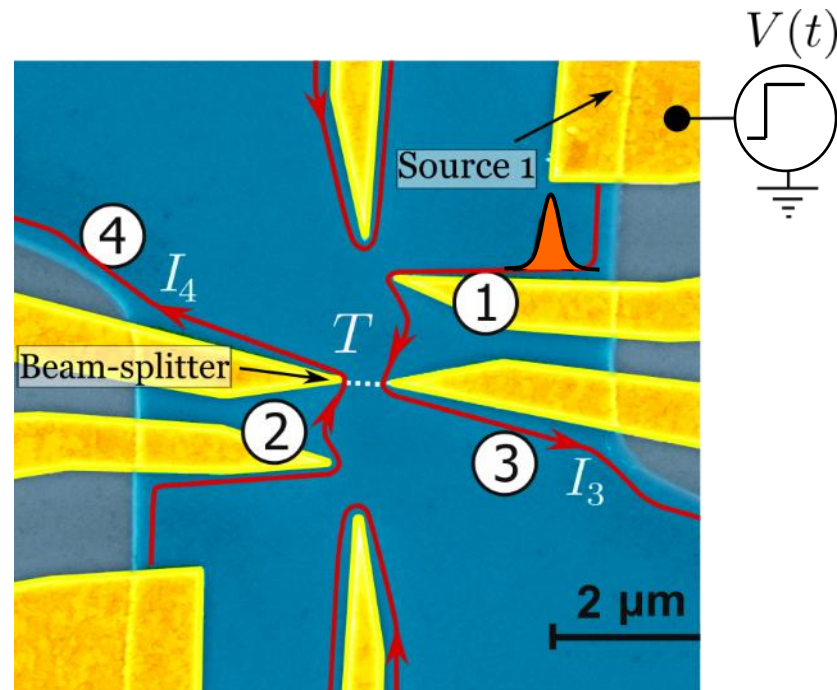
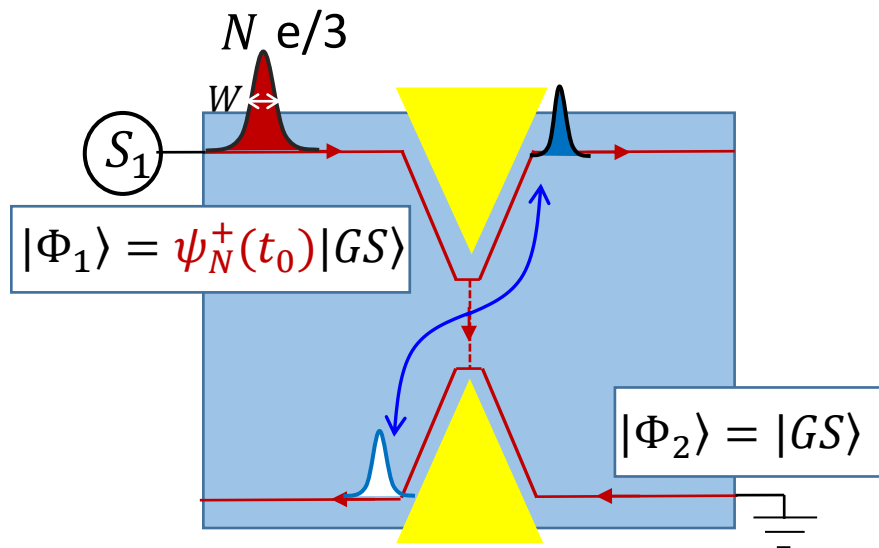


$$\sin(\theta N_1(t, t'))$$



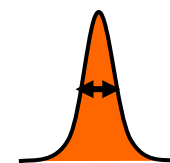
Brings information on both θ and δ

T. Jonckheere, et al., PRL **130**, 186203 (2023).



Current pulse carrying charge $Ne/3$
emulates N –anyon source

$N e/3$



$W \approx 100 \text{ ps}$

$$N = CV/e^*$$

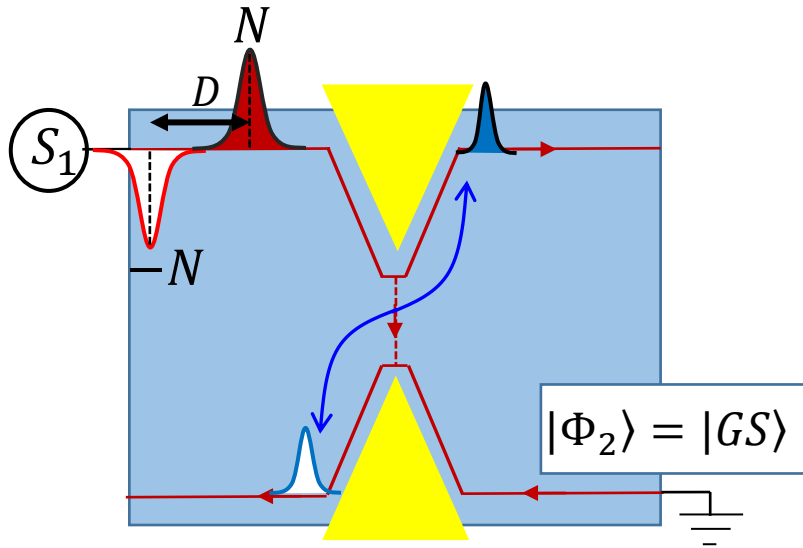
Calibrated capacitance...

J.Y. Lee, C. Han, H.S. Sim, PRL **125**, 196802 (2020)

C. Mora, arXiv:2212.05123 (2022)

M. Ruelle et al., Science **389**, eadm7695 (2025).

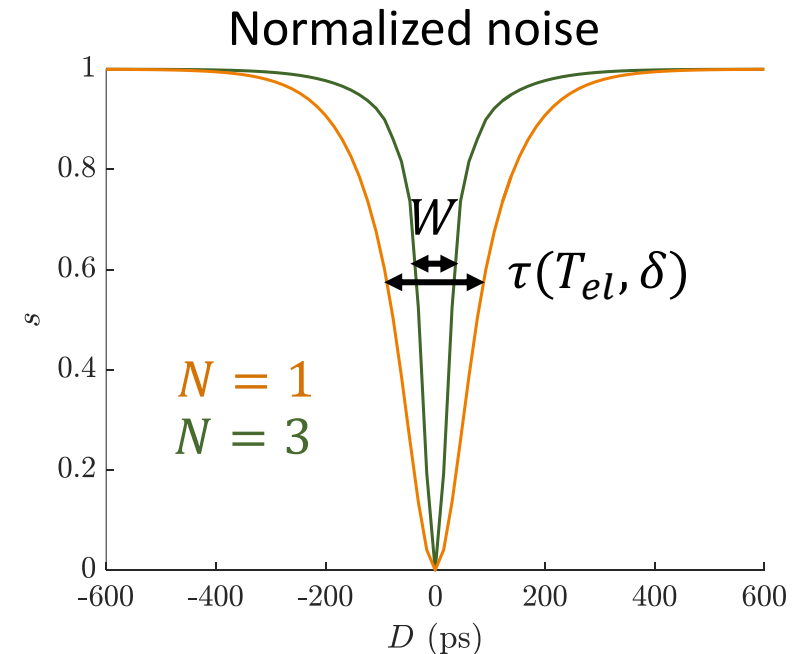
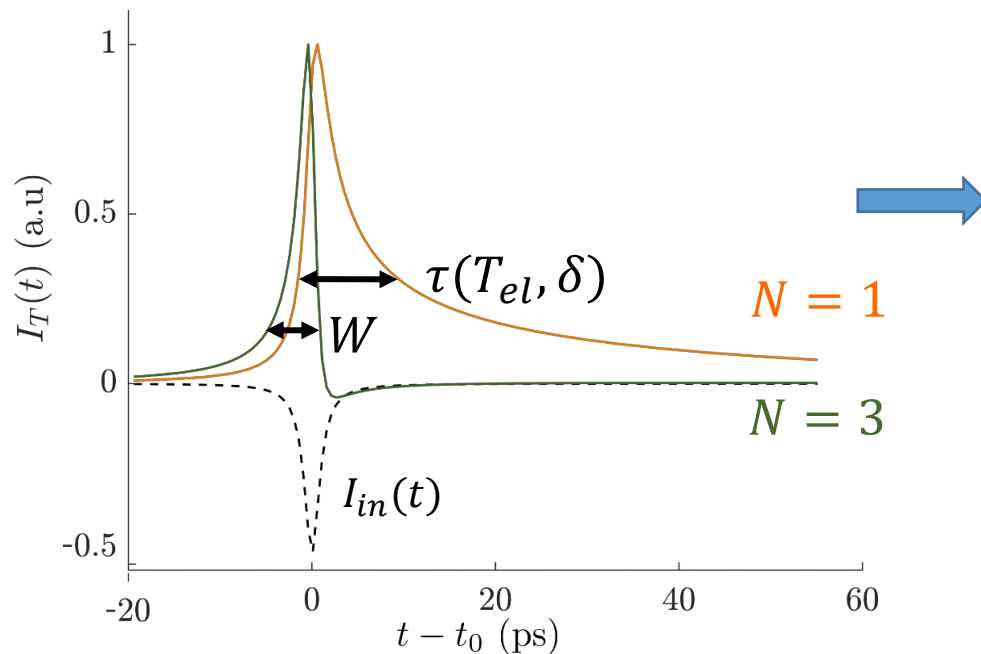
T. Jonckheere, et al., PRL **130**, 186203 (2023).

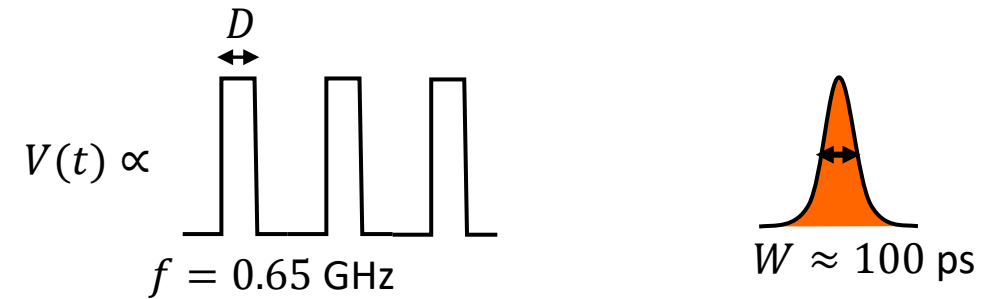
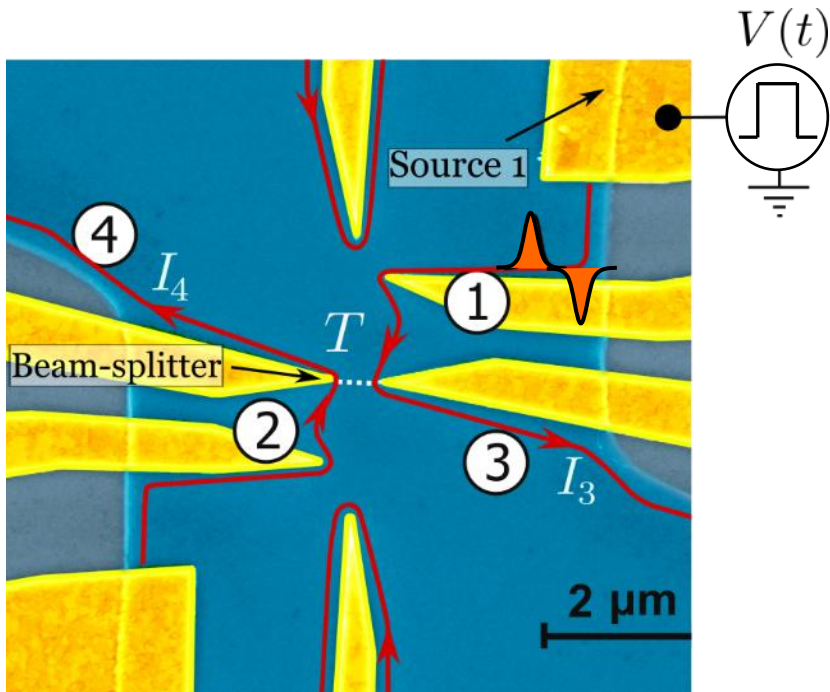


A. Latyshev et al., arXiv:2509.07875

S. Varada et al., PRB **111** (20), L201407 (2025)

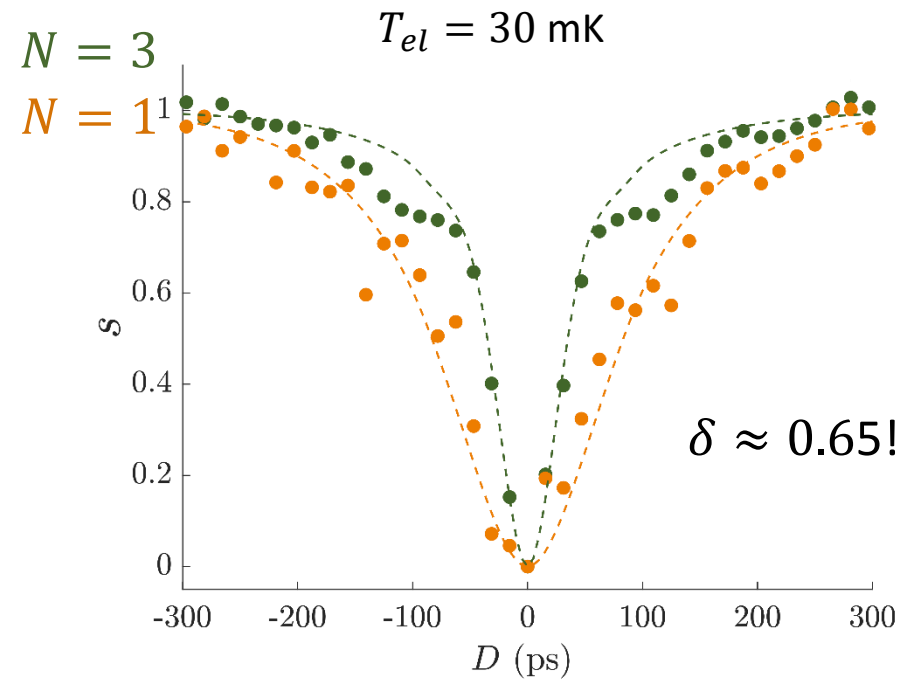
$N = 1$: single anyon emission at $\nu = 1/3$
 $N = 3$: single electron emission at $\nu = 1/3$



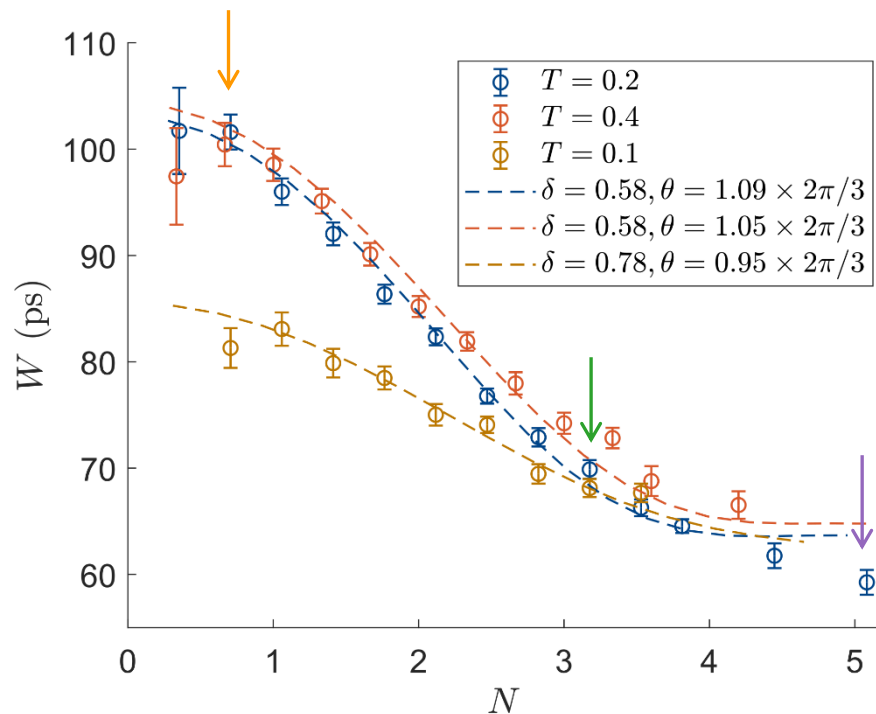
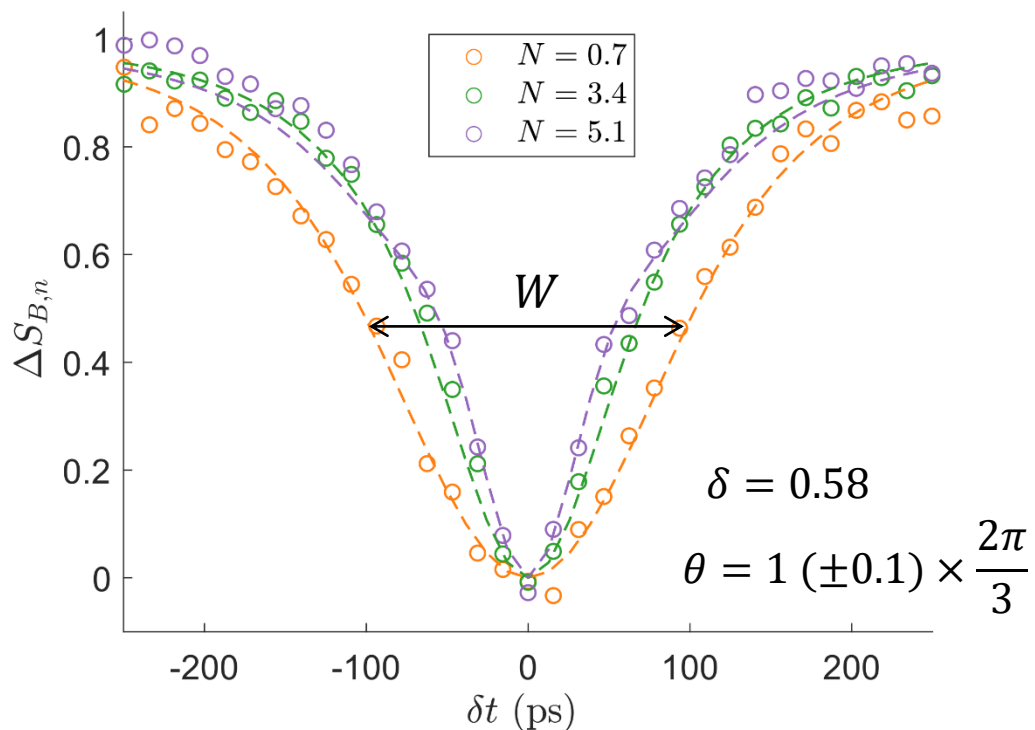


Current pulse carrying charge $Ne/3$
emulates N –anyon source

M. Ruelle et al., Science **389**, eadm7695 (2025).



$$T_{el} = 45mK, W = 65 ps$$

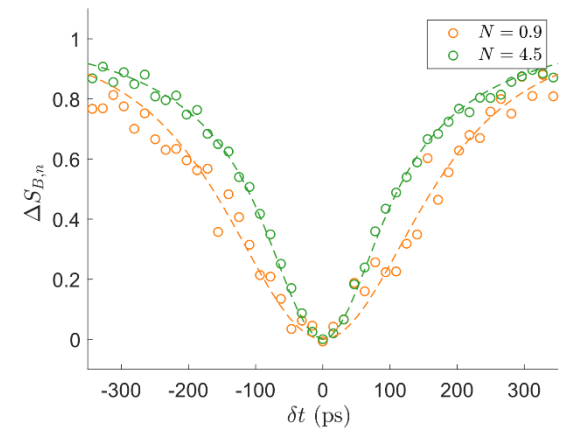
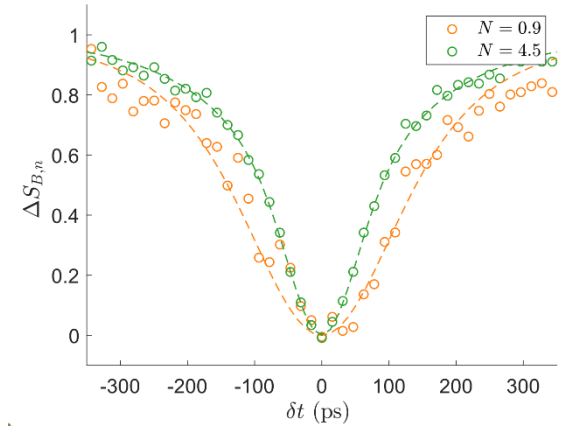
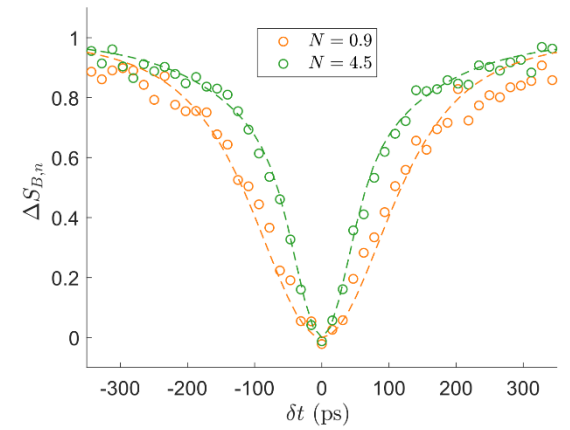
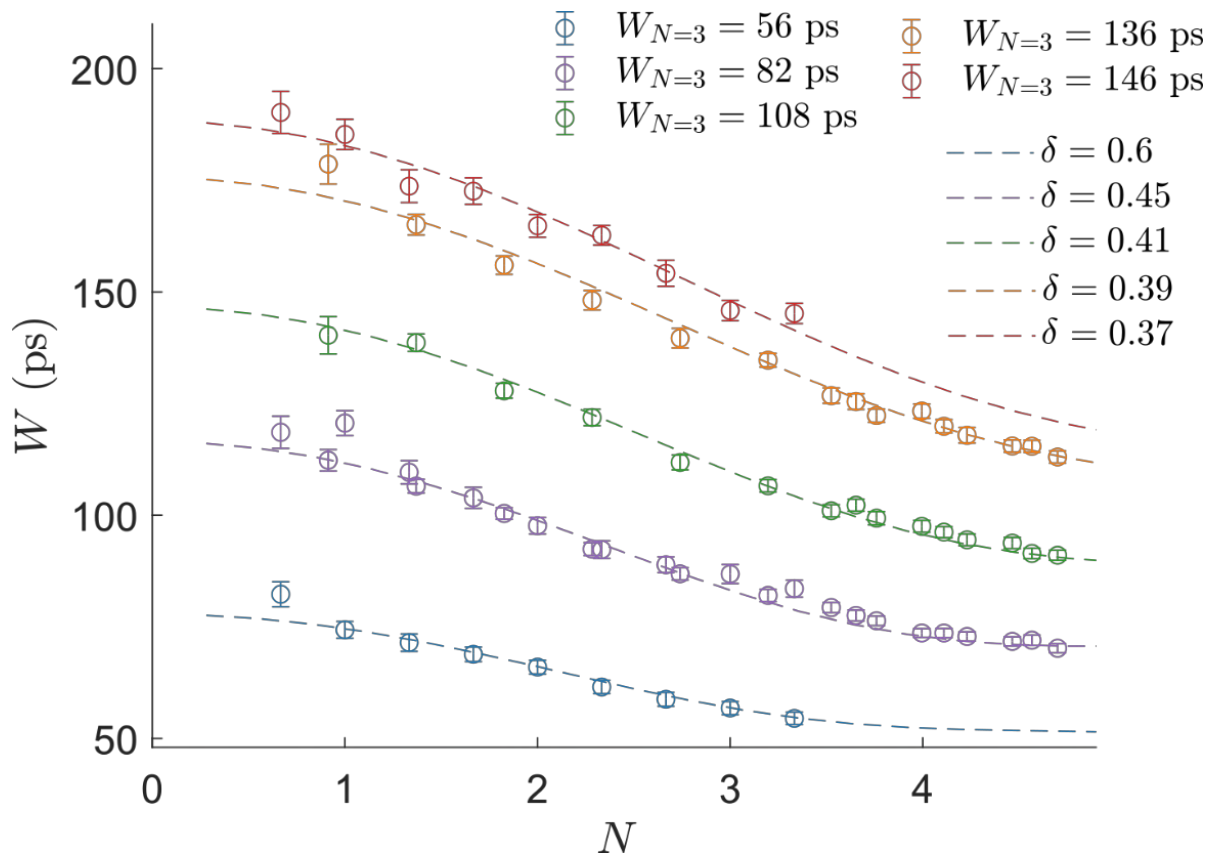


Braiding phase has the expected value $\theta = \frac{2\pi}{3}$ for $\nu = 1/3$

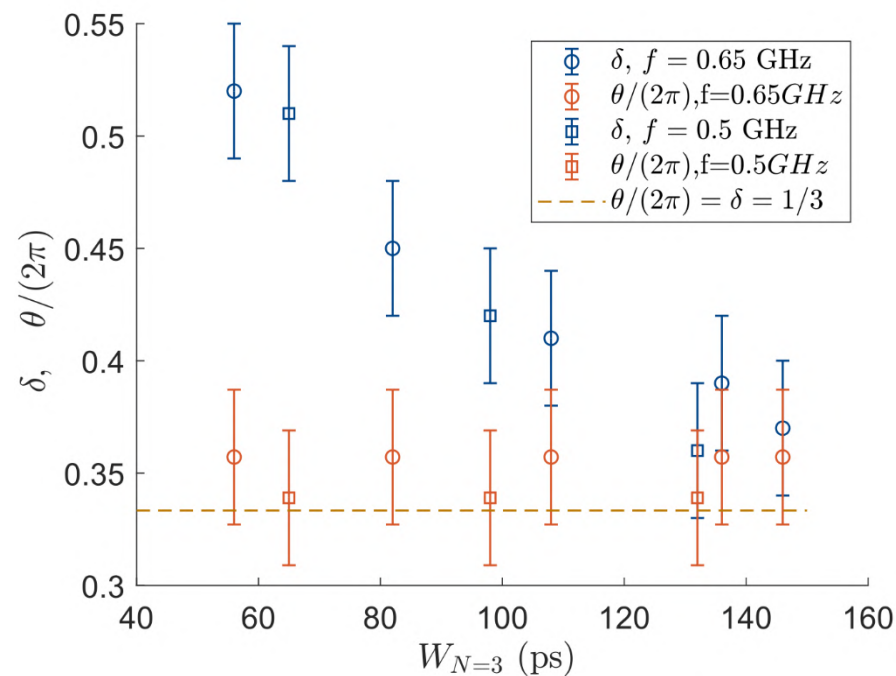
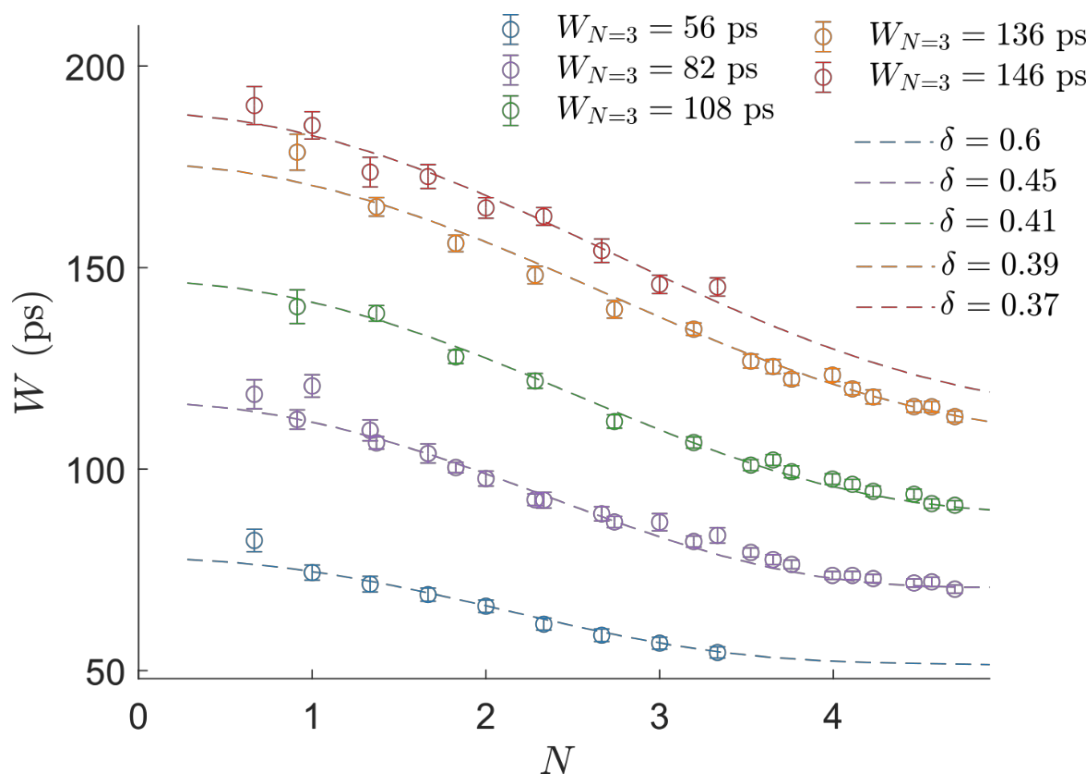
Not the scaling dimension! $\delta = 0.58$ instead of $\delta = 0.33$ expected

Varying the width of the pulse

$T_{el} = 60$ mK $T = 0.45$



$T_{el} = 60$ mK $T = 0.45$

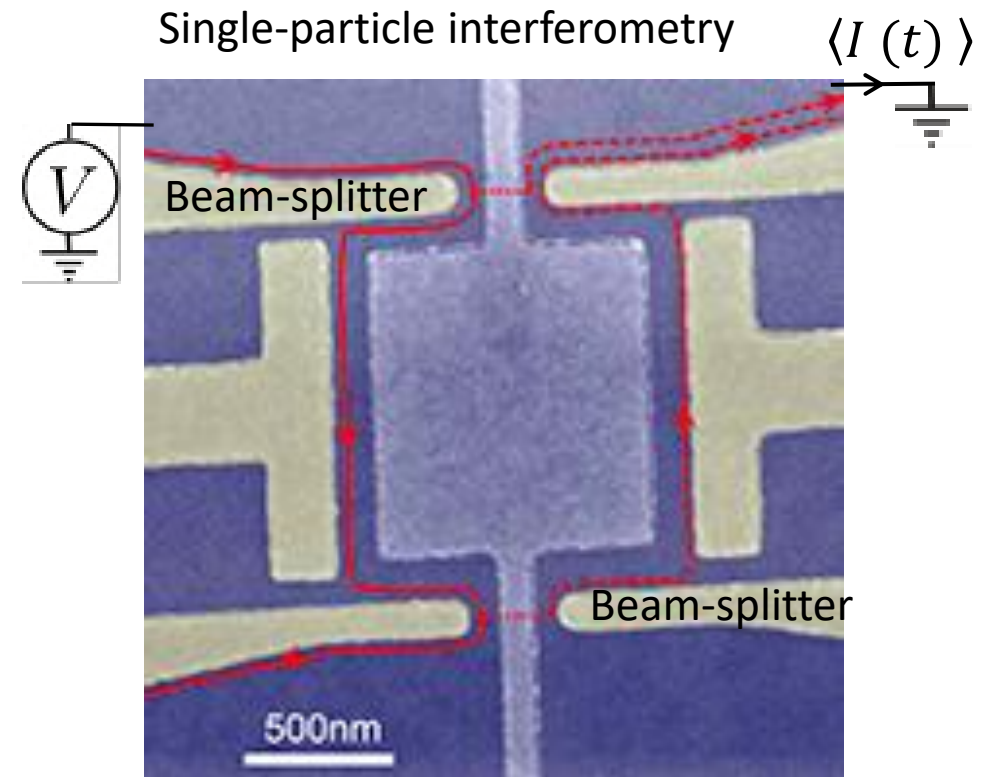
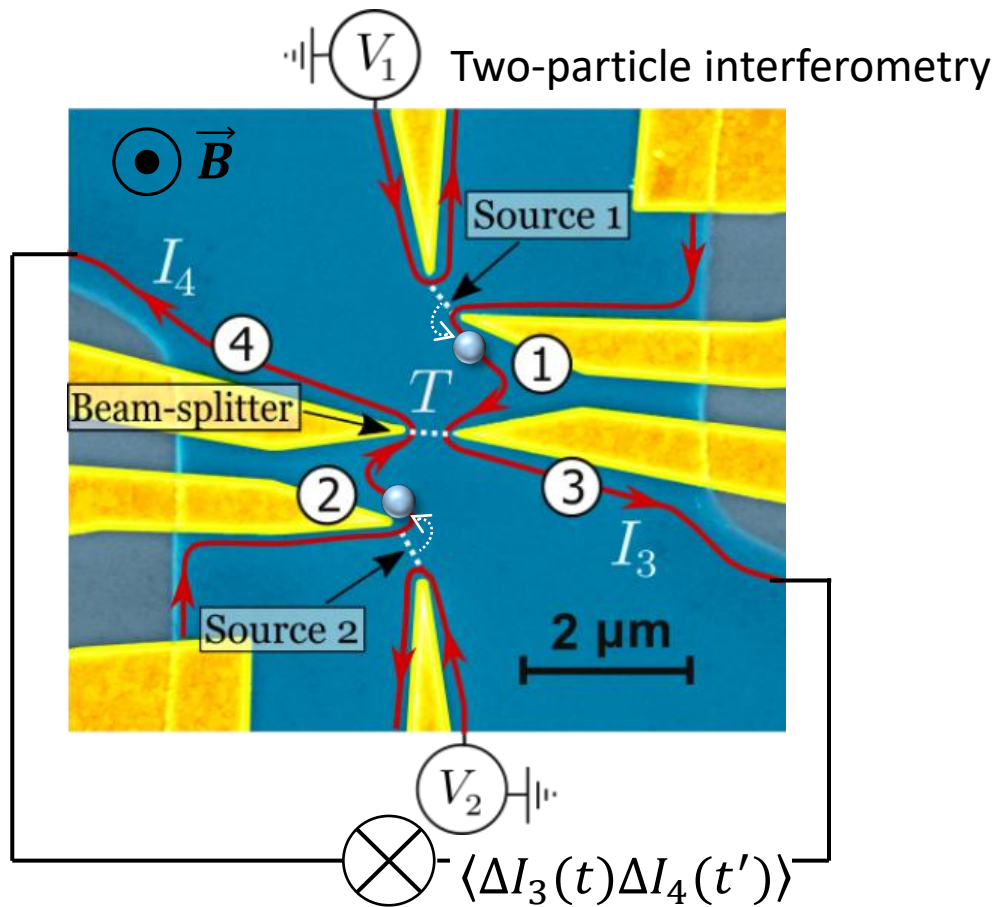


Deviations appear on short timescales

$\delta \approx 0.37$ recovered for larger pulses

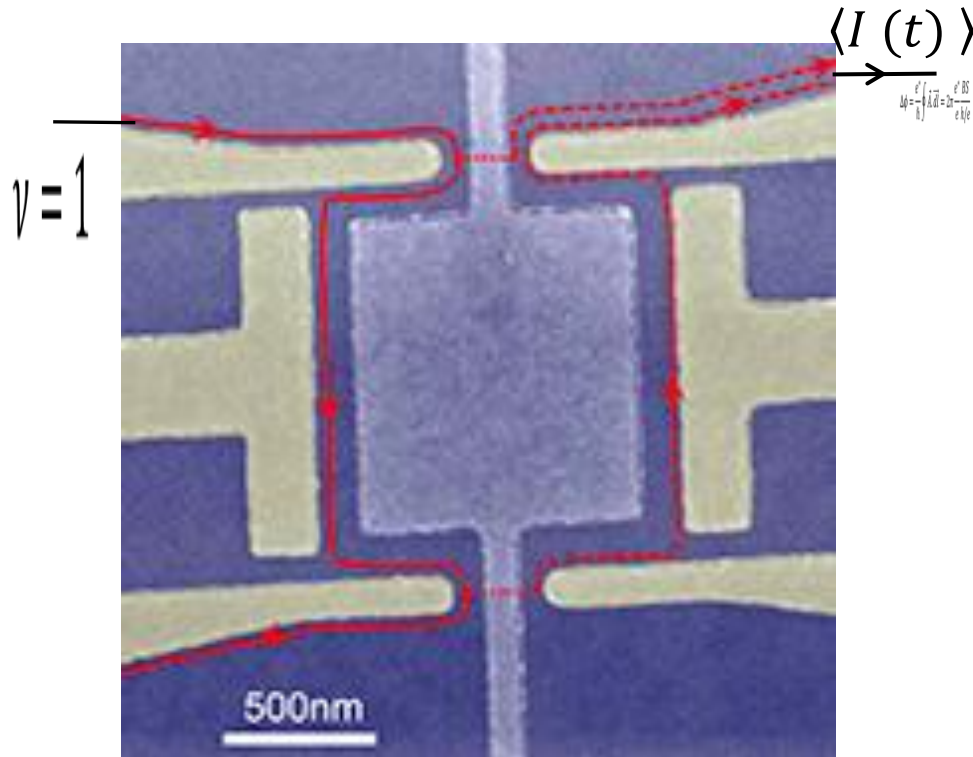
Single and two particle interferometers in quantum Hall conductors

Electron optics experiments in quantum Hall conductors



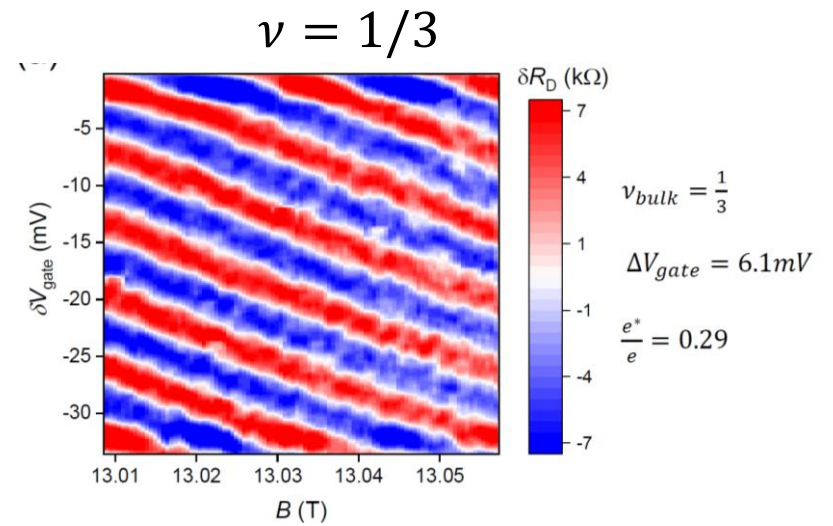
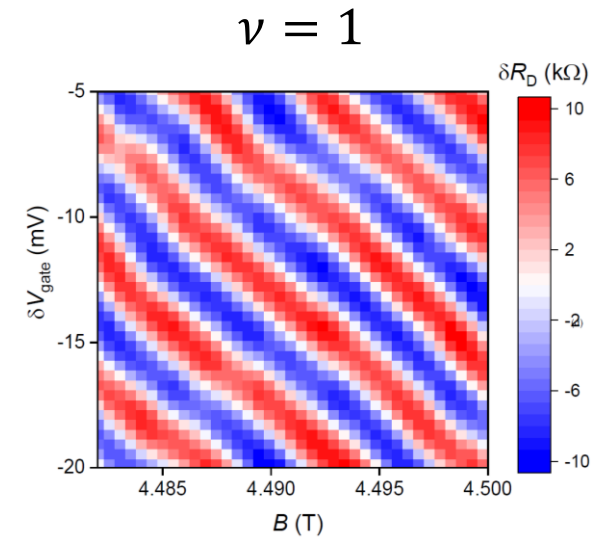
J. Nakamura, et al., Nature Physics **16** 931 (2020).

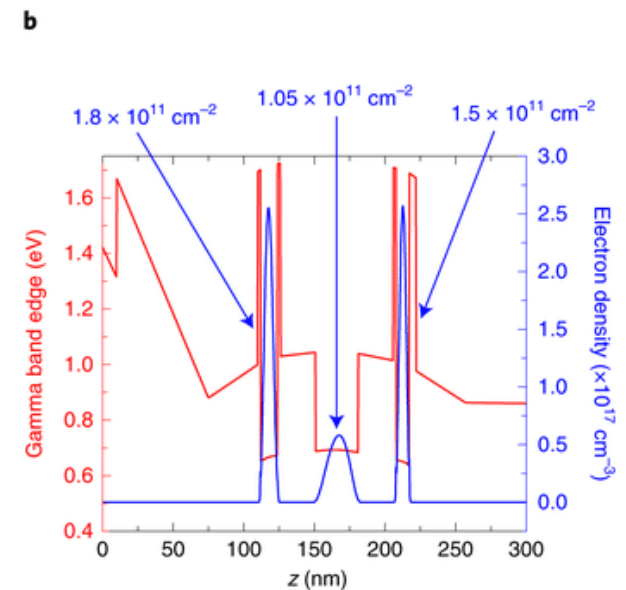
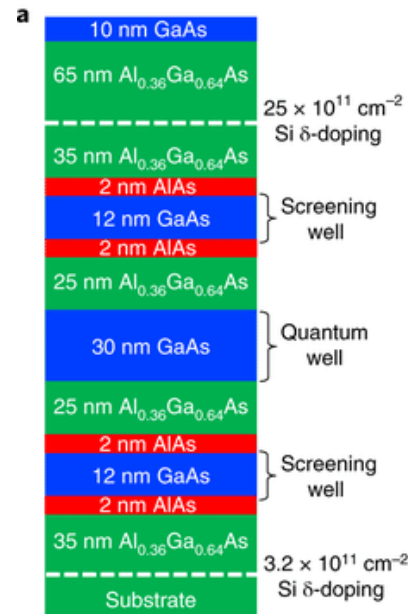
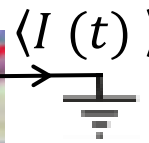
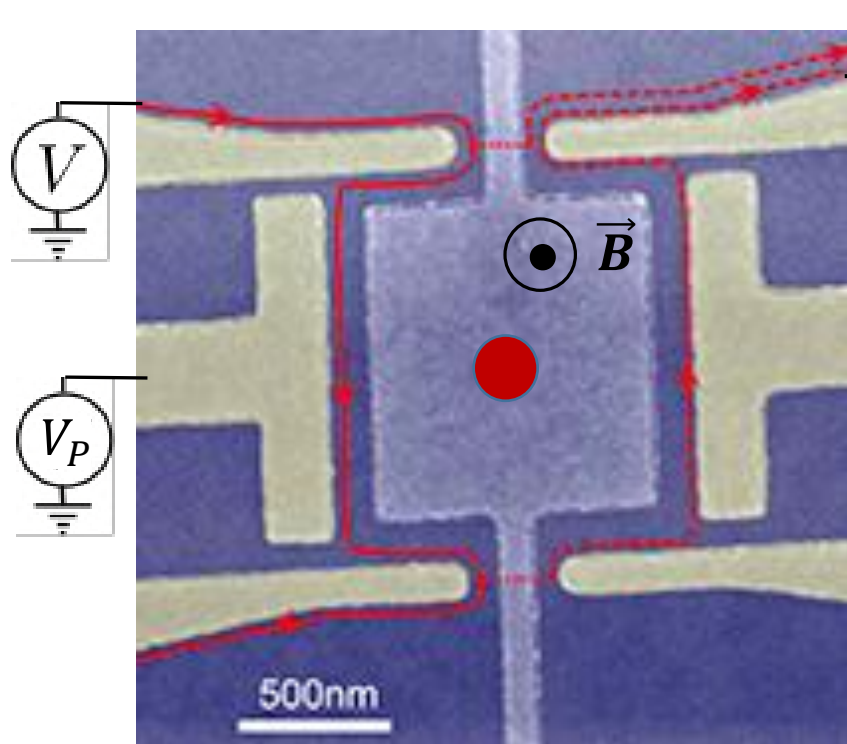
H. Bartolomei, M. Kumar et al., Science **368** 173 (2020)



Interferometry of anyons of charge e^*

$$\Delta\phi = \frac{e^*}{\hbar} \oint \vec{A} \cdot d\vec{l} = 2\pi \frac{e^* BS}{e h/e}$$

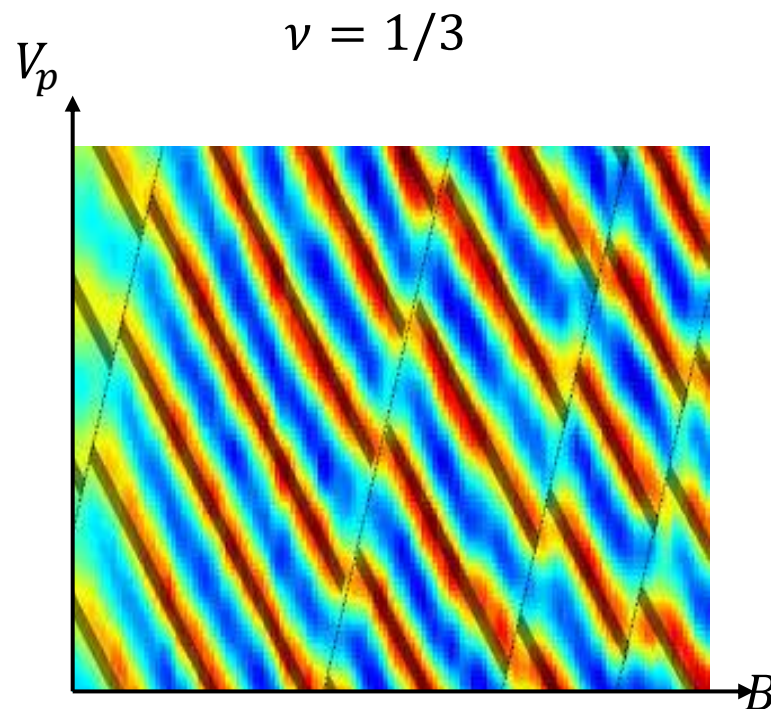
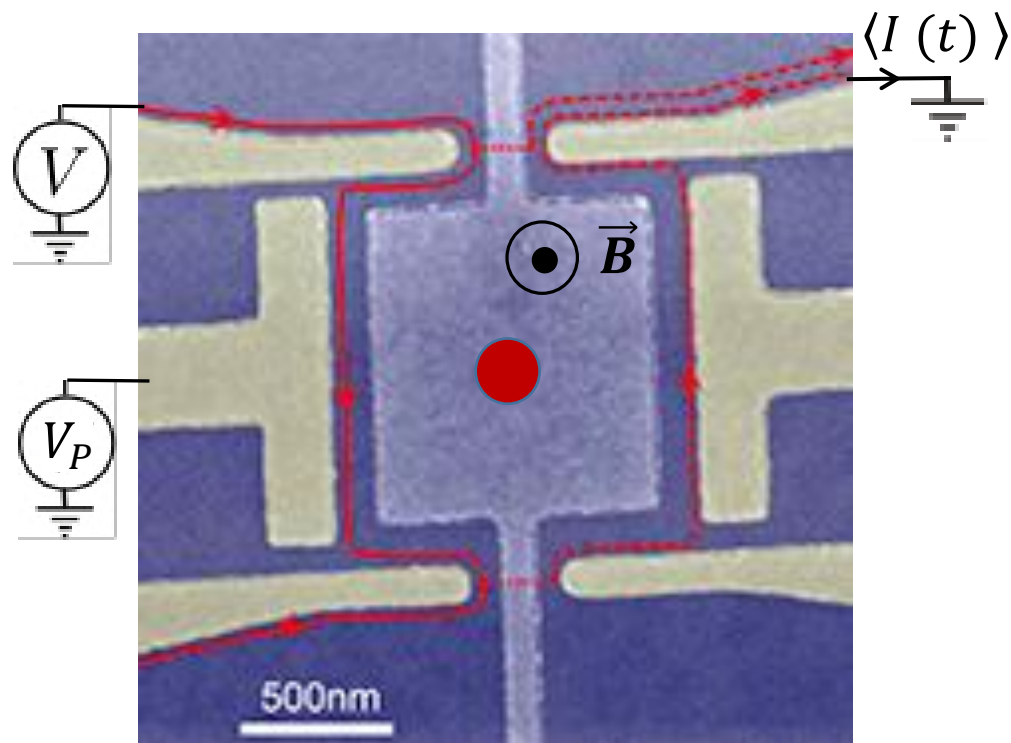




J. Nakamura, S. Liang, G.C. Gardner, M.J. Manfra, Nature Physics **16** 931 (2020).

$$\Delta\phi = 2\pi \frac{e^* BS}{e h/e} + \theta \Delta N$$

Pb: real situation much more complex due to Coulomb interaction



$$\Delta\phi = 2\pi \frac{e^*}{e} \frac{BS}{h/e} + \theta \Delta N$$

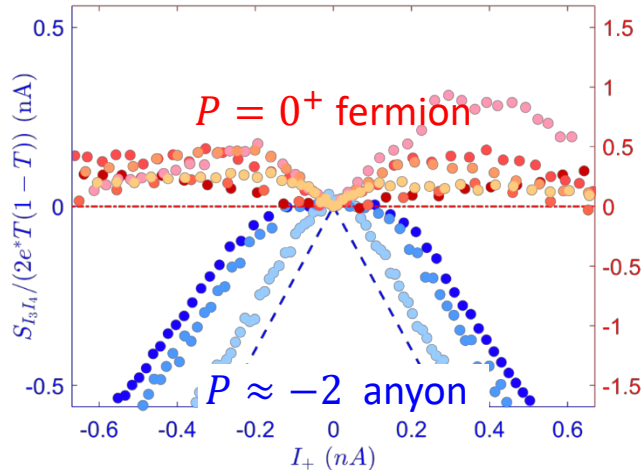
Pb: real situation much more complex due to Coulomb interaction

J. Nakamura, S. Liang, G.C. Gardner, M.J. Manfra, Nature Physics **16** 931 (2020).

$$\theta = 2\pi/3$$

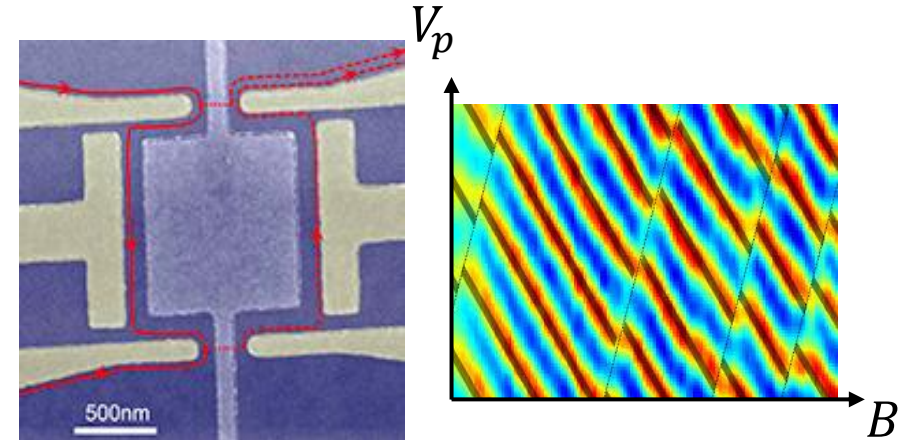
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- Two-particle interferometry



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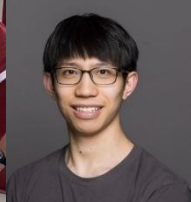
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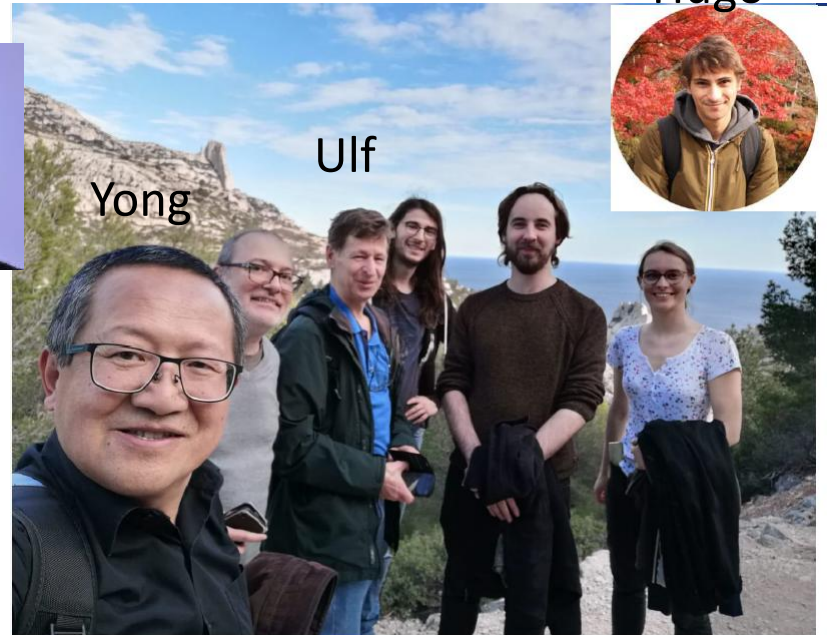


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